Search for doubly charged scalar bosons decaying into same-sign $W$ boson pairs with the ATLAS detector

ATLAS Collaboration*  
CERN, 1211 Geneva 23, Switzerland

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Abstract A search for doubly charged scalar bosons decaying into $W$ boson pairs is presented. It uses a data sample from proton–proton collisions corresponding to an integrated luminosity of 36.1 fb$^{-1}$ collected by the ATLAS detector at the LHC at a centre-of-mass energy of 13 TeV in 2015 and 2016. This search is guided by a model that includes an extension of the Higgs sector through a scalar triplet, leading to a rich phenomenology that includes doubly charged scalar bosons $H^\pm\pm$. Those bosons are produced in pairs in proton–proton collisions and decay predominantly into electroweak gauge bosons $H^\pm\pm \rightarrow W^\pm W^\pm$. Experimental signatures with several leptons, missing transverse energy and jets are explored. No significant deviations from the Standard Model predictions are found. The parameter space of the benchmark model is excluded at 95% confidence level for $H^\pm\pm$ bosons with masses between 200 and 220 GeV.

1 Introduction

An extension of the scalar sector of the Standard Model (SM) is possible in the context of type II seesaw models [1], originally conceived to explain the smallness of the neutrino masses. In the model investigated in this paper, the scalar sector includes a hypercharge $Y = 2$ scalar triplet, $\Delta$, in addition to the SM scalar doublet $H$ [2,3]. Electroweak symmetry breaking (EWSB) is achieved if the neutral components of $H$ and $\Delta$ acquire vacuum expectation values, $v_d$ and $v_t$ respectively. After the EWSB, the mixing between these fields results in seven scalar bosons: $H^\pm\pm$, $H^\pm$, $A^0$ (CP odd), $H^0$ (CP even), $h^0$ (CP even). A small mixing between the CP-even scalars allows $h^0$ to have the expected properties of the SM Higgs boson. In addition, the triplet-neutrino Yukawa term provides non-zero neutrino masses proportional to the vacuum expectation value of the triplet $v_t$. Constraints from electroweak precision measurements lead to an upper bound on $v_t$ of around 1 GeV. This range is significantly lower than the electroweak scale and matches the need for small values suggested by the natural association of $v_t$ with the neutrino masses.

The assumption of a non-zero $v_t$, of the order of a hundred MeV, opens the possibility for the doubly charged boson to decay into a pair of same-sign $W$ bosons, $H^\pm\pm \rightarrow W^\pm W^\pm$, while the leptonic decays $H^\pm\pm \rightarrow \ell^\pm\ell^\pm$ are suppressed with increasing $v_t$ [4,5]. Extensive searches for leptonic decays $H^\pm\pm \rightarrow \ell^\pm\ell^\pm$ have been performed at various colliders [6–11], where $H^\pm\pm$ bosons with masses up to about 800 GeV have been excluded. Moreover, searches for $H^\pm\pm \rightarrow W^\pm W^\pm$ decays have been performed by the CMS Collaboration in the context of single $H^\pm\pm$ production through vector-boson fusion at large $v_t$ (of order of tens of GeV) [12,13] for a model with two Higgs triplets [14]. For that model, a custodial symmetry avoids large contributions to the electroweak precision observables [15]. In contrast, the $H^\pm\pm \rightarrow W^\pm W^\pm$ decay mode has not been directly searched for so far for small values of $v_t$, where the vector-boson fusion is suppressed.

The present paper focuses on the phenomenology of doubly charged scalar bosons $H^\pm\pm$ that can be produced in pairs at the Large Hadron Collider (LHC) and decay into $W$ bosons. The triplet vacuum expectation value is taken to be $v_t = 0.1$ GeV such that only the $H^\pm\pm \rightarrow W^\pm W^\pm$ decays are relevant, leading to final states with four $W$ bosons. The mixing between the CP-even scalars is taken to be $10^{-4}$ and the remaining five Yukawa parameters in the potential are adjusted to obtain a given $H^\pm\pm$ mass hypothesis while requiring $h^0$ to have a mass of 125 GeV. The corresponding cross-section calculation is performed for on-shell $W$ bosons, and therefore only the region $m_{H^\pm\pm} > 200$ GeV is considered in the present analysis.

The four-boson final states are identified by the presence of light charged leptons (electrons or muons), missing transverse momentum, and jets. The analysis uses three final states defined according to the number of light leptons: same-sign (SS) dilepton channel ($2\ell^\pm$), trilepton channel ($3\ell$) and four-lepton channel ($4\ell$). Similar final states were used for other searches for new phenomena in ATLAS [16–18]. However,
the previously searched signal topologies differ significantly from those targeted in the present analysis and a dedicated event selection optimisation is therefore applied.

This paper includes a description of the experimental set-up in Sect. 2, followed by a description of the simulation used in the analysis in Sect. 3. The event selection and background estimations for the three explored signatures are described in Sect. 4. The signal region optimisation is described in Sect. 5. The systematic uncertainties are presented in Sect. 6. The results are shown in Sect. 7, followed by the conclusions in Sect. 8.

2 ATLAS detector

The ATLAS experiment [19] at the LHC is a multipurpose particle detector with a forward–backward symmetric cylindrical geometry and a near 4π coverage in solid angle. It consists of an inner tracking detector surrounded by a superconducting solenoid providing a 2 T axial magnetic field, electromagnetic and hadronic calorimeters, and a muon spectrometer. The inner tracking detector, covering the pseudorapidity range |η| < 2.5, consists of silicon pixel and silicon microstrip tracking detectors inside a transition-radiation tracker that covers |η| < 2.0. It includes, for the √s = 13 TeV running period, a newly installed innermost pixel layer, the insertable B-layer [20]. Lead–liquid-argon (LAr) sampling calorimeters provide electromagnetic (EM) energy measurements for |η| < 2.5 with high granularity and longitudinal segmentation. A hadronic calorimeter consisting of steel and scintillator tiles covers the central pseudorapidity range (|η| < 1.7). The endcap and forward regions are instrumented with LAr calorimeters for EM and hadronic energy measurements up to |η| = 4.9. The muon spectrometer surrounds the calorimeters and is based on three large air-core toroid superconducting magnets with eight coils each. It includes a system of precision tracking chambers (|η| < 2.7) and fast detectors for triggering (|η| < 2.4). A two-level trigger system is used to select events [21]. The first-level trigger is implemented in hardware and uses a subset of the detector information to reduce the accepted rate to a design maximum of 100 kHz. This is followed by a software-based trigger with a sustained average accepted event rate of about 1 kHz.

3 Data and simulation

The data sample collected by the ATLAS Collaboration at √s = 13 TeV during 2015 and 2016 was used. After the application of beam and data quality requirements, the integrated luminosity is 36.1 fb⁻¹.

Monte Carlo (MC) simulation samples were produced for signal and background processes using the full ATLAS detector simulation [22] based on GEANT4 [23] or, for selected smaller backgrounds and some of the signal samples, a fast simulation using a parameterisation of the calorimeter response and GEANT4 for the tracking system [24]. To simulate the effects of additional pp collisions in the same and nearby bunch crossings (pile-up), additional interactions were generated using PYTHIA 8.186 [25, 26] with a set of tuned parameters for the underlying event, referred to as the A2 tune [27], and the MSTW2008LO set of parton distribution functions (PDF) [28], and overlaid on the simulated hard-scatter event. The simulated events were reweighted to match the distribution of the number of interactions per bunch crossing observed in the data and were reconstructed using the same procedure as for the data.

The signal events containing \(H^{±±}\) pairs were simulated with the CalCEP generator version 3.4 [29], which is at leading order in QCD, using the Lagrangian described in Ref. [3] and the PDF CTEQ6.1 [30, 31]. The modelling of the parton showering and hadronisation of these events was performed using PYTHIA 8.186 [25, 26] with the A14 tune [32]. Event samples for the process \(pp \rightarrow H^{±±} H^{±±} \rightarrow W^± W^± W^± W^±\) were simulated for \(m_{H^{±±}}\) in the range from 200 to 700 GeV with steps of 100 GeV. The production cross-section decreases rapidly with \(m_{H^{±±}}\) and is 80.7 fb for \(m_{H^{±±}} = 200\) GeV, 5.0 fb for \(m_{H^{±±}} = 400\) GeV, and 0.35 fb for \(m_{H^{±±}} = 700\) GeV. Next-to-leading order (NLO) corrections [33] in QCD were applied, which increase these cross-sections by a factor 1.25. The fast detector simulation was used for the samples corresponding to \(m_{H^{±±}} > 500\) GeV.

