Abstract A search for neutral heavy resonances is performed in the $WW \rightarrow e\nu\mu\nu$ decay channel using $pp$ collision data corresponding to an integrated luminosity of $36.1 \text{ fb}^{-1}$, collected at a centre-of-mass energy of 13 TeV by the ATLAS detector at the Large Hadron Collider. No evidence of such heavy resonances is found. In the search for production via the quark–antiquark annihilation or gluon–gluon fusion process, upper limits on production via the quark–antiquark annihilation or gluon–gluon fusion process, constraints are also obtained on these resonances, as well as on a Higgs boson in the Georgi–Machacek model and a heavy tensor particle coupling only to gauge bosons.

1 Introduction

The measured properties [1–4] of the Higgs boson discovered in 2012 by the ATLAS [5] and CMS [6] collaborations at the Large Hadron Collider (LHC) are, within experimental uncertainties, consistent with those predicted for the Standard Model (SM) Higgs boson, $h$. Nevertheless, the SM is thought to be an incomplete theory and many scenarios beyond the SM (BSM) predict an extended Higgs sector [7,8]. Diboson vector and tensor resonances are also predicted in several other extensions to the SM, such as in composite Higgs models [9,10] and models with warped extra dimensions [11–14].

This article reports on the results of a search for heavy neutral resonances decaying into two $W$ bosons, which then decay into the $e\nu\mu\nu$ final state, either directly or via leptonic tau decays with additional neutrinos. The analysis is based on the full $pp$ collision dataset collected by the ATLAS detector in 2015 and 2016 at the centre-of-mass energy of $\sqrt{s} = 13$ TeV, corresponding to an integrated luminosity of 36.1 $\text{ fb}^{-1}$.

The results are interpreted in terms of different benchmark models. For the case of a scalar resonance produced by gluon–gluon fusion (ggF) or vector-boson fusion (VBF), two scenarios with different intrinsic widths are considered. Constraints on the heavy neutral scalar in two-Higgs-doublet models (2HDM) are also obtained. The neutral member of the fiveplet in the Georgi–Machacek (GM) model [15,16] also serves as a reference model in the VBF production mode. The parameterisation of heavy vector triplet (HVT) Lagrangians [17,18] permits the interpretation of searches for spin-1 resonances in a generic way. The bulk Randall–Sundrum (RS) model [11,19] features a spin-2 Kaluza–Klein (KK) graviton excitation ($G_{KK}$) decaying into $WW$, while a tensor resonance signal in the VBF production mode is based on an effective Lagrangian model (ELM) [20].

A previous search for a heavy Higgs boson in the $e\nu\mu\nu$ final state was performed by ATLAS [21] based on a data sample with an integrated luminosity of $20.3 \text{ fb}^{-1}$ at $\sqrt{s} = 8$ TeV. The CMS Collaboration also published a search for a high-mass scalar decaying into two $W$ bosons in the fully leptonic final state [22], using datasets at $\sqrt{s} = 7$ and 8 TeV with integrated luminosities of 5.1 and 19.5 $\text{ fb}^{-1}$, respectively. A search for heavy resonances in the RS models in the leptonic decays of the $WW$ channel, using a dataset of 4.7 $\text{ fb}^{-1}$ at 7 TeV [23], was reported by the ATLAS Collaboration. The ATLAS and CMS collaborations have obtained constraints on the HVT and bulk RS models, based on other decay modes of the $VV$ channels, with $V$ being either a $W$ or a $Z$ boson [24–36]. The search in the $e\nu\mu\nu$ decay mode is complementary to searches performed in other decay modes. In particular, the sensitivity to low mass resonances is higher in the fully leptonic final state than in final states that include jets due to background from jet production.

The article is organised as follows. Section 2 presents the various models used in this analysis. Section 3 describes the ATLAS detector. The data and simulated event samples are discussed in Sect. 4. The event reconstruction and selection...
are described in Sects. 5 and 6, respectively, followed by the background estimation techniques in Sect. 7. Systematic uncertainties are discussed in Sect. 8 and the results are presented in Sect. 9. Finally, the conclusions are given in Sect. 10.

2 Theoretical models

The different signal models studied are presented in Table 1. One scenario for the heavy scalar assumes that the scalar has a width much smaller than the detector resolution. This is referred to as the narrow-width approximation (NWA). Larger widths (large-width assumption, LWA) of 5, 10 and 15% of the heavy Higgs boson mass, are also considered. The choice of the width range for the heavy Higgs boson is motivated by the fact that, for several of the most relevant BSM models, widths above 15% are already excluded by indirect limits [37].

The 2HDM comes in different types [38], defined by assumptions about the couplings of each of the Higgs doublets and the discrete symmetries imposed. This analysis considers Type I, where one Higgs doublet couples to vector bosons while the other couples to fermions, and Type II of the minimal supersymmetric (SUSY)-like model in which one Higgs doublet couples to up-type quarks and the other to down-type quarks and charged leptons. This analysis uses a generic charge-conjugation- and parity-conserving (CP-conserving) 2HDM with a softly broken $Z_2$ symmetry [38] which has several free parameters: (i) four masses $m_h$, $m_H$, $m_A$ and $m_{H^\pm}$ for the two CP-even neutral states, the pseudo-scalar and the charged Higgs boson pair, respectively, (ii) a mixing angle $\alpha$ between the CP-even neutral Higgs fields, and (iii) the ratio of the vacuum expectation values of the two Higgs doublets $\tan\beta = \nu_2/\nu_1$. The benchmark model is defined by setting $m_h = 125$ GeV and the masses of the supersymmetric particles heavy enough so that Higgs boson decays into SUSY particles are kinematically forbidden. The cross sections and branching fractions are calculated with SusHi and 2HDMC [39,40].

The GM model extends the Higgs sector with the addition of a real and a complex triplet of SU(2)$_L$ in a way which preserves the SM value of $\rho = M_W^2/(M_Z^2 \cos^2 \theta_W) = 1$ at tree level, with $m_W$, $m_Z$ and $\theta_W$ being the $W$ and $Z$ boson mass and the weak mixing angle, respectively. The physical states include a fermiophobic fiveplet, $H^0_5$, $H_5^{\pm}$ and $H^+_5$, of custodial SU(2) symmetry which couples preferentially to vector bosons [41]. For that reason, the GM model is less constrained [42], when produced by the VBF process, than other standard benchmark models of a triplet Higgs field, such as the little Higgs model [43] or the left–right symmetric model [44]. The model has many parameters [45,46], but, if the other new Higgs bosons are heavier than those of the $H_5$ multiplet, the only production mode is via the VBF process. The cross section and decay width into $VV$ are then proportional to a single parameter, $\sin^2 2\beta_H$, which characterises the fraction of the gauge boson masses generated by the triplet Higgs fields.

The HVT Lagrangian [18] parameterises the couplings of the new spin-1 heavy bosons to SM particles in a generic manner and allows their mixing with SM gauge bosons. The $s$-channel production mechanism of the heavy gauge bosons is primarily via $q\bar{q}$ annihilation ($qqA$). The HVT bosons couple to the Higgs boson and SM gauge bosons with coupling strength $c_{h,gV}$ and to the fermions with coupling strength $g^2 c_F/g_V$, where $g$ is the SM SU(2)$_L$ gauge coupling, $c_h$ and $c_F$ are multiplicative factors that modify the couplings to the Higgs boson and to the fermions, and $g_V$ represents its coupling strength to the $W$ and $Z$ bosons. For the case of vector-boson fusion, it is assumed that there is no coupling to fermions so that non-VBF production processes are suppressed.

