Measurement of Dijet Azimuthal Decorrelations in $pp$ Collisions at $\sqrt{s} = 7$ TeV

The ATLAS Collaboration

Azimuthal decorrelations between the two central jets with the largest transverse momenta are sensitive to the dynamics of events with multiple jets. We present a measurement of the normalized differential cross section based on the full dataset ($\int Ldt = 36$ pb$^{-1}$) acquired by the ATLAS detector during the 2010 $\sqrt{s} = 7$ TeV proton-proton run of the LHC. The measured distributions include jets with transverse momenta up to 1.3 TeV, probing perturbative QCD in a high energy regime.

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The production of events containing high transverse-momentum ($p_T$) jets is a key signature of quantum chromodynamic (QCD) interactions between partons in $pp$ collisions at large center-of-mass energies ($\sqrt{s}$). The Large Hadron Collider (LHC) opens a window into the dynamics of interactions with high-$p_T$ jets in a new energy regime of $\sqrt{s} = 7$ TeV. QCD predicts the decorrelation in the azimuthal angle between the two most energetic jets, $\Delta\phi$, as a function of the number of partons produced. Events with only two high-$p_T$ jets have small azimuthal decorrelations, $\Delta\phi \sim \pi$, while $\Delta\phi \ll \pi$ is evidence of events with several high-$p_T$ jets. QCD also describes the evolution of the shape of the $\Delta\phi$ distribution, which narrows with increasing leading jet $p_T$. Decorrelations in $\Delta\phi$ therefore test perturbative QCD (pQCD) calculations for multiple jet production without requiring the measurement of additional jets. Furthermore, a detailed understanding of events with large azimuthal decorrelations is important to searches for new physical phenomena with dijet signatures, such as supersymmetric extensions to the Standard Model [4].

In this Letter, we present a measurement of dijet azimuthal decorrelations with jet $p_T$ up to 1.3 TeV as measured by the ATLAS detector, beyond the reach of previous colliders. The normalized differential cross section $$(1/\sigma)\frac{d\sigma}{d\Delta\phi}$$ is based upon an integrated luminosity $\int Ldt = (36 \pm 4)$ pb$^{-1}$ [2]. The $\Delta\phi$ distribution is normalized by the inclusive dijet cross section, $\sigma$, integrated over the same phase space. This construction minimizes experimental and theoretical uncertainties. Previous measurements of $\Delta\phi$ from the D0 [3] and CMS [4] collaborations are extended here to higher jet $p_T$ values.

Jets are reconstructed using the anti-$k_t$ algorithm [5] (implemented with FASTJET [6]) with radius $R = 0.6$, and the jet four-momenta are constructed from a sum over its constituents, treating each as an $(E, \vec{p})$ four-vector with zero mass. The anti-$k_t$ algorithm is well-motivated since it is infrared-safe to all orders, produces geometrically well-defined cone-like jets, and is used for pQCD calculations (from partons), event generators (from stable particles), and the detector (from energy clusters [7]). The azimuthal decorrelation, $\Delta\phi$, is defined as the absolute value of the difference in azimuthal angle between the jet with the highest $p_T$ in each event, $p_T^{\text{max}}$, and the jet with the second-highest $p_T$ in the event. There are nine analysis regions in $p_T^{\text{max}}$, where the lowest region is bounded by $p_T^{\text{max}} > 110$ GeV and the highest region requires $p_T^{\text{max}} > 800$ GeV [7]. Only jets with $p_T > 100$ GeV and $|y| < 2.8$, where $y$ is the jet rapidity [8], are considered. The two leading jets that define $\Delta\phi$ are required to satisfy $|y| < 0.8$, restricting the measurement to a central $y$ region where the momentum fractions $(x)$ of the interacting partons are roughly equal and the experimental acceptance for multijet production is increased. In this region where $0.02 \lesssim x \lesssim 0.14$, the parton distribution function (PDF) uncertainties are typically $\pm 3\%$ (at fixed factorization scale) [9]. The cross sections, measured over the range $\pi/2 \leq \Delta\phi \leq \pi$ and normalized independently for each analysis region, are compared with expectations from a pQCD calculation [10] that is next-to-leading order (NLO) in three-parton production. The perturbative prediction for the cross section is $\mathcal{O}(\alpha_s^3)$, where $\alpha_s$ is the strong coupling constant.

