Ultra-High-Energy Cosmic Rays from Radio Galaxies

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Abstract. Radio galaxies are intensively discussed as the sources of cosmic rays observed above about $3 \times 10^{18}$ eV, called ultra-high energy cosmic rays (UHECRs). We present a first, systematic approach that takes the individual characteristics of these sources into account, as well as the impact of the extragalactic magnetic-field structures up to a distance of 120 Mpc. We use a mixed simulation setup, based on 3D simulations of UHECRs ejected by observed, individual radio galaxies taken out to a distance of 120 Mpc, and on 1D simulations over a continuous source distribution contributing from beyond 120 Mpc. Additionally, we include the ultra-luminous radio galaxy Cygnus A at a distance of about 250 Mpc, as its contribution is so strong that it must be considered as an individual point source. The implementation of the UHECR ejection in our simulation setup, both that of individual radio galaxies and the continuous source function, is based on a detailed consideration of the physics of radio jets and standard first-order Fermi acceleration. This allows to derive the spectrum of ejected UHECRs as a function of radio luminosity, and at the same time provides an absolute normalization of the problem involving only a small set of parameters adjustable within narrow constraints.

We show that the average contribution of radio galaxies taken over a very large volume cannot explain the observed features of UHECRs measured at Earth. However, we obtain excellent agreement with the spectrum, composition, and arrival-direction distribution of UHECRs measured by the Pierre Auger Observatory, if we assume that most UHECRs observed arise from only two sources: The ultra-luminous radio galaxy Cygnus A, providing a mostly light composition of nuclear species dominating up to about $6 \times 10^{19}$ eV, and the nearest radio galaxy Centaurus A, providing a heavy composition dominating above $6 \times 10^{19}$ eV. Here we have to assume that extragalactic magnetic fields out to 250 Mpc, which we did not include in the simulation, are able to isotropize the UHECR events arriving from Cygnus A. We discuss the plausibility for this scenario, and also its consequences if this is indeed the explanation for the origin of UHECR. Most notably, it would mean that the UHECR flux in our local cosmic environment might be up to about one order of magnitude above the average throughout the universe.

Keywords: ultra high energy cosmic rays, magnetic fields, radio galaxies

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1 Introduction

Fully ionized nuclei that penetrate Earth’s atmosphere with an energy above $3 \text{ EeV} \simeq 10^{18.5} \text{ eV}$ are commonly defined as Ultra-High-Energy Cosmic Rays (UHECRs). The energy of 3 EeV hereby stems from a long known feature in the cosmic-ray spectrum called “the ankle” at...
about this energy.\textsuperscript{1} Observed up to several hundred \text{EeV}, they belong to the most enigmatic phenomena of astrophysics, and for some time they were even assumed to be products of some new-physics processes \cite{2}.

The situation has somewhat changed in recent years, where the Pierre Auger Observatory (Auger) in the southern hemisphere, and the High Resolution Fly’s eye (HiRes) as well as the Telescope Array (TA) experiments in the northern hemisphere, have provided measurements with unprecedented statistics \cite{3,4,5,6}. While confirming the known spectral features with high precision, besides the ankle also a cutoff above about $10^{19.5}\text{eV}$ \cite{7,8}, the main new result of these experiments was the discovery of an increase of the fraction of heavier elements above $10^{18.3}\text{eV}$ among observed \text{UHECRs} \cite{4,9,10}. This conclusion is based on observations of the depth of the maximal energy release, $X_{\text{max}}$, in the particle shower (commonly called “air shower”) unleashed by the primary \text{UHECR}. The translation of $X_{\text{max}}$ into the nuclear mass $A$ of the primary \text{UHECR} is impeded by large fluctuations among individual air showers and the incomplete knowledge of the hadronic interaction model, but in the light of the current data the formerly common assumption that \text{UHECRs} are mostly protons is hardly tenable.

This discovery has significantly eased the explanation of the origin of \text{UHECRs}, and at the same time caused a new problem: while before there were no sources known in the universe which could accelerate particles to the highest observed energies, now there are so many different source types that it becomes a challenge to find out the real responsible ones. As noted by Hillas already in 1984 \cite{11}, there are objects in the universe spreading over 20 orders of magnitude in size and magnetic field strength, which all fulfill the condition $Z\epsilon\beta\mu R \sim 100\text{EeV}$ (hereafter referred to as the \text{Hillas-condition}) for $Z \sim 10$, as typical for the most abundant metals in the universe. Hillas grouped these sources into a diagram known after his name, and indeed, for all sources in this diagram a theory has been presented that they can not only reach the required energies, but also the observed flux — from compact sources like pulsar winds \cite{12,13}, tidal disruption of neutron stars or white dwarfs \cite{14–18}, gamma-ray bursts \cite{19–21}, to the largest structures in the universe \cite{22–24}.

The objects in the center of the Hillas diagram are dominated by structures related to active galactic nuclei (AGN). The most powerful of them sit in the centers of so-called \text{FR-II} radio galaxies, which have also been shown to be able to produce the energy, flux and spectrum observed in \text{UHECRs} \cite{25,26}. The problem with these sources is that they are very rare, and none of them is closer to Earth than 100 Mpc — which poses a problem as \text{UHECRs} arriving from such distances are constrained to observed energies below 60 EeV due to the so-called Greisen-Zatsepin-Kuzmin (GZK) effect \cite{27,28}, while \text{UHECRs} are observed at significantly higher energies. It has therefore been suggested to also include less luminous AGN, which are related to so-called \text{FR-I} radio galaxies containing less powerful and less extended jets. The consequences of this have been investigated in relation to multi-messenger signals of \text{UHECRs} \cite{29–32} or a potential anisotropy arising from the nearest \text{FR-I} galaxy, Centaurus A, at about 3 Mpc \cite{33–36}.

In fact, if the observed radio brightness of radio galaxies were a reliable tracer for their expected contribution to the \text{UHECR} flux, there should be a significant anisotropy in the \text{UHECR} arrival directions. At about 1 GHz, the brightest three radio galaxies, i.e., Virgo A (M87), Centaurus A (NGC5198), and Cygnus A, contribute about 75\% of the total flux.

\begin{footnote}
This denotation stems from the similarity of the cosmic-ray spectrum with a human leg, with a spectral steepening at about $10^{15}\text{eV}$ being called the “knee”, followed by a flattening at about $3\text{EeV}$ called the “ankle”. This anatomic analogy has been stretched by the recent discovery of a “second knee” \cite{1}, but is still commonly used.
\end{footnote}
most of it coming from Cygnus A alone. A strong anisotropy in UHECR arrival directions, however, is not observed. Early claims for anisotropy correlated with extragalactic structure [37, 38] were never confirmed to a sufficient level of statistical significance, and a joined data analysis of 2130 + 8259 events above $10^{19}$ eV from TA and Auger are on a 99% confidence level compatible with isotropy [39]. Recently, Auger reported a $5\sigma$ detection of a dipole with an amplitude of only $\sim 6.5\%$ in the UHECR arrival directions, while higher-order multipoles are still consistent with isotropy [40]. This discrepancy between expectation and observation is often used to disregard radio galaxies or other powerful, steady sources as an explanation for the origin of UHECRs, and instead numerous transient sources are proposed which would trivially explain the observed isotropy [14, 41].

An important unknown in this conclusion is the deflection of UHECRs in magnetic fields pervading extragalactic space and our Galaxy. The effect of extragalactic magnetic fields (EGMFs) can be simply estimated as follows: If we assume that a magnetic field with strength $B_0$ is homogeneous over a coherence length $\ell$, a particle passing with energy $E$ and charge $Z$ will be deflected on average by an angle $\theta_0 \approx Z e B_0 \ell / 2 E$, and after a path of length $D$ undergoing random deflections the total deflection angle will be $\theta \sim \theta_0 \sqrt{D/\ell}$. For a typical UHECR as derived from observations, i.e., $E \sim 30$ EeV, $Z \sim 10$, and a coherence length $\ell \sim 1$ Mpc comparable to the typical distance between galaxies, and $B_0 \sim 1\,\text{nG}$, which is the current upper limit on primordial magnetic fields [42], we find $\theta_0 \sim 3^\circ$. Thus, a particle traversing a distance of 100 Mpc would be deflected by about $30^\circ$, a quite significant amount. In fact, if these magnetic deflections are combined with the various energy-loss processes of UHECRs during propagation, one finds that for this simple setup and reasonable values of $B_0$, no UHECRs can reach us from beyond a "magnetic horizon" of about 300–500 Mpc, which is nearly constant over the UHECR energy regime [43]. That means, sources outside this horizon will not contribute to the observed UHECR flux, and UHECRs from point sources near to it will reach us nearly isotropized. On the other hand, if UHECR protons exist at the highest energies $E \sim 100$ EeV, and considering that they would need to arrive from not more than $D \sim 30$ Mpc due to the GZK effect, one would expect deflections of not more than a few degrees.

This simple scenario demonstrates that deflections in EGMFs are expected to be significant. For a quantitative estimates, however, it is not useful, because EGMFs are expected to be correlated with the observed large-scale structure of matter. Various work has been done to estimate the effect of realistic EGMF structure on UHECR propagation [44–46]. One of the currently most sophisticated descriptions of the EGMF in the local universe is given by Dolag et al. [47], where constrained simulations of large-scale structure formation have been used to derive the EGMF from cosmological magnetic fields grown in a magnetohydrodynamic amplification process. The results of this simulation are straightforward to implement in the publicly available code CRPropa3 [48], and therefore easy to apply to UHECR propagation.

Apart from EGMFs, another important impact on arrival directions comes from the Galactic magnetic field (GMF). Unlike EGMFs, which are supposed to just scatter UHECR events around the line of sight to their sources, the largely coherent structure of the GMF can deflect them systematically, in addition to a random scattering induced by a small-scale turbulent field of similar strength. Simulations of UHECR trajectories through one of the currently most advanced GMF models [49, 50] have shown that, for fiducial values of UHECRs used above, both systematic and random deflections in the GMF should be again of the order of $10–30^\circ$ [51, 52]. This effect depends on direction, but is obviously independent of the distance of an extragalactic source, thus affects very nearby sources in the same way.
These considerations show that a model proposing dominant and powerful sources like radio galaxies cannot be trivially abandoned on the base of observed UHECR isotropy. Rather, it appears to be time to review radio galaxies as putative UHECR sources in the light of current data on the UHECR spectrum, composition and arrival-direction distribution. In this paper, we present results of UHECR propagation simulations comprising three components: (a) A 3D structure of radio galaxies and EGMF within a radius of 120 Mpc; (b) a continuous source function derived from a luminosity function of radio galaxies for the contributions from beyond 120 Mpc; (c) a contribution of the powerful radio galaxy Cygnus A, which is near the magnetic horizon and is expected to deliver the dominant contribution to UHECRs due to its extreme radio power and brightness. All simulations are carried out with CRPropa3.

The paper is organized as follows: In Sect. 2 we outline the physics of radio galaxies as cosmic-ray accelerators and the derivation of the continuous source function. In Sect. 3 we explain the simulation setup and describe various tests performed to ensure its consistency. In Sect. 4 we present the simulation results for various choices of best-fit parameters. In Sect. 5 we discuss the implications of our results, and we conclude in Sect. 6, and given an outlook on future work on the subject.

2 Radio galaxies as UHECR accelerators

2.1 Phenomenology of radio galaxies

The term radio galaxy (RG) is commonly used for a subclass of AGN which form extended plasma jets, called radio-loud AGN [53], most likely due to accretion onto a strongly rotating supermassive black hole [54]. It has to be distinguished from the more general term of radio-emitting galaxies, to which also other objects, like starburst galaxies, belong [55].

According to a classification introduced in 1974 by Fanaroff and Riley [56], radio galaxies fall into two major classes: FR-I RGs, in which the jets are terminating within the galactic environment on scales of a few kiloparsec, and which are observed as center brightened radio objects; and FR-II RGs, where the jets extend on scales of \( \gtrsim 100 \) kpc deep into extragalactic space, and terminate there in two extremely bright hot spots, causing a double-source appearance in radio observations. This morphological distinction strongly correlates with radio power\(^2\): Sources with \( P_{178} \lesssim 2 \times 10^{25} \) W Hz\(^{-1}\) sr\(^{-1}\) tend to be FR-I galaxies, while sources with \( P_{178} \gtrsim 2 \times 10^{25} \) W Hz\(^{-1}\) sr\(^{-1}\) usually have FR-II morphology — although there are notable exceptions to this, like the very powerful FR-I galaxy Hydra A at about 230 Mpc distance.

Plasma jets from AGN are usually assumed to start with relativistic velocities (bulk Lorentz factor \( \Gamma \sim 10 \)), which is required for the explanation of highly-variable blazar emission in either leptonic or hadronic scenarios [57]. At larger distances from the AGN, however, jets become sub-relativistic, with estimates for the dimensionless velocity of the jet \( \beta_{jet} = v_{jet}/c \) ranging from \( \lesssim 0.1 \) for FR-I RGs [58], to \( \gtrsim 0.3 \) for FR-II RGs [59]. Although cosmic-ray acceleration may take place in the relativistic part of the jet [31, 60], leading to highly anisotropic UHECR ejection via neutrons produced in photohadronic interactions [29, 34, 61, 62], we will assume here that the main sites of UHECR acceleration located in or at the

\(^2\)Here and in the following we denote the spectral power of a radio source at a frequency \( \nu \), measured in Watt/Hz or Watt/Hz/sr, with \( P_\nu \), where a three digit number without decimal point in the subscript refers to a frequency in MHz, while a number with decimal point refers to a frequency in GHz. With \( L_\nu = \nu P_\nu \), we denote a radio luminosity, measured in erg/s, with the same conventions for frequency notation.
end of the mildly relativistic extended part of the jets, which would cause a largely isotropic emission.

The distribution of radio galaxies in the universe is described by the radio-luminosity function (RLF) of radio-loud AGN. An expression for the local RLF has been given by Mauch and Sadler [63] — hereafter referred to as MS07 — as\(^3\) \(\phi(P_{1.4}) = dN_{RG}/dM_K dV\), which we translate to the luminosity at \(L_{1.1} \equiv 1.1 \text{ GHz} P_{1.1}\) via the spectral behavior of \(P_\nu \propto \nu^{-0.6}\), i.e., \(P_{1.4} = 0.79 L_{1.1} \text{ GHz}^{-1}\). Since the absolute \(K\)-band magnitude depends on the radio luminosity \(L\) according to \(dM_K = 1.09 \ln L\), the total number \(N_{RG}\) of radio-loud AGN per volume element \(dV\) and luminosity interval \(dL\) can be written as

\[
\frac{dN_{RG}}{dV dL_{1.1}} = \frac{C_{RG}}{L_{1.1} \left( (L_{1.1}/L_\star)^\alpha + (L_{1.1}/L_\star)^\beta \right)^3},
\]

with the parameters \(C_{RG} \approx 2.9 \times 10^{-6} \text{ Mpc}^{-3}\) and \(L_\star \approx 4.9 \times 10^{40} \text{ erg/s}\), both determined within a factor of 2, and the power-law indices \(\alpha = 1.27 \pm 0.18\) and \(\beta = 0.49 \pm 0.04\). This is essentially a broken power law with a break at \(L_{1.1} = L_\star\), where most of the radio power is emitted by RGs with \(L_{1.1} \sim L_\star\), which also corresponds quite well to the radio luminosity separating FR-I and FR-II RGs. The small value of \(C_{RG}\) makes clear that powerful radio galaxies are very rare: For sources like Centaurus A, at a distance of \(\sim 3 \text{ Mpc}\) the closest radio galaxy with \(L_{1.1} \sim 5 \times 10^{39} \text{ erg/s}\), we would expect on average 1 within a sphere of 30 Mpc radius, and for the brightest radio galaxy in the sky, Cygnus A \(L_{1.1} \approx 2 \times 10^{44} \text{ erg/s}\) we have a number density of \(\sim 0.3 \text{ Gpc}^{-3}\), thus roughly a few hundred similarly powerful RGs in the entire observable universe!