The spin-2 $G_{KK}$ is the first Kaluza–Klein excitation of the graviton in the RS model with a warped extra dimension [11,19], where the SM fields are localised in the bulk [12–14]. This model is characterised by the dimensionless coupling constant $k/M_{Pl} \sim O(1)$ where $k$ determines the curvature of the space, and where $M_{Pl} = M_{Pl}/\sqrt{8\pi}$ is the reduced Planck scale.

For the VBF production mode, the spin-2 signal is based on an effective Lagrangian approach with $\Lambda$ as a characteristic energy scale of the underlying new physics [20],

$$\mathcal{L} = \frac{1}{\Lambda} T_{\mu\nu} \left( f_1 B_{\mu\nu}^a B^a_{\mu\nu} + f_2 W_{\mu\nu}^a W^a_{\mu\nu} + 2f_3 (D^\mu \Phi)(D^\nu \Phi) \right).$$

Here, $f_i$ are variable coupling parameters, $T_{\mu\nu}$ is the spin-2 singlet field, $B_{\mu\nu}^a$ and $W_{\mu\nu}^a$ are the electroweak field strength...
tensors, and $\Phi$ is the scalar Higgs field. The covariant derivative $D^\mu = \partial^\mu - ig W_i^\mu \sigma^i / 2 - ig' Y B^\mu$, where $\sigma^i$ are the Pauli matrices, $Y$ the weak hypercharge, and $g$ and $g'$ the corresponding gauge coupling constants. The model differs from the RS model in that the couplings to fermions or gluons are not included in the Lagrangian. Also, the BSM amplitude is multiplied by a form factor which is a function of a cut-off scale $\Lambda_{\text{ff}}$ and a suppression power $n_{\text{ff}}$ in order to preserve unitarity at high energies:

$$f(p_1^2, p_2^2, k_{\text{sp}2}^2) = \left( \frac{\Lambda_{\text{ff}}}{|p_1^2| + \Lambda_{\text{ff}}^2}, \frac{\Lambda_{\text{ff}}^2}{|p_2^2| + \Lambda_{\text{ff}}^2}, \frac{\Lambda_{\text{ff}}^2}{|k_{\text{sp}2}^2| + \Lambda_{\text{ff}}^2} \right)^{n_{\text{ff}}}$$

where $p_1^2$ and $p_2^2$ are the squared invariant masses of the incoming electroweak bosons and $k_{\text{sp}2}^2$ is the squared invariant mass of the sum of the initial boson momenta, equivalent to that of an $s$-channel spin-2 particle. The specific parameter settings for the signal models used are given in Sect. 4.

3 ATLAS detector

The ATLAS detector [47,48] is a general-purpose particle detector used to investigate a broad range of physics processes. It includes an inner tracking detector (ID) surrounded by a thin superconducting solenoid, electromagnetic and hadronic calorimeters and a muon spectrometer (MS) incorporating three large superconducting toroidal magnets with eight coils each. The ID consists of fine-granularity silicon pixel and microstrip detectors, and a straw-tube tracker. It is immersed in a 2 T axial magnetic field produced by the solenoid and provides precision tracking for charged processes. It includes an inner tracking detector (ID) surrounded by a thin superconducting solenoid, electromagnetic and hadronic calorimeters and a muon spectrometer (MS) incorporating three large superconducting toroidal magnets with eight coils each. The ID consists of fine-granularity silicon pixel and microstrip detectors, and a straw-tube tracker. It is immersed in a 2 T axial magnetic field produced by the solenoid and provides precision tracking for charged particles in the range $|\eta| < 2.5$, where $\eta$ is the pseudorapidity of the particle. The straw-tube detector also provides transition radiation measurements for electron identification. The calorimeter system covers the pseudorapidity range $|\eta| < 4.9$. It is composed of sampling calorimeters with either liquid argon (LAr) or scintillator tiles as the active medium, and lead, steel, copper, or tungsten as the absorber material. The MS provides muon identification and momentum measurements for $|\eta| < 2.7$. The ATLAS detector has a two-level trigger system [49] to select events for further analysis.

4 Data and simulation samples

The data used in this analysis were collected with a single-electron or single-muon trigger. These triggers have a transverse energy or momentum threshold, $E_T$ or $p_T$, that depends on the data-taking period, with the lowest threshold varying between 20 and 26 GeV. The trigger efficiency for $WW$ events passing the offline event selection (Sect. 6) is greater than 99%. Data quality criteria are applied to ensure that events are recorded with stable beam conditions and with all relevant subdetector systems operational.

Samples of simulated signal and background events are used to optimise the event selection and to estimate the signal acceptance and the background yields from various SM processes.

The sample for the NWA heavy Higgs boson signal was produced with POWHEG-BOX2.0 [50–52] which calculates separately the ggF [53] and VBF [54] production mechanisms with matrix elements up to next-to-leading order (NLO) in quantum chromodynamics (QCD). It uses the CT10 NLO parton distribution function (PDF) set [55] and is interfaced with PYTHIA 8.186 [56] for the $H \to WW$ decays, for parton showering and hadronisation. A set of tuned parameters called the AZNLO tune [57] is used to describe the underlying event. The NWA Higgs boson is generated with a width of 4 MeV. This event sample is also used to constrain the 2HDM. The LWA heavy Higgs boson signal was simulated at NLO using the MADGRAPH5_AMC@NLO 2.3.2 event generator [58] with the NNPDF23LO PDF set [59]. The generated particles at matrix element level are showered by PYTHIA 8.186 with the A14 tune [60] for the underlying event. The mass of the heavy Higgs boson consid- ered in this analysis spans the range between 200 GeV and $4 (3)$ TeV for the ggF-induced (VBF-induced) signals. Both NWA and LWA samples were generated in steps of 100 GeV up to 1 TeV, and in steps of 200 GeV thereafter.

The POWHEG-BOX samples describe the production of a ggF-induced heavy Higgs boson in association with one jet at leading-order (LO) precision, while further jets are emu- lated by the parton shower generator, PYTHIA. A more precise calculation of higher jet multiplicities is provided by using MADGRAPH5_AMC@NLO 2.3.2 to simulate $gg \to H$ events in association with up to two jets at NLO precision. Here, the overlap between identical final states generated at the matrix element (ME) and the parton shower (PS) stage is removed using FxFx merging [61]. The fraction of ggF events passing the event selection requirements of the $N_{\text{jet}} = 1$ and $N_{\text{jet}} \geq 2 VBF$ categories (defined later in Sect. 6) predicted by the POWHEG-BOX event generator is reweighted to match that of the MADGRAPH5_AMC@NLO FxFx samples. The corresponding scale factors are calculated for several hypothetical heavy Higgs boson masses. It is the largest, 1.14, for the 200 GeV mass point, and decreases with increasing res-
onance mass to a value of 0.85 for the 4 TeV mass point, for the \( N_{\text{jet}} = 1 \) VBF category. The corresponding numbers are 0.91 and 0.73 for the \( N_{\text{jet}} \geq 2 \) VBF category.