The angular decorrelation is sensitive to multijet configurations such as those produced by event generators like SHERPA [11], which matches higher-order tree-level pQCD diagrams with a dipole parton-shower model [12]. Samples for $2 \to 2, 3$ jet production are combined using an improved CKKW matching scheme [13]. The progression of the parton shower is vetoed to avoid double counting of emissions. Event generators such as PYTHIA [14] and HERWIG [15] use $2 \to 2$ leading order pQCD matrix elements matched with phenomenological parton-cascade models to simulate higher-order QCD effects. Such models have been successful at reproducing other QCD processes measured by the ATLAS collaboration [7, 10].

The ATLAS detector [17, 18] consists of an inner tracking system surrounded by a thin superconducting solenoid providing a 2T magnetic field, electromagnetic and hadronic calorimeters, and a muon spectrometer based on large superconducting toroids. Jet measurements depend most heavily on the calorimeters. The electromagnetic calorimeter is a lead liquid-argon (LAr) detector with an accordion geometry. Hadron calorimetry is based on two different detector technologies, with scintillator tiles or LAr as the active medium, and with either steel, copper, or tungsten as the absorber material. The pseudorapidity ($\eta$) [8] and $\phi$ segmentations of...
the calorimeters are sufficiently fine to ensure that angular resolution uncertainties are negligible compared to other sources of systematic uncertainty.

A hardware-based calorimeter jet trigger identified events of interest; the decision was further refined in software [17, 18]. Events with at least one jet that satisfied a minimum transverse energy (E_T) requirement were recorded for further analysis. The events in each p_T^max range are selected by a single trigger with a given E_T threshold, and the lower end of the range is chosen above the jet p_T at which that trigger is ≈ 100% efficient. Three sets of triggered events with different integrated luminosity are considered: 2.3 pb^{-1} for 110 < p_T^max ≤ 160 GeV, 9.6 pb^{-1} for 160 < p_T^max ≤ 260 GeV, and 36 pb^{-1} for p_T^max > 260 GeV [2]. Events are also required to have a reconstructed primary vertex within 15 cm in z of the center of the detector; each vertex had ≥ 5 associated tracks. The inputs to the anti-kt jet algorithm are clusters of calorimeter cells seeded by cells with energy that is significantly above the measured noise [3]. Jets reconstructed in the detector, whether in data or the GEANT4-based simulation [19, 20], are corrected for the effects of hadronic shower response and detector-material distributions using a p_T- and η-dependent calibration [7] based on the detector simulation and validated with extensive test-beam [17] and collision data [21] studies. Jets likely to have arisen from detector noise or cosmic rays are rejected [22].

The resulting Δφ distribution is shown in Fig. 1 for jets with p_T > 100 GeV. There are 146788 events in the data sample, 85 of which have at least five jets with p_T > 100 GeV. Also shown is the PYTHIA sample with MRST 2007 LO* PDF [23] and ATLAS MC09 underlying event tune [24], processed through the full detector simulation and normalized to the number of events in the data sample. Two- and three-jet production primarily populates the region 2π/3 < Δφ < π while smaller values of Δφ require additional activity such as soft radiation or more jets in an event. Fig. 1 illustrates that the decorrelation increases when a third high-p_T jet is also required. Events with additional high-p_T jets widen the overall distribution.