### 2.2 From radio luminosity to cosmic-ray luminosity and maximal rigidity

Radio galaxies have been proposed UHECR accelerators because their jets are suitable sites for first-order Fermi acceleration [64–68] at their discontinuities, i.e., in internal or external shocks [25, 60, 69, 70], or by shear at the jet boundary [71]. First-order Fermi acceleration is a self-similar process producing power-law spectra with a canonical spectral index of \(a = 2\) at ideal discontinuities (e.g., strong shocks with large Mach-numbers). Relativistic effects can modify this to indices in a range of \(1.5 < a < 2.3\) [72, 73]. Acceleration takes place on a timescale

\[
\tau_{acc} = \frac{f_{\text{diff}} r_L}{c \beta_{sh}^2}
\]

for cosmic-ray particles with a Larmor radius \(r_L = E/Z eB = R/B\), where \(R\) is called the particle rigidity, \(c \beta_{sh}\) is the shock or shear velocity, and \(f_{\text{diff}} \sim 3\) encapsulates all details of the diffusion process in a strongly turbulent magnetic field for standard geometries [72].

As the Fermi acceleration process is only dependent on rigidity, abundances \(f_i\) of different nuclear species with charge \(Z_i\), measured at a given rigidity \(R\), remain unchanged in the acceleration process. The abundances at injection, however, will depend on the environment or details of the injection process. Regarding the environment, we would expect that the interior of radio-galaxy jets contains a similar material mix as the accretion disk around the supermassive black hole, from which the jet likely feeds [74, 75], and which can be assumed to have abundances close to solar (denoted \(f_\odot\)). For FR-II galaxies where these jets end far outside their host galaxy, it should be just this material which is injected into the acceleration. For FR-I galaxies, however, the jets dissipating inside the galactic environment may pick

\(^3\)We use the RLF of Eq. (6) but revise the wrong notation \((P_\nu/P_\star)\) to \((P/P_\star)\).
up material from gaseous shells and galactic cosmic rays, which can lead to a much heavier composition\cite{76}. Additionally, the injection of a particle into the Fermi process may depend on other parameters than just the rigidity of the fully ionized nucleus. As a working hypothesis, we may assume that particles are injected at some minimal rigidity $\hat{R}$ with abundances $f_i = f_{\odot}Z_i^q$, which reduces the heaviness of the composition to a single parameter $q$.

To estimate the maximal rigidity which can be obtained in radio galaxies, we need to estimate the strength of the magnetic field $B$. For this we need to derive the total kinetic power of the jets from observational quantities. A relation between the jet power $Q_{\text{jet}}$ and the extended radio luminosity $P_{151}$ at 151 MHz has been given by Willott et al.\cite{77} using the narrow emission-line data from the 7C Redshift Survey. In general, the jet power depends on the projected linear size $D$ of the source according to $Q_{\text{jet}} \propto D^{-4/7+\gamma/2}$ in the case of an electron number density profile that scales with the distance to the power of $-\gamma$ at radii $\sim 100$ kpc. Further, it is pointed out that a very weak $D$-dependence is obtained in the case of a realistic value of $\gamma \approx 1.5$, so that the jet power can be estimated by

$$Q_{\text{jet}} = f^{3/2} Q_{\text{jet},0} \approx 3 \times 10^{45} f^{3/2} \left( \frac{P_{151}}{10^{28} \text{W Hz}^{-1} \text{sr}^{-1}} \right)^{6/7} \text{erg s}^{-1}. \quad (2.3)$$

with $1 \lesssim f \lesssim 20$ resulting from the uncertain efficiency of converting jet internal energy into observable radio luminosity (see also\cite{78, 79} for additional discussion on the normalization factor). Furthermore, the minimal energy condition for $Q_{\text{jet}}$ is given by\cite{80}

$$Q_B \approx \frac{3}{4} (Q_e + Q_{\text{cr}}), \quad (2.4)$$

where $Q_B$, $Q_e$ and $Q_{\text{cr}}$ denote the total power in the magnetic field, relativistic electrons, and cosmic rays, respectively. Assuming that these components dominate the energy density of the jet, we then have $Q_{\text{jet}} \approx Q_B + Q_e + Q_{\text{cr}}$, hence $Q_{\text{jet}} \approx \frac{7}{5} Q_B$. The magnetic field power of the jet with a bulk velocity of $c\beta_{\text{jet}}$ and a size $r$ is determined by $Q_B = c \beta_{\text{jet}} \pi r^2 B^2/8\pi$, so that we obtain

$$B r \approx \sqrt{\frac{24 Q_{\text{jet}}}{7c \beta_{\text{jet}}}} \quad (2.5)$$

and by equating the acceleration time (2.2) to a characteristic loss time scale $t_{\text{loss}} = f_{\text{loss}} r/\beta_{\text{sh}} c$, where $f_{\text{loss}} \sim 1$ is a fudge factor adapting the speed of expansion or diffusive particle escape downstream of the shock to the shock velocity; essentially we suppose that there is one dominant velocity scale in the system. Absorbing this factor into $f_{\text{diff}} \sim 3$, the maximal cosmic ray rigidity yields

$$\hat{R} = \frac{E_{\text{max}}}{Ze} = \beta_{\text{sh}} f_{\text{diff}} B r = g_{\text{acc}} \sqrt{g_{cr} Q_{\text{jet},0}} / c = g_{\text{acc}} e \sqrt{\alpha f g_{cr} Q_{\text{jet},0}}. \quad (2.6)$$

The last writing is useful when calculating energies, $Ze\hat{R}$, as all electromagnetic constants disappear in the dimensionless Sommerfeld constant $\alpha = e^2/\hbar c \approx 1/137$\cite{34}. Here, we introduced the dimensionless coefficients

$$g_{cr} = \frac{4}{7} f^{3/2} \quad \text{and} \quad g_{\text{acc}} = \sqrt{\frac{6 \beta_{\text{sh}}^2}{f_{\text{diff}}^2 \beta_{\text{jet}}^2}}. \quad (2.7)$$
The cosmic ray luminosity of a source in steady state is equal to the cosmic ray fraction of the jet power, \( Q_{\text{cr}} \), so that we can use these terms synonymous and write for minimal energy conditions with \( Q_{\text{cr}} \gg Q_{\text{c}} \),

\[
Q_{\text{cr}} \simeq 1.3 \times 10^{42} g_{\text{cr}} \left( \frac{L_{1,1}}{L_{\ast}} \right)^{6/7} \frac{\text{erg}}{\text{s}}, \tag{2.8}
\]

where we converted to the quantities used to describe the RLF, again assuming a typical \( \nu^{-0.6} \) spectrum; note that \( Q_{\text{cr}} = \frac{4}{7} Q_{\text{jet}} \) and \( \dot{R}(Q_{\text{cr}}) = g_{\text{acc}} \sqrt{Q_{\text{cr}}/c} \). Supposing the typical shock and jet velocities \( \beta_{d\gamma} \sim \beta_{\text{jet}} \sim 0.1 \) in extended jets of radio galaxies, we obtain \( 0.01 \lesssim g_{\text{acc}} \lesssim 1 \), and as we assume \( Q_{\text{cr}} \gg Q_{\text{c}} \), the jet power normalization factor \( f \) will rather be at the high end of its range, hence we assume \( 5 \lesssim g_{\text{cr}} \lesssim 50 \). If we distinguish between species of cosmic ray nuclei of charge \( Z_i \) and assume that they represent a number fraction \( f_i \) of mixed composition cosmic rays accelerated all over the same rigidity regime, the fractional contribution of this species to the total cosmic ray power will be \( f_i Z_i \).

Consequently, even for optimal conditions (\( g_{\text{acc}} = 1 \) and \( g_{\text{cr}} = 50 \)), only sources with radio luminosity \( L_{1,1} > 0.028 L_{\ast} Z^{-7/3} \) can contribute cosmic rays with charge number \( Z \) to the UHECR regime, i.e., an energy of 3 EeV or higher. To reach the highest energies, \( E \gtrsim 100 \text{ EeV} \), this requirement becomes \( L_{1,1} \gtrsim 100 L_{\ast} Z^{-7/3} \). Thus, only the most powerful radio galaxies can accelerate protons to the highest energies, and for those radio galaxies emitting the most power per volume (\( L_{1,1} \approx L_{\ast} \)), we need to require \( Z > 7 \) to obtain this. Interesting is also the relation to \( Q_{\text{cr}} \), which implies

\[
Q_{\text{cr}} \sim 10^{44} \text{erg/s} \quad \text{for} \quad \dot{R} \sim 10^{20} \text{ V} \tag{2.9}
\]

This relation has been first pointed out by Lovelace [81] for a dynamo model of radio galaxies, but as we see here, it follows in general from combining the Hillas limit with minimal energy conditions in radio galaxy jets. We will refer to the general correspondence between cosmic ray power and maximal rigidity expressed by Eq. (2.6) as the Lovelace-Hillas relation henceforth.

### 2.3 Continuous source function of UHECR from radio galaxies

Lacking the possibility to do efficient Monte Carlo calculations, early work on UHECR origin, e.g. [19, 25], used continuous source functions, derived from average properties of the sources considered. In this approach, cosmic rays are injected with a homogeneous rate per unit volume throughout the universe, where potentially dependences cosmological evolution (i.e., redshift) are allowed. The flux is then obtained by integration over the comoving volume, using modification factors to consider propagation effects [82, 83]. Radio-galaxy source functions, however, did so far not consider the Lovelace-Hillas relation, which certainly has a major impact on the result as it predicts a power-dependent maximal rigidity of radio galaxies.

The local continuous source function, \( \Psi_{i,0}(R) \), derives from the local luminosity function of radio galaxies and the spectrum of individual sources as

\[
\Psi_{i,0}(R) \equiv \frac{\text{d}N_{\text{cr}}(Z_i)}{\text{d}V \text{d}R \text{d}t} = \int_{Q_{\text{cr}}}^{Q_{\text{c}}} \frac{\text{d}N_{\text{RG}}}{\text{d}V \text{d}Q_{\text{cr}}} \frac{\text{d}Q_{\text{cr}}}{\text{d}Q_{\text{c}}} \tag{2.10}
\]

where \( \text{d}N_{\text{cr}} \) is the number of cosmic ray particles of charge number \( Z \) emitted at rigidity \( R \), per energy bin \( Z \in [R, \dot{R}(Q_{\text{cr}})] \), volume element \( \text{d}V \), time interval \( \text{d}t \). \( S_i(R, \dot{R}(Q_{\text{cr}})) \equiv \text{d}N_{\text{cr}}(Z_i)/\text{d}R \text{d}t \) is the spectrum of cosmic ray nuclei species \( i \) with charge number \( Z_i \), emitted by a source with
total cosmic ray power per charge number, \(Q_{cr,i} \equiv Q_{cr}(Z_i) = f_i Z_i Q_{cr}/\bar{Z}\), up to a maximal rigidity \(\tilde{R}(Q_{cr})\), and \(dN_{RG}/dV\,dQ_{cr}\) is the number of radio galaxies per volume per power bin. The limits of integration are the smallest, \(\tilde{Q}_{cr}\), respectively largest, \(Q_{cr}\), CR powers we need to consider. To convert rigidity into energy we use the average cosmic ray charge number at a given rigidity, \(\bar{Z} = \sum_i f_i Z_i\), where the sum goes over all elements we consider with modified solar injection abundances \(f_i\), at a constant injection rigidity \(\tilde{R} \approx 10\, G\)V corresponding to the peak of the Galactic CR flux.

To solve this integral analytically, we make the simplifying assumption that individual source spectra are power laws with a sharp cutoff at \(\tilde{R}(Q_{cr}) = g_{acc}\sqrt{Q_{cr}/c}\), i.e.

\[
S_i(R, \tilde{R}(Q_{cr})) = \nu_i(a) Q_{cr} \left(\frac{R}{\tilde{R}}\right)^{−a} \Theta(\tilde{R}(Q_{cr}) − R),
\]

where \(\Theta(x)\) is the Heaviside step function. From the requirement

\[
Q_{cr,i} = \frac{f_i Z_i Q_{cr}}{\bar{Z}} = c Z_i \int_{\tilde{R}}^{\tilde{R}(Q_{cr})} dR\,R\,S_i(R, \tilde{R}(Q_{cr})),
\]

we obtain the spectral normalization correction \(\nu_i(a)\) as

\[
\nu_i(a) = \frac{f_i}{c Z_i R^2} \times \begin{cases} 
2 − a / (2 − a) − 1, & \text{for } a \neq 2 \\
1/\ln \rho_{cr}, & \text{for } a = 2
\end{cases}
\]

where we introduced the cosmic ray dynamical range \(\rho_{cr} \equiv \tilde{R}(Q_{cr})/\tilde{R}\). The radio luminosity function \(\Phi_{RG}(Q)\) provides the number of radio galaxies per volume per logarithmic power bin, so that \(dN_{RG}/dV\,dQ_{cr} = Q_{cr}^{−1} \Phi_{RG}(Q_{cr})\). Using the MS07 luminosity function dependent on the CR power where

\[
P_{1.4} = 3.31 \times 10^{28} \left(\frac{Q_{cr}}{3 \times 10^{28} \, g_{cr} \, W}\right)^{7/6} \, \text{W Hz}^{−1},
\]

for a fiducial radio spectrum \(\propto \nu^{−0.6}\), we obtain

\[
dN_{RG}/dV\,dQ_{cr} = \frac{7}{6} \Phi_{RG}(P_{1.4}(Q_{cr})) = \frac{7}{6} \frac{C_{RG}}{g_{cr} \, Q_*} \left[\left(\frac{Q_{cr}}{g_{cr} \, Q_*}\right)^{α'} + \left(\frac{Q_{cr}}{g_{cr} \, Q_*}\right)^{β'}\right]^{−1},
\]

where \(C_{RG} \approx 3 \times 10^{-6}\,\text{Mpc}^{−3}\) is the source density parameter and \(α' = \frac{7}{6} α + 1 \approx 2.48\) as well as \(β' = \frac{7}{6} β + 1 \approx 1.55\) are related to the indices for the radio luminosity function given below Eq. (2.1). Further, we introduced \(Q_* \approx 1.3 \times 10^{42}\,\text{erg/s}\), according to Eq. (2.8), as the CR power related to the break in the radio luminosity function. We can then find an approximate analytical solution to Eq. (2.10) as

\[
Ψ_{i,0}(R) \approx \frac{7}{6} \frac{C_{RG}}{g_{acc}} f_i \tilde{ν}(a) c Z_i \left(\frac{R}{\tilde{R}_*}\right)^{−a} \times \left\{F_{αβ} \left(\frac{R^2}{\tilde{R}_*^2}\right) - \tilde{F}_{αβ} \left(\frac{Q_{cr}}{g_{cr} \, Q_*}\right)\right\},
\]

for the three simplifying cases

- \(\tilde{ν}(a) = 2 − a\); \(\tilde{α} = α' − a/2\); \(\tilde{β} = β' − a/2\); for \(a < 2\), \(a \neq 2\)
- \(\tilde{ν}(a) = 1/\ln \rho_{cr}\); \(\tilde{α} = α' − 1\); \(\tilde{β} = β' − 1\); for \(a \approx 2\)
- \(\tilde{ν}(a) = (a - 2) ρ_{cr}^{2 - a}\); \(\tilde{α} = α' − 1\); \(\tilde{β} = β' − 1\); for \(a > 2\), \(a \neq 2\)
where we introduced $R_* = g_{\text{acc}} \sqrt{g_{\text{cr}} Q_*/c}$, as well as the corresponding dynamic range $\rho_* = R_*/\dot{R}$, and
\[
F_{\alpha,\beta}(x) = \frac{x^{1-\alpha}}{\alpha - 1} \, {}_2F_1\left[1, \frac{\alpha - 1}{\alpha - \beta}; \frac{\alpha - \beta}{\alpha - \beta}; -x^{\beta - \alpha}\right]
\]
with $_2F_1$ denoting the hypergeometric function. The first respectively last case corresponds to $\rho_*^2 - a \gg 1$ for $a < 2$ (and is thus included in the integration), and $\rho_*^2 - a \ll 1$ for $a > 2$, which can both be assumed even for moderate deviations of the spectral index from $a = 2$ as $\rho_* \approx 10^9$. For $a \approx 2$, we simply set $\dot{R}(Q_{\text{cr}}) \approx R_*$ constant in the slowly varying logarithm. $Q_{\text{cr}}$ is the highest jet power considered, and should correspond to the most powerful FR-II galaxies, like Cygnus A; we set it to $10^{45} \, g_{\text{cr}} \, \text{erg/s}$. Assuming higher values of $Q_{\text{cr}}$ would not significantly change the resulting production rate.