The SM background processes were simulated using the MC event generator programs and configurations shown in Table 1. The production of \(VV\), \(VVqq\), and \(VVV\) (where \(V\) denotes a vector boson \(W\) or \(Z\) and \(qq\) labels the vector-boson fusion production mechanism) was simulated with a NLO QCD matrix element computed by SHERPA and matched to the SHERPA parton shower. The main background contribution in the \(2ℓ±±\) and \(3ℓ±\) channels is from \(WZ\) production, for which the total cross-section prediction is 48.2 ± 1.1 pb [44]. The main contribution to the \(4ℓ\) topology is from \(ZZ\) production with a total cross-section of 16.9 ± 0.6 pb [45, 46], which is suppressed by requiring significant missing transverse momentum in these events. The MC samples used to simulate \(t\bar{t}H\), \(t\bar{t}V\), \(VV\) and \(t\bar{t}\) are described in more detail in Refs. [47–49].

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1 ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z-axis along the beam pipe. The x-axis points from the IP to the centre of the LHC ring, and the y-axis points upwards. Cylindrical coordinates \((r, φ)\) are used in the transverse plane, \(φ\) being the azimuthal angle around the z-axis. The pseudorapidity is defined in terms of the polar angle \(θ\) as \(η = −\ln \tan(θ/2)\). Angular distance is measured in units of \(ΔR ≡ \sqrt{Δη^2 + (Δφ)^2}\).
Table 1 Configurations used for event generation of background processes. If only one PDF is shown, the same is used for both the matrix element (ME) and parton shower generators; if two are shown, the first is used for the matrix element calculation and the second for the parton shower. $V$ refers to the production of an electroweak boson ($W$ or $Z^*$). "Tune" refers to the underlying-event tune of the parton shower generator. "MG5_aMC" refers to MADGRAPH5_AMC@NLO 2.2.1; "Pythia 6" refers to version 6.427; "Pythia 8" refers to version 8.1; "Herwig++" refers to version 2.7. The samples have heavy flavour hadron decays modelled by EVTGEN1.2.0 [34], except for samples generated with SHERPA.

<table>
<thead>
<tr>
<th>Process</th>
<th>Event generator</th>
<th>ME order</th>
<th>Parton shower</th>
<th>PDF</th>
<th>Tune</th>
</tr>
</thead>
<tbody>
<tr>
<td>$VV, qqVV, VVV$</td>
<td>SHERPA 2.1.1</td>
<td>MEPS NLO</td>
<td>SHERPA 2.1.1</td>
<td>CT10 [36]</td>
<td>SHERPA 2.1.1 default</td>
</tr>
<tr>
<td>$\tilde{t}\tilde{t}H$</td>
<td>MG5_AMC [37]</td>
<td>NLO</td>
<td>PYTHIA 8</td>
<td>NNPDF 3.0 NLO [38]</td>
<td>A14 [32]</td>
</tr>
<tr>
<td>$VH$</td>
<td>PYTHIA 8</td>
<td>LO</td>
<td>PYTHIA 8</td>
<td>NNPDF 2.3 LO</td>
<td>A14</td>
</tr>
<tr>
<td>$t\bar{t}q\bar{b}$</td>
<td>MG5_AMC</td>
<td>LO</td>
<td>PYTHIA 8</td>
<td>CT10</td>
<td>A14</td>
</tr>
<tr>
<td>$t\bar{t}W$</td>
<td>MG5_AMC</td>
<td>NLO</td>
<td>Herwig++ [39]</td>
<td>CT10</td>
<td>UE-EE-5 [40]</td>
</tr>
<tr>
<td>$t\bar{t}(Z^<em>/\gamma^</em>)$</td>
<td>MG5_AMC</td>
<td>NLO</td>
<td>PYTHIA 8</td>
<td>NNPDF 3.0 NLO</td>
<td>A14</td>
</tr>
<tr>
<td>$t(Z^<em>/\gamma^</em>)$</td>
<td>MG5_AMC</td>
<td>LO</td>
<td>PYTHIA 6 [25]</td>
<td>CTEQ6L1 [30, 31]</td>
<td>Perugia2012 [41]</td>
</tr>
<tr>
<td>$tW(Z^<em>/\gamma^</em>)$</td>
<td>MG5_AMC</td>
<td>NLO</td>
<td>PYTHIA 8</td>
<td>NNPDF 2.3 LO</td>
<td>A14</td>
</tr>
<tr>
<td>$\tilde{t}\tilde{t}, \tilde{t}\bar{t}$</td>
<td>MG5_AMC</td>
<td>LO</td>
<td>PYTHIA 8</td>
<td>NNPDF 2.3 LO</td>
<td>A14</td>
</tr>
<tr>
<td>$\tilde{t}\bar{t}W^{+}W^{-}$</td>
<td>MG5_AMC</td>
<td>LO</td>
<td>PYTHIA 8</td>
<td>NNPDF 3.0 NLO</td>
<td>A14</td>
</tr>
<tr>
<td>$VV\gamma$</td>
<td>SHERPA 2.2</td>
<td>MEPS NLO</td>
<td>SHERPA 2.2</td>
<td>NNPDF 3.0 NLO</td>
<td>SHERPA 2.2 default</td>
</tr>
<tr>
<td>$s\rightarrow t$-channel, $Wt$ single top</td>
<td>POWHEG-BOX v 2 [42, 43]</td>
<td>NLO</td>
<td>PYTHIA 6</td>
<td>CT10/CTEQ6L1</td>
<td>Perugia2012</td>
</tr>
</tbody>
</table>
The simulated SM contributions in each of the channels considered are separated into prompt-lepton and fake-lepton contributions, depending on the source of the reconstructed leptons at generator level. The processes that contain only reconstructed charged leptons originating from prompt leptonic decays of $W$ and $Z$ bosons are classified as a prompt-lepton contribution, while processes with at least one of the reconstructed leptons being a misidentified hadron or photon, or a lepton from hadron decays constitute the fake-lepton contribution. The simulated events are not used to evaluate the background originating from charge-misidentified leptons for the $2\ell^\text{es}$ channel and fake leptons for the $2\ell^\text{es}$ and $3\ell$ channels. These are estimated in general using data-driven methods because they are not well modelled by simulations. This is the case in particular for $Z \rightarrow \ell^+\ell^-$, $W \rightarrow \ell\nu$ and $t\bar{t}$ processes. The background process $V\gamma$ can contribute if electrons originating from the photon conversion are selected. This contribution is found to be small and adequately modelled, so it is estimated using the MC simulation. For the $4\ell$ channel, the background from fake leptons is small and the data-driven methods are not applicable due to the low number of events available, so the MC simulation is used to estimate both the prompt-lepton and the fake-lepton contributions.

4 Event selection and background estimates

4.1 Event reconstruction

Interaction vertices originating from $pp$ collisions are reconstructed using at least two tracks with transverse momentum $p_T > 0.4$ GeV, and required to be consistent with the beamspot envelope. The primary vertex is identified as the vertex with the largest sum of squares of the transverse momenta from associated tracks [50].

Electrons are reconstructed as tracks in the inner detector matched to clusters in the electromagnetic calorimeter, within the region of pseudorapidity $|\eta| < 2.47$ [51]. The candidates in the transition region between the barrel and the endcap calorimeters ($1.37 < |\eta| < 1.52$) are removed. Only those electron candidates with transverse momentum greater than 10 GeV are considered. The electron identification is based on a multivariate likelihood-based discriminant that uses the shower shapes in the electromagnetic calorimeter and the associated track properties measured in the inner detector. In particular, the loose and tight identification working points, described in Ref. [51], are used, providing electron identification efficiencies of approximately 95% and 78–90% (depending on $p_T$ and $\eta$), respectively.  