Benchmark samples for the GM, HVT and bulk RS models were generated at LO using MADGRAPH5_AMC@NLO interfaced to PYTHIA 8.186 with the NNPDF23LO PDF set. A value of \( \sin \theta_H = 0.4 \) is chosen for the GM benchmark model. For the HVT interpretation in the \( q \bar{q} \) annihilation mode, samples were generated according to the extended gauge symmetry model \( A \) [18] with \( g_V = 1 \). In the VBF mode, samples were generated using the same \( g_V \) value but setting the couplings to the fermions to zero so that the new vector boson couples only to the SM vector and Higgs bosons. For the bulk RS model, a curvature scale parameter \( k/\tilde{M}_p \) of either 0.5 or 1 is considered. The ELM VBF spin-2 signals were generated at LO with VBFNLO 3.0.0 beta 2 [62] with the NNPDF30LO PDF set [63] and using the following parameter setting [20]: \( \Lambda_{fj} = 3 \text{ TeV}, \, n_{ff} = 4 \), \( \Lambda = 1.5 \text{ TeV} \) and \( f_1 = f_2 = f_3 = 1 \). The mass range considered is between 200 GeV and 5 TeV for the KK graviton signal, between 250 GeV and 5 TeV for the HVT qqA signal, between 200 GeV and 1 TeV for the GM and ELM VBF signals, and between 300 GeV and 1 TeV for the HVT VBF signal.

The main sources of SM background include events from the production of single top quarks, \( t \bar{t} \), dibosons \( (WW, WZ \text{ and } ZZ) \), \( Z/\gamma^{*}+\text{jets} \) and \( W+\text{jets} \). Single-top-quark simulated events were generated with POWHEG-BOX 2.0 [64,65] using the CT10 NLO PDF set interfaced to PYTHIA 6.428 [66] for parton showering and hadronisation, with the PERUGIA2012 tune [67] and CTEQ6L1 PDF [68] to describe the underlying event. The \( t \bar{t} \) events were generated with POWHEG-BOX 2.0 [69] using the NNPDF30NLO PDF set [63] interfaced to PYTHIA 8.186 for parton showering and hadronisation, with the A14 tune and CTEQ6L1 PDF to describe the underlying event. The sample was generated by setting the resummation damping parameter \( h_{\text{damp}} \) to 1.5 times the top-quark mass, \( m_{\text{top}} \), which was set to 172.5 GeV. The \( h_{\text{damp}} \) parameter controls the ME/PS matching and effectively regulates the high-\( p_T \) radiation. The EVTGEN 1.2.0 [70] package was used to model the properties of the bottom and charm hadron decays. Diboson samples were generated with SHERPA 2.1.1 [71–75] for the \( gg \) production processes and SHERPA 2.2.1 for the \( g \bar{g} \) production processes, using the CT10 NLO and NNPDF30NNLO PDF sets, respectively. The SHERPA event generator for the latter processes produces up to one additional parton at NLO and up to three additional partons at LO. Production of \( W \) and \( Z \) bosons in association with jets was also simulated using SHERPA 2.1.1 with the CT10 NLO PDF set, where \( b \)- and \( c \)-quarks are treated as massive particles. The \( gg \rightarrow WW \) production also includes the contribution of the SM Higgs boson at 125 GeV and the interference effects between the continuum and Higgs resonance processes. The VBF part of SM Higgs boson production was generated with POWHEG-BOX [54] interfaced to PYTHIA 8.186 for parton showering and hadronisation.

The effect of multiple \( pp \) interactions in the same and neighbouring bunch crossings (pile-up) was included by overlaying minimum-bias collisions, simulated with PYTHIA 8.186, on each generated signal and background event. The number of overlaid collisions is such that the distribution of the average number of interactions per \( pp \) bunch crossing in the simulation matches the pile-up conditions observed in the data, which is about 25 interactions per bunch crossing on average. The generated samples were processed through a GEANT4-based detector simulation [76,77], followed by the standard ATLAS reconstruction software used for collision data.

5 Event reconstruction

Events used in this analysis are required to have at least one primary vertex with a minimum of two associated tracks, each with transverse momentum \( p_T > 400 \text{ MeV} \). If there is more than one vertex reconstructed in an event that meets these conditions, the one with the highest sum of track \( p_T^2 \) is chosen as the primary vertex.

Electrons are reconstructed from clusters of energy deposits in the electromagnetic calorimeter that match a track reconstructed in the ID. They are identified using the likelihood identification criteria described in Ref. [78]. The electrons used in this analysis are required to pass the “MediumLH” selection for \( p_T > 25 \text{ GeV} \) or the “TightLH” selection for \( p_T < 25 \text{ GeV} \) and be within \( \left| \eta \right| < 2.47 \), excluding the transition region between the barrel and endcaps in the LAr calorimeter (1.37 < \( \left| \eta \right| < 1.52 \)). These “MediumLH” and “TightLH” selection categories have identification efficiencies of 84 and 74%, respectively, for electrons with \( p_T \) of 25 GeV. The corresponding probabilities to misidentify hadrons as electrons are approximately 0.5 and 0.3%, respectively.

Muons are reconstructed by combining ID and MS tracks that have consistent trajectories and curvatures [79]. The muon candidates used in this analysis are required to have \( \left| \eta \right| < 2.5 \) and pass the “Medium” selection for \( p_T > 25 \text{ GeV} \) or the “Tight” selection for \( p_T < 25 \text{ GeV} \), defined on the basis of the quality of the reconstruction and identification. These selections have a reconstruction efficiency of approximately 96 and 92%, respectively, for muons originating from the decay of \( W \) bosons [80]. The corresponding

For electrons, \( p_T \) is defined as the magnitude of the transverse component of the electron momentum as measured using the electromagnetic calorimeter.

 Springer
probabilities to misidentify hadrons as muons are approximately 0.2 and 0.1%, respectively.

To ensure that leptons originate from the interaction point, a requirement of $|d_0|/\sigma_{d_0} < 5$ (3) is imposed on the electrons (muons) and $|z_0\sin\theta| < 0.5$ mm is applied to both lepton types. Here $d_0$ and $z_0$ are the transverse and longitudinal impact parameters of the lepton with respect to the primary vertex, respectively, and $\sigma_{d_0}$ is the uncertainty in the measured value of $d_0$. In addition, electrons and muons are required to be isolated from other tracks and calorimetric activities by applying $p_T$- and $\eta$-dependent isolation criteria. For muons, the calorimeter isolation is based on energy deposits in the calorimeter within a cone $\Delta R$ of 0.2 around the muons. The muon track isolation uses a variable cone size starting at $\Delta R = 0.3$ and shrinking with increasing $p_T$ of the muon [81]. The same calorimeter isolation is used for electrons, and the electron track isolation uses a variable cone size starting at $\Delta R = 0.2$. The efficiency of these isolation requirements is 90% for both lepton types with $p_T$ of 25 GeV, increasing to 99% at 60 GeV.

Jets are reconstructed from three-dimensional clusters of energy deposits in the calorimeters using the anti-$k_T$ algorithm [82] with a radius parameter of $R = 0.4$ implemented in the FastJet package [83]. The four-momenta of the jets are calculated as the sum of the four-momenta of their constituents, which are assumed to be massless. Jets are corrected for energy from pile-up using the pile-up subtraction based on jet areas [84]. The jet energy scale is estimated in Ref. [85]. Jets are required to have $p_T > 30$ GeV and $|\eta| < 4.5$.

For jets with $p_T < 60$ GeV and $|\eta| < 2.5$, the multivariate “jet vertex tagger” algorithm [86] is used to suppress jets from pile-up interactions. To avoid double counting, jets of any transverse momentum are discarded if they are within a cone of size $\Delta R = 0.2$ around an electron candidate or if they have fewer than three associated tracks and are within a cone of size $\Delta R = 0.2$ around a muon candidate. However, if a jet with three or more associated tracks is within a cone of size $\Delta R < 0.4$ of a muon candidate, or the separation between an electron and any jet is within 0.2 $< \Delta R < 0.4$, the corresponding muon or electron candidate is rejected.