The function $\Psi_{i,0}(R)$ is the local continuous source function as it is derived from a radio luminosity function determined in the local universe ($z < 0.3$). To extend it to larger redshifts, we use the approximation
\[
\Psi_i(R, z) \approx \Psi_{i,0}(R) (1 + z)^{m-1},
\]
where $m$ is called the source evolution index. This writing disregards the (likely) dependence on characteristic parameters like $L_*$, $\alpha$ and $\beta$ on $z$. However, it has been shown that different redshift dependent RLFs with roughly the same mean index $m$ produce very similar UHECR spectra [25, 84], so we adopt Eq. (2.17) with the range $2 < m < 4$.

With this, the free parameters of the model are $f_i$, $a$, $g_{\text{cr}}$, $g_{\text{acc}}$, and $m$, with quite limited plausibility ranges as discussed the previous sections (note that $R_*$ and $\rho_*$ depend on $g_{\text{cr}}$ and $g_{\text{acc}}$). This means, that there is little margin to adjust the total normalization of the model. Setting canonical values of $a = 2$, $g_{\text{cr}} = 50$, $g_{\text{acc}} = 0.5$ and using solar abundances $f_i = f_\odot$, we get a total UHECR production rate density of
\[
(\Psi)_{\text{UHE}} = \sum_i \int_{10^{19}\text{eV}}^{10^{20}\text{eV}} \text{d}E \frac{\Psi_{i,0}(E/Z_i e)}{Z_i e} \approx 4 \times 10^{44} \, \text{erg Mpc}^{-3} \text{year}^{-1}
\]
comparable to the value obtained by Waxman [85] as required to explain the observed UHECR flux at Earth. As the total normalization at $R \gtrsim R_*$ is approximately $\propto g_{\text{cr}}^{1.5} g_{\text{acc}}$ (see below), this estimate presents and upper limit regarding these parameters. The source evolution parameter $m$ is mostly relevant below the ankle and cannot significantly modify the UHECR normalization. Looking at $\nu(a)$, up to an order of magnitude increase is possible for $a < 2$, while for $a > 2$ every increase of $a$ by 0.1 drops the normalization by almost one order of magnitude. Thus, spectral indices significantly steeper than $a = 2$ are strongly disfavored.

Let us now analyze the shape and characteristic energies connected to this spectrum. The term in curly brackets of Eq. (2.16) consists of two modified hypergeometric functions, where the second is simply a (usually very small) constant. In the case of a small value of $R$, the first term runs asymptotically to a constant $\sim 1$, so at low rigidities we see an undistorted $R^{-a}$ power law. Towards large values of $R$, it approaches asymptotically a power law $R^{2-2\tilde{a}}$, so that in total we obtain a production rate density $\propto R^{4-2\tilde{a}-a}$ for $a \gtrsim 2$, which approaches a limit $\propto R^{2-2\tilde{a}}$ for $a < 2$; thus for flat spectra and rigidities above the break, the spectral index is only determined by the spectral slope of the RLF, and is close to $E^{-3}$. This is a result of the sliding-cutoff effect introduced by the Lovelace-Hillas relation, and a

\[^{4}\text{This solution is only valid for } 0 < \tilde{\beta} < 1 < \tilde{\alpha}, \text{ which is the case here.}\]
Figure 1: UHECR production rate density (colored, solid lines) and corresponding \( \langle \ln(A) \rangle \) (black, dashed line) using \( a = 2, \ g_{ct} = 50, \ g_{acc} = 0.5 \), as well as an initial solar abundance \( f_i = f_{\odot} \). Note that the sharp cutoffs in the spectra and also the “stairs” in the \( \langle \ln(A) \rangle \) curve are artefacts from using a Heaviside cutoff in the derivation of the continuous source function.

Putting everything together, this means that for light nuclei the spectral index in the UHECR regime is about \( E^{-3} \), thus too steep to explain the observed UHECR spectrum (as propagation will make the spectrum even steeper). For heavy nuclei, \( \epsilon Z R_s \) may be large enough to allow UHECR spectra consistent with the data, but this would mean that they have to dominate in the entire UHECR regime, which is inconsistent with the measured composition. We can therefore already get to the preliminary conclusion that the cosmic ray spectrum averaged over the radio luminosity function cannot explain the known UHECR data. The only way to save the day is therefore to consider the contribution of powerful individual sources in the local universe. The estimation of their impact is subject of our simulation setup, which is explained in the following.

3 Simulation setup

The Monte Carlo based algorithm of CRPropa has been designed to describe the propagation of UHECR nuclei in extragalactic background light and cosmic magnetic fields [86, 87]. The recently published version CRPropa3 has been redesigned for better performance, and includes
improved methods and new functionalities [48]. It thus represents the ideal framework to explain the energy spectrum, composition and arrival directions of UHECRs (e.g. [45, 46, 88]).

In CRPropa3, the trajectory of the particles is determined by the Lorentz force according to the Cash-Karp algorithm, an adaptive Runge-Kutta algorithm of the fourth order. In doing so, energy loss due to photonuclear interactions (e.g. pair production, photo-pion production and photo-disintegration) with the cosmic microwave background (CMB) and the UV/optical/IR background (IRB) is taken into account. Furthermore, the decay of unstable nuclei is considered as well. The necessary photon fields of the CMB and the IRB, where the IRB model of Gilmore et al. [89] is used, are already implemented. The uncertainties in the simulation results due to poorly known physical quantities as the IRB and the photo-disintegration cross sections should not affect the results in a significant way [90].

In order to obtain suitable statistics, an observer sphere with a radius of \( r_{\text{obs}} = 100 \text{kpc} \) is chosen. This size also ensures, that no significant deformation of the higher multipole moments of the arrival directions occurs, as recently shown by Dundović and Sigl [91]. Thus, the UHECR candidates are only propagated on extragalactic scales at first. Afterward, the Galactic deflection of the candidates in the Milky Way is easily applied using the galactic lensing module. Here, the GMF model of Jansson & Farrar [49] — hereafter referred to as the JF12 model — is adopted to trace the particles from the observer sphere through the Galactic magnetic field to Earth. This treatment neglects energy losses inside the Galaxy, which is justified as Galactic trajectory lengths are small compared to the mean free path of interactions. Further details and explanations of the simulation setup are given in Appendix B. Consequently, the physical constituents that still need to be applied to the simulation setup are (i) the EGMF and (ii) the source setup (e.g. spatial distribution, energy spectrum, composition).

3.1 The extragalactic magnetic field

We apply the local EGMF model from Dolag et al. [47] to define the magnetic field up to a distance of about 120 Mpc from Earth. Here, a uniform magnetic seed field with \( B_0 = 2 \times 10^{-12} \text{G} \) is treated in the framework of ideal MHD in a constrained structure-formation simulation of the local universe (see [92] for more details). The arbitrarily chosen initial orientation of the uniform seed magnetic field gets randomized in dense environments, e.g. galaxy clusters, but in low density regions the initial orientation of the seed field is conserved. Nevertheless, it is shown by Dolag et al. that there is no statistical difference in the distribution of deflected UHE protons with respect to the chosen orientation of the seed field. Magnetic-field strengths of the order of a few nG and higher are only found in dense environments like galaxy clusters, so significant deflections are predominantly expected from UHECRs that pass through dense environments. On large scales these environments are rare and not surrounding the Milky Way, so that the sky is dominated by areas showing small deflections [47]. We note, however, that other simulations obtain higher cumulative filling factors of strong fields [46].

The advantage of using the Dolag et al. field is that it represents the true structure of the local universe, which is essential when comparing with a source distribution based on astronomical catalogs. Therefore, the Dolag et al. EGMF is sampled in this simulation with a high resolution of 14.6 kpc and subsequently stored in a multi-resolution grid that has been developed by Müller [93] using a relative error of 0.4 and an absolute error of \( 10^{-14} \).
3.2 Characteristics of the source sample of UHECRs

As our implemented EGMF structure extends only to 120 Mpc, we need to distinguish two groups of sources: Sources up to a distance of 120 Mpc from Earth are called local RGs, and are treated as individual sources in 3D simulations making full use of the EGMF structure. Sources beyond 120 Mpc are called non-local RGs, and are described by the continuous source function derived in Sect. 2.3 which is used as input for 1D simulations. Continuity requirements imply that the average source features of local and non-local RGs are the same in particular in the transition regime around 120 Mpc, which is carefully checked (see Appendix A.2). For all sources we assume a temporally constant and isotropic outflow of UHECRs described by a single power law with the same spectral index $a$, limited to a range $1.7 \leq a \leq 2.2$ according to the expectations for first order Fermi acceleration. We explicitly simulate the propagations of the nuclear species $^1$H, $^4$He, $^{12}$C, $^{14}$N, $^{16}$O, and $^{56}$Fe, i.e., $Z_i \in \{1, 2, 6, 7, 8, 26\}$. As nuclei in the range $Z = 3, 4, 5$ have low abundances and are quite unstable to spallation in high-energy propagation, this range can be considered complete up to oxygen. Reducing the heavy elements beyond CNO to iron, however, is a simplification chosen just for convenience, and we need to discuss potential artifacts introduced by it in our results.

For the chemical composition at a given rigidity $R$ we distinguish two cases: (1) the light composition scenario assuming solar abundances (in particle number, e.g. [94]), given by

$$f_\odot = \begin{cases} 
0.922, & \text{for } H, \\
0.078, & \text{for } He, \\
2.3 \times 10^{-4}, & \text{for } C, \\
6.7 \times 10^{-5}, & \text{for } N, \\
4.9 \times 10^{-4}, & \text{for } O, \\
2.7 \times 10^{-5}, & \text{for } Fe, 
\end{cases} \quad (3.1)$$

with a strong dominance of $^1$H and $^4$He nuclei; (2) the heavy composition scenario, where we assume modification of the abundances at injection to Fermi acceleration as $f_i \propto f_\odot Z^q$ with $1 \lesssim q \lesssim 2$. For the limiting case of $q = 2$ we obtain abundances

$$f_\bullet = Z^2 f_\odot = \begin{cases} 
0.706, & \text{for } H, \\
0.238, & \text{for } He, \\
0.006, & \text{for } C, \\
0.003, & \text{for } N, \\
0.033, & \text{for } O, \\
0.014, & \text{for } Fe. 
\end{cases} \quad (3.2)$$

In this case, the iron contribution to the UHECR source spectrum is comparable to that of hydrogen (note that the abundances compared at a given energy $E = ZeR$ are further enhanced by $Z^{a-1}$). We treat these two cases as extremes in many explanatory examples, but for the fit to the data we optimize the value of $q$ in its given range. We note that other composition enhancements would be possible depending on the unknown details of the injection process, but they are not considered here unless explicitly noted.

Based on the characteristics of the Monte Carlo method, most source characteristics are changeable after the simulation by re-weighting (e.g. [86]). Further details on the weights we use are given in Sect. 3.4.
3.2.1 Local sources

To obtain a complete sample of radio galaxies powerful enough to contribute to the UHECR spectrum we use the catalog of van Velzen et al. [95] — hereafter referred to as vV12 catalog. In total it contains 575 radio-emitting galaxies with a flux greater than 213 mJy at 1.4 GHz. We note, however, that the catalog is based 2MASS redshift survey [96], which introduces other selection criteria, in particular a limiting infrared $K$-band magnitude of 11.75. This excludes preferentially distant powerful radio galaxies for which the radio-to-infrared luminosity ratio is larger than average. An example is the powerful FR-II galaxy Pictor A at 125 Mpc distance, which is among the ten brightest radio galaxies in the sky but not contained in the vV12 catalog. Within 120 Mpc, however, where we actually use individual sources, the catalog is fairly complete at the bright end, as cross checks with searches in the NASA/IPC extragalactic database showed. Excluding 52 star-forming galaxies, we are finally left with 121 radio galaxies, predominantly of FR-I type, which will subsequently be called the visible, local RGs. In addition to the spatial coordinates $(x_0, y_0, z_0)$, the catalog also provides the radio luminosity $L$ at 1.1 GHz of each RG and the angular distance $\delta_d$ between the galaxy and the radio matches. But still, the catalog only covers 88% of the sky and misses a multitude of low-luminous, distant RGs due to its flux limit.

To complete the source sample, we use the MS07 RLF (2.1) and consider spherical shell elements with a thickness $\bar{d} - \bar{d}$ at small redshift and a radio luminosity between $\bar{L}$ and $\bar{L}$. Hence, we obtain the expected total number of radio-loud AGNs in the fraction $P_V$ of the volume by

$$N_{RG} = C_{RG} \frac{4\pi}{3} \left( d^3 - \bar{d}^3 \right) P_V \int_{\bar{L}}^{\bar{L}_1} dL L^{-1} \left[ (L/L_*)^\alpha + (L/L_*)^\beta \right]^{-1}. \tag{3.3}$$

with the best-fit parameter $C_{RG} \approx 3.162 \times 10^{-6}$ Mpc$^{-3}$, which has been determined by fitting the Eq. (3.3) with $P_V = 0.88$ to the known number of radio-loud sources from the vV12 catalog. All other parameters are taken as the best-fit values given by MS07 [63]. Since the integrand of Eq. (3.3) hardly depends on the upper integration limit $\bar{L}_1$, we will set $\bar{L}_1 = 10^{45}$ erg/s in the following. Due to the flux limitation of the catalog we had to increase the lower luminosity limit $\bar{L}_1$ with distance according to $\bar{L}_1^{\text{obs}} = 3.241 \times 10^{35} (d/\text{Mpc})^3$ erg/s. Fig. 2 shows the best-fit result with a $p$-value from a Kolmogorov-Smirnov test of about 0.9. Note that the resulting fit slightly depends on the number $N_{\text{bins},d}$ of distance bins, but in order to obtain a sufficient number of known sources in each distance bin we suggest $100 \lesssim N_{\text{bins},d} \lesssim 300$. More important, the constant of proportionality $C_{RG}$ does not change significantly with $N_{\text{bins},d}$.