The missing transverse momentum, with magnitude $E_T^\text{miss}$, is defined as the negative vector sum of the transverse momenta of all identified leptons and jets and the remaining unclustered energy of the event, which is estimated from tracks associated with the primary vertex but not assigned to any physics object [58].
4.2 Event preselection

Candidate events are selected using triggers that require at least one electron or one muon to pass various thresholds of $p_T$ [21]. The higher thresholds are applied with looser lepton identification and/or isolation requirements in order to ensure efficiencies close to 100% for leptons with transverse momentum above 30 GeV.

The signal topologies studied in this search involve the presence of at least two leptons of the same charge and are classified as explained above in three mutually exclusive categories: $2\ell^{ss}$, $3\ell$ and $4\ell$ channels. The $2\ell^{ss}$ channel targets signal events where the two same-sign $W$ bosons from one of the doubly charged Higgs boson decays leptotonically, while the other two $W$ bosons from the other doubly charged Higgs boson decay hadronically. In the $3\ell$ channel, one $W$ boson decays hadronically and in the $4\ell$ channel, all $W$ bosons decay leptonically. All channels present significant $E_T^{miss}$ corresponding to the neutrinos from leptonic $W$ boson decays. In the $2\ell^{ss}$ and $4\ell$ channels, jets from $W$ boson decays originate from the first- and second-generation quarks, and therefore lead to events without $b$-jets. The event selection is divided into two steps: the preselection and the signal region selection.

The preselection requirements are summarised in Table 2. The electrons (muons) are selected in the pseudorapidity range $|\eta| < 2.47$ ($2.5$) with a transverse momentum of at least 10 GeV, satisfying the type $L$ requirement. Events are selected only if the absolute value of the sum of charges of the leptons is two, one and zero for the $2\ell^{ss}$, $3\ell$ and $4\ell$ channels, respectively. At least one of the leptons is required to have $p_T > 30$ GeV to ensure a high trigger efficiency. To reduce the fake-lepton contamination in the $2\ell^{ss}$ channel, the second highest $p_T$ (subleading) lepton is required to have $p_T > 20$ GeV and both leptons are required to be of type $T$.

Similarly in the $3\ell$ channel, each lepton in the pair of leptons of the same sign, which is expected to suffer more from fake-lepton contamination, is required to have $p_T > 20$ GeV and to both be of type $T$. In the $2\ell^{ss}$ and $4\ell$ channels, the leptons are labelled by descending $p_T$, and are denoted by $\ell_{1,2}$. The ranking follows a different logic for the $3\ell$ channel: the lepton that has a charge opposite to the total lepton charge is denoted as $\ell_0$, while the same-sign leptons are denoted by $\ell_1$ and $\ell_2$, ranked by increasing distance to $\ell_0$ in the $\eta$-$\phi$ plane.

Further preselection requirements are based on $E_T^{miss}$, the jet multiplicity $N_{jets}$ and the number of jets tagged as $b$-jets $N_{b-jet}$. Moreover, in order to reduce the background from $Z$ bosons and neutral mesons decaying into same-flavour opposite-sign leptons (SFOS), the invariant mass of such lepton pairs is required to be greater than 12 (15) GeV for the $3\ell$ ($4\ell$) channel and to have an invariant mass that is not compatible with the $Z$ boson. For the $2\ell^{ss}$ channel, the $Z$ boson invariant mass veto is also applied to $e^\pm e^\pm$ events, in order to reduce the contributions originating from electron charge misidentification.

After this preselection, 562 data events are selected in the $2\ell^{ss}$ channel, 392 events in the $3\ell$ channel, and 44 events in the $4\ell$ channel.

Table 2 The preselection criteria for the three analysis channels. The leptons are ordered by decreasing $p_T$ ($\ell_1, \ell_2, \ldots$) in the $2\ell^{ss}$ and $4\ell$ channels, while for the $3\ell$ channel $\ell_1, \ell_2$ denote the same-sign leptons and $\ell_0$ the lepton with a charge opposite to the total lepton charge. $Q_\ell$ denotes the charge of each lepton.

<table>
<thead>
<tr>
<th>Selection criteria</th>
<th>$2\ell^{ss}$</th>
<th>$3\ell$</th>
<th>$4\ell$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Trigger</td>
<td>At least one lepton with $p_T &gt; 30$ GeV that fulfills the requirements of single-lepton triggers</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$N_{jets}(L\text{-type})$</td>
<td>$2$</td>
<td>$3$</td>
<td>$4$</td>
</tr>
<tr>
<td>$N_{jets}(T\text{-type})$</td>
<td>$2$</td>
<td>$2 (\ell_1, \ell_2)$</td>
<td>$-1$</td>
</tr>
<tr>
<td>$</td>
<td>\sum Q_\ell</td>
<td>$</td>
<td>$2$</td>
</tr>
<tr>
<td>Lepton $p_T$ threshold</td>
<td>$p_T^{\ell_1, \ell_2} &gt; 30$, $20$ GeV</td>
<td>$p_T^{b-jet} &gt; 30$, $20$ GeV</td>
<td>$p_T^{\ell_1, \ell_2, \ell_3, \ell_4} &gt; 10$ GeV</td>
</tr>
<tr>
<td>$E_T^{miss}$</td>
<td>$\geq 3$</td>
<td>$\geq 2$</td>
<td>$-1$</td>
</tr>
<tr>
<td>$b$-jet veto</td>
<td>$N_{b-jet} = 0$</td>
<td>$N_{b-jet} = 0$</td>
<td>$N_{b-jet} = 0$</td>
</tr>
<tr>
<td>Low SFOS $m_{\ell\ell}$ veto</td>
<td>$m_{e^\pm e^\pm} - m_Z &gt; 10$ GeV</td>
<td>$m_{\ell^+\ell^-} &gt; 15$ GeV</td>
<td>$m_{\ell^+\ell^-} &gt; 12$ GeV</td>
</tr>
<tr>
<td>$Z$ boson decays veto</td>
<td>$</td>
<td>m_{e^\pm e^\pm} - m_Z</td>
<td>&gt; 10$ GeV</td>
</tr>
</tbody>
</table>
4.3.1 Charge misidentification

In the $2\ell^{ss}$ channel, a background contribution is expected from events with opposite-sign lepton pairs when the charge of one of the leptons is misidentified, while the background contribution from charge misidentification is negligible for $3\ell$ and $4\ell$ channels. In the transverse momentum domain relevant for this analysis, charge misidentification is only significant for electrons and is due mainly to bremsstrahlung interactions with the inner detector material. The radiated photon produces an $e^+e^-$ pair near the original electron trajectory leading to a charge identification confusion.

The misidentification rate is measured using a large data sample of dilepton events originating mainly from $Z \rightarrow e^+e^-$ decays selected by two type $T$ electrons with an invariant mass between 80 and 100 GeV. The sample contains mostly opposite-sign dileptons, with a small fraction of same-sign dileptons. The fraction of same-sign dilepton events is estimated using sidebands around the peak. Its impact on the charge misidentification rate is about 2% and is included in the systematic uncertainty.

The background from fake leptons in both the opposite-sign and same-sign samples is estimated using sidebands around the $Z$ boson mass peak. Its impact on the charge misidentification rate is about 2% and is included in the systematic uncertainty.

The background from charge misidentification in a given region is estimated using a data control sample selected with the same criteria as the nominal sample but with opposite-sign dilepton pairs, where at least one lepton is an electron, weighted by the probability that the charge of the electron(s) is misidentified.

4.3.2 Fake-lepton contributions

The composition of the fake-lepton background varies considerably among the analysis channels. Therefore, the methods to estimate the fake-lepton contributions are different for the $2\ell^{ss}$, $3\ell$ and $4\ell$ channels. The contribution from fake leptons for the $2\ell^{ss}$ and $3\ell$ channels are estimated using the fake-factor method, while the simulation prediction corrected with data-driven scale factors is used for the $4\ell$ channels. Those methods involve various fake-enriched control samples that are summarised in Table 3 and described below.

