ABSTRACT: Coinage metal clusters are of great importance for a wide range of scientific fields, ranging from microscopy to catalysis. Despite their clear fundamental and technological importance, the experimental structural determination of copper clusters has attracted little attention. We fill this gap by elucidating the structure of cationic copper clusters through infrared (IR) photodissociation spectroscopy of $\text{Cu}_n^{+}$–$\text{At}_m$ complexes. Structures of $\text{Cu}_n^{+}$ ($n = 3–10$) are unambiguously assigned based on the comparison of experimental IR spectra in the $70–280$ cm$^{-1}$ spectral range with spectra calculated using density functional theory. Whereas $\text{Cu}_5^{+}$ and $\text{Cu}_6^{+}$ are planar, starting from $n = 5$, $\text{Cu}_n^{+}$ clusters adopt 3D structures. Each successive cluster size is composed of its predecessor with a single atom adsorbed onto the face, giving evidence of a stepwise growth.

Transition-metal clusters are of great interest and considerable importance to the fields of heterogeneous catalysis, solid-state physics, surface chemistry, and organometallic chemistry due to their frequently strongly size-dependent properties of magnitudes that are often not present in small molecular or bulk systems. Part of the fascination of clusters stems from their often enhanced reactivity with respect to the bulk, typically attributed to the larger number of undercoordinated atoms, involving free d electrons to actively participate in the bonding. The coinage metals’ (Cu, Ag, Au) special place in this stems from their closed d shell, forcing s electrons to become involved in bonding and often leading to much more gentle activation of feedstock molecules.$^{1,2}$

Of the coinage metal clusters, gold has, not unreasonably, attracted the most attention, especially after the groundbreaking experiments by Haruta and coworkers demonstrating catalytic CO oxidation at low temperatures.$^{3}$ Whereas gold certainly has laid claim to being a “special” element, the chemical importance of other coinage metal clusters has to a certain point been neglected by this proverbial gold rush. The catalytic activity of, for instance, the industrial CO$_2$ hydrogenation catalyst has been largely attributed to copper nanoparticles on an Al$_2$O$_3$ surface with ZnO as a cocatalyst.$^{4,5}$ It has been shown that the CO$_2$ hydrogenation is structure-sensitive, where smaller Cu nanoparticles exhibit larger turnover frequencies.$^7$ Yet the mechanism for hydrogenation remains elusive, even for the very first step. Using clusters as well-defined model systems, proposed reaction pathways can be tested with high confidence, allowing much needed insight. Beside this role as a model system, clusters may also act as a catalyst under technical conditions themselves. Recently, it was shown that deposited Cu$_4^{+}$ can lower the activation barrier for the methanol formation from CO$_2$,$^7$ thereby illustrating the potential of nanotailored catalysts.

For both roles, however, knowledge of the clusters themselves and, in particular, of their structures, is imperative. In the model system scenario, it is of great importance to know the cluster morphology prior to exposing it to reactants, whereas for studies on deposited clusters, size-selected prior to soft landing, the influence of the substrate on the cluster geometry must be taken into account.$^{10}$ Despite the fact that Cu clusters were among the first clusters produced,$^9$ the available structural information is limited to photoelectron spectroscopy for anions,$^{10–12}$ mass spectrometric studies using H$_2$O as the molecular probe to reveal the number of adsorption sites,$^{14}$ photodissociation spectra in the visible range,$^{15,16}$ and an ion-mobility mass spectrometry (IMS) study$^{17}$ for cations.

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Infrared multiple photon dissociation (IRMPD) spectroscopy has a solid track record in determining the molecular structure of metal clusters and can do so mass-selectively using weakly bound messenger atoms or molecules.\textsuperscript{18,19} We recorded IRMPD spectra of Cu\textsuperscript{+}−Ar\textsubscript{m} \((n = 3−10, m = 1−4)\) in a molecular beam environment using IR light produced by the free-electron laser (FEL) FELIX in the 70−280 cm\textsuperscript{-1} spectral range\textsuperscript{20} and mass-selective detection. The experimental IR spectra for four of the cluster sizes studied are shown in Figure 1, whereas a complete overview of spectra for Cu\textsuperscript{+}−Ar\textsubscript{m} can be found in the Supporting Information (SI). The experimental IR spectra are complemented by spectra calculated for trial structures using density functional theory (DFT) calculations at the PBE-D3/TZVP level. Isomers for pure copper clusters as well as clusters complexed with argon are considered to assess the role of the messenger.