To complete the catalog at the low luminosity end, we use a strategy as follows. According to Eq. (2.6) the minimal total luminosity at 1.1 GHz that enables an acceleration of UHECRs yields

$$\bar{L}_{1,1}(Z) = 1.4 \times 10^{39} Z^{-7/3} \text{erg/s} \tag{3.4}$$

Since the chemical composition around the ankle shows that $Z < 6$ (e.g. [97]), there cannot be a significant contribution from sources with $L_{1,1} \leq 10^{38}$ erg/s. Therefore, the luminosity of M49 at a distance of 15.8 Mpc, which is the least luminous source of the van Velzen et al. catalog, provides a reasonable lower luminosity limit of $\bar{L}_{1,1} = 9.7 \times 10^{37}$ erg/s. Thus, as a

\begin{footnotesize}
5http://ned.ipac.caltech.edu/

6The most luminous RG in the vV12 catalog, Cygnus A, has a radio luminosity $L_{1,1} = 1.58 \times 10^{44}$ erg/s.
\end{footnotesize}
sufficiently luminous radio galaxy at lower distances would likely be contained in the catalog, we distribute additional RGs only beyond a distance of \( \hat{d} = 15.8 \) Mpc. So, the Eq. (3.3) is used to determine the unknown RGs up to a maximal distance of \( \hat{d} = 120 \) Mpc:

(i) On the one hand, the invisible, local RGs with a luminosity \( L_{\text{inv}} \in [\hat{L}_{1.1}, \hat{L}_{1.1}] \) within the distance \( d_{\text{un}} = [\hat{d}, \hat{d}] \), that are located in the non-covered part of the sky\(^7\), i.e. \( P_V = 0.12 \). Here, we obtain a total number \( N_{\text{inv RG}} = 100 \) of invisible sources.

(ii) On the other hand, the low luminous, local RGs with a luminosity \( L_{\text{low}} \in [\hat{L}_{1.1}, \hat{L}_{d_{\text{obs}}}] \) within the distance \( d_{\text{un}} \in [\hat{d}, \hat{d}] \), that are located in the covered part of the sky, i.e. \( P_V = 0.88 \), but emit a radio flux below the observable flux limit. Here we obtain a total number \( N_{\text{low RG}} = 617 \) of low luminous sources.

Next, we need to determine different realizations of a distribution \( \Phi \) of these invisible and low luminous sources in space and luminosity. For this, each distribution has to comply with the following expectations.

First, there is a certain radial distribution within the distance interval \( d_{\text{un}} \). Up to a distance of about \( 40 \, h^{-1} \) Mpc, where \( h = 0.678 \) denotes the dimensionless Hubble parameter, we expect the sources to be predominantly located within the supergalactic plane [98, 99]. However, a uniform distribution in space is expected beyond 59 Mpc, so that the distance-

---

\(^7\)The vV12 catalog includes only sources with \(|b| > 5^\circ\) for \(30^\circ < l < 330^\circ\), \(|b| > 8^\circ\) for \(l < 30^\circ\) or \(l > 330^\circ\), and \(|b| > 10^\circ\) for \(\delta < -40^\circ\), where \(l\) and \(b\) denote the Galactic longitude and latitude, respectively, and \(\delta\) refers to the declination of the source.
dependent source distribution yields

\[ \phi_1(d) \propto \begin{cases} d, & \text{for } d \leq 59 \text{ Mpc}, \\ d^2, & \text{for } d > 59 \text{ Mpc}. \end{cases} \]  

(3.5)

Second, the RLF

\[ \phi_2(L_{1,1}) = \frac{dN_{\text{RG}}}{dL_{1,1} \, dV} = C_{\text{RG}} \, L_{1,1}^{-1} \left[ \left( \frac{L_{1,1}}{L_*} \right)^\alpha + \left( \frac{L_{1,1}}{L_*} \right)^\beta \right]^{-1} \]  

(3.6)

is applied to distribute the luminosity of the invisible and low-luminous sources in the corresponding luminosity interval \( L_{\text{inv}} \) and \( L_{\text{low}} \), respectively.

Finally, the invisible and low-luminous RGs need to be placed at a certain position \((x_0, y_0, z_0)\) at the previously derived distance. Here, we use the magnetic-field strength \( B \) from Dolag et al., since \( B \) is to first order proportional to the gas-particle density [47]. The unknown spatial distribution of RGs is also expected to comply with \( \phi_3(x_0, y_0, z_0) \propto B(x_0, y_0, z_0) \), so that the total RG distribution function is described by

\[ \Phi(d, L_{1,1}, x_0, y_0, z_0) \propto \phi_1(d) \, \phi_2(L_{1,1}) \, \phi_3(x_0, y_0, z_0), \]  

(3.7)

where \( d = \sqrt{x_0^2 + y_0^2 + z_0^2} \). Fig. 3 shows the resulting total number of RGs up to a distance \( d \) with a radio luminosity above \( \tilde{L}_{1,1} \). Extending the source catalog of the vV12 catalog by the invisible and low luminous local RGs obviously leads to a very good agreement with the expected total number \( N_{\text{RG}} \) of radio-loud AGNs in the local Universe, i.e. \( d = 120 \text{ Mpc} \) and \( P_V = 1 \).
The local RG distribution is eventually composed of 121 RGs from the vV12 catalog, 100 invisible RGs and 617 low luminous RGs with individual positions and radio luminosities as displayed in Fig. 4. Here, ‘ESO 075-G 041’ is the most distant RG in the sample with \( d \sim d = 120 \text{Mpc} \) and also the most luminous one with \( L_{1.1} = 7.1 \times 10^{41} \text{erg/s} \). The most luminous sources at a small distance are Centaurus A (with \( L_{\text{CenA}} = 2.6 \times 10^{40} \text{erg/s}, d_{\text{CenA}} = 3.6 \text{Mpc} \)), M87 (with \( L_{\text{M87}} = 7.6 \times 10^{40} \text{erg/s}, d_{\text{M87}} = 18.4 \text{Mpc} \)) and Fornax A (with \( L_{\text{ForA}} = 8.3 \times 10^{40} \text{erg/s}, d_{\text{ForA}} = 20.9 \text{Mpc} \)). Due to these three sources we obtain a local excess in luminosity up to a distance of about 30 Mpc, that can also be seen in Fig. 3. If the deflection of the UHECRs by the EGMF is rather small and the CR luminosity is roughly proportional to radio luminosity (which is both expected), we expect a dominant contribution to the observed UHECR flux by Centaurus A, M87 and Fornax A. As derived in Sect. 2.1, the radio luminosity also provides an estimate of the corresponding maximal rigidity given by

\[
\hat{R}(L_{1.1}) = 0.139 \, g_{\text{acc}} \sqrt{g_{\text{cr}}} \left( \frac{L_{1.1}}{10^{38} \text{erg s}^{-1}} \right)^{3/7} \text{EV}.
\]  

(3.8)

The parameters \( g_{\text{cr}} \) and \( g_{\text{acc}} \) are used as free fit parameters with their value ranges \( 5 \leq g_{\text{CR}} \leq 50 \) and \( 0.01 \leq g_{\text{acc}} \leq 0.5 \), as discussed in Sect. 2.2. Assuming maximal values for these parameters and a canonical \( a = 2 \), Fig. 5 shows that among the local sources only very few sources can contribute to the explanation of the highest observed energies, and only for a composition significantly heavier than solar. It further confirms our expectation that Centaurus A, M87 and/or Fornax A are the prime candidates to explain the CR flux these energies, and we see that at least one of them must have enhanced metal abundances. Up to energies of about \( 10^{19.9} \text{eV} \), the total contribution by all the other local sources is similar to the one from M87 or Fornax A, however, at higher energies its contribution is negligible.

Although, at energies around the ankle all local sources are able to contribute to the observed CR flux, even for solar composition, but still the flux is expected to be dominated
by very few, bright sources.

An explanation of the whole UHECR spectrum by the local sources requires an individual adjustment of the initial abundance, since a continuous spectral behavior — without significant jumps — is needed where the contribution of light elements dominates around the ankle and the heavy elements at the highest energies. However, such a scenario of an ad hoc fine-tuning of the individual, initial abundances has a low plausibility, as it is highly unlikely to be by accident realized by nature.

### 3.2.2 The average non-local contribution

From simple geometrical considerations one can see that the flux provided by sources inside 120 Mpc cannot give a complete picture of the UHECR flux: assuming a homogeneous source distribution and neglecting propagation effects, the contribution of every distance bin $\delta d$ to the flux would be equal. As UHECR around $10^{18}\text{eV}$ can reach us in principle from 1 Gpc distance, the contribution from beyond 120 Mpc should be an order of magnitude higher at this energy than from inside this boundary. This is somewhat reduced by propagation effects, but it is enhanced by positive source evolution which we need to assume for radio galaxies.

In fact, the main reason for limiting our 3D simulations to a volume of 120 Mpc radius is of rather technical nature:

1. The EGMF structure available for our local universe is limited to this distance.

2. The vV12 catalog is known to become incomplete for distances larger than 120 Mpc, expressed best by the omission of the very powerful radio galaxy Pictor A close to this distance.

3. With the current setup, the necessary CPU time of 3D simulations with the same precision will become unmanageable in the case of a substantial increase of the source sample or the trajectory lengths. This argument becomes even more restrictive if we consider that for such propagation distances cosmological-evolution effects cannot be neglected, requiring 4D simulations which are even more CPU demanding [48, 88].

Therefore we follow a different approach to include distant sources: We perform 1D simulations based on the continuous source function for radio galaxies derived in Sect. 2.3. This
simulation is carried out from 120 Mpc up to a distance corresponding to a redshift \( z = 2 \). An advantage of this treatment is that it is fast, and cosmological evolution effects are easy to include. The disadvantage is that we neglect all potential anisotropy provided by distant, but extremely powerful sources. We return to this issue in the next section.

![Figure 6](image)

Figure 6: UHECR flux from the non-local contribution at the observer using \( a = 2 \), \( g_{cr} = 50 \), \( g_{acc} = 0.5 \). The left plot shows the influence of the different initial elements, where \( f_i = f_{\odot} \) (solid line) and \( f_i = f_{\odot} \) (dashed line). The right plot exhibits the change of the spectrum dependent on the different distance ranges the continuous source function is adopted.

Fig. 6 demonstrates two important aspects of the non-local contribution. The left panel shows that, as for local sources, only a heavy composition would be needed to reach out to the highest energy data. However, as the corresponding mean logarithm of the mass number at around \( 10^{18.5} \) eV is already above 2.5 in this scenario, it is in disagreement with the observations of the chemical composition. It thus confirms our prior expectation outlined in Sect. 2.3. The right panel shows the spectral steepening that results from adopting the continuous source function for an increasing distance range due to the increasing impact of interactions. Hence, the flux contribution from the average local sources is about a magnitude smaller and provides a significantly flatter spectral behavior compared to the one of the average non-local sources. We show in Appendix A.2 that indeed the continuous source function in the limited range of 5 to 120 Mpc perfectly agrees with the sum of contributions from all local source described in the previous section.

Hence, we can already draw the conclusion that the average non-local contribution needs to be subdominant above the ankle. As shown in Fig. 7, for the case of a heavy composition the parameter values of \( g_{cr} \) and \( g_{acc} \) need to be at the lower end of the possible range in order to make an initial spectral index \( a < 2 \) possible. In the case of \( f_i = f_{\odot} \), the range of possible values enlarges, but still the lower end of the parameter space is preferred. In the contrary, the previous Sect. has shown that at least a few, powerful local sources need a high value of \( g_{cr} \) to explain the highest energies observed. Thus we conclude, that we need to assume an exceptional parameter value for Centaurus A, M87 and/or Fornax A in order to explain the observational data.

### 3.2.3 Cygnus A

Cygnus A, with a CR luminosity \( Q_{cr}(L) = 1.3 \times 10^{45} g_{cr} \) erg/s and a corresponding maximal rigidity \( \hat{R}(L) = 62.6 g_{acc} \sqrt{g_{cr}} \) EV at a distance of 255 Mpc, represents an ideal UHECR accelerator. It can accelerate protons to the highest energies observed, and nuclei to energies
Figure 7: Minimal spectral index of the average non-local contribution dependent on $g_{cr}$ and $g_{acc}$ to obtain a negligible contribution to the observed UHECR flux. Here, the limiting cases of a solar-like initial composition with $f_i = f_\odot$ (left plot) and a heavy initial composition with $f_i = f_\ast$ (right plot) are shown.

Far beyond. Due to its large distance, its UHECR spectrum arriving at the observer is almost entirely shaped by the GZK process, unlike local FR-I galaxies, where the resulting spectrum is mostly determined by source properties. A slight residual effect of the source parameters is given through the strength of its GZK bump [82, 83], as shown in figure 8. This shows that Cygnus A is in principle able to explain the observed UHECR flux up to a few tens of EeV, irrespective of its composition, where a better fit to the spectrum obtained for a light composition scenario. It could therefore provide a source of light elements, balancing the required heavy composition of FR-I galaxies. In fact, the different morphology of FR-
II galaxies like Cygnus A gives rise to assume not only a different composition, but also a
different parameter set of $g_{\text{CygA}}$ and $g_{\text{acc}}$. 

As we cannot perform a full 3D or even 4D simulation for Cygnus A due to the lack of
a reliable EGMF structure, we use a simplified approach: We simulate Cygnus A as a point
source in 3D neglecting deflections in the EGMF. In doing so, we are still able to normalize
the flux — as it will be introduced in Sect. 3.4.1 — in the same manner as we do it for
the local sources. GMF deflections can be still applied. As this is a completely unrealistic
scenario, we also consider the case of a full isotropization of UHECRs arriving from Cygnus
A in the EGMF, a scenario which is even possible in limiting EGMF scenarios [46]. Together,
they give limits on the anisotropy introduced by this source, and the results for spectrum and
composition are largely independent on deflections, if we neglect the propagation-time delay
which could in principle lead to a slight increase of energy losses or photo-disintegration.

As Cygnus A can easily provide more than half of all UHECRs observed, it is justified
to treat it as the only point source outside 120 Mpc. There are two other FR-II radio galaxies
which could make an individual contribution on a $\sim 10\%$ level, Pictor A and 3C353, both at
about 125 Mpc distance, which are expected to show very similar properties in spectrum and
composition as Cygnus A. Unlike the northern source Cygnus A, however, they are visible by
Auger which could be of some relevance (see discussion in Sect. 5). All weaker FR-II galaxies
can, with some restrictions, be considered to be properly treated by the continuous source
function approach, and in fact they are on average included in the high luminosity part of
the source function.

3.3 Simulation parameters

As mentioned in Sect. 3.2, the free parameters of the model are optimized after the simulation
by applying weights to the simulated candidates — we describe this procedure in detail in
Sect. 3.4. Nevertheless, we need to consider the plausibility ranges of these parameters already
during the simulation to keep the statistics of all simulated sources, energies and particle
species in an optimal relation.