**Table 3** The selection criteria defining the fake-enriched control regions used to determine the fake factors for the $2\ell^{ss}$ and $3\ell$ channel and the MC scale factors for the $4\ell$ channel. The symbol “...” means no requirement. The transverse mass $m_T$, used for the 4\ell-Z region to reduce the $WZ$ contributions, is calculated as the invariant mass of the vector sum of transverse momentum of the fake-lepton candidate and the missing transverse momentum.

| Sample | $2\ell^{ss}$ | $3\ell$ | $4\ell$-$Z$ | $4\ell$-$T$
|--------|-------------|---------|-------------|-------------
| $N_t$ (type $L$) | 2 | 3 | 3 | 3 |
| $\sum |Q_t|$ | 2 | 1 | 1 | 1 |
| $p_T^T$ | $> 30$, 20 GeV | $> 10$, 20, 20 GeV | $> 10$, 10, 10 GeV | $> 10$, 10, 10 GeV |
| $N_{jets}$ | $\geq 3$ | 1 | 1 or 2 | 1 or 2 |
| $N_{b jet}$ | 0 | - | - | - |
| $Z$-window | $|m_{\ell\ell} - m_Z| > 10$ GeV | $|m_{\ell\ell}^{os} - m_Z| > 10$ GeV | $|m_{\ell\ell}^{os} - m_Z| < 10$ GeV | No same-flavour opposite-sign lepton pair |
| $m_{\ell\ell}^{os}$ | - | - | - | - |
| $E_T^{miss}$ | $< 70$ GeV | - | $< 50$ GeV | - |
| $m_T$ | - | - | $< 50$ GeV | - |
Fig. 1 Distribution of variables used for the signal region optimisation of the $2\ell_{ss}$ final state. The events are selected with the preselection requirements listed in Table 2. The data (dots) are compared with the predictions (histograms) that include the contributions from the dominant prompt-lepton background ($WZ$), other prompt-lepton backgrounds, processes where a fake lepton is reconstructed, and electrons with misidentified charge (QMisdID). The expected signal distributions corresponding to two $H^{\pm\pm}$ masses are also shown, scaled up for visibility. The last bin includes overflows. In each figure the bottom panel shows the ratio of data to the prediction, where the band around unity represents the total uncertainty of the SM prediction.

\[
\theta_{\mu}^{2\ell_{ss}} = \left( \frac{N_{\text{Data}}^{C} - N_{\text{Prompt}}^{C}}{N_{\text{Data}}^{C} - N_{\text{Prompt}}^{C}} \right)_{\mu\mu},
\]

where $N_{\text{Data}}^{C}$ and $N_{\text{Prompt}}^{C}$ are the number of events where both muons are of type $T$, and where one is of type $T$ and the other of type $T'$, respectively. The prompt-lepton contributions $N_{\text{Prompt}}^{C}$, which are estimated using MC simulation, are subtracted from data event yields to obtain a pure estimate of the fake-lepton contributions in the $\mu\mu$ and $\mu\mu'$ regions. The superscript $C$ indicates the fake-enriched control region.

The electron fake factor is computed using the fake-enriched region where a same-sign $e\mu$ pair was selected:

\[
\theta_{e\mu}^{2\ell_{ss}} = \left( \frac{N_{\text{Data}}^{C} - N_{\text{Prompt}}^{C} - N_{\text{fake-}e\mu}^{C}}{N_{\text{Data}}^{C} - N_{\text{fake-}e\mu}^{C}} \right)_{e\mu}.
\]

In addition to the prompt-lepton contribution, the electron charge-misidentification contribution, denoted by $N_{e\ell}^{C}$, needs to be subtracted. It is computed using the method described in Sect. 4.3.1. Furthermore, the fake-muon contribution in the $e\mu$ sample is subtracted from the numerator. It is computed as:

\[
N_{e\mu}^{\text{fake-}e\mu} = \theta_{e\mu}^{2\ell_{ss}} (N_{\text{Data}}^{C} - N_{\text{Prompt}}^{C})_{e\mu}.
\]

The muon fake contribution is not considered in the denominator of the electron fake factor, in Eq. (2), because it is negligible.

The muon fake factor is measured to be $0.14 \pm 0.03$, while the electron fake factor is $0.48 \pm 0.07$, where the uncertainties are statistical only. A systematic uncertainty of 35% (56%) in the electron (muon) fake factor is estimated from complementary control samples with low jet multiplicity or by applying a different selection to vary the fraction of jets containing heavy-flavour hadrons. The uncertainty in the muon fake factor is larger due to lower number of data events available for those checks. The fake-lepton contributions in the nominal region (signal or preselection, denoted collectively by the superscript $R$) are obtained by multiplying the fake factors by the number of events in a region with the same selection as the nominal region, but where at least one lepton is of type $T$.
Fig. 2 Distribution of variables used for the signal region optimisation of the 3ℓ channel (a detailed description can be found in the caption of Fig. 1)

\[
N_{\text{fake},R} = \theta_2^{2\ell ss} \times (N^{\text{Data}} - N^{\text{Prompt}} - N^{\ell_0} R_{\mu \mu}) \\
+ \theta_2^{2\ell ss} \times (N^{\text{Data}} - N^{\text{Prompt}} - N^{\ell_0} R_{ee}^{\ell_0}),
\]

\[
N_{\text{fake},R} = \theta_2^{2\ell ss} \times (N^{\text{Data}} - N^{\text{Prompt}} - N^{\ell_0} R_{\mu \mu}) \\
+ \theta_2^{2\ell ss} \times (N^{\text{Data}} - N^{\text{Prompt}} - N^{\ell_0} R_{ee}^{\ell_0}),
\]

(3)

where the prompt-lepton and the charge misidentifications contributions are subtracted as explained above.

**Fake-lepton contribution estimate for the 3ℓ channel**

A method similar to that employed for the 2ℓss channel is applied for the 3ℓ channel. Here the opposite-sign lepton \( \ell_0 \) is assumed to be prompt, an assumption that was found to be valid in MC simulation. The fake-enriched region used to calculate the fake factors for the 3ℓ channel, which is described in Table 3, follows the 3ℓ preselection conditions except that the jet multiplicity is required to be exactly one. The fake factors for electrons and muons are both calculated by applying a formula analogous to Eq. (1) to the \( \ell_0 ee/\ell_0 e\mu \) and \( \ell_0 \mu e/\ell_0 \mu \mu \) regions, respectively. The muon fake factor is found to be 0.17 ± 0.06 and the electron fake factor is found to be 0.39 ± 0.07, where the errors are statistical only. The values are compatible with those obtained for the 2ℓss channel. Additional control samples, defined such that the content is enriched in either Z+jets or tt events, are used to test the method and to estimate systematic uncertainties of 55% and 81% for the electron and muon fake factors, respectively. The fake-lepton contributions to the nominal regions are then calculated using relations analogous to Eq. (3).

**Fake-lepton contribution estimate for the 4ℓ channel**

There are too few data events to apply the fake-factor method in the 4ℓ channel. Instead, the fake-lepton contribution is estimated from the yields predicted by the MC simulation but corrected using process-dependent scale factors that are extracted in two fake-enriched control regions. The fake-lepton contribution in this channel comes mainly from \( t\bar{t}V \) processes, where the fake lepton originates from a b-jet. A small component from light quarks is also present. Two data samples designed to contain fake leptons originating from Z+jets and \( t\bar{t} \) events are used to study the capability of the simulation to describe fake leptons originating from light- and heavy-flavour jets, respectively. The two control samples are labelled Z and T and are defined in Table 3. The samples are required to have three identified leptons. For the Z region, the fake-lepton candidate is assumed not to be part of the lepton pair forming the Z boson candidate. For the T region, the fake lepton is assumed to be the lepton with the lower \( p_T \) in the same-sign lepton pair. The scale fac-
Fig. 3 Distribution of variables used for the signal region optimisation of the 4ℓ channel (a detailed description can be found in the caption of Fig. 1)

tors are derived independently for fake electrons and fake muons. Four scale factors $\lambda_\ell^X$ (with $\ell = e, \mu$ and $X = Z, T$) are obtained by solving the system of equations

$$N_{\text{Data}}^\ell|X - N_{\text{Prompt}}^\ell|X = \lambda_\ell^T N_T^\ell|X + \lambda_\ell^Z N_Z^\ell|Z+\text{jets}|X,$$

where the event yields $N^\ell$ are labelled by the nature of the contribution, data (Data) or simulation (Prompt), $T$ and $Z+\text{jets}$, and the equations are derived in each of the respective control region $X$ (Z or T). The obtained scale factors are $\lambda_e^T = 1.12 \pm 0.05$, $\lambda_\mu^T = 1.02 \pm 0.07$, $\lambda_e^Z = 1.11 \pm 0.05$ and $\lambda_\mu^Z = 0.94 \pm 0.07$, where the errors are statistical only. Alternative trilepton control samples, where the jet multiplicity and the lepton $p_T$ threshold are varied, are used to estimate a systematic uncertainty of 50% in these scale factors. The scale factors are used as weights to the simulated events that contain a fake lepton according to the fake-lepton flavour and the presence of heavy-flavour jets in the event.