Complexation of copper clusters with Ar atoms is strongly influenced by temperature and Ar concentration; the mass distribution with which IR spectra were recorded forms a compromise between low Ar coverage and the mass range over which IR spectra could be recorded. To reduce contamination due to ingrowth from Cu\textsuperscript{+}−Ar\textsubscript{m+1} complexes, for each cluster size \(n\), we display the spectrum for the complex Cu\textsuperscript{+}−Ar\textsubscript{m} with the highest intensity in the mass spectrum. As a result, spectra for the lower cluster masses have been recorded at relatively high Ar coverage, illustrated by Cu\textsubscript{4}+−Ar\textsubscript{4}. Its IRMPD spectrum exhibits four well-resolved bands \((83, 103, 132, \text{ and } 207 \text{ cm}^{-1})\) with line widths (full width at half-maximum (fwhm)) of \(\sim 11 \text{ cm}^{-1}\). To assign this spectrum, it is compared with the calculated spectra of three isomers of Cu\textsubscript{4}+; comparisons with further isomers can be found in the SI. Only doublet spin isomers are considered here (just as for further even-sized clusters, where singlet states are considered for odd-sized clusters) because they were calculated to be lower in energy than other spin state isomers, in agreement with what has been reported at the LCGTO-GGA theoretical level.\textsuperscript{21} The assignment of the IRMPD spectrum based on the calculated spectra of bare Cu\textsubscript{4}+ is not possible, as it is evident

\begin{figure}[h]
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\includegraphics[width=\textwidth]{figure1.png}
\caption{Comparison of IRMPD spectra of Cu\textsubscript{4}+−Ar\textsubscript{4}, Cu\textsubscript{5}+−Ar\textsubscript{3}, Cu\textsubscript{7}+−Ar\textsubscript{3}, and Cu\textsubscript{10}+−Ar (black dots; the blue line represents a five-point adjacent average) with calculated vibrational spectra of three low-energy isomers (black lines) and their complexes with Ar (red shading). The corresponding structures and relative energies of the bare clusters are reported above each spectrum.}
\end{figure}
that each of the structures proposed does not have sufficient bands to explain the experimental spectrum. Sure enough, the 132 and 207 cm\(^{-1}\) bands match the calculated spectrum for the lowest energy \(D_{3h}\) rhombic structure 4A very well, but this leaves the lower frequency bands unaccounted for. The inclusion of the Ar messenger atoms in the DFT calculations quite drastically changes the predicted spectrum. Now all four observed bands are predicted, agreeing in both frequency and relative intensity. The two lowest-frequency bands (observed at 83 and 103 cm\(^{-1}\)) that were not predicted by the calculations for the bare Cu\(_4^+\) originate from motions of the Ar relative to cluster. Also, slight shifts for the higher frequency bands are calculated, from 130 and 212 cm\(^{-1}\) for the bare clusters to 138 and 215 cm\(^{-1}\) for the complex. These bands are associated with elongations of the cluster along the Cu–Cu axes, which are stericly hindered by the presence of the Ar. However, they do not disturb the original structure dramatically, as illustrated by the relatively low calculated binding energies of 0.20 and 0.27 eV for binding to the copper atoms positioned along the long and short axes, respectively. It is of further interest that the presence of Ar atoms increases the IR intensity of the bands to over ten times, suggesting that the Ar acts as an “antenna”.