To estimate the number of $b$-tags in the event, jets with $p_T > 20$ GeV and within $|\eta| < 2.5$ are considered to contain a $b$-hadron if they yield a $b$-tagging algorithm discriminant value exceeding a reference value. The MV2c10 algorithm [87, 88] is chosen at the 85% $b$-tagging efficiency benchmark point, estimated from $b$-jets in simulated $t\bar{t}$ events. The misidentification rate for jets which originate from a light quark or gluon is less than 1%, while it is approximately 17% for $c$-jets.

The missing transverse momentum, with magnitude $E_T^{\text{miss}}$, is calculated as the negative vectorial sum of the transverse momenta of calibrated electrons, muons, and jets originating from the primary vertex, as well as tracks with $p_T > 500$ MeV compatible with the primary vertex and not associated with any of these [89].

### 6 Event selection

As a first step, $WW$ candidate events are selected by requiring two oppositely charged, different-flavour leptons ($e$ or $\mu$). Both leptons must satisfy the minimal quality criteria discussed in Sect. 5. When ordered in $p_T$, these leptons are called the leading and subleading ones, $p_T^L$,sublead. In order to suppress the background from diboson processes, a veto is imposed on events with an additional lepton with $p_T^{\text{other}} > 15$ GeV.

Table 2 summaries the selections and the definition of signal regions (SRs). The variables used in the selections are the most discriminating ones chosen by a boosted decision tree (BDT) [90], based on the NWA signal samples. These are $p_T^L$,lead, the invariant mass of the leading and subleading leptons, $m_{\ell\ell}$, and the pseudorapidity difference between the two leptons, $\Delta \eta_{\ell\ell}$. The first two variables provide good separation between a heavy resonance signal and the $WW$ and top-quark background. The separation of signal from background based on the $\Delta \eta_{\ell\ell}$ distribution is found to have a reasonable efficiency and allows, at the same time, a control region to be defined for the $WW$ background (Sect. 7.2). For each selected variable, the selection criterion is set by maximising the signal significance in the presence of background. The optimised selection is checked to be applicable to the LWA signals.

In order to further suppress the top-quark background, events with at least one $b$-tagged jet ($N_{b\text{-tag}} \geq 1$) are rejected.

Table 2. Selection conditions and phase space definitions used in the ggF and VBF signal regions

<table>
<thead>
<tr>
<th>SR_{ggF}</th>
<th>SR_{VBF1J}</th>
<th>SR_{VBF2J}</th>
</tr>
</thead>
<tbody>
<tr>
<td>Common selections</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$N_{b\text{-tag}} = 0$</td>
<td>$</td>
<td>\Delta\eta_{\ell\ell}</td>
</tr>
<tr>
<td>$m_{\ell\ell} &gt; 50$ GeV</td>
<td>$p_T^L$,lead &gt; 45 GeV</td>
<td>$p_T^L$,sublead &gt; 30 GeV</td>
</tr>
<tr>
<td>veto if $p_T^L$,other &gt; 15 GeV</td>
<td>$\max(m_T^{b}) &gt; 50$ GeV</td>
<td></td>
</tr>
<tr>
<td>ggF phase space</td>
<td>VBF1J phase space</td>
<td>VBF2J phase space</td>
</tr>
<tr>
<td>Inclusive in $N_{jet}$ but excluding $</td>
<td>\eta_{j}</td>
<td>&gt; 2.4$ and $\min(</td>
</tr>
</tbody>
</table>

\[\text{Page 5 of 34}\]
from the signal regions. To reduce the Z+jets and W+jets contributions, two other variables are used: \( p_T^\ell \) and the maximum value of the transverse mass calculated with either of the two leptons and the missing transverse momentum, \( m_T^W \). The latter variable is defined as:

\[
m_T^W = \sqrt{2 p_T^\ell E_T^{\text{miss}} (1 - \cos(\phi^\ell - \phi^{E_T^{\text{miss}}}))},
\]

where \( p_T^\ell \) and \( \phi^\ell \) are the transverse momentum and azimuthal angle of a given lepton and \( \phi^{E_T^{\text{miss}}} \) is the azimuthal angle of the missing transverse momentum vector.

Three event categories are defined: two disjoint categories optimised for the VBF production, VBF \( N_{\text{jet}} \geq 2 \) (SR\(_{\text{VBF1J}}\) and SR\(_{\text{VBF2J}}\)), and one quasi-inclusive category (excluding the VBF phase space) dedicated to the ggF or qqA signal (SR\(_{\text{ggF}}\)). For the VBF \( N_{\text{jet}} = 1 \) category, two discriminating variables are used to minimise the contribution of the ggF signal: the pseudorapidity of the jet, \( \eta \), and the minimum value of the pseudorapidity difference between the jet and either of the leptons, \( \min(|\Delta\eta_{\ell j}|) \). For the VBF \( N_{\text{jet}} \geq 2 \) category, the invariant mass, \( m_{jj} \), and the rapidity difference, \( \Delta\eta_{jj} \), of the two leading jets are used to select the VBF signal.

The NWA and LWA signal acceptance times the efficiency, after all selection requirements for a 700 GeV ggF signal, is approximately 50% in the quasi-inclusive ggF category and 5% or less in the VBF \( N_{\text{jet}} = 1 \) and VBF \( N_{\text{jet}} \geq 2 \) categories. For a 700 GeV VBF signal, it is between 15 and 25% for the three event categories. The acceptance times efficiency for the three event categories combined, as a function of resonance mass, is shown in Fig. 1 for the different signals. For the spin-1 and spin-2 signals, the range up to 1 TeV is considered in the case of VBF model processes. For samples with lower resonance masses, the acceptance times efficiency is lower because the leptons are softer. This is also the reason why the search is limited to signal mass values greater than about 200 GeV. The same selection is applied to all models and the different selection efficiencies between the models are mainly due to different \( \Delta\eta_{\ell\ell} \) distributions for the different spin states.

The discriminating variable used for the statistical analysis (Sect. 9) in this search is the transverse mass defined as

\[
m_T = \sqrt{(E_{T\ell\ell}^\ell + E_T^{\text{miss}})^2 - |p_T^\ell + E_T^{\text{miss}}|^2},
\]

where

\[
E_{T\ell\ell}^\ell = \sqrt{|p_T^\ell|^2 + m_{\ell\ell}^2},
\]

and \( p_T^\ell \) is the transverse momentum vector of the leading and subleading leptons.

### 7 Background estimation

The dominant background for the \( e^+e^-\mu^+\mu^- \) final state is due to events with top quarks and due to SM W W events. Additional contributions to the background arise from V+jets and the diboson processes ZV, V\( \gamma \) and V\( \gamma^* \). Since the discriminating variable used for this search is the transverse mass, \( m_T \), both the normalisation and the shape of the background \( m_T \) distribution must be estimated. The shape of the background is modelled using simulated events while the top-quark and W W background normalisations are determined by a simultaneous fit (Sect. 9) to the data in \( m_T \)-binned distributions in the signal regions and the total event yields in control regions. The normalisation factors of the fit, named “post-fit” normalisation factors\(^3\) hereafter, provide the best overall matching

\(^3\) The post-fit normalisation factors are checked to be consistent within the quoted uncertainties with the pre-fit ones obtained using the control regions only.
Table 3 Summary of all the selections used in the ggF and VBF WW and top-quark control regions. The common selection “veto if $p_T^{\text{other}} > 15 \text{ GeV}” applied to all the regions is not explicitly shown.