5 Signal region optimisation

The hypothetical signal produces four $W^\pm$ bosons in each event. Since at least two leptonic $W$ boson decays are needed to lead to the multi-lepton topologies considered in this analysis, all signal events are expected to feature significant $E_T^{\text{miss}}$, while jets are expected from hadronic $W$ boson decays for $2\ell^\mp$ and $3\ell$ channels. Moreover, when the mass of the doubly charged Higgs boson is in the range of 200–300 GeV, each $H^{\pm\pm}$ is produced with a significant momentum and the two subsequent $W$ bosons are emitted close to each other in the laboratory frame. Consequently, the two same-sign leptons from the decays of the two $W$ bosons tend to be close in the $\eta\phi$ plane. The decay products of the other doubly charged Higgs boson are generally well-separated from the two same-sign leptons.

The analysis channels face different background contributions from the SM. The $2\ell^\mp$ category is populated with events containing one prompt lepton from a $W$ boson, or to a lesser extent from a $Z$ boson, and one fake lepton from the hadronic final state produced. The $2\ell^\mp$ events with two same-sign electrons can also originate from Drell–Yan and $t\bar{t}$ production, where the charge of one of the electrons is misidentified, as explained above. In the $2\ell^\mp$ and $3\ell$ channels, most of the expected prompt-lepton contribution is due to the production of $WZ$ associated with jets, with both bosons subsequently decaying into leptons. This process also produces other features of the signal, such as significant $E_T^{\text{miss}}$ and the absence
### Table 4

The selection criteria used to define the signal regions. The variables are described in Sect. 5

<table>
<thead>
<tr>
<th>Selection criteria</th>
<th>$2\ell^\pm$</th>
<th>$e^\pm\mu^\pm$</th>
<th>$\mu^\pm\mu^\pm$</th>
<th>$3\ell$</th>
<th>$4\ell$</th>
</tr>
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<tr>
<td>$m_{H^{\pm\pm}} = 200$ GeV</td>
<td></td>
<td></td>
<td></td>
<td>SFOS 0</td>
<td>SFOS 1.2</td>
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<td>$&gt; 100$</td>
<td>$&gt; 100$</td>
<td>$&gt; 100$</td>
<td>$&gt; 45$</td>
<td>$&gt; 45$</td>
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<tr>
<td>$m_{\ell\ell}$ [GeV]</td>
<td>[25, 130]</td>
<td>[15, 150]</td>
<td>[35, 150]</td>
<td>$&gt; 160$</td>
<td>$&gt; 170$</td>
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<td>$\Delta R_{\ell\ell}$ [rad.]</td>
<td>$&lt; 0.8$</td>
<td>$&lt; 1.8$</td>
<td>$&lt; 0.9$</td>
<td>$[0.15, 1.57]$</td>
<td>$[0.00, 1.52]$</td>
</tr>
<tr>
<td>$\Delta \phi(\ell \ell, E_T^{\text{miss}})$ [rad.]</td>
<td>$&lt; 1.1$</td>
<td>$&lt; 1.3$</td>
<td>$&lt; 1.3$</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S$ [rad.]</td>
<td>$&lt; 0.3$</td>
<td>$&lt; 0.3$</td>
<td>$&lt; 0.2$</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$m_{\text{jets}}$ [GeV]</td>
<td>[140, 770]</td>
<td>[95, 330]</td>
<td>[95, 640]</td>
<td></td>
<td></td>
</tr>
<tr>
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<td></td>
<td></td>
<td>$[0.08, 1.88]$</td>
<td>$[0.07, 1.31]$</td>
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<td>$p_T^{\text{leading jet}}$ [GeV]</td>
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<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\Delta R_{\ell\ell}$ [rad.]</td>
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<td></td>
<td></td>
<td>$[0.16, 1.21]$</td>
<td></td>
</tr>
<tr>
<td>$m_{H^{\pm\pm}} = 300$ GeV</td>
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<td></td>
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</tr>
<tr>
<td>$E_T^{\text{miss}}$ [GeV]</td>
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<td>$&gt; 200$</td>
<td>$&gt; 200$</td>
<td>$&gt; 65$</td>
<td>$&gt; 55$</td>
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<tr>
<td>$m_{\ell\ell}$ [GeV]</td>
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<td>[80, 320]</td>
<td>[80, 320]</td>
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<td>$&lt; 1.8$</td>
<td>$&lt; 1.8$</td>
<td>$[0.18, 2.23]$</td>
<td>$[0.08, 2.23]$</td>
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<tr>
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<td></td>
</tr>
<tr>
<td>$S$ [rad.]</td>
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<td>$&lt; 0.4$</td>
<td>$&lt; 0.4$</td>
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<td></td>
</tr>
<tr>
<td>$m_{\text{jets}}$ [GeV]</td>
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<td>[130, 640]</td>
<td>[130, 640]</td>
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<tr>
<td>$\Delta R_{\ell\ell}$ [rad.]</td>
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<td></td>
<td>$[0.27, 2.37]$</td>
<td>$[0.21, 2.08]$</td>
</tr>
<tr>
<td>$p_T^{\text{leading jet}}$ [GeV]</td>
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<td>$p_T^{\ell\ell}$ [GeV]</td>
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<td></td>
<td></td>
</tr>
<tr>
<td>$\Delta R_{\ell\ell}$ [rad.]</td>
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<td></td>
<td></td>
<td>$[0.09, 1.97]$</td>
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</tr>
<tr>
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<td>$&gt; 200$</td>
<td>$&gt; 200$</td>
<td>$&gt; 200$</td>
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<td>$&gt; 85$</td>
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<td>$&lt; 1.8$</td>
<td>$&lt; 1.8$</td>
<td>$[0.22, 2.39]$</td>
<td>$[0.29, 2.69]$</td>
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<tr>
<td>$\Delta \phi(\ell \ell, E_T^{\text{miss}})$ [rad.]</td>
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<td>$&lt; 2.4$</td>
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<td></td>
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<td>$&lt; 0.6$</td>
<td>$&lt; 0.5$</td>
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<td>$m_{\text{jets}}$ [GeV]</td>
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<td></td>
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<tr>
<td>$\Delta R_{\ell\ell}$ [rad.]</td>
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<td></td>
<td>$[0.30, 2.59]$</td>
<td>$[0.31, 2.30]$</td>
</tr>
<tr>
<td>$p_T^{\text{leading jet}}$ [GeV]</td>
<td></td>
<td></td>
<td></td>
<td>$&gt; 120$</td>
<td>$&gt; 100$</td>
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<td>$p_T^{\ell\ell}$ [GeV]</td>
<td></td>
<td></td>
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<tr>
<td>$\Delta R_{\ell\ell}$ [rad.]</td>
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<td></td>
<td>$[0.39, 2.22]$</td>
<td></td>
</tr>
<tr>
<td>$\Delta R_{\ell\ell}$ [rad.]</td>
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<td></td>
<td></td>
<td>$[0.55, 2.90]$</td>
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</table>
of $b$-jets for most of the production cross-section. For the $WZ$ events, the mass of the same-flavour opposite-sign lepton pair is close to the $Z$ boson mass, while no such resonant distribution is expected for the signal. In the $4\ell$ channel, the dominant background originates from $t\bar{t}V$ and $ZZ$ production. Processes containing top quarks ($t\bar{t}$, $t\bar{t}V$) can lead to events with multiple leptons in the final state. A noticeable feature of those processes is the presence of $b$-jets.