The measured differential Δφ distributions in data are corrected in a single step with a bin-by-bin unfolding method [7] to compensate for trigger and detector inefficiencies and the effects of finite experimental resolutions. These correction factors, evaluated using the PYTHIA sample, lie within ±9% of unity. The leading sources of systematic uncertainty on the normalized cross sec-
tion are the jet energy scale calibration (2 – 17%) [7], the bin-by-bin unfolding method (1 – 19%), and the jet energy and position resolutions (0.5 – 5%). The ranges in parentheses represent the magnitude of the uncertainties near π and π/2, respectively, and correspond to the analysis region with the smallest statistical uncertainty (160 < $p_T^{\text{max}}$ ≤ 210 GeV). Uncertainties due to multiple $pp$ interactions in the same beam crossing (< 0.8% on the cross section for all analysis regions) are included in the evaluation of the jet energy scale uncertainties.

The normalized differential cross section is shown for each of the nine $p_T^{\text{max}}$ analysis regions as a function of $\Delta \phi$ in Fig. 2. As $p_T^{\text{max}}$ increases, and the probability for the emission of a hard third jet is reduced, the fraction of events near π becomes larger. Overlaid on the data are the results from a NLO pQCD [$\mathcal{O}(\alpha_s^3)$] calculation, NLOJET++ [10] with fastNLO [25] and using the MSTW 2008 PDF [3]. The factorization and renormalization scales are set to $p_T^{\text{max}}$ and are varied independently up and down by a factor of two to determine the scale uncertainties. The scale uncertainties are larger between $\pi/2 < \Delta \phi < 2\pi/3$ where the pQCD calculation is effectively leading order in four-parton production. The PDF uncertainties are treated as the envelope of the 68% CL uncertainties from MSTW 2008 [3], NNPDF 2.0 [24], and CTEQ 10 [27], and are combined with the uncertainties resulting from an $\alpha_s$ variation of ±0.004; the $\alpha_s$ contributions dominate. The calculation is corrected for non-perturbative effects due to hadronization and the underlying event [28, 29]; the correction is smaller than 3%. The fixed-order calculation fails near $\Delta \phi \to \pi$ where soft processes dominate and contributions from logarithmic terms are enhanced. Figure 3 displays the ratio of the cross section with respect to the NLO calculation. In most regions, the theory is consistent with the data. However, the prediction in the range 110 < $p_T^{\text{max}}$ < 160 GeV is relatively low in the central region of $\Delta \phi$ where the scale uncertainties are small.

The data are also compared with predictions from SHERPA, PYTHIA, and HERWIG in Fig. 3. The leading-logarithmic approximations used in these event generators’ parton-shower models effectively regularize the divergence at $\Delta \phi \to \pi$; all three provide a good description of the data in this region. In the region $\pi/2 < \Delta \phi < 5\pi/6$, where multijet contributions are significant, this observable distinguishes between the three generators. SHERPA, which explicitly includes higher-order tree-level diagrams, performs well in most $\Delta \phi$ and $p_T^{\text{max}}$ regions. Having phenomenological parameters that have been adjusted to previous ATLAS measurements, PYTHIA [28] and HERWIG [24] also describe the data.

In summary, we present a measurement of dijet azimuthal decorrelations in events produced in $pp$ collisions at $\sqrt{s} = 7$ TeV. The normalized differential cross sections...
are based on the full dataset (∫L dt = 36 pb−1) collected by the ATLAS collaboration during the 2010 run of the LHC. Expectations from NLO pQCD [O(αs 4)] and those of several event generators successfully describe the general characteristics of our measurements, including the increasing slope of the ∆φ distribution with pT max and the shape near ∆φ ∼ π/2 where events with multiple jets make a considerable contribution. Our data, which include jets with pT values that significantly exceed earlier measurements, explore QCD in a new kinematic region.

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[8] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point in the center of the detector. Cylindrical coordinates (r, φ) are used in the transverse plane, φ being the azimuthal angle around the beam axis. The pseudorapidity is defined in terms of the polar angle θ as η = −ln tan(θ/2). Rapidity is defined as y = ±1/2 ln[(E + p_z)/(E − p_z)], where E is the energy and p_z is the longitudinal component of the momentum along the beam direction.
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