<table>
<thead>
<tr>
<th>$WW\text{ CR}_{\text{ggF}}$</th>
<th>Top CR$_{\text{ggF}}$</th>
<th>$WW\text{ CR}_{\text{VBF1J}}$</th>
<th>Top CR$_{\text{VBF}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$N_{b\text{-tag}} = 0$</td>
<td>$N_{b\text{-tag}} = 1$</td>
<td>$N_{b\text{-tag}} = 0$</td>
<td>$N_{b\text{-tag}} \geq 1$</td>
</tr>
<tr>
<td>$</td>
<td>\Delta \eta_{\ell \ell}</td>
<td>&gt; 1.8$</td>
<td>$</td>
</tr>
<tr>
<td>$m_{\ell \ell} &gt; 55 \text{ GeV}$</td>
<td>$p_T^{\ell,\text{lead}} &gt; 45 \text{ GeV}$</td>
<td>$p_T^{\ell,\text{lead}} &gt; 25 \text{ GeV}$</td>
<td>$p_T^{\ell,\text{sublead}} &gt; 25 \text{ GeV}$</td>
</tr>
<tr>
<td>$p_{T,\text{sublead}} &gt; 30 \text{ GeV}$</td>
<td>max$(m_W^W)$ &gt; 50 GeV</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Excluding VBF1J and VBF2J phase space</td>
<td>phase space</td>
<td>VBF1J and VBF2J</td>
<td></td>
</tr>
</tbody>
</table>

between the number of observed data events and the corresponding SM background expectations in all the signal and control regions. The control regions are defined by criteria similar to those used for the signal regions, but with some requirements loosened or reversed to obtain signal-depleted samples, enriched in the relevant background. These criteria are summarised in Table 3.

The following subsections describe the methods used to estimate the most important background processes, namely top quark, $WW$, and $W+\text{jets}$. The $Z/\gamma^*+\text{jets}$ and non-$WW$ diboson background contributions are small. The $Z/\gamma^*+\text{jets}$ Monte Carlo (MC) samples are normalised using NNLO cross sections [91] and the non-$WW$ ones with NLO cross sections from the SHERPA event generator. The small background from the $m_h \simeq 125 \text{ GeV}$ Higgs boson resonance and its off-shell component is included and its interference with the continuum $WW$ background is taken into account.

7.1 Top-quark background

Events containing top quarks can be produced as a $t\bar{t}$ pair or as a single top quark in association with either a $W$ boson or a quark of another flavour. In this analysis, contributions from $t\bar{t}$ and single-top-quark events are estimated together, with their relative contributions determined by their predicted cross sections and by their relative acceptances obtained from MC simulation. The single-top-quark contribution varies from about 10 to 30% depending on the signal event category.

The normalisation of the top-quark background for the quasi-inclusive ggF category is determined in a control region (Top CR$_{\text{ggF}}$) where one jet is required to be $b$-tagged in addition to the signal region selection. The purity of the top-quark background in this CR is high (97%) and thus allows the modelling of the MC simulation to be validated. The distribution of the simulated leading lepton $p_T$ in the Top CR$_{\text{ggF}}$ is found to disagree with the data and the ratio between the data and the simulation decreases with increasing $p_T^{\ell,\text{lead}}$. The simulated distribution is corrected in the SR$_{\text{ggF}}$ and corresponding CRs with factors obtained by fitting the ratio with a linear function. The correction varies between $+4$ and $-10\%$ as $p_T^{\ell,\text{lead}}$ increases from 50 to 200 GeV.

The top-quark background control regions for the VBF categories (Top CR$_{\text{VBF}}$) have a small number of data events and are therefore merged. At least one jet is required to be $b$-tagged. In addition, the selection thresholds imposed on $m_{\ell \ell}$ and $p_T^{\ell,\text{sublead}}$ are relaxed to 10 and 25 GeV, respectively, and the selection on $|\Delta \eta_{\ell \ell}|$ and max$(m_W^W)$ is removed. The threshold value on $m_{\ell \ell}$ of 10 GeV is used to suppress background contributions from low-mass resonances decaying into different-flavour final states via $\tau^+\tau^-$. In this control region, the purity of the top-quark background is 96%, and no mis-modelling of the $p_T^{\ell,\text{lead}}$ distribution is observed.

The post-fit normalisation factors from the simultaneous fit are $0.96 \pm 0.05$ and $1.12^{+0.13}_{-0.12}$ in the ggF and the VBF control regions, respectively, where the uncertainty quoted corresponds to the combined statistical and systematic uncertainties.

Figure 2 shows the $m_T$ distributions in the ggF and VBF top-quark CRs. The different background components are scaled according to the event yields obtained from the simultaneous fit. In the control regions the fit uses only the integrated event yields. The shape of the distributions is compared between data and MC predictions and found to be in good agreement after the application of the $p_T^{\ell,\text{lead}}$ correction described above for the ggF top-quark CR. The shapes of the $m_T$ distribution for 700 GeV and 2 TeV NWA Higgs boson signals are also shown, normalised to the expected limits on $\sigma_H \times B(H \rightarrow WW)$ from this analysis. The ggF contribution from the SM Higgs boson is included in the $WW$ component. The SM Higgs boson VBF contribution is negligibly small and is not shown in this and following figures.

7.2 WW background

The WW CR for the quasi-inclusive ggF category ($WW\text{ CR}_{\text{ggF}}$) uses the same selection as for the SR except for $|\Delta \eta_{\ell \ell}|$ which is reversed so that the CR and SR are orthogonal. The selection conditions are shown in Table 3. The $m_T$ distributions of the $q\bar{q} \rightarrow WW$ SHERPA MC sample in the

---

**Note:** The table in the original text is not properly formatted, and the content is partially obscured. The corrected and unobscured content is presented above. The corrections and unobscured content are not shown in the original table.
SR\textsubscript{ggF} and WW CR\textsubscript{ggF} are compared at MC generator level with corresponding predictions combining NNLO QCD calculations \cite{92} with NLO electroweak (EW) corrections \cite{93}. While the integrated yields of the distributions agree within 3\% in both the SR\textsubscript{ggF} and the WW CR\textsubscript{ggF}, a small \(m_T\) shape difference is observed, particularly in the SR. The \(m_T\) distributions of the \textsc{Sherpa} samples are thus reweighted to the combined NNLO QCD and NLO EW predictions. The post-fit normalisation factor obtained from the simultaneous fit for the WW contributions in the quasi-inclusive ggF categories is \(1.14 \pm 0.09\), where the uncertainty quoted corresponds to the combined statistical and systematic uncertainties. The post-fit purity of the WW background in the control region is 51\%.

In order to select more data events, the WW CR for the \(N_{\text{jet}} = 1\) VBF category (WW CR\textsubscript{VBF1J}) uses a slightly different selection (shown in Table 3) from the one in the SR, but still disjoint from the SR. The normalisation factor obtained from the same simultaneous fit for the WW contribution in the WW CR\textsubscript{VBF1J} is \(1.0 \pm 0.2\), where the uncertainty quoted corresponds to the combined statistical and systematic uncertainties. The post-fit purity of the WW background in the control region is 44\%.

The WW contribution in the \(N_{\text{jet}} \geq 2\) VBF category is about 20\%, and its prediction is taken from simulation because it is difficult to isolate a kinematic region with a sufficient number of WW events and with a small contamination from the top-quark background.

Figure 3 shows the \(m_T\) distributions in the \(WW\) CR\textsubscript{ggF} and CR\textsubscript{VBF1J}. The different background contributions are scaled according to the event yields obtained from the simultaneous fit. For the WW control regions only integrated event yields are used in the fit, like in the fits of the top control regions.