Given these properties of the signal and of the expected background, the following discriminating variables, in addition to $E_T^{miss}$, are considered:

- $m_{x\ell}$, the invariant mass of the system composed of all selected leptons in the event, where $x$ can be 2, 3 or 4.
- $\Delta R_{\ell+\ell}$, the distance in $\eta$--$\phi$ between two same-sign leptons. This variable is used for the $2\ell^{ss}$ and $3\ell$ channels. In the $4\ell$ channel, two such variables can be calculated per event, $\Delta R_{\ell+\ell}^{\text{min}}$ and $\Delta R_{\ell+\ell}^{\text{max}}$, denoting the minimum and maximum values, respectively.
- $m_{\text{jets}}$, the invariant mass of the system composed of all jets in the event. When there are more than four jets in the event, only the leading four jets are used. This variable is used only for the $2\ell^{ss}$ channel.
- $p_T^{\text{leading jet}}$, the transverse momentum of the highest-$p_T$ jet.
- $\Delta\phi(\ell\ell, E_T^{\text{miss}})$, the difference in azimuth between the dilepton system and $E_T^{miss}$. This variable is used in the $2\ell^{ss}$ channel.
- $\Delta R_{\ell-jet}$, the minimal distance in $\eta$--$\phi$ between any lepton and its closest jet. This variable is used in the $3\ell$ channel.
- $S$, is a variable used for the $2\ell^{ss}$ channel to describe the event topology in the transverse plane, and defined using the spread of the $\phi$ angles of the leptons, $E_T^{miss}$, and jets as follows:

$$S = \frac{R(\phi_{\ell_1}, \phi_{\ell_2}, \phi_{E_T^{miss}}) \cdot R(\phi_{\ell_1}, \phi_{\ell_2}, \phi_{E_T^{miss}}, \phi_{j_1}, \phi_{j_2}, \ldots)}{R(\phi_{\ell_1}, \phi_{\ell_2}, \phi_{E_T^{miss}})}$$

where the $R$ is the root mean square that quantifies the spread, $R(\phi_1, \ldots, \phi_n) = \sqrt{\frac{1}{n} \sum_{i=1}^{n} (\phi_i - \bar{\phi})^2}$. The azimuthal angles $\phi$ are bounded in $(-\pi, \pi]$, and the bound is considered in the calculation. The $S$ variable is expected to be on average smaller for the signal than for the background for low $H^{\pm\pm}$ mass values.

The distributions of the selected variables for the $2\ell^{ss}$, $3\ell$ and $4\ell$ channels are shown at preselection level in Figs. 1, 2 and 3, respectively. The data are compared with the sum of the prompt lepton, fake lepton and charge-misidentified lepton background predictions. The prompt-lepton backgrounds are estimated with simulations while the background from fake leptons and charge-flipped leptons are measured with the methods described in the previous section. Good agreement is observed in both normalisation and shape, demonstrating that the background contributions are well modelled. The expected signal distributions for various $H^{\pm\pm}$ masses are also shown to illustrate the discriminating power of the selected variables.

The strategy used to extract the signal is based on rectangular cut optimisation using the TMVA package [59]. For each $m_{H^{\pm\pm}}$ hypothesis, six signal regions are defined using the following lepton flavour content: in the $2\ell^{ss}$ channel, three signal regions are optimised separately for $ee$, $e\mu$ and $\mu\mu$ channels; in the $3\ell$ channel, the signal regions are optimised...
Fig. 4 Event yields in the signal regions optimised for the $m_{H^{±±}} = 200, 300, 400$ and 500 GeV searches. The bottom panel shows the ratio of the data to the total background prediction, where the band illustrates the total uncertainty of the SM background. The error bars attributed to data are estimated assuming a Poisson distribution with the average equal to the respective yields. The signal prediction is represented as a dotted histogram, stacked on the SM background.

In order to verify the background estimate reliability for the signal region, three further checks were performed: the optimised cuts were applied individually, the cuts were applied successively, or each cut was inverted while the other cuts were applied. The agreement between data and prediction remains adequate for all those cases.

6 Systematic uncertainties

The theoretical uncertainties associated with the signal prediction originate from the PDFs, the matrix element calculation and the parton shower simulation. The uncertainties related to PDFs are evaluated using the Hessian method provided in LHAPDF6 [60] and are found to be in the range separately for events with no same-flavour opposite-sign lepton pairs (SFOS 0, for which the SM background is small) and for events with one or two such pairs (SFOS 1,2); the $4\ell$ channel is treated globally, with no further lepton flavour distinction. The selection criteria used to define the signal regions are shown in Table 4. The optimisation is performed as a function of the $H^{±±}$ mass for $m_{H^{±±}} = 200, 300, 400$ and 500 GeV, and seeks the best expected signal significance. The last optimisation point is applied to $m_{H^{±±}} \geq 600$ GeV as well, since the signal discrimination power does not vary significantly in this regime.

In order to verify the background estimate reliability for the signal region, three further checks were performed:
Table 5: Event yields in the signal regions of corresponding targeted masses of $H^{±±}$. The signal yield is for the corresponding mass point and is normalised to the luminosity of 36.1 fb$^{-1}$. The dominant background from prompt-lepton sources is from the $WZ$ process in the $2ℓ^{±}$ channel. For the $3ℓ$ and $4ℓ$ channels, the dominant background from prompt-lepton sources is from $WZ$ and $t\bar{t}V$ processes. The overall signal acceptance $A$ and the upper limit of extra contribution to each signal region at 95% confidence level $n_{95}$ are also presented. The data and SM prediction yields obtained for $m_{H^{±±}} = 500$ GeV are also valid for $m_{H^{±±}} = 600$ and 700 GeV.