All 3D simulations of individual sources are carried out up to the maximal rigidity of
the source, $\hat{R} = g_{\text{acc}} \sqrt{g_{\text{cr}} Q_{\text{jet}}/c}$ for the maximal parameters $g_{\text{cr}} = 50$ and $g_{\text{acc}} = 0.6$, where
$Q_{\text{jet}}$ is determined from the source luminosity $L_{1,1}$ as given in the vV12 catalog. We use
a canonical spectral index $\tilde{\alpha} = 2$, thus the number of particles simulated per logarithmic
energy bin follows an $E^{-1}$ power law. To obtain a sufficient number of candidates from each
initial type of element, we simulate a similar number of H, He, CNO and Fe elements. Here,
CNO is composed of 29% $^{12}$C, 8% $^{14}$N and 63% $^{16}$O, like the abundances of these elements
in the solar system [94]. In the case of 3D simulations that include CR particle interactions
and deflections, the possible statistics is mostly determined by the available CPU time. With
respect to the available computational power, we aimed to reach sufficient statistics to
obtain a statistical Poisson error per energy bin of less than 1% at low energies, i.e. up to
40 EeV. This way, the resulting statistical accuracy from the simulation becomes similar to
the error of the observed UHECR flux. This translates to the requirement that about $10^4$
candidates per low energy bin need to arrive at the observer sphere. However, this statistical
accuracy is neither needed nor realizable in less than $10^6$ CPU hours for each individual
source. Due to the absolute normalization of the CR flux, as introduced in Sect. 3.4.1, the

\[\text{In our case, we used about } 2 \times 10^5 \text{ CPU hours for the simulation of the total local source sample.}\]

\[\text{To obtain these numbers in the high energy bins, we would have to increase the total number of simulated}
\text{events by a factor of } 10^2 \text{ or more.}\]
low-luminous sources contribute to the result only in their total sum, so that in the end a comparable statistics to the powerful sources is obtained. Hence, it is a reasonable choice to simulate a total number $\tilde{N}_{i,j}$ of CR candidates per element $i$ and source $j$, that is proportional to its radio luminosity $L_j$, i.e. $\tilde{N}_{i,j} \propto L_j$.

The simulation of CR candidates from Cygnus A is designed in the same manner as the simulation of the local sources, but disregarding the influence of the EGMF, which speeds up the simulations significantly. Here, we simulate the same total number for each element type $i$ and by using $\tilde{N}_{i,CygA} = 6.1 \cdot 10^{11}$ we obtain statistics that are comparable to ones from the powerful, local sources.

For the continuous non-local source function (2.17), 1D simulations are carried out using a maximal rigidity of $\hat{R}_{nl} = 100$ EV and a flat initial spectrum of $\tilde{a}_{nl} = 1$. Although 1D simulations are so fast that CPU time is not a limiting factor, the computational time of subsequent analysis (the fitting procedure) strongly depends on size of the output data. In order to obtain an appropriate analyzable output with sufficient statistics, we simulate about $10^6$ candidates in total with some more light than heavy elements, due to the dominant contribution of these elements according to the solar abundances.

The final numbers of simulated particles together with the initial simulation parameters that characterize the energy distribution of the UHECR outflow at the sources are summarized in Tab. 1. Further details on the simulation program and the used parameters are given in Appendix B.

### Table 1: Initial simulation parameters.

<table>
<thead>
<tr>
<th></th>
<th>$\tilde{N}_H$</th>
<th>$\tilde{N}_{He}$</th>
<th>$\tilde{N}_C$</th>
<th>$\tilde{N}_N$</th>
<th>$\tilde{N}_O$</th>
<th>$\tilde{N}_{Fe}$</th>
<th>$\tilde{a}_l$</th>
<th>$g_{ct}$</th>
<th>$g_{acc}$</th>
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<tr>
<td>local</td>
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<td>$1.3 \cdot 10^{11}$</td>
<td>$3.7 \cdot 10^{10}$</td>
<td>$1.1 \cdot 10^{10}$</td>
<td>$8.0 \cdot 10^{10}$</td>
<td>$1.2 \cdot 10^{11}$</td>
<td>2</td>
<td>50</td>
<td>0.6</td>
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<tr>
<td>CygA</td>
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<td>$6.1 \cdot 10^{11}$</td>
<td>$6.1 \cdot 10^{11}$</td>
<td>$6.1 \cdot 10^{11}$</td>
<td>$6.1 \cdot 10^{11}$</td>
<td>$6.1 \cdot 10^{11}$</td>
<td>2</td>
<td>50</td>
<td>0.6</td>
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<tr>
<th></th>
<th>$\tilde{N}_H$</th>
<th>$\tilde{N}_{He}$</th>
<th>$\tilde{N}_C$</th>
<th>$\tilde{N}_N$</th>
<th>$\tilde{N}_O$</th>
<th>$\tilde{N}_{Fe}$</th>
<th>$\tilde{a}_{nl}$</th>
<th>$\hat{R}_{nl}$ [EV]</th>
</tr>
</thead>
<tbody>
<tr>
<td>non-local</td>
<td>$5 \cdot 10^5$</td>
<td>$2.5 \cdot 10^5$</td>
<td>$8.3 \cdot 10^4$</td>
<td>$7.1 \cdot 10^4$</td>
<td>$6.3 \cdot 10^4$</td>
<td>$2 \cdot 10^4$</td>
<td>1</td>
<td>100</td>
</tr>
</tbody>
</table>

### 3.4 Absolute normalization of the CR flux

As our model is absolutely normalized due to the connection between CR power and maximal rigidity of each source, we have no freedom to fit a spectrum obtained in "arbitrary units" to the measured CR spectrum. Instead, we need to transfer the absolute normalization to the simulation setup. In general, the task is to achieve absolute weights $W_k = dN_k/dS\, dt$ to each CR candidate $k$ which arrived at the observer with energy $E_k^*$ and charge $Z_k^*$ in the physical dimension of particles per square-cm per second. Here we need the information on each candidates initial rigidity $R_k = E_k^*/(eZ_k^*)$ at the source, which is provided in the CRPropa output. Due to the different treatments of the local and non-local flux contribution, their absolute weights differ as shown in the following.

#### 3.4.1 Cosmic ray flux from local sources and Cygnus A

As rigidity is the natural measure for this problem, we use the simulated spectrum $d\tilde{N}/dR \propto R^{-\tilde{a}}$ dependent on the initial rigidity $R$ of the candidates and re-weight it to obtain

$$\frac{d\tilde{N}_i}{dR} = \nu_0 f_i R^{-\tilde{a}} \exp(-R/\hat{R})$$

(3.9)
with an arbitrary normalization constant $\nu_0$. In order to derive a characteristic emission time scale, we use the CR power of the local source $j$, of the individual charge type $i$ given by

$$Q_{cr,i,j} = Q_{cr,j} \frac{f_i Z_i}{Z},$$

(3.10)

as well as the CR energy of

$$E_{cr,i} = \nu_0 f_i Z_i e \int_{R_0}^{R_{cut}} dR \frac{R^{1-a}}{R} \exp(-R/\hat{R})$$

(3.11)

that is provided by the re-weighted CR spectrum (3.9). Hence, the CR emissivity of the local source $j$ with a given total CR power $Q_{cr}$ yields

$$\frac{dN_{i,j}}{dR} = \frac{dN_{i,j}}{dR} \frac{Q_{cr,i,j}}{E_{cr}} = \frac{f_i}{Z} \frac{Q_{cr,j} R^{-a}}{e \int_{R_0}^{R_{cut}} dR \frac{R^{1-a}}{R} \exp(-R/\hat{R})}$$

(3.12)

The surface of the observer$^{10}$ is given by $S = 4\pi r_{obs}^2$. The mean rigidity transfer from the source to the observer sphere is determined by

$$\bar{R}_k = \frac{\int_{R_k}^{\hat{R}_k} dR \frac{dN}{dR}}{\int_{R_k}^{\hat{R}_k} dR} \frac{dN}{dR},$$

(3.13)

where $R_k$ denotes the initial rigidity of an individual CR candidate that arrived at the observer, and $\hat{R}_k = 3E_{cr}/Z_i$ is the minimal initial rigidity of the simulated candidates. Knowing the total number $N_{i,j}$ of simulated CRs of a given charge type $i$, we are able to determine the flux weight of an individual CR candidate that managed to arrive at the observer

$$W_{k,i,j} = \frac{\bar{R}_k}{S N_{i,j}} \frac{dN_{i,j}}{dR_k} \frac{dR_k}{dt} = \frac{f_i Q_{cr,j} (\hat{R}_1^{1-\hat{a}} - \hat{R}_1^{1-\hat{a}})}{4\pi r_{obs}^2 \bar{N}_{i,j} Z e (1 - \hat{a}) (G_a(\hat{R}_0) - G_a(\hat{R}_{cut}))} R_k^{\hat{a}-a} \exp(-R_k/\hat{R}),$$

(3.14)

with

$$G_a(x) = \hat{R}^{2-a} \Gamma \left(2 - a, x/\hat{R} \right)$$

(3.15)

where $\Gamma(c,x)$ denotes the incomplete gamma function. Thus, the total, observed UHECR flux from all local sources yields

$$W_\kappa = \sum_{k,i,j} \delta_{k\kappa} W_{k,i,j}.$$

(3.16)

Here, the Kronecker Delta accounts for the final binning of the observed flux in energy bins $E_\kappa$ (denoting the central energy of the bin) of equal logarithmic width $\Delta = d \log(E_\kappa)$, so that the linear width of the bin $E_\kappa$ is $dE_\kappa = \Delta E_\kappa$. Hence, $\delta_{k\kappa} = 1$ for $E_k^* \in dE_\kappa$ and $\delta_{k\kappa} = 0$ otherwise. The absolute particle number flux density in the energy bin $E_\kappa$ (meaning in CGS units, particles per square-cm per erg per second) is given by

$$j_\kappa \equiv W_\kappa / dE_\kappa = \sum_{k,i,j} \delta_{k\kappa} \frac{f_i Q_{cr,j} (\hat{R}_1^{1-\hat{a}} - \hat{R}_1^{1-\hat{a}}) R_k^{\hat{a}-a} \exp(-R_k/\hat{R})}{4\pi r_{obs}^2 \bar{N}_{i,j} Z e (1 - \hat{a}) (G_a(\hat{R}_0) - G_a(\hat{R}_{cut}))}.$$

(3.17)

It is shown in the Appendix A.1 that the previously derived absolute normalization is in good agreement with the theoretical expectations.

Note that negligible energy losses are supposed between the observer sphere in the simulation and the ‘real observer’, i.e. the Earth.
### 3.4.2 Cosmic ray flux from the non-local source function

In the case of the 1D simulation of non-local sources, all CR candidates with a final energy \( E_k^* \geq 3 \text{EeV} \) are stored by the observer. Thus, the absolute normalization of the CR flux cannot be derived as before by using the rigidity transfer to a geometrical observer sphere with a certain surface. However, the production rate density Eq. (2.16) enables an absolute normalization, where we have to consider source evolution effects according to Eq. (2.17), fixing the evolution parameter to \( m = 3 \) as it has only a minor influence on our results. We define an effective comoving volume element

\[
\Delta V_k \approx 4\pi d_k^2 \Delta d_i
\]

(3.18)

with the propagation distance \( d_k \) of the observed CR candidate and the thickness \( \Delta d_i = \hat{d}/\tilde{N}_i \) of the spherical shell that is represented by each simulated candidate of type \( i \). Here \( \hat{d} \) denotes the maximal comoving distance that has been adopted in the 1D simulation. Thus, the CR production rate yields

\[
\dot{P}_{cr,i,k} \equiv \frac{dN_k}{dR_k \, dt} = \int dV \Psi_{i,0}(R_k) \simeq \Psi_{i,0}(R_k) (1 + z)^{m-1} 4\pi d_k^2 \Delta d_i .
\]

(3.19)

Note, that the result of the integral only holds, when a homogeneous distribution of sources is used in the 1D simulation. Since the non-local candidates are simulated with a flat energy distribution, we obtain a mean rigidity transfer of

\[
\bar{R}_k = \frac{\int \hat{R}_{nl} \, d\tilde{N}/dR}{\int \tilde{N}/dR} = R_k \ln \left( \frac{\hat{R}_{nl}}{\hat{R}_1} \right) ,
\]

(3.20)

so that the total CR production rate yields

\[
\frac{dN_k}{dt} = R_k \dot{P}_{cr,i,k} = R_k \ln \left( \frac{\hat{R}_{nl}}{\hat{R}_1} \right) \dot{\rho}_{cr,i}(R_k) (1 + z)^{m-1} 4\pi d_k^2 \Delta d_i .
\]

(3.21)

In three spatial dimensions, an isotropically emitting source with a total rate \( dN/dt \) provides a particle flux of \( \frac{dN}{dA \, dt} = \frac{dN}{4\pi d^2} \) at a distance \( d \). The CR production rate (3.19) corresponds to \( dN/dt \) for an isotropically emitting source at distance \( d_k \) in three spatial dimensions, which can easily be shown by the transformation of \( \dot{P}_{cr,i,k} \) into three spatial dimensions.\(^{11}\) However, we still need to take into account that the cross section of an extended, spherical observer is reduced by a factor of \( 1/4 \) by switching from a 1D to a 3D treatment, i.e. \( (dN_k/dt)_{3D} = (dN_k/dt)/4 \). Hence, the flux contribution by a CR candidate with a propagation distance \( d_k \) is given by

\[
W_{k,i} \equiv \frac{dN_k}{dt} \, dA_k = \left( \frac{dN_k}{dt} \right)_{3D} \frac{1}{4\pi d_k^2} = \frac{R_k \ln \left( \frac{\hat{R}_{nl}}{\hat{R}_1} \right) \hat{d}}{4\tilde{N}_i} \dot{\rho}_{cr,i}(R_k) (1 + z)^{m-1} \]

(3.22)

Like in the previous section, the total, observed UHECR flux from all non-local sources yields

\[
W_\kappa = \sum_{k,i} \delta_{k\kappa} W_{k,i}
\]

(3.23)

\(^{11}\)First \( \dot{P}_{cr,i,k} \) needs to be integrated over the corresponding surface area, in order to obtain the total CR rate that is emitted in direction of the observer. Since, we expect the sources to emit CRs rather isotropically, we also need to divide by \( 4\pi d_k^2 \), so that the surface area from the previous integration is canceled out.
where, the Kronecker Delta once again accounts for the final binning of the observed flux into energy bins $E_\kappa$ with a linear width of $dE_\kappa$. So, the absolute particle number flux density in the energy bin $E_\kappa$ (meaning in CGS units, particles per square-cm per erg per second) yields

$$ j_\kappa = \frac{W_\kappa}{dE_\kappa} = \sum_{k,i} \delta_{k\kappa} \frac{R_k}{4N_i} \ln(\hat{R}_{nl}/\hat{R}_1) \hat{d} \rho_{\text{cr},i}^n R_k (1+z)^{m-1}. \quad (3.24) $$

A sanity check of the previously derived absolute normalization is shown in the Appendix A.2.

Of course, the total particle number flux density from local and nonlocal sources is given by the sum of Eq. (3.17) and (3.24). In addition, the normalization weights $W_{k,i,j}$ and $W_{k,i}$ from the local and nonlocal contribution, respectively, enable to determine the mean logarithm of the mass number $A_\eta$ within the energy bin $E_\eta$ according to

$$ \langle \ln(A_\eta) \rangle = \frac{\sum_{k,i,j} \delta_{k\eta} (\ln(A_{k,i,j}^* W_{k,i,j} + \ln(A_{k,i}^i) W_{k,i})}{\sum_{k,i,j} \delta_{k\eta} (W_{k,i,j} + W_{k,i})}. \quad (3.25) $$

### 3.5 Fit parameters of the model

As introduced in Sect. 2, the main free parameters of our model and their plausibility ranges are the spectral index $a \in [1.7, 2.2]$, the cosmic-ray load $g_{\text{cr}} \in [1, 50]$, and the acceleration efficiency $g_{\text{acc}} \in [0.01, 0.5]$. The abundances of nuclear species are connected, according to Sect. 2.2, by the dependence $f_i = Z_i^q$, with an abundance parameter $q \in [0, 2]$. Finally, the source-evolution parameter was fixed to $m = 3$ as its influence is mainly important below the ankle [83].