### 7.3 W+jets background

Events with W bosons produced in association with jets may enter the SR when a jet is misidentified as a lepton. Due to the difficulties in accurately modelling the misidentification process in the simulation, the W+jets background contribution is estimated using the data-driven method developed for the SM \(h \rightarrow WW\) analysis \cite{94}. A sample of events is used which satisfies all event selection criteria, except that one of the two lepton candidates fails to meet the quality criteria for being an identified lepton but satisfies a less restrictive selection, referred to as “anti-identified”. Anti-identified muons (electrons) have loosened isolation and impact parameter (likelihood identification) selection criteria as compared to the identified selection. From this data sample the non-W+jets contribution, dominated by top-quark and WW background processes, is subtracted on the basis of MC predictions. The W+jets purity of the samples is 46, 59 and 22\% for the quasi-inclusive ggF, \(N_{\text{jet}} = 1\) and \(N_{\text{jet}} \geq 2\) VBF categories, respectively.

The W+jets contamination in the signal region is then determined by scaling the number of events in the background-subtracted data sample by an extrapolation factor, which is the ratio of the number of identified leptons to the number of anti-identified leptons in a data sample.
of dijet events in bins of lepton $p_T$ and $\eta$. The dijet sample is collected using prescaled low-$p_T$ single-lepton triggers with thresholds of 12 GeV for electrons and 14 GeV for muons. Events are selected with exactly one candidate lepton, back-to-back with the leading jet. Electroweak processes in the dijet event sample, dominated by $W$+jets and $Z/\gamma^*$ background contributions, are subtracted. The dominant systematic uncertainty in the estimation of the $W$+jets background is due to the differences between dijet and $W$+jets sample characteristics. All systematic uncertainties associated with this background estimate are listed in Sect. 8.1.

8 Systematic uncertainties

In this section, experimental and theoretical uncertainties in the normalisation and shape of the $m_T$ distributions of the background and the signal are described. Except for those explicitly mentioned here, the shape uncertainties are small and thus neglected. Overall, the systematic uncertainty dominates, except in the tails of the $m_T$ distributions where the statistical uncertainty is larger.

8.1 Experimental uncertainties

The dominant sources of experimental uncertainty in the signal and background yields are the jet energy scale and resolution (Jet) [85], the $b$-tagging efficiency ($b$-tag) [87], and the pile-up modelling [86]. Other systematic uncertainties such as those associated with trigger efficiencies, lepton reconstruction and identification efficiencies, lepton momentum scales and resolutions [78,80], missing transverse momentum reconstruction [89] and the jet vertex tagger [86] are also considered when evaluating systematic effects on the shape and normalisation of the background, or the shape and efficiency of the signal yield. The uncertainty associated with the pile-up modelling is assessed by performing a variation of $\pm 9\%$ in the number of simulated pile-up interactions to cover the uncertainty in the ratio of the predicted and measured cross sections of non-diffractive inelastic events producing a hadronic system of mass $m_X,_{\text{had}}>13$ GeV [95].

For the main background from top-quark and $WW$ processes, the impact of the most important experimental systematic uncertainties is summarised in Tables 4 and 5 together with dominant theoretical uncertainties. The maximum changes in yield for the up and down variations are shown in the various signal and control regions. The correlation between the SRs and CRs is taken into account in the simultaneous fit.

Systematic effects due to lepton identification efficiencies, momentum and scale resolutions, are found to be approximately 1%. They are not shown in the tables. The last column in the tables shows the total uncertainty, including these small uncertainty sources.

The data-driven $W$+jets background estimate is subject to several sources of systematic uncertainty. The subtraction of the subdominant electroweak processes (Sect. 7.3) has a
The uncertainty in the total 2015 and 2016 integrated luminosity is 2.1%. It is derived, following a methodology similar to that detailed in Ref. [96], from van der Meer scans performed in August 2015 and May 2016, calibrated at high luminosity by various luminosity detectors.

8.2 Theoretical uncertainties of the background

For background sources which are normalised using control regions, theoretical uncertainties are evaluated for the extrapolation from the control region to the signal region.

For the top-quark and WW background, theoretical uncertainties in the extrapolation are evaluated according to the prescription from the LHC Higgs Cross Section Working Group [97]. The uncertainties include the impact of missing higher-order corrections, PDF variations and other MC modelling. The dominant theoretical uncertainties are shown in Tables 4 and 5.

For the top-quark background, the uncertainty from the event generator and parton shower modelling (ME+PS) is estimated by comparing the nominal POWHEG-BOX+PYTHIA8 generated samples with those from an alternative event generator, SHERPA 2.2.1. The uncertainty named “Scale” corresponds to variations of the renormalisation \( \mu_R \) and factorisation \( \mu_F \) scales as well as \( h_{\text{damp}} \). The variations for \( \mu_R \) and \( \mu_F \) are between 0.5 and 2 from their nominal scale of \( \sqrt{m_t^2+p_T^2} \), with \( p_T \) being the top-quark transverse momentum. The parameter \( h_{\text{damp}} \) is varied between \( m_{\text{top}}^2 \) and \( 2 \cdot m_{\text{top}}^2 \) from its nominal scale \( h_{\text{damp}} = 1.5 \cdot m_{\text{top}} \). In the analysis the single-top-quark and \( t\bar{t} \) processes are studied together. An uncertainty of 20% [98,99] is assigned to the relative contribution of the single-top-quark processes, corresponding to the source “Single top” in Table 4. The PDF uncertainty is obtained by taking the envelope of the uncertainty of the NNPDF30NLO PDF set and its differences in central value with the CT14 [100] and MMHT 2014 [101] PDF sets, following the recommendations of Ref. [55]. The PDF uncertainties are \( m_T \) dependent and increase from 2 to 10% with \( m_T \). This \( m_T \) dependence is taken into account in the signal regions. In the ggF quasi-inclusive category, two additional shape systematic uncertainties associated with the scale variations and the \( p_T \) reweighting for the leading lepton in the top-quark background are applied, the latter corresponding to \( \pm 50\% \) of the reweighting correction. These two uncertainties are comparable and vary from a few percent at low \( m_T \) to about 10% at \( m_T \approx 1 \text{ TeV} \), without affecting the integrated event yield of the top-quark background in the category.

For the WW background, the ME+PS modelling uncertainty is obtained by comparing the nominal SHERPA 2.2.1 sample with an alternative sample generated with POWHEG-BOX+PYTHIA8. The renormalisation, factorisation, and resummation scales are varied separately by factors of 0.5

---

Table 4 Relative impact (in %) of dominant experimental and theoretical uncertainties in the event yields for the top-quark background processes in the three signal regions (SR_\text{ggF}, SR_\text{VBF1J} and SR_\text{VBF2J}) and the top-quark and WW control regions (\text{Top} CR_\text{ggF/VBF1J} and the WW CR_\text{ggF/VBF1J}). Jet and \( b \)-tag sources dominate the experimental uncertainty while ME+PS, Scale, Single top and PDF are the dominant theoretical uncertainties. The last column shows the total uncertainty including those not listed here.