<table>
<thead>
<tr>
<th>Subchannel</th>
<th>$2ℓ^{±}$</th>
<th>$μ^{±}μ^{±}$</th>
<th>$3ℓ$</th>
<th>$4ℓ$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$m_{H^{±±}} = 200$ GeV</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>
| Prompt lepton | $0.5 \pm 0.2$ | $0.3 \pm 0.2$ | $1.3 \pm 0.6$ | $0.3 \pm 0.1$ | $1.4 \pm 0.5$ | $0.07 \pm 0.03$
| QMisID | $0.6 \pm 0.2$ | $0.4 \pm 0.1$ | $-$ | $-$ | $-$ | $-$
| Fake lepton | $1 \pm 1$ | $< 0.4$ | $0.4 \pm 0.3$ | $0.2 \pm 0.1$ | $0.2 \pm 0.1$ | $0.03 \pm 0.02$
| Total background | $2 \pm 1$ | $0.6 \pm 0.3$ | $1.7 \pm 0.7$ | $0.5 \pm 0.1$ | $1.7 \pm 0.6$ | $0.11 \pm 0.05$
| Signal | $1.1 \pm 0.2$ | $2.3 \pm 0.4$ | $2.4 \pm 0.4$ | $1.8 \pm 0.3$ | $5.0 \pm 0.9$ | $1.1 \pm 0.2$
| $A$ [%] | $0.037$ | $0.080$ | $0.082$ | $0.061$ | $0.17$ | $0.038$
| $n_{95}$ | $12.3$ | $7.1$ | $7.5$ | $4.1$ | $7.7$ | $3.8$
| Data | $3$ | $2$ | $2$ | $1$ | $2$ | $0$
| $m_{H^{±±}} = 300$ GeV | | | | |
| Prompt lepton | $0.1 \pm 0.1$ | $0.9 \pm 0.4$ | $0.02 \pm 0.02$ | $0.4 \pm 0.1$ | $4 \pm 1$ | $0.3 \pm 0.1$
| QMisID | $0.1 \pm 0.1$ | $0.07 \pm 0.04$ | $-$ | $-$ | $-$ | $-$
| Fake lepton | $0.4 \pm 0.5$ | $< 0.2$ | $< 0.4$ | $0.3 \pm 0.2$ | $0.8 \pm 0.4$ | $0.2 \pm 0.2$
| Total background | $0.7 \pm 0.5$ | $1.0 \pm 0.5$ | $0.02 \pm 0.02$ | $0.8 \pm 0.2$ | $5 \pm 2$ | $0.5 \pm 0.2$
| Signal | $0.16 \pm 0.03$ | $0.6 \pm 0.1$ | $0.29 \pm 0.05$ | $0.6 \pm 0.1$ | $1.8 \pm 0.3$ | $0.43 \pm 0.08$
| $A$ [%] | $0.027$ | $0.10$ | $0.049$ | $0.11$ | $0.30$ | $0.071$
| $n_{95}$ | $4.0$ | $9.6$ | $3.0$ | $3.1$ | $22.7$ | $3.8$
| Data | $0$ | $3$ | $0$ | $0$ | $11$ | $0$
| $m_{H^{±±}} = 400$ GeV | | | | |
| Prompt lepton | $0.7 \pm 0.3$ | $1.0 \pm 0.4$ | $0.2 \pm 0.1$ | $0.3 \pm 0.1$ | $4 \pm 1$ | $0.3 \pm 0.1$
| QMisID | $0.3 \pm 0.1$ | $0.2 \pm 0.1$ | $-$ | $-$ | $-$ | $-$
| Fake lepton | $0.4 \pm 0.5$ | $< 0.3$ | $< 0.4$ | $0.3 \pm 0.2$ | $0.2 \pm 0.1$ | $0.05 \pm 0.04$
| Total background | $1.4 \pm 0.6$ | $1.2 \pm 0.5$ | $0.3 \pm 0.1$ | $0.6 \pm 0.2$ | $4 \pm 1$ | $0.4 \pm 0.1$
| Signal | $0.20 \pm 0.04$ | $0.38 \pm 0.07$ | $0.19 \pm 0.03$ | $0.23 \pm 0.04$ | $0.6 \pm 0.1$ | $0.17 \pm 0.03$
| $A$ [%] | $0.11$ | $0.21$ | $0.11$ | $0.13$ | $0.36$ | $0.092$
| $n_{95}$ | $10.4$ | $18.3$ | $6.4$ | $3.1$ | $10.4$ | $4.3$
| Data | $2$ | $6$ | $1$ | $0$ | $4$ | $1$
| $m_{H^{±±}} = 500$ GeV | | | | |
| Prompt lepton | $1.0 \pm 0.4$ | $0.7 \pm 0.3$ | $0.3 \pm 0.2$ | $0.4 \pm 0.1$ | $3 \pm 1$ | $0.2 \pm 0.1$
| QMisID | $0.3 \pm 0.1$ | $0.2 \pm 0.1$ | $-$ | $-$ | $-$ | $-$
| Fake lepton | $0.2 \pm 0.5$ | $0.3 \pm 0.5$ | $< 0.4$ | $0.11 \pm 0.06$ | $0.10 \pm 0.05$ | $0.2 \pm 0.2$
| Total background | $1.6 \pm 0.6$ | $1.2 \pm 0.6$ | $0.3 \pm 0.2$ | $0.5 \pm 0.1$ | $3.0 \pm 0.8$ | $0.4 \pm 0.2$
| Signal | $0.10 \pm 0.02$ | $0.16 \pm 0.03$ | $0.07 \pm 0.01$ | $0.09 \pm 0.02$ | $0.24 \pm 0.04$ | $0.06 \pm 0.01$
| $A$ [%] | $0.16$ | $0.25$ | $0.11$ | $0.14$ | $0.37$ | $0.098$
| $A$ [%] $m_{H^{±±}} = 600$ GeV | $0.22$ | $0.36$ | $0.16$ | $0.17$ | $0.44$ | $0.11$
| $A$ [%] $m_{H^{±±}} = 700$ GeV | $0.26$ | $0.38$ | $0.17$ | $0.19$ | $0.48$ | $0.12$
| $n_{95}$ | $8.6$ | $12.7$ | $3.8$ | $3.0$ | $7.9$ | $4.9$
| Data | $4$ | $3$ | $0$ | $0$ | $2$ | $3$
from 2.5% to 4.5%. The uncertainty of the parton shower simulation is assessed by comparing PYTHIA (with A14 tune) and Herwig++ (with UEE5 tune [61]), and is found to be 2.4%, 1.7%, and 3.8% for the 2\ell^{\pm\mp}, 3\ell, and 4\ell channels, respectively. The higher-order corrections are assumed to induce an additional 15% uncertainty in the cross-section calculations [33]. Combining those uncertainties in quadrature, an overall uncertainty of 17% is obtained for the signal normalisation.

The theoretical uncertainties associated with the largest SM backgrounds, VV [48] (including same-sign WWqq and ZZ processes) and t\bar{t}V [62], are estimated using dedicated MC samples, where the factorisation and renormalisation scales are varied independently by factors of 2 and 0.5 and the parton shower parameters are varied within the given model uncertainties. The theoretical uncertainties obtained are 24% for VV and 17% for t\bar{t}V. The uncertainties related to PDFs are found to be negligible. The uncertainty associated with the V\gamma contributions is taken to be 25%, as indicated by dedicated studies using converted photons. The uncertainties related to the VVV and rZ process predictions are taken from the respective inclusive cross-section measurements [63,64]. For other rare backgrounds which have no dedicated measurements yet (t\bar{t}f, t\bar{t}W^{+}W^{-}), uncertainties of 50% are assumed and are found to have a negligible impact on the sensitivity.

The experimental uncertainties arise from the accuracy of the detector simulation and from the uncertainties associated with the data-driven methods that are used to estimate the instrumental backgrounds. These uncertainties originate from the following sources:

- The uncertainties related to event reconstruction include the lepton [52,65] and the jet [66] energy scales and resolutions and the uncertainties in the reconstruction of \( E_{T}^{\text{miss}} [/itex] [58]. The impact of this type of uncertainty on the signal and background yields is in the range 3–8% and 10–30%, respectively.

- The uncertainties related to the efficiencies of electron [51] and muon [52] reconstruction and identification, including the uncertainties in the trigger efficiency are estimated in dedicated studies. The impact of this type of uncertainty on the signal and background yields is found to be in the range 4–6% and 2–5%, respectively. The uncertainties related to the b-jet identification algorithms, used in the analysis to veto events containing b-jets, are found to be negligible.

- The uncertainties originating from data-taking conditions include the luminosity measurement and the pile-up simulation procedure. The uncertainty of the integrated luminosity is 2.1%, determined using a methodology similar to that detailed in Ref. [67]. The uncertainty from the pile-up simulation is about 5%.

- The uncertainties related to the background contributions from electron charge misidentification are 22–28% for the 2\ell^{\pm\mp} ee and e\mu channels. The uncertainties of the fake-lepton contributions range from 50 to 250%, and mainly originate from the fake factors (2\ell^{\pm\mp} and 3\ell channels) and the scale factors for MC simulation (4\ell channel) described in Sect. 4.3.2 and the statistics of the control samples. The uncertainties exceed 100% in some cases due to the subtraction of the prompt-lepton contributions in the fake-lepton control regions.

The theoretical and experimental systematic uncertainties described above are assumed to be correlated amongst the various signal regions in the interpretation of the final results. Overall, the sensitivity of the search is dominated by the statistical uncertainty of the event yield in the signal regions.

7 Results

The expected and observed event yields in the signal regions are shown in Fig. 4 and Table 5. For a \( H^{\pm\mp} [/itex] mass of 200 GeV, substantial signal yield is expected in all channels, and the analysis sensitivity is found to be comparable across the 2\ell^{\pm\mp}, 3\ell and 4\ell channels. No significant excess has been observed. Table 5 also includes the overall signal acceptance \( A [/itex], defined as the number of selected events selected in a given channel divided by the total number of \( pp \rightarrow H^{\pm\mp} H^{\pm\mp} \rightarrow W^{\pm} W^{\pm} W^{\pm} W^{\mp} [/itex] events and representing the signal reduction due to phase space acceptance, branching ratio and detector efficiency.

![Graph](image-url)
The statistical analysis of the results is based on a likelihood-ratio test [68] using the CLs method [69]. The parameter of interest is the signal strength, defined as the cross-section of the hypothetical contribution from physics beyond the SM in units of the cross-section of the benchmark model. The likelihood function is constructed from Poisson probability distributions of counting experiments for each of the six channels in each signal region. The systematic uncertainties are treated as nuisance parameters implemented in the likelihood functions with Gaussian constraints.

The expected and observed upper limits of the $H^{\pm\pm} \to W^\pm W^\pm$ cross-section at 95% confidence level (CL), obtained from the combination of 2$\ell^{ss}$, 3$\ell$ and 4$\ell$ channels for the six $H^{\pm\pm}$ mass hypotheses are shown in Fig. 5. Assuming a linear interpolation of the sensitivity between neighbouring mass hypotheses, and the cross-section of the benchmark model, the observed (expected) lower limit on the mass of the $H^{\pm\pm}$ boson is 220 GeV (250 GeV) at 95% CL.

8 Conclusion

A search for the pair production of doubly charged Higgs scalar bosons with subsequent decays into $W$ bosons is performed in proton–proton collisions at a centre-of-mass energy of 13 TeV. The data sample was collected by the ATLAS experiment at the LHC and corresponds to an integrated luminosity of 36.1 fb$^{-1}$. The search for the $H^{\pm\pm} \to W^\pm W^\pm$ decay mode, not considered in previous analyses at colliders, is motivated by a model with an extended scalar sector that includes a triplet in addition to the Standard Model scalar doublet. The analysis proceeds through the selection of multi-lepton events in three channels (a pair of same-sign leptons, three leptons and four leptons) with missing transverse momentum and jets. The signal region is optimised as a function of the $H^{\pm\pm}$ mass. The data are found to be in good agreement with the Standard Model predictions for all channels investigated. Combining those channels, the model considered is excluded at 95% confidence level for $H^{\pm\pm}$ boson masses between 200 and 220 GeV.

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References

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1 Department of Physics, University of Adelaide, Adelaide, Australia  
2 Physics Department, SUNY Albany, Albany, NY, USA  
3 Department of Physics, University of Alberta, Edmonton, AB, Canada  
4 (a) Department of Physics, Ankara University, Ankara, Turkey; (b) Istanbul Aydin University, Istanbul, Turkey; (c) Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey  
5 LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy, France  
6 High Energy Physics Division, Argonne National Laboratory, Argonne, IL, USA  
7 Department of Physics, University of Arizona, Tucson, AZ, USA  
8 Department of Physics, University of Texas at Arlington, Arlington, TX, USA  
9 Physics Department, National and Kapodistrian University of Athens, Athens, Greece  
10 Physics Department, National Technical University of Athens, Zografou, Greece  
11 Department of Physics, University of Texas at Austin, Austin, TX, USA  
12 (a) Faculty of Engineering and Natural Sciences, Bahcesehir University, Istanbul, Turkey; (b) Faculty of Engineering and Natural Sciences, Istanbul Bilgi University, Istanbul, Turkey; (c) Department of Physics, Bogazici University, Istanbul, Turkey; (d) Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey  
13 Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan  
14 Institut de Física d’Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain  
15 (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China; (b) Physics Department, Tsinghua University, Beijing, China; (c) Department of Physics, Nanjing University, Nanjing, China; (d) University of Chinese Academy of Science (UCAS), Beijing, China  
16 Institute of Physics, University of Belgrade, Belgrade, Serbia  
17 Department for Physics and Technology, University of Bergen, Bergen, Norway  
18 Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, USA  
19 Institut für Physik, Humboldt Universität zu Berlin, Berlin, Germany  
20 Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland  
21 School of Physics and Astronomy, University of Birmingham, Birmingham, UK  
22 Centro de Investigaciones, Universidad Antonio Nariño, Bogotá, Colombia  
23 (a) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy; (b) INFN Sezione di Bologna, Bologna, Italy  
24 Physikalisches Institut, Universität Bonn, Bonn, Germany  
25 Department of Physics, Boston University, Boston, MA, USA  
26 Department of Physics, Brandeis University, Waltham, MA, USA  
27 (a) Transilvania University of Brasov, Brasov, Romania; (b) Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania; (c) Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi, Romania; (d) Physics Department, National Institute for Research and Development of Isotopic and Molecular Technologies, Cluj-Napoca, Romania; (e) University Politehnica Bucharest, Bucharest, Romania; (f) West University in Timisoara, Timisoara, Romania  
28 (a) Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovakia; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic  
29 Physics Department, Brookhaven National Laboratory, Upton, NY, USA  
30 Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina  
31 Cavendish Laboratory, University of Cambridge, Cambridge, UK  
32 (a) Department of Physics, University of Cape Town, Cape Town, South Africa; (b) Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg, South Africa; (c) School of Physics, University of the Witwatersrand, Johannesburg, South Africa  
33 Department of Physics, Carleton University, Ottawa, ON, Canada  
34 (a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies-Université Hassan II, Casablanca, Morocco; (b) Centre National de l’Énergie des Sciences Techniques Nucléaires (CNESTEN), Rabat, Morocco; (c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Marrakech, Morocco; (d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda, Morocco; (e) Faculté des sciences, Université Mohammed V, Rabat, Morocco  
35 CERN, Geneva, Switzerland
125 Department of Physics, Oklahoma State University, Stillwater, OK, USA
126 Palacký University, RCP TM, Joint Laboratory of Optics, Olomouc, Czech Republic
127 Center for High Energy Physics, University of Oregon, Eugene, OR, USA
128 LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France
129 Graduate School of Science, Osaka University, Osaka, Japan
130 Department of Physics, University of Oslo, Oslo, Norway
131 Department of Physics, Oxford University, Oxford, UK
132 LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris, France
133 Department of Physics, University of Pennsylvania, Philadelphia, PA, USA
134 Konstantinov Nuclear Physics Institute of National Research Centre “Kurchatov Institute”, PNPI, St. Petersburg, Russia
135 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, USA
136 (a) Laboratório de Instrumentação e Física Experimental de Partículas-LIP, Lisbon, Portugal; (b) Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisbon, Portugal; (c) Departamento de Física, Universidade de Coimbra, Coimbra, Portugal; (d) Centro de Física Nuclear da Universidade de Lisboa, Lisbon, Portugal; (e) Departamento de Física, Universidade do Minho, Braga, Portugal; (f) Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada, Spain; (g) Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal
137 Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
138 Czech Technical University in Prague, Prague, Czech Republic
139 Faculty of Mathematics and Physics, Charles University, Prague, Czech Republic
140 State Research Center Institute for High Energy Physics, NRC KI, Protvino, Russia
141 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, UK
142 IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France
143 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, USA
144 (a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile; (b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
145 Department of Physics, University of Washington, Seattle, WA, USA
146 Department of Physics and Astronomy, University of Sheffield, Sheffield, UK
147 Department of Physics, Shinshu University, Nagano, Japan
148 Department Physik, Universität Siegen, Siegen, Germany
149 Department of Physics, Simon Fraser University, Burnaby, BC, Canada
150 SLAC National Accelerator Laboratory, Stanford, CA, USA
151 Physics Department, Royal Institute of Technology, Stockholm, Sweden
152 Departments of Physics and Astronomy, Stony Brook University, Stony Brook, NY, USA
153 Department of Physics and Astronomy, University of Sussex, Brighton, UK
154 School of Physics, University of Sydney, Sydney, Australia
155 Institute of Physics, Academia Sinica, Taipei, Taiwan
156 (a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia; (b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
157 Department of Physics, Technion, Israel Institute of Technology, Haifa, Israel
158 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
159 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
160 International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan
161 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
162 Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
163 Tomsk State University, Tomsk, Russia
164 Department of Physics, University of Toronto, Toronto, ON, Canada
165 (a) TRIUMF, Vancouver, BC, Canada; (b) Department of Physics and Astronomy, York University, Toronto, ON, Canada
166 Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan
167 Department of Physics and Astronomy, Tufts University, Medford, MA, USA
168 Department of Physics and Astronomy, University of California Irvine, Irvine, CA, USA
169 Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
Also at Near East University, Nicosia, North Cyprus, Mersin, Turkey

Also at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany

Also at School of Physics, Sun Yat-sen University, Guangzhou, China

Also at The City College of New York, New York, NY, USA

Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing, China

Also at Tomsk State University, Tomsk, and Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia

Also at TRIUMF, Vancouver, BC, Canada

Also at Universita di Napoli Parthenope, Naples, Italy

* Deceased