EVIDENCE FOR SOFT MODES AT $\mathbf{q} \neq 0$ IN THE FAR-INFRARED SPECTRUM OF $\text{Rb}_2\text{ZnBr}_6$

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The possibility to observe optically inactive and $\mathbf{q} \neq 0$ soft modes in an indirect way by using far-infrared techniques is investigated. Using the fact that, due to anharmonicity, all modes of lattice vibrations are coupled with each other, the softening of one particular mode can be seen indirectly in the damping of other modes. By applying this idea to the modulated structure $\text{Rb}_2\text{ZnBr}_6$, the infrared inactive soft mode at the normal-incommensurable transition, the soft phason mode at the lock-in transition as well as a soft mode behaviour at an intermediate temperature have been observed.

Incommensurable structures have attracted a great deal of interest during the last years. In these systems, below a transition temperature $T_c$, the atomic positions are shifted periodically, with a characteristic wavevector which is incommensurable with the underlying lattice. In most cases, this modulation becomes commensurable at a temperature $T < T_c$. New types of excitations appear in this incommensurable phase, originating from spatial variations of the amplitude or the phase of the modulation wave, called the amplitude- and phason-modes. The transition to this phase is very often associated with the softening of a lattice mode with a wavevector $\mathbf{q} \neq 0$ situated within the Brillouin-zone. Obviously, this mode is optically inactive, and therefore the softening of this lattice mode will not be seen directly by simple far-infrared techniques. It is expected from theoretical considerations, that also one of phason modes should soften at the lock-in transition temperature $T_c$. Although this mode is in principle optical active, it has not been observed directly due to its small oscillator strength.

In this paper we will show that it is possible to study these soft modes at both phase transitions in a rather indirect way. The method is based on the fact that all lattice modes are coupled with each other due to the anharmonicity of the interaction potential. Not the soft mode itself, but the effect of its softening on the other optical active modes is investigated.

The electrodynamic response of a system of damped harmonic oscillators of frequency $\omega_j$ and damping constant $\gamma_j$ is characterized by the frequency dependent dielectric function:

$$
\varepsilon(\omega) = \varepsilon(\infty) + \sum_j f_j (\omega_j^2 - \omega^2 + i\omega\gamma_j),
$$

where $f_j$ is the oscillator strength of the mode $j$. The coupling of the modes is reflected in the damping constants $\gamma_j$, which are the only entries in Eq. (1) containing information of properties at non-zero wavevectors. Using the usual perturbation theory, these damping constants are given in first order by:

$$
\gamma_j = \frac{\hbar}{16\pi k_B T} \sum_{k,l} \left\{ n_{k} n_{l} + (n_{k} + n_{l} + 1) \{ 2\delta(\omega_j - \omega_k - \omega_l) - \delta(\omega_k + \omega_l) - \delta(\omega_j - \omega_k) \} + \frac{1}{2} \right\},
$$

where the summation includes all pairs of nodes which fulfill momentum conservation. $N$ is the number of unit cells, the temperature dependent occupation number is given by $n_i = [\text{exp}(\hbar\omega_i/kT) - 1]$ and $\gamma_{ij}$ are the coefficients of the third order term of the anharmonic potential.

If at a $q \neq 0$, a mode $\omega_Q(q)$ softens, then one gets for the temperature dependence close to the phase transition

$$
\omega_Q = \sqrt{T_c - T},
$$

where $\beta$ usually equals to $\frac{1}{3}$. Here $T_c$ is the transition temperature. Near $T_c$, we have $\hbar\omega_i << kT$ and $n_i >> 1$, for all $i \neq s$. Therefore, the leading term in Eq. (2) will be $\propto n_i \gamma_i$, which gives in first approximation:

$$
\gamma_1(T) = T / [T_c - T].
$$

The softening of the mode at $q$, thus results in an increase of the damping constants $\gamma_i$ of all other modes $i \neq s$ near the phase transition. This can indirectly be observed by measuring $\gamma_1$ as a function of temperature.

As an example for a normal structural phase transition accompanied by the softening of an infrared inactive mode, in Fig. 1a a part of the far-infrared reflection spectrum of $\text{CH}_3\text{NH}_2\text{CuCl}_4$ is presented. For the experimentally clearly distinguishable resonances in-
Fig. 1a: Reflectivity R of (CH₃NH₃)₂CuCl₄, relative to the reflectivity R₀ of a gold sample, as a function of frequency for several temperatures.

1b: Inverse of the normalized damping constant γ(0)/γ(T) as a function of the inverse temperature T⁻¹ for the peaks indicated in Fig. 1a.

Dicated by numbers in the figure, the width of the line was taken as a measure for the damping. The damping constants obtained in this way are plotted as a function of the inverse temperature in Fig. 1b. As can be seen from the figure, close to Tₜ, they all show a linear dependence with the inverse temperature, in accordance with Eq. (4).