In derivation of the continuous source function, we assumed that these parameters are the same for all radio galaxies to allow an analytical solution. This, however, is a simplification and by no means implied. The scatter in the correlation plots showed by Willott et al. suggests that $g_{\text{cr}}$ can vary from source to source within an order of magnitude centered at the best fit value for the sample. There is no reason to assume that a similar scatter does not apply to $g_{\text{acc}}$ and $q$, as these parameters depend on individual source properties as jet speeds and the material in which the jet dissipates. Only for the spectral index $a$, we have good reasons to consider it as a parameter which is, in narrow ranges, determined by fundamental physics and thus universal.

Hence, it would be in principle required for our treatment of individual sources to give each source its own parameter set. However, this would increase the number of free parameters far beyond the number of data constraints, making the problem ambiguous and ill-defined. Luckily, from the discussion in the previous section, we have good reasons to assume that only four sources, Centaurus A, M87, Fornax A, and Cygnus A, make a significant individual contribution to the cosmic ray flux. All other individual sources can be given the same parameter values, and continuity implies that they should be identical to those chosen for the continuous source function. The contribution by latter needs to be subdominant, as already exposed in Sect. 3.2.2, hence, there is no benefit in using different parameter values of $g_{\text{cr}}$ and $g_{\text{acc}}$ in the case of these sources. Here, we only use an individual value of $g_{\text{cr}}$, subsequently called $\bar{g}_{\text{cr}}$, but stick to the same parameter value of $g_{\text{acc}}$ that is applied for Centaurus A, M87 and Fornax A, respectively. Regarding the abundance parameter $q$, we assume from the discussion in Sect. 2.2 that $q = 0$ for Cygnus A (meaning solar abundances), while for FR-I
galaxies an enrichment of heavy elements is allowed by fitting \( q \) within its plausibility limits. To keep the parameter set manageable, we assume hereby that this parameter is the same for all local and non-local radio galaxies. Thus, our model has a total number of seven free parameters, \( \mu, g_{\text{cr}}^{\text{CygA}}, g_{\text{acc}}^{\text{CygA}}, g_{\text{cr}}^{\mu}, g_{\text{acc}}^{\mu}, g_{\text{cr}}^{\text{M87}}, q \), where \( \mu \) refers to either Centaurus A, M87 or Fornax A.

4 Fit Results

To obtain a best fit parameter set, we use the `minimize()` function provided by the python package `lmfit` [100]. This function minimizes the residual \( R \) of the weighted and binned results \( j_\kappa \) and \( \langle \ln(A_\eta) \rangle \) from the simulation and the 24 observational data points above \( 5 \times 10^{18} \text{eV} \) from Auger [4, 97]. We perform the minimization algorithm 10 times using Powell’s method [101] with a random set of initial weights within the allowed range, to increase the probability that the obtained minimum is indeed a global minimum. The proposed fitting method tends to lead to different local minimums, especially for an unfitting model, however, in the case of a suitable fit model the algorithm converges most times to the same minimum. As usual the goodness of the fit is quantified by the minimal reduced chi-square \( \chi^2 = \sum R^2/n_{\text{dof}} \), where \( n_{\text{dof}} \) refers to the corresponding degree of freedom of the fits. The normalization of the observed \( \langle \ln(A) \rangle \) data depends pretty much on the used hadronic interaction model. There are three contemporary models that were either tuned to recent LHC data (QGSJetII-04 [102], Epos-LHC [103, 104]) or found in good agreement with these measurements (Sibyll2.1 [105]). The Auger collaboration processed the measured \( X_{\text{max}} \) data with these interaction models exposing no difference in the spectral behavior but a systematic shift in the normalization, whereof the QGSJetII-04 (Epos-LHC) model yields the most light (heavy) composition [97]. If there are no basic changes in the fit result, we will predominantly display the results that refer to the Sibyll2.1 model.

In the case of an accurate fit, we further analyze the corresponding arrival directions of the best-fit models, although, we neither included the arrival directions into the fit nor are deflection effects for CR candidates from non-local sources, and Cygnus A in particular, taken into account. However, there are two extreme scenarios we like to discuss: (a) UHECRs from Cygnus A do not suffer from any EGMF deflections, and (b) UHECRs from Cygnus A are completely isotropized by the EGMF. It needs to be kept in mind that Cygnus A was used as a representative of the powerful nearby FR-II sources, hence, the scenario (b) could also be realized by replacing Cygnus A with a multitude of FR-II galaxies with slightly less radio power. This analysis is in particular interesting in the light of the recent observation [106] of a dipol moment above \( 8 \text{EeV} \).

In addition, the resulting best-fit parameters also enable to draw further conclusions on the origin and the deflection of the observed UHECR candidates from local sources. Here, the corresponding best-fit weights \( W_{k,i,j}^{\text{best}} \) according to the absolute normalization of the local flux (3.14), have to be taken into account. Thus, the contribution of the local source \( \mu \) is given by

\[
N_\mu = \frac{\sum_{k,i,j} \delta_{\mu j} W_{k,i,j}^{\text{best}}}{\sum_{k,i,j} W_{k,i,j}^{\text{best}}} \tag{4.1}
\]

where the Kronecker delta \( \delta_{\mu j} \) is used to ensure that the observed CR candidate originates from this source. Further, the mean deflection angle of the UHECR candidates from the...
source $\mu$ is determined by

$$\bar{\theta}_{\text{def},\mu} = \frac{\sum_{k,i,j} \delta_{\mu j} W^\text{best}_{k,i,j} \theta_{\text{def},i,k}}{\sum_{k,i,j} W^\text{best}_{k,i,j}}. \quad (4.2)$$

Here, the deflection angle of an individual candidate $i$ of type $k$ is estimated by

$$\theta_{\text{def},i,k} = \arccos \left( \frac{\vec{p}_{i,k} \cdot \vec{p}_{i,k}^*}{|\vec{p}_{i,k}| |\vec{p}_{i,k}^*|} \right) \quad (4.3)$$

where we only use its initial and final momentum $\vec{p}_{i,k}^*$ and $\vec{p}_{i,k}$, respectively. Note, that in general $\theta_{\text{def},i,k}$ only determines the minimal deflection angle, since a total deflection $> \theta_{\text{def},i,k}$ can not be reconstructed\footnote{In principle, a detailed tracking of the candidates is possible with CRPropa. However, further developments on the propagation module are needed in order to handle the vast amount of simulated data.}. However, in the case that the minimal deflection angles of the candidates are predominantly small, most likely small-angle deflections take place and $\theta_{\text{def},i,k}$ is also a good approximation for the total deflection angle.

### 4.1 Case 1: All local sources equal

First, we discuss the case that none of the powerful individual local sources offer deviating parameters, i.e. $g^\mu_{\text{cr}} = \bar{g}_{\text{cr}}$, and only Cygnus A provides solar abundances while for all other sources heavy elements get enhanced by fitting $q$. The resulting best-fit parameters, using the different hadronic interaction models, are summarized in Table 2. As shown in Fig. 9, no satisfactory fit to the observational data is possible, even in the case of the EPOS-LHC.

<table>
<thead>
<tr>
<th>Model</th>
<th>$a$</th>
<th>$\bar{g}_{\text{cr}}$</th>
<th>$CygA_{\text{CygA}}$</th>
<th>$\bar{g}_{\text{acc}}$</th>
<th>$CygA_{\text{acc}}$</th>
<th>$q$</th>
<th>$\chi^2$</th>
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</table>

**Figure 9:** Best-fit results to energy spectrum (left) and chemical composition (right) using the EPOS-LHC model and no local source exception.
model which leads to the heaviest observed composition. Apart from the decreasing heaviness
with increasing energy, a distinct disagreement with the observations, also the CR flux at the
highest energies cannot be explained. Hence, at least Centaurus A, M87 or Fornax A need to
obtain a higher value of $g_{\text{cr}}$ than the average sources.

We want to point out one additional aspect here which is relevant also for the following
two cases. Our fit procedure finds for Cygnus A a rather low acceleration efficiency, $g_{\text{acc}}^{\text{CygA}}$. This is counter-intuitive, as from its jet power and radio structure one would expect that there can hardly be a more efficient accelerator among radio galaxies than Cygnus A. As we expound in Sect. 5, the very large angle between the Auger zenith and Cygnus A is expected to cause a steepening of the spectrum from Cygnus A measured by Auger, an effect which is not considered here. To compensate this, we would need an originally flatter spectrum of Cygnus A, which would require a larger value of $g_{\text{acc}}$ as seen from Fig. 8. We cannot quantify this effect here as we are not able in our setup to calculate the detailed propagation of UHECR from Cygnus A, but qualitatively it should be noted that our findings for $g_{\text{acc}}^{\text{CygA}}$ should not be regarded a robust result.

4.2 Case 2: Powerful Centaurus A

In this case we assume that candidates arriving from Centaurus A can be weighted with
an individual parameter value of $g_{\text{cr}}^{\text{CenA}}$. We consider the two limiting cases where (a) only
Centaurus A has an enhanced heavy composition according to the parameter $q$ — the so-
called light composition scenario; and (b) the case where all sources except Cygnus A eject an
enhanced heavy composition — the heavy composition scenario. Cygnus A gets an individual
terminator value of $g_{\text{cr}}^{\text{CygA}}$ and $g_{\text{acc}}^{\text{CygA}}$ as well, but all other sources (including Fornax A and
M87) get an average parameter value of $g_{\text{cr}}$ and $g_{\text{acc}}$. The resulting best-fit parameters
for both composition scenarios are summarized in Table 3 and 4. Independent of the

| Table 3: Best fit parameters (light composition scenario with powerful Centaurus A). |
|---|---|---|---|---|---|---|---|
| $a$ | $\bar{g}_{\text{cr}}$ | $g_{\text{cr}}^{\text{CenA}}$ | $g_{\text{cr}}^{\text{CygA}}$ | $g_{\text{acc}}^{\text{CenA}}$ | $g_{\text{acc}}^{\text{CygA}}$ | $q$ | $\chi^2$ |
| EPOS-LHC | 1.85 | 7.73 | 41.54 | 43.94 | 0.127 | 0.059 | 2 | 1.1 |
| QGSJetII-04 | 1.82 | 6.31 | 21.20 | 48.84 | 0.220 | 0.056 | 2 | 1.4 |
| Sibyll2.1 | 1.83 | 6.67 | 24.77 | 47.90 | 0.19 | 0.056 | 2 | 1.3 |

| Table 4: Best fit parameters (heavy composition scenario with powerful Centaurus A). |
|---|---|---|---|---|---|---|---|
| $a$ | $\bar{g}_{\text{cr}}$ | $g_{\text{cr}}^{\text{CenA}}$ | $g_{\text{cr}}^{\text{CygA}}$ | $g_{\text{acc}}^{\text{CenA}}$ | $g_{\text{acc}}^{\text{CygA}}$ | $q$ | $\chi^2$ |
| EPOS-LHC | 1.77 | 1.02 | 46.73 | 30.48 | 0.117 | 0.068 | 1.77 | 1.3 |
| QGSJetII-04 | 1.82 | 1.06 | 27.42 | 49.34 | 0.211 | 0.057 | 1.81 | 1.8 |
| Sibyll2.1 | 1.83 | 1 | 38.77 | 50 | 0.159 | 0.057 | 1.8 | 1.6 |

preferred hadronic interaction model, we are able to explain the observational data of the
energy spectrum and the chemical composition in the case of the light composition scenario,
as shown in Fig. 10 for Sibyll2.1. The heavy composition scenario leads to slightly less
accurate fits, predominantly due to a too heavy composition around $10^{18.8}$ eV, and the full
range of possible values for the parameter $g_{\text{cr}}$ is used in order to keep the average non-local
contribution subdominant. Concentrating on the light composition scenario, the significantly smaller parameter value of $g_{cr}$ compared to the values of Centaurus A or Cygnus A indicates that not every RG with a certain minimal radio luminosity has a high cosmic ray load. We could probably account for at most $15–25\%$ of the RGs with a radio luminosity at 1.1 GHz above $\tilde{L} \approx 10^{38}$ erg/s to have parameters comparable to Centaurus A and Cygnus A, while the contribution of all other RGs then would have to be irrelevant (i.e., $g_{cr} \ll \bar{g}_{cr}$). The same would have to be the case for the abundance enhancement, which needs to be stronger in Centaurus A (or a small subset of sources) than in most other sources. In total, we would need to assume that both Cygnus A and Centaurus A, which are anyway unusual due to their proximity given their power, are also preferred as cosmic-ray accelerators. In this scenario, almost all UHECRs observed would originate in one of those two sources. For this scenario, the isotropically ejected CR power above the ankle of these sources would be $Q_{\text{CenA}}^{\text{uhecr}} \approx 2 \times 10^{43}$ erg/s and $Q_{\text{CygA}}^{\text{uhecr}} \approx 2 \times 10^{47}$ erg/s, with only minor changes caused by the chosen scenario or the hadronic interaction model.

Figure 10: Best-fit results to energy spectrum (left) and chemical composition (right) using Sibyll2.1 and the light composition scenario with powerful Centaurus A.

Figure 11: Best-fit results to energy spectrum (left) and chemical composition (right) using Sibyll2.1 and the heavy composition scenario with powerful Centaurus A.

The corresponding arrival directions for the different best-fit models or initial composition scenarios do not differ significantly. Figs. 12-14 show the results for the limiting Cygnus A scenarios (point source vs. complete isotropisation) in the case of the Sibyll2.1 best-fit model.
for the light scenario. The angular power spectra shown also include the case where initial H and He candidates from Centaurus A are excluded, which does not affect the total energy spectrum or the chemical composition, but lowers the level of anisotropy and improves consistency with the data set at $4 - 8 \text{EeV}$. The motivation for this is that models for a particular heavy enhancement of Centaurus A would expect a pure heavy composition [76]. We note, however, that at $4 - 8 \text{ EeV}$ an impact of cosmic rays from the Galaxy or other extragalactic components cannot be excluded, and that at $E > 8 \text{ EeV}$ our results are in good agreement with the data, for the case that UHECR from Cygnus A are completely isotropized.

**Figure 12:** Skymap with non-deflected Cygnus A events for $4 \text{EeV} \leq E \leq 8 \text{EeV}$ (left), and $E > 8 \text{EeV}$ (right) using Sibyll2.1 and the light composition scenario with a powerful Centaurus A.

**Figure 13:** Skymap with isotropized Cygnus A events for $4 \text{EeV} \leq E \leq 8 \text{EeV}$ (left), and $E > 8 \text{EeV}$ (right) using Sibyll2.1 and the light composition scenario with a powerful Centaurus A.

Fig. 15 shows the resulting mean deflection $\bar{\theta}_{\text{def}, \mu}$ as well as $N_\mu$ of the individual local RG that contribute to the best-fit flux. It is seen that the dominant contribution from Centaurus A does not suffer significant deflections in the EGMF, rather it is the Galactic magnetic field which randomizes the arrival directions from Centaurus A to a large degree as they are mostly iron nuclei and thus their rigidity is small. Other radio galaxies which still have some significant contribution mostly show EGMF deflections up to about $10^\circ$. 