The IRMPD spectrum recorded for Cu\(_4^+\)−Ar\(_3\) is simpler than that for Cu\(_5^+\)−Ar\(_4\), exhibiting just two well-resolved bands and suggesting a higher symmetry structure. In the literature, the structure of Cu\(_5^+\) has been the subject of considerable debate, with both planar and 3D structures proposed to be most stable.\(^{1,17,22,23}\) Here it is compared with the calculated spectra of three isomers of Cu\(_5^+\) within 0.1 eV relative to each other. The comparison for the bare structures suggests the assignment to a \(D_{3h}\) trigonal bipyramid isomer 5A, with a reasonable agreement for the bands observed at 188 cm\(^{-1}\) (177 cm\(^{-1}\) calculated) and at 121 cm\(^{-1}\) (93 cm\(^{-1}\) calculated). Calculated modes correspond to in-plane and out-of-plane distortions of the triangular base plane. Ar binding to atoms of the triangular base is energetically the most favorable, where the steric hindrance caused by the Ar shifts the calculated bands to 189 and 121 cm\(^{-1}\), respectively, resulting in an almost perfect match with the experimental bands. Ar complexation of the planar structure 5C is not energetically favorable (0.06 eV higher than the lowest energy complex of 5A), nor does it provide a better agreement with the observed spectrum. IMS measurements by Weis suggest a planar structure 5C is not energetically favorable (0.06 eV higher than that for Cu4\(^+\)) and suggesting a higher symmetry structure. In the literature, measurements by Weis suggest a planar structure 5C for Cu5\(^+\). Because of the increase in the system size, the relative influence of the Ar messenger atom on the spectrum is further reduced in comparison with the Cu5\(^+\) case: DFT calculations for the Ar-tagged system predict an almost identical spectrum as that for the bare cluster, with (apart from a cluster−Ar vibration at 70 cm\(^{-1}\)) minimal modifications in band frequencies and only a two-fold intensity increase.

In this way, we recorded IRMPD spectra for Cu\(_n^+\)−Ar\(_3\) (\(n = 3–10\)) and unambiguously assigned each spectrum to a specific structure. Further experimental and calculated spectra, including spectra for Ar-complexed clusters, can be found in the SI. Beside the clusters discussed above, we find a triangular structure for Cu\(_5^+\), a capped triangular bipyramid structure for Cu\(_4^+\), and a pentagonal bipyramid structure capped with one and two Cu atoms for Cu\(_7^+\) and Cu\(_9^+\), respectively. It is now of interest to evaluate the stepwise evolution of the cluster structures found, as shown in Figure 2. Clearly, the two smallest clusters Cu\(_4^+\) and Cu\(_5^+\) are planar. However, from \(n = 6\), all structures found are 3D. What is particularly interesting is that every structure found is formed by the addition of a Cu atom onto the structure of the previous cluster size. From Cu\(_7^+\), the pentagonal motif dominates, forming a template onto which larger clusters form. We thus see no signs of the emergence of fcc-like structures, as one would expect from the bulk structure. All structures found here are in agreement with most of the previous theoretical studies\(^{21–26}\) except for the structure of Cu\(_5^+\) discussed above.

When comparing the structures found with those reported from calculations for neutral and anionic copper clusters, it is remarkable that cationic structures undergo the 2D to 3D transformation at smaller sizes (\(n = 5\) for cations and 6 and 7 for anions and neutrals, respectively).\(^{24–25}\) The same trend was shown for Ag\(_n^+\) (\(n = 6\) versus \(n = 6/7\) (anions/neutrals)) and Au\(_n^+\) (\(n = 8\) versus 12/11),\(^{34}\) illustrating the importance of
charge for the cluster geometry: The structure is a fine balance between interatomic stabilization and surface energies, and as a consequence, the loss of an electron leads to the weakening of the interatomic bonds, and the cluster needs to rearrange to reduce the surface energy.22

We can further compare the structures found for Cu_{n}^{+} with those for other coinage metals. Interestingly, for Ag_{n}^{+}, the same structures for the small cluster sizes (n = 3–10) are found, except for n = 5.55 It is further noteworthy that the structure found for Cu_{n}^{+} is the same as that used for computational studies on deposited Cu_{n}^{+} clusters. Such calculations are thus consistent with the form the clusters had prior to deposition on the Al_{2}O_{3} support.7

In conclusion, we established the structure of small cationic copper clusters based on a combination of IR photofragmentation spectroscopy of Cu_{n}^{+}–Ar_{n} complexes and DFT calculations, which convincingly account for the influence of the Ar.

ASSOCIATED CONTENT

Supporting Information
The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcllett.9b00539.

Experimental and computational details, IRMPD spectra of clusters with n = 3–10 and their comparison to different bare isomers, and the isomers of Ar-tagged complexes of the matching isomer (PDF)

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Notes
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