For systems with high transition temperatures, the sharp lines will coalesce into broad bands, making a direct measurement of γ from the line width impossible. However, it is nevertheless possible to detect a change in γ by carefully measuring the temperature dependence of the transmission of monochromatic far-infrared radiation in this broad band. The relative transmission t* of electromagnetic radiation through a dielectric slab of thickness d may be approximated, for the case of normal incidence, neglecting multiple reflections and assuming k << n, by

\[ t^* = (1-R)^2 t, \]

where t = \exp(-4πkdωk), R = (n-1)^2/(n+1)^2. Here n and k are the real and imaginary parts of the complex index of refraction n = n + ik = n ⅈk. As can be seen from Eq. (5), the dominating term for the transmission is given by k in the exponential leading to t. Therefore we can take n as a temperature independent constant and study the effect of the temperature dependence of γ(T) on k only. For frequencies \( ω \ll ω_j \) (i.e. \( ω = ω_j \) and \( ω \ll ω_j \) for all other \( j ≠ 1 \)), k(ω) can be written as

\[ k(ω) = k_1(ω) + k_0(ω) \]

with

\[ k_1(ω) = \frac{1}{2n} \frac{∫_0^1{ωγ_j}^2}{ω^2 + ω_1^2} \]

and

\[ k_0 = \frac{1}{2n} ∑_{j ≠ 1} ∫_0^1{ωγ_j} \frac{ω_j^2}{ω_j^2} \]

According to Eq. (4), \( γ_j \) will increase strongly near the phase transition and the denominator in Eq. (7) will be dominated by \( γ_j^2ω_j^2 \), leading to

\[ k(T) = k_1(T) + k_0 \]

Therefore, k(T) will show a relative minimum at the transition temperature Tₜ, if k₀ is only weakly temperature dependent in this region. As a result, according to

\[ t^* = t_0(1-4πdωk_1)^2 = t_0(1-4πdω_1^2T/T_0)^2 \]

the transmission will show a maximum at Tₜ. The k₁ term has been included in t₂. So by carefully measuring the transmission as a function of temperature at a fixed frequency, possible optical inactive soft modes can be observed indirectly.
We have applied this method to Rb₂ZnBr₄, known to be a modulated structure with \( T_1 = 353 \) K and \( T_c = 201 \) K. As the effect is expected to be rather small, the experimental conditions for measuring the transmission have to be controlled rather accurately. Therefore, a very stable source of monochromatic far-infrared radiation is needed. We used a harmonic generator of microwaves, which is continuously tunable over the range 2 to 25 cm⁻¹. As a detector, a He cooled silicon bolometer was used.

In Fig. 2 the transmission at 7.8 cm⁻¹ of Rb₂ZnBr₄ relative to the transmission of the system without sample is plotted against inverse temperature. The figure shows a pronounced maximum at 201 K and a very small one at 353 K at the two known transition temperatures \( T_c = 201 \) K and \( T_1 = 353 \) K. In addition, at 300 K another well-defined maximum is found, indicating an extra phase transition within the incommensurable phase. In Fig. 3 more detailed measurements for temperatures close to \( T_1 \) are presented, using different frequencies, and all showing the linear \( T^{-1} \) dependence of the transmission around \( T_1 \), in accordance with Eq. (10).

These results indicate that one can observe the softening of the modes at both \( T_c \) and \( T_1 \) indirectly, which seemed to be impossible in a direct way. As the damping has already increased considerably with temperature around \( T_1 \), the soft mode contribution is probably relatively smaller here, leading to a smaller effect in the transmission at \( T_1 \). As the amplitude is expected to decrease only slightly in frequency at the temperature \( T_1 \), the \( q \) phason mode should soften at \( T_1 \). Therefore, the observed effect at \( T_1 \) can tentatively be ascribed to this soft \( q_0 \) phason mode. The maximum at 300 K indicates a possible other phase transition within the incommensurable phase, a fact that is also suggested recently by other experiments. This could be the formation of a domain-like structure, consisting of commensurate microdomains, before the transition to the real commensurate phase at 200 K takes place.

In conclusion we have shown the possibility of indirect investigation of \( q \neq 0 \), infrared inactive soft modes through their contribution to the damping of the other optically active lattice modes. The experiment involves the measuring with very high accuracy of the temperature dependence of the transmission, using very stable monochromatic far-infrared radiation. In this way the optically inactive soft mode at the normal-incommensurable transition as well as the soft mode at the lock-in transition at \( T_c \) could be seen in Rb₂ZnBr₄. At \( T_1 \), the observed effect can be ascribed to a soft phason. This gives new evidence for the existence of these excitations, characteristic for
incommensurable structures but rather hard to observe. At 300 K an intermediate transition within the incommensurable phase has been seen, which is probably a transition to a domain-like structure.

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References