– 29 –
Figure 14: Angular power spectrum with isotropized (solid and dash-dotted line) and non-deflected (dashed line) Cygnus A events for $4 \text{EeV} \leq E \leq 8 \text{EeV}$ (left), and $E > 8 \text{EeV}$ (right) using Sibyll2.1 and the light composition scenario with a powerful Centaurus A.

Figure 15: The mean deflection angle of CR candidates from the individual local sources and its contribution (given by the color) to the best-fit CR flux. Here for the Sibyll2.1 model and the light composition scenario with a powerful Centaurus A.

4.3 Case 3: Powerful M87 and Fornax A

Finally, we discuss the case where M87 and Fornax A, instead of Centaurus A, have an individual parameter value of $g_{\text{M87/Fornax}}$. The resulting best-fit parameters for the light and the heavy composition scenarios are summarized in table 5 and 6, respectively.

In general, the resulting fit is less accurate than in the case 2, as shown in Fig. 16 and 17: Dependent on the chosen initial composition scenario the resulting composition at the observer is either too light or too heavy (with a rather decreasing heaviness with increasing energy). The resulting CR flux shows a better agreement with the observations, although,
Table 5: Best-fit parameters (light composition scenario, powerful M87 and Fornax A).

<table>
<thead>
<tr>
<th></th>
<th>$a$</th>
<th>$\bar{g}_c$</th>
<th>$g_{cr}$</th>
<th>$M/F$</th>
<th>$g_{CygA}$</th>
<th>$g_{acc}$</th>
<th>$q$</th>
<th>$\chi^2$</th>
</tr>
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<tr>
<td>EPOS-LHC</td>
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<td>50</td>
<td>32.04</td>
<td>0.100</td>
<td>0.065</td>
<td>2</td>
<td>4.8</td>
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<td>QGSJetII-04</td>
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<td>9.83</td>
<td>50</td>
<td>41.61</td>
<td>0.133</td>
<td>0.061</td>
<td>1.87</td>
<td>2.8</td>
</tr>
<tr>
<td>Sibyll2.1</td>
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<td>10.02</td>
<td>48.91</td>
<td>34.05</td>
<td>0.121</td>
<td>0.064</td>
<td>1.87</td>
<td>3.1</td>
</tr>
</tbody>
</table>

Table 6: Best-fit parameters (heavy composition scenario, powerful M87 and Fornax A).

<table>
<thead>
<tr>
<th></th>
<th>$a$</th>
<th>$\bar{g}_c$</th>
<th>$g_{cr}$</th>
<th>$M/F$</th>
<th>$g_{CygA}$</th>
<th>$g_{acc}$</th>
<th>$q$</th>
<th>$\chi^2$</th>
</tr>
</thead>
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<tr>
<td>EPOS-LHC</td>
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<td>0.078</td>
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<tr>
<td>QGSJetII-04</td>
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<td>44.20</td>
<td>0.170</td>
<td>0.065</td>
<td>1.84</td>
<td>2.4</td>
</tr>
<tr>
<td>Sibyll2.1</td>
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<td>2.14</td>
<td>44.96</td>
<td>31.50</td>
<td>0.147</td>
<td>0.072</td>
<td>1.78</td>
<td>2.4</td>
</tr>
</tbody>
</table>

Figure 16: Best-fit results to energy spectrum (left) and chemical composition (right) using QGSJetII-04 and the heavy composition scenario with a powerful M87 and Fornax A.

Figure 17: Best-fit results to energy spectrum (left) and chemical composition (right) using QGSJetII-04 and the light composition scenario with a powerful M87 and Fornax A.

the data point at $10^{19.85}$ eV is missed. Even though the fit-procedure has already chosen the
Figure 18: Angular power spectrum of CR events with $4 \leq E \leq 8$ EeV (left plot) and $E > 8$ EeV (right plot) for the light composition scenario with a powerful M87 and Fornax A using the QGSJetII-04 model. We show the cases for isotropized (solid and dash-dotted line) and non-deflected (dashed line) Cygnus A events.

maximal possible value of $g_{cr}^{MF}$, M87 and Fornax A together are still too weak to explain the CR flux below the three highest data points, as Centaurus A could. Hence, there is obviously no need in generating multiple realizations of the distribution function $\Phi$ of the unknown local sources.

Aside from the fit accuracy the resulting spectra expose features that refer to the transition from helium to CNO at $\sim 10^{18.9}$ eV and from CNO to iron at $\sim 10^{19.4}$ eV which have not been seen in the previous case due to the different $q$ value (note that for this single source spectrum we did use exponential cut-offs for all nuclear components, so these features are no mathematical artifacts). Thus, for $q > 1.9$ the observed CR spectra become completely dominated by the initial iron ejecta of the sources. The corresponding arrival directions of the best-fit model yields a similar but slightly worse — with respect to the observed dipole — picture compared to the case of a powerful Centaurus A (case 2): The resulting dipole of the multipole moments for CR energies $> 8$ EeV is slightly above the observed one — for completely isotropized Cygnus A events — but at CR energies $< 8$ EeV we obtain a very good agreement with isotropy when we exclude the initially light elements (H, He) from M87 and Fornax A (see Fig. 18).

The previously discussed disagreements with the data in terms of energy spectrum and composition only get worse when we suggest to provide either M87 or Fornax A with an individual parameter value of $g_{cr}$. Here, a good value for $g_{cr}$ within its plausibility limits will no longer be found, and the fit-procedure will try to compensate this by increasing $g_{acc}$ to the maximal value. This will lead to a wrong spectral behavior of the CR flux at the highest data points, as well as a stronger contribution by the average non-local contribution at the lowest energies. Since Fornax A and M87 possess an almost equal radio luminosity and distance to Earth, energy spectrum and chemical composition will expose the same deficits independent of which source is suggested to become powerful. The main difference between both scenarios is shown in the resulting arrival directions of the CRs and its corresponding angular power spectrum, which finds its explanation that the two sources have different positions and are thus UHECR trajectories are differently influenced by EGMF and GMF structures. As shown in Fig. 19, a powerful M87 source is in the case of the heavy composition scenario and completely isotropized Cygnus A events able to explain the observed dipole above 8 EeV, whereas the
scenario with a powerful Fornax A source is not.

Figure 19: Angular power spectrum with isotropized (solid line) and non-deflected (dashed line) Cygnus A events for $E > 8$ EeV using Sibyll2.1 and the light composition scenario. The left plot shows the case of a powerful CR source by M87 and the right plot the case of a powerful CR source by Fornax A.

Certainly one could discuss whether the results could be improved by further loosening parameter constraints, e.g., allowing $g_{\text{acc}}$ and $g_{\text{M/F}}$ to be fitted independently, but in total we consider the scenario of M87 and/or Fornax A being the main contributor(s) of the cosmic ray flux beyond the GZK cutoff for Cygnus A, as disfavored. Certainly, both can make a significant individual contribution which might become important when comparing UHECR results in the northern and southern hemisphere (like we discuss also for Cygnus A, see Sect. 4.1 and 5 — note that M87 is a northern source while Centaurs A and Fornax A are in the south), but we do not consider an accurate explanation of the Auger data plausible without a dominant contribution by Centaurus A.

5 Discussion

Radio galaxies are highly individual sources; both their intrinsic luminosities and their radio brightness as seen from Earth spread over many orders of magnitude. Given the physical constraints explained in Sect. 2 regarding their potential role as UHECR accelerators, these highly individual properties should reflect on total power and spectrum of their UHECR emission. Nevertheless, we present here results from a simplified model describing UHECR emission from radio galaxies, which involves maximally seven free parameters, adjustable in narrow ranges. But in spite of the simplifications we used, we think we can make three tentative conclusions regarding the possibility of radio sources to be the sources of UHECR, which we outline and discuss as follows.

The most solid conclusion we can draw is that an average contribution of all radio galaxies cannot explain both spectrum and composition of UHECR. This is a result of what we called the Lovelace-Hillas relation, which sets the maximal rigidity cosmic rays can reach in the acceleration process in relation to the radio luminosity of the source. Here of course, we need to note that our continuous source function has been derived by applying average values of cosmic-ray load, acceleration efficiency, and composition to all radio galaxies. We disregarded both the spread, plus potential correlations of these parameters between themselves or to observed properties of the radio galaxies. What would happen if we release this constraint?
Clearly, the main effect would be it would smear out our continuous source function to something which is no longer a sharply broken power law, but a smooth function with slowly steepening spectrum. How strong this effect can be is difficult to say on a qualitative base, but as only the flat end of this function would be able to provide good fits to the observed UHECR spectrum, this assumption would not evade our second conclusion.

The second conclusion is, that we need some outstanding, individual sources which explain the cosmic ray flux above the ankle. Here, two sources recommend themselves: Cygnus A, the by far brightest radio galaxy in the sky, and Centaurus A, the second-brightest radio galaxy and also the one nearest to Earth. The fact that we have such bright sources in our vicinity is already a cosmic coincidence, but additionally we need to assume that they have a cosmic ray load significantly above the average of the bulk of radio galaxies. While this can easily be argued for Cygnus A, which is the by far the most powerful radio galaxy in our cosmic environment, it is difficult to argue for Centaurus A — this is just a regular radio galaxy, and is only distinguished by its proximity to Earth. Clearly, we do not need to assume that Centaurus A is the only local source with such special properties, but it would need to be one of maximally 25% of the sources — a number which stems from assuming the extreme scenario that all other radio galaxies fall in two strict classes, in which one class has the cosmic-ray load as Centaurus A, and the other not producing any UHECR. Among these 25%, there could be the other two exceptionally bright local radio galaxies, M87 and Fornax A, but we showed that they cannot solve the issue alone — Centaurus A is needed.

The third conclusion is that Cygnus A needs to provide a predominantly light (meaning: near to solar) composition, while Centaurus A would need to have a heavy composition, with an iron fraction comparable to protons at a given cosmic-ray energy. Only this way, the increase of heaviness seen in the Auger data can be explained. We argued already that if the injection process into first-order Fermi acceleration is purely rigidity dependent (as the acceleration process itself), and since FR-II galaxies like Cygnus A can be expected to feed their powerful jets exclusively from the accretion disk of the supermassive black hole, which again are expected to have abundances not far from solar, one would indeed expect a light composition from Cygnus A. On the other hand, for FR-I galaxies where the much weaker jets dissipate inside the galactic environment, one would assume that they mix much stronger with heavy material present in their ISM. Indeed, such scenarios have been proposed for Centaurus A based on observations of the jet environment [76].

A necessary condition for all this to work is that UHECRs from Cygnus A are largely isotropized by the EGMF. This may be considered reasonable, at least for EGMF models which propose rather high filling factors of field strengths close to 1 nG or above [17]. But it is not only necessary to assume this to make our predictions consistent with the anisotropy limits seen in the arrival-direction distribution, rather, it is the condition to allow Auger actually to see UHECR from Cygnus A. The angle between the zenith of the Auger experiment and Cygnus A is minimally $\approx 75^\circ$, so it is beyond the limit of zenith angles Auger usually takes into account. Only deflected cosmic rays from Cygnus A can therefore be included in the Auger spectrum. With our simulation setup, we cannot quantify how strong these deflections are, but certainly they will be to some degree energy dependent — higher energy particles are less strongly deflected as less energetic particles. This will lead to a steepening of the spectrum from Cygnus A measured by Auger towards higher energies, which, as we discussed in Sect. 4.1, may be compensated by choosing a higher value of $g_{\text{acc}}^{\text{CygA}}$ as also expected from theory. An interesting question would therefore be how this affects the spectra which would be measured by a northern observatory, like TA. First, at some sufficiently
high energy \( \gtrsim 30 \text{ EeV} \), where a complete isotropization of light nuclei in the EGMF becomes unlikely, TA should see more light nuclei from Cygnus A than Auger, which may cause subtle differences in spectrum and composition seen by these experiments. Second, we would expect that TA sees some anisotropy in the direction of Cygnus A in this energy regime, but note that EGMF and GMF effects can certainly not be completely neglected at \( \sim 30 \text{ EeV} \). We also note that, although there is no significant sign of Cygnus A in the arrival-direction analysis already published by TA [107], this does not contradict our expectation, as this analysis was performed at energies \( > 57 \text{ EeV} \) where Cygnus A is already in its GZK cutoff and thus hardly contributes.

Certainly, there is lot of room to discuss the plausibility of the scenario we present here. The interesting point is, however, that we manage with our simplified 7-parameter model (a) to provide an excellent fit to the UHECR spectrum, (b) to reproduce the heaviness of chemical composition and their trend as a function of energy, (c) to match the strength of the dipole in arrival directions, and (d) to remain consistent with the limits on anisotropy for higher multipoles — all measured by the Pierre Auger Observatory. If we take this success as a reason to assume that radio galaxies, with the properties as assumed in our best-fit scenario, are indeed the sources of UHECR, two interesting corollaries would follow.

First, the contribution of the rest of the radio galaxies (i.e., those except Cygnus A and Centaurus A), although irrelevant as CR sources at ultra-high energies, would have a spectrum so steep that towards lower energies it has the potential to reach the level of the measured cosmic-ray flux at a break energy around \( 10^{17} \text{ eV} \). Thus, it would have a shape which perfectly matches the feature known as the "second knee". It would thus add radio galaxies to the various scenarios which predict a strong contribution of extragalactic cosmic rays below the ankle [108]. We did not check this scenario further as it is out of the scope of this paper, but it is an intriguing possibility to follow.

Second, as we have shown that indeed the predictions of the continuous source scenario are consistent with the sum of individual sources in our local universe, we need to conclude that if we are able to determine the average UHECR flux over a very large volume, it should match the one predicted by the continuous source function. But as its contribution lies below the UHECR flux measured at Earth (at least in the energy range above the ankle), it means in turn that this local UHECR flux is roughly one order of magnitude higher than the average UHECR flux throughout the universe. One immediate conclusion would be that predicted cosmogenic neutrino fluxes (e.g. [109–111]) would be overestimated by an order of magnitude, as they all rely on the assumption that our measured UHECR flux is typical for the universe. Furthermore, it would also reduce the slight tension that predicted cosmogenic photon fluxes show in the case of models with a strong source evolution [112].

Finally, we want to note two conclusions which affect the modeling of the UHECR origin in general. Firstly, our model is based on what we call the Lovelace-Hillas relation, which puts the highest rigidity of accelerated particles in relation to the radio luminosity of the source, and which has been used to argue that UHECR sources must be very numerous transients. The simple reason why we manage to remain within the isotropy limits although we propose that essentially all measured UHECR come from just two steady sources, is that one of these sources is far away enough that extragalactic magnetic fields can reasonable isotropize its contribution [46], and the other one is heavy enough that the Galactic magnetic field is able to do the job.

Secondly, we obtain all our fits for a power-law spectrum approximately \( \propto E^{-1.8} \), thus for an index which is quite reasonable and expected from first-order Fermi acceleration at
weakly relativistic shocks. This questions the necessity of the currently fashionable Peter’s cycle models, and shows that there is no need to go to extremely hard injection spectra (e.g. [113]) to still get a good fit of the data. Our result is even slightly softer (closer to theoretical expectations for first-order Fermi acceleration) than the already good results found by Wittkowski and Kampert [88].

6 Conclusions and Outlook

We examined the possibility that radio galaxies are the source of the observed ultra-high energy cosmic rays (UHECRs). We performed 3D simulations using the CRPropa3 code [48] together with the extragalactic magnetic field obtained in constrained structure-formation simulations by Dolag et al. [47], which represents the structure of our local universe out to a distance of 120 Mpc. Within this volume, we used individual radio galaxies taken from the catalog of van Velzen et al. [95], for which we derived a parametrized model for their UHECR emission based on their radio luminosity. For the contribution outside this volume, we derived a continuous source function for UHECRs from radio galaxies using the same physical relations as for the individual sources and a radio-luminosity function provided by Mauch and Sadler [63], and determined the measured UHECR flux by 1D simulations with CRPropa3. To complete the model, we needed further to include the ultra-luminous radio galaxy Cygnus A, which is expected from its radio brightness to provide about half of all observed UHECRs within our model. As our setup did not allow a detailed propagation calculations for Cygnus A, it was treated as a point source neglecting scattering in the extragalactic magnetic field. In the final analysis of the results, we allowed then that the contribution by Cygnus A is isotropized to mimic the likely effect of extragalactic magnetic fields on a source at this distance [46]. As we use a physical model to link UHECR emission power and spectrum to the radio luminosity of the sources, our model is absolutely normalized in flux for all UHECR components, and we present an approach how to implement this absolute normalization into a Monte Carlo simulation.

The results of our simulations were then fitted to data published by the Pierre Auger Observatory for spectrum, composition, and arrival-direction distribution of UHECRs, using a re-weighting mechanism common for fitting Monte Carlo results with the residual minimizer contained in the lmfit python package. To keep the parameter space manageable, we constrained our fits to several plausible scenarios, in which usually the parameters for all sources except a few were set equal. Adopting this necessarily simplified approach, we could obtain two main conclusions:

(1) The continuous source function of UHECRs from radio galaxies cannot explain the spectrum and composition of cosmic rays above the “ankle” of the spectrum at $\sim 5 \times 10^{18}$ eV. However, we found from our fits that this component, which represents a broken power law with a break at about $3 \times 10^{17}$ eV might have the potential to explain part of the cosmic-ray spectrum between the “second knee” (at $\sim 10^{17}$ eV) and the ankle. This requires further investigation.

(2) The spectrum and composition above the ankle can be explained by the contribution of the two brightest radio galaxies seen from Earth, Cygnus A at a distance of about 250 Mpc, and Centaurus A at a distance of about 4 Mpc. Here we have to assume that (a) both sources have a baryonic load in their jets which is significantly above the average of radio galaxies; (b) that Cygnus A provides a predominately light composition,
not far from solar abundances of accelerated cosmic rays compared at the same rigidity, while Centaurus A provides a heavy composition, in which the abundance of iron nuclei is enhanced roughly to the level of hydrogen.

A requirement for this scenario not to violate isotropy constraints is that extragalactic magnetic fields are strong enough to largely isotropize UHECR arriving from Cygnus A. We did check several other scenarios, for example for the radio galaxies M87 and/or Fornax A being the dominant local source(s) instead of Centaurus A, but found none of them more plausible than assuming a dominant contribution of Centaurus A. We deliver a critical discussion of the assumptions required for this scenario to work, and also mention some intriguing corollaries which might follow from it, in particular a possible inhomogeneity of the UHECR density throughout the universe, with the locally measured flux being about one order of magnitude above average.

Eventually, we presented a physics-based, absolutely normalized, 7-parameter model which is able to explain all currently-known features of UHECR. We consider this result a clear demonstration that radio galaxies still present very good candidates for an explanation of the origin of UHECR. To further substantiate this claim, we need to overcome several simplifications which we applied to obtain our result within the technical limitations present. In particular, we see the need to set up a simulation and optimization procedure which fully accounts for the astrophysical properties of the model, which means (i) performing 3D, or better 4D simulations at least out to the distance of Cygnus A (250 Mpc), ideally out to the magnetic horizon at about 500 Mpc [43], (ii) investigating several possible scenarios of magnetic-field structure out to this distance, and (iii) applying a Bayesian optimization method which works for high-dimensional parameter spaces, and which would thus allow to treat at least all radio galaxies which could make a measurable, individual contribution with an own parameter set. While by the time this work was started, such a setup still seemed illusionary, it will become possible in the near future thanks to several developments:

(a) The availability of detailed reconstructions of the matter distribution of the local universe out to about 400 Mpc distance [114, 115], which could be used for a phenomenological, parametrized construction of extragalactic magnetic fields with various field strengths and volume filling factors as demonstrated by Alves Batista et al. [46].

(b) The development of efficient targeting methods in CRPropa which could speed up 3D and 4D simulations by factors of 10 to $10^7$ [116] compared to the still common $4\pi$ emission.

(c) The development and computational implementation of very efficient Bayesian Markov Chain Monte Carlo methods in the last years (e.g. PyMC3 [117]), which are proven to be able to handle very high dimensional parameter spaces.

The next, straightforward step, however, will be to test the model developed here as a possible extragalactic explanation of cosmic rays below the ankle, for which we expect to be able to present results very soon.

Acknowledgments

This work most notably benefits from the development of CRPropa3 and in particular from the support by David Walz, Gero Müller and Nils Nierstenhöfer. We are grateful to Günter
Validation of the flux normalization

Within this work the radio luminosities are used to normalize the observed UHECR flux from the simulations. To validate the normalization suggested in Sect. 3.4 we first made a simplified simulation with a well known theoretical expectations to show that the local flux contribution according to Eq. (3.17) provides the expected behavior. Afterward, the non-local flux contribution according to Eq. (3.24) is determined in the local region with \( d \leq 120 \) Mpc and compared with the previously validated local contribution.

A.1 Local contribution

A three-dimensional simulation without any particle interactions or deflections is performed using a random subset of the luminous sources of the vV12 catalog. Here, the source \( j \) provides a radio luminosity \( L_j \) and a corresponding CR power \( Q_{cr,j} \) between \( \tilde{R}_0 \) and \( \tilde{R}_1 \) according to Eq. 2.8. Within the UHE regime between \( \tilde{R}_1 \) and \( \tilde{R}_ \) we obtain

\[
Q_{\text{uhecr},i,j} = \frac{f_i Z_i Q_{\text{cr},j} G_a(\tilde{R}_1) - G_a(\tilde{R}_\text{cut})}{G_a(\tilde{R}_0) - G_a(\tilde{R}_\text{cut})}
\]

with

\[
G_a(x) = -\int dR \, R^{1-a} \exp(-R/\tilde{R}) = \tilde{R}^{2-a} \Gamma \left( 2-a, x/\tilde{R} \right)
\]

where \( \Gamma(c, x) \) denotes the incomplete gamma function. The first term in Eq. (A.1) accounts for the distribution of the total CR power between the different types of nuclei \( i \) with respect to the initial abundances \( f_i \). The second term determines the proportion of the CR power that goes into UHECRs with a spectral behavior according to the re-weighted energy spectrum under the assumption that the spectral behavior does not change between \( \tilde{R}_0 \) and \( \tilde{R}_1 \). Thus, the theoretically expected total energy flux \( J_{i,j}^{(th)} \) for an extended, spherical observer with a radius \( r_{\text{obs}} \), surface \( S_{\text{obs}} \) and cross-section \( \sigma_{\text{obs}} \) at a distance \( d_j \gg r_{\text{obs}} \) yields

\[
J_{i,j}^{(th)} = \frac{Q_{\text{uhecr},i,j} \sigma_{\text{obs}} S_{\text{obs}}}{4\pi d_j} \approx \frac{1}{16\pi d_j^2} \frac{f_i Z_i Q_{\text{cr},j} G_a(\tilde{R}_1) - G_a(\tilde{R}_\text{cut})}{G_a(\tilde{R}_0) - G_a(\tilde{R}_\text{cut})}.
\]

It should be kept, that the used cross-section \( \sigma_{\text{obs}} = \pi r_{\text{obs}}^2 \) is only valid for parallel incident CRs, which is almost the case for \( d_j \gg r_{\text{obs}} \).

The total energy flux \( J_{i,j}^{(sim)} \) due to the normalization (3.17) is given by

\[
J_{i,j}^{(sim)} = \sum_{k,i,j} E_k^* W_{k,i,j} = \sum_{k,i,j} \frac{E_k^* f_i Q_{\text{cr},j} \left( \tilde{R}_1^{1-a} - \tilde{R}_0^{1-a} \right) R_{k}^{1-a} \exp(-R_k/\tilde{R})}{4\pi r_{\text{obs}}^2 \tilde{N}_{i,j} Z \left( 1-\tilde{a} \right) \left( G_a(\tilde{R}_0) - G_a(\tilde{R}_\text{cut}) \right)}.
\]
Using 100 random sets of the parameters \(a, f_i, g_{cr}, g_{acc}\) the corresponding relative error

\[
\Delta_{i,j} = 100 \times \frac{|J_{i,j}^{(sim)} - J_{i,j}^{(th)}|}{J_{i,j}^{(th)}}
\]  (A.5)

for a single source as well as the total relative error \(\sum_j \Delta_{i,j}\) for multiple sources is shown in the Fig. 20. Here, the error increases in the case of smaller statistics but predominantly stay below a few percentage level yielding a good validation of the normalization of the local contribution.

**A.2 Non-local contribution**

The contribution of non-local sources is only taken into account for candidates with a propagation distance \(> 120\, \text{Mpc}\). Nevertheless, we simulated CR candidates in this 1D approach with \(0 < d \leq 120\, \text{Mpc}\) in order to compare their resulting flux with the local flux contribution from the 3D approach. Since we cannot account for deflections in 1D and the cosmic variance can cause a local excess or lack of CRs, we do not expect the results to be match exactly, but in a reasonable range. In order to avoid differences by the different types of cutoff — exponential in the 3D case and by a Heaviside function in the 1D case — we use \(g_{cr} = 50, g_{acc} = 0.6\) so that \(\hat{R}_{cut} = \hat{R}\) and both spectra have a sharp cutoff. Fig. 21 displays the resulting energy flux from candidates that originate between \(5\, \text{Mpc} - 120\, \text{Mpc}\) using the 3D approach and a subsequent normalization with Eq. (3.17) and the 1D approach normalized with Eq. (3.24), respectively.

Here, the shown energy flux

\[
J_{\nu,\kappa}^{(1D)} = \frac{E_{\kappa}}{dE_{\kappa}} \sum_k \delta_{\kappa\nu} W_{k,i}
\]  (A.6)
from the 1D simulation and
\[
J_{i,\kappa}^{(3D)} = \frac{E_{\kappa}}{dE_{\kappa}} \sum_{k,j} \delta_{k\kappa} W_{k,i,j} \quad (A.7)
\]
from the 3D simulation of the different initial nuclei type \(i\) in the energy bin \(E_{\kappa}\) have the same spectral behavior as well as a similar absolute magnitude. The selected distance range is chosen so that the contribution by Centaurus A, at a distance of 3.6 Mpc, gets excluded, otherwise \(J_{i,\kappa}^{(3D)}\) will significantly exceed \(J_{i,\kappa}^{(1D)}\). But due to the close distance to Earth and the high radio luminosity it is most likely that Centaurus A causes a local excess of UHECRs. To ensure, that the relative error
\[
\Delta_{i,\kappa} = 100 \times \frac{|J_{i,\kappa}^{(1D)} - J_{i,\kappa}^{(3D)}|}{J_{i,\kappa}^{(3D)}} \quad (A.8)
\]
is independent of the chosen parameters \(a\), \(f_i\) and \(m\), we draw 100 random sets of these parameters and determine the corresponding \(\Delta_{i,\kappa}\). In addition, we also change the spatial region the candidates need to originate between \(\tilde{d} \in [5, 30] \text{ Mpc}\) and \(\tilde{d} \in [100, 120] \text{ Mpc}\), to account for cosmic variance. As shown in the Fig. 22 the 1D flux hardly differs from the 3D flux by more than a few hundred percent and the mean is well below. So we conclude that the normalization (3.24) yields the correct CR flux for the non-local source contribution.

B Further details of the simulation setup

The results shown in the Sect. 4 are based on the following simulation setup:

(i) The observer is located at \(\vec{r}_{\text{obs}} = (118.34, 117.69, 119.2) \text{ Mpc}\) with a radius of 0.1 Mpc.
Figure 22: Difference of the 1D and the 3D flux normalization in the local regime using 100 random sets of parameters: (i) Relative error $\Delta = \sum_i \Delta_{i,\kappa}$ dependent on the energy bin $E_\kappa$ (left plot); and (ii) relative error $\Delta = \sum_\kappa \Delta_{i,\kappa}$ dependent on the initial nucleus type $i$ (right plot).

(ii) The sources are located according to the catalog of RGs from van Velzen et al. [95] with a maximal distance to the observer of 120 Mpc. In total, $N_{RG} = 838 = 121 + 100 + 617$ different source positions are used within 120 Mpc where all RGs are considered as a point source. In addition, we take the average CR contribution by the non-local sources beyond 120 Mpc into account as well the the impact of an extraordinarily powerful radio source, like Cygnus A, at a distance of 255 Mpc.

(iii) The ‘extragalactic’ propagation in the case of the local sources is performed with the DeflectionCK() module from CRPropa3, where the equation of motion is solved by a Runge-Kutta integration method. Thereby, we used a local accuracy of $10^{-3}$ between the fourth and the fifth order of the algorithm, a minimal step of 1 kpc and a maximal step of 1 Mpc. Further, we use the magnetic field data from Dolag et al. [47], which is stored in a multi-resolution grid according to the QuimbyMagneticField() module [93]. These EGMF data are intended to become part of the CRPropa program and are currently available at https://forge.physik.rwth-aachen.de/public/quimby/mhd/. In addition, we account for nuclear decay as well as electron pair production, photo-pion production and photo-disintegration due to the cosmic microwave background (CMB) and the UV/optical/IR background (IRB).

To decrease the necessary CPU time for the propagation, we use the EmissionMap() module from CRPropa3. Here, three prior simulations are made, where each simulation generates a so-called emission map for the individual sources and the initial candidate types with a certain angular resolution in the spherical coordinates $\theta$ and $\phi$. Each emission map provides information on the number of candidates that hit the observer sphere dependent on its initial momentum. By using this emission map (in the subsequent simulation) a CR candidate will immediately be rejected from the propagation, if its initial momentum has not provided any hit according to the map. In order to optimize the resolution of the final emission map, we run three simulation with a gradually decreasing observer size, in which each simulation uses the previously generated emission map and generates a new one with a higher resolution. The final emission map of each source and candidate type has an angular resolution of $10^\circ$ in $\theta$ and $\phi$. 
In the case of the non-local sources, no realistic EGMF structure is available, so that the SimplePropagation() module is used, where CR candidates keep their initial direction of momentum. CR candidates from the single non-average source at 255 Mpc are simulated in three spatial dimensions like the local sources, whereof each source emits UHECR candidates isotropically. However, the average non-local contribution is simulated in one spatial dimension, where the sources are homogeneously distributed in 1D according to the SourceUniform1D module. Here, all CR candidates will reach the observer, if their energy does not drop below the minimal energy of 3 EeV that is used for all simulations. Further, the effects due to the redshift evolution of the photon target are taken into account.

(iv) The ‘galactic’ propagation is performed with the PARSEC code [122] in order to obtain the arrival directions of the UHECR candidates at the Earth. Here, only the deflection by the galactic magnetic field from the JF12 model, its regular and turbulent component, is taken into account. The magnetic lensing according to the PARSEC code is already part of the CRPropa program.

References