<table>
<thead>
<tr>
<th>Source</th>
<th>Jet</th>
<th>( b )-tag</th>
<th>ME+PS</th>
<th>Scale</th>
<th>Single top</th>
<th>PDF</th>
<th>Total</th>
</tr>
</thead>
<tbody>
<tr>
<td>SR_\text{ggF}</td>
<td>5.2</td>
<td>17</td>
<td>1.3</td>
<td>3.0</td>
<td>4.2</td>
<td>2.5</td>
<td>19</td>
</tr>
<tr>
<td>SR_\text{VBF1J}</td>
<td>9.6</td>
<td>7.8</td>
<td>1.0</td>
<td>1.6</td>
<td>5.9</td>
<td>2.6</td>
<td>15</td>
</tr>
<tr>
<td>SR_\text{VBF2J}</td>
<td>9.7</td>
<td>14</td>
<td>9.5</td>
<td>5.0</td>
<td>2.1</td>
<td>3.4</td>
<td>21</td>
</tr>
<tr>
<td>\text{Top} CR_\text{ggF}</td>
<td>2.2</td>
<td>4.8</td>
<td>0.34</td>
<td>0.21</td>
<td>2.6</td>
<td>3.0</td>
<td>6.6</td>
</tr>
<tr>
<td>WW CR_\text{ggF}</td>
<td>5.3</td>
<td>18</td>
<td>1.1</td>
<td>6.3</td>
<td>4.0</td>
<td>3.2</td>
<td>20</td>
</tr>
<tr>
<td>\text{Top} CR_\text{VBF}</td>
<td>8.2</td>
<td>3.5</td>
<td>10</td>
<td>1.5</td>
<td>1.3</td>
<td>3.7</td>
<td>14</td>
</tr>
<tr>
<td>WW CR_\text{VBF1J}</td>
<td>9.9</td>
<td>8.3</td>
<td>9.4</td>
<td>3.9</td>
<td>5.3</td>
<td>2.7</td>
<td>18</td>
</tr>
</tbody>
</table>

Table 5 Relative impact (in %) of dominant experimental and theoretical uncertainties in the event yields for the WW background processes in the three signal regions (SR_\text{ggF}, SR_\text{VBF1J} and SR_\text{VBF2J}) and the WW control regions (WW CR_\text{ggF/VBF1J}). Jet and Pile-up sources dominate the experimental uncertainty while ME+PS, \( \mu_R \), Resummation and PDF are the dominant theoretical uncertainties. The last column shows the total uncertainty including those not listed here.

<table>
<thead>
<tr>
<th>Source</th>
<th>Jet</th>
<th>Pile-up</th>
<th>ME+PS</th>
<th>( \mu_R )</th>
<th>Resummation</th>
<th>PDF</th>
<th>Total</th>
</tr>
</thead>
<tbody>
<tr>
<td>SR_\text{ggF}</td>
<td>1.2</td>
<td>1.8</td>
<td>2.4</td>
<td>1.7</td>
<td>3.1</td>
<td>2.7</td>
<td>5.5</td>
</tr>
<tr>
<td>SR_\text{VBF1J}</td>
<td>17</td>
<td>2.8</td>
<td>11</td>
<td>7.3</td>
<td>5.0</td>
<td>2.3</td>
<td>23</td>
</tr>
<tr>
<td>SR_\text{VBF2J}</td>
<td>18</td>
<td>3.1</td>
<td>38</td>
<td>18</td>
<td>1.4</td>
<td>2.1</td>
<td>47</td>
</tr>
<tr>
<td>WW CR_\text{ggF}</td>
<td>1.1</td>
<td>1.8</td>
<td>2.6</td>
<td>0.95</td>
<td>2.9</td>
<td>3.6</td>
<td>5.9</td>
</tr>
<tr>
<td>WW CR_\text{VBF1J}</td>
<td>16</td>
<td>4.5</td>
<td>12</td>
<td>11</td>
<td>2.3</td>
<td>2.8</td>
<td>23</td>
</tr>
</tbody>
</table>

significant impact on the extrapolation factor calculation at high lepton \( p_T \). The subtraction is varied, as described in Ref. [94], and the variation of the event yield in the signal region is taken as the uncertainty. The method assumes that the extrapolation factors of the dijet and W+jets samples are equal. Differences in the jet flavour composition between dijet and W+jets events introduce an additional systematic uncertainty. This is evaluated as the sum in quadrature of two contributions: differences between the extrapolation factors calculated with dijet samples and Z+jets samples in data, and differences between the extrapolation factors evaluated with W+jets and Z+jets MC samples. Finally, the statistical uncertainties of the different data and MC samples used to evaluate the extrapolation factors are taken as an additional source of systematic uncertainty. The overall relative systematic uncertainty of the W+jets background is found to be approximately 35% for each of the three signal event categories, with the dominant uncertainty being associated with the jet flavour composition.
and 2. The uncertainty corresponding to the factorisation scale variation is smaller than the other uncertainties and is not shown. The PDF uncertainty for the $W W$ background is obtained and treated in the same way as for the top-quark background. In the ggF quasi-inclusive category, an additional shape uncertainty from ME+PS is applied. It varies from a few percent at low $m_T$ to about 20% at $m_T \simeq 1$ TeV. There are no significant shape uncertainties in the $m_T$ distributions in the VBF categories.

In addition to the scale uncertainties described above, a relative uncertainty of $\pm 50\%$ is assigned to the reweighting corrections of the $q\bar{q} \rightarrow WW$ SHERPA sample to the combined NNLO QCD and NLO EW predictions in the ggF SR and $W W$ CR.

The $gg \rightarrow (h^*) \rightarrow W W$ process, where the SM 125 GeV Higgs boson is off-shell, is modelled at leading order with the SHERPA event generator with a $K$-factor of 1.7 that is used to account for higher-order cross-section corrections with an uncertainty of 60%, following the studies in Refs. [102–105].

Other small background processes, such as $W Z$, $Z Z$, $Z/\gamma^* +\text{jets}$ and $W W$ in the $N_{\text{jet}} \geq 2$ VBF category, do not have their own control regions. They are normalised to the theoretical predictions. The uncertainties in their yields due to the uncertainties in the predictions are evaluated with the same prescription as described above. The impact of these uncertainties is small (see Tables 6, 7 in Sect. 9).

8.3 Theoretical uncertainties in the signal predictions

Theoretical uncertainties in the signal acceptance include effects due to the choice of QCD renormalisation and factorisation scales, the PDF set as well as the underlying-event modelling, the parton shower model and the parton shower tune. These uncertainties are evaluated separately in each of the three event categories as a function of the resonance mass and independently for ggF- and VBF-induced resonances.

The effect of missing higher-order corrections in QCD on the signal acceptance is estimated by varying the renormalisation and factorisation scales independently by factors of 0.5 and 2 from the nominal scale of $\sqrt{m_H^2 + p_{T,H}^2}$, with $m_H$ and $p_{T,H}$ being the mass and the transverse momentum of the heavy Higgs boson, respectively. The acceptance values obtained with these modified MC samples are compared to the signal acceptance of the nominal sample. For resonances produced via ggF, these uncertainties are found to be negligible in the quasi-inclusive ggF and $N_{\text{jet}} = 1$ VBF categories, while in the $N_{\text{jet}} \geq 2$ VBF category they range between 2.5 and 0.2% for a resonance mass varying from 200 GeV to 4 TeV (unless stated otherwise, the following uncertainties are quoted for the same mass range). For resonances produced via vector-boson fusion, these uncertainties range from 0.9 to 2.8% in the quasi-inclusive ggF category, from 1.9 to 3.6% in the $N_{\text{jet}} = 1$ VBF category and from 1.0 to 7.3% in the $N_{\text{jet}} \geq 2$ VBF category.

The PDF-induced uncertainties in the signal acceptance are determined in the same way as for the top-quark and $W W$ background processes. For the ggF-induced (VBF-induced) signal, these uncertainties reach 0.4% (1.7%), 1.5% (1.2%) and 1.6% (1.5%) for the quasi-inclusive ggF, $N_{\text{jet}} = 1$ and $N_{\text{jet}} \geq 2$ VBF event categories, respectively.

The uncertainties corresponding to the parton shower tune and the underlying event are derived by moving independently, up or down, the PYTHIA internal parameters that are associated with final-state radiation or the multiple parton interactions to study separately their influence on the signal acceptance of the various signal mass points. These uncertainties are compared for each event category and mass point to the uncertainties from the choice of parton shower model, which are estimated by comparing the results obtained for the nominal parton shower generator to those obtained using

### Table 7 Event yields in the signal and control regions for the quasi-inclusive ggF category. The predicted background yields and uncertainties are calculated after the same simultaneous fit to the data in all the event categories as in Table 6. The statistical and systematic uncertainties are combined. The notation “$V V$” represents non-$W W$ diboson background

<table>
<thead>
<tr>
<th>Event Category</th>
<th>$SR_{\text{ggF}}$</th>
<th>$Top_{\text{CR}_{\text{ggF}}}$</th>
<th>$WW_{\text{CR}_{\text{ggF}}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$W W$</td>
<td>11,500 ± 800</td>
<td>820 ± 120</td>
<td>3360 ± 220</td>
</tr>
<tr>
<td>Top quark</td>
<td>11,800 ± 600</td>
<td>52,550 ± 330</td>
<td>2610 ± 180</td>
</tr>
<tr>
<td>$Z/\gamma^*$</td>
<td>1420 ± 110</td>
<td>111 ± 20</td>
<td>20.9 ± 2.0</td>
</tr>
<tr>
<td>$W +\text{jets}$</td>
<td>1180 ± 320</td>
<td>710 ± 190</td>
<td>280 ± 70</td>
</tr>
<tr>
<td>$VV$</td>
<td>866 ± 34</td>
<td>101 ± 12</td>
<td>250 ± 11</td>
</tr>
<tr>
<td>Background</td>
<td>26,740 ± 170</td>
<td>54,290 ± 250</td>
<td>6510 ± 80</td>
</tr>
<tr>
<td>Data</td>
<td>26,739</td>
<td>54,295</td>
<td>6515</td>
</tr>
</tbody>
</table>
Fig. 4 Post-fit distributions of the transverse mass $m_T$ in the SR$_{ggF}$ (top left), SR$_{VBF1J}$ (top right) and SR$_{VBF2J}$ (bottom) categories. In each plot, the last bin contains the overflow. The hatched band in the upper and lower panels shows the total uncertainty of the fit. The top-quark and $WW$ background event yields are scaled using the indicated normalisation factors obtained from the simultaneous fit to all signal and control regions. The heavy Higgs boson signal event yield is normalised to the expected limits on $\sigma H \times B(H \rightarrow WW)$ and is shown for masses of 700 GeV and 2 TeV in the NWA scenario.

**HERWIG++** [106,107]. The tune uncertainties are found to be smaller than the shower uncertainties for all mass points. Thus only the latter uncertainties are considered in the final results. The corresponding uncertainties for ggF-induced signals increase from 1.3 to 3.1%, from 13 to 28%, and from 2.3 to 15% for increasing resonance masses in the quasi-inclusive ggF, $N_{\text{jet}} = 1$ and $N_{\text{jet}} \geq 2$ VBF categories, respectively. The uncertainties for VBF-induced signals increase from 4.3 to 19%, from 5.1 to 9.0%, and from 3.3 to 8.0% in the three categories.

In addition, uncertainties due to missing higher-order corrections in QCD are evaluated for ggF-induced processes for each event category, considering also event migration effects between different event categories. This follows the method proposed by Stewart and Tackmann [108]. The corresponding uncertainties range from 3 to 10% for the quasi-inclusive ggF category and from 4 to 30% (30–60) for the $N_{\text{jet}} = 1$ ($N_{\text{jet}} \geq 2$) VBF event categories.

9 Results

The statistical method used to interpret the results of the search is described in Ref. [109]. A likelihood function $\mathcal{L}$
is defined as the product of Poisson probabilities associated with the number of events in bins of the \( m_T \) distributions in the signal regions and of the total yields in the control regions. Each source of systematic uncertainty is parameterised by a corresponding nuisance parameter \( \theta \) constrained by a Gaussian function.

The \( m_T \) distributions in the signal regions are divided into 18 (8) bins for the ggF quasi-inclusive (\( N_{\text{jet}} = 1 \) and \( \geq 2 \) VBF) categories. The bins are of variable size to reflect the increasing width of the \( m_T \) distribution of the expected signal with increasing mass, while keeping the statistical precision of the background contributions in each bin sufficiently high.

The numbers of events predicted and observed in the signal and control regions are shown for the quasi-inclusive ggF categories in Table 6 and for the VBF \( N_{\text{jet}} = 1 \) and \( \geq 2 \) categories in Table 7. These yields are obtained from a simultaneous fit to the data in all the SRs and the CRs. The fitted signal event yield is consistent with zero. The background compositions depend strongly on the event categories: the top-quark and \( W \bar{W} \) processes are comparable in \( \text{SR}_{\text{ggF}} \) and \( \text{SR}_{\text{VBF1J}} \) while the top-quark events dominate in \( \text{SR}_{\text{VBF2J}} \).

The large reduction of the total background uncertainty is due to strong anti correlations between some of the uncertainty sources of the top-quark and \( W \bar{W} \) background. The \( m_T \) distributions in \( \text{SR}_{\text{ggF}} \), \( \text{SR}_{\text{VBF1J}} \) and \( \text{SR}_{\text{VBF2J}} \) are shown in Fig. 4. As no excess over the background prediction is observed, upper limits at 95% confidence level (CL) are set on the production cross section times the branching fraction, \( \sigma_X \times B(\mu \rightarrow WW) \), for signals in each benchmark model.

The 95% CL upper limits are computed using the modified frequentist method known as CL_{s} [110], using the asymptotic approximation of the distribution of a test statistic [111], \( q_{\mu} \), a function of the signal strength \( \mu \), defined as the ratio of the measured \( \sigma_X \times B(X \rightarrow WW) \) to that of the prediction:

\[
q_{\mu} = -2 \ln \left( \frac{\mathcal{L}(\mu; \hat{\theta}_\mu)}{\mathcal{L}((\hat{\mu}); \hat{\theta})} \right).
\]

The quantities \( \hat{\mu} \) and \( \hat{\theta} \) are those values of \( \mu \) and \( \theta \), respectively, that unconditionally maximise \( \mathcal{L} \). The numerator depends on the values \( \hat{\theta}_\mu \) that maximise \( \mathcal{L} \) for a given value of \( \mu \).

Limits are obtained separately for ggF and VBF production for the NWA and LWA signal hypotheses. To derive the expected limits on the ggF (VBF) production modes, the VBF (ggF) production cross section is set to zero so that the...
Fig. 7 Exclusion contours at 95% CL in the plane of $\tan \beta$ and $\cos(\beta - \alpha)$ for Type I (left) and Type II (right) 2HDM signals with three mass values of 200 GeV (top), 300 GeV (middle) and 500 GeV (bottom). The inner and outer bands show the $\pm 1 \sigma$ and $\pm 2 \sigma$ ranges around the expected limit and the hatched regions are excluded.
expected limits correspond to the background-only hypothe-

sis. To derive the observed limits on the ggF (VBF) produc-
tion mode, the VBF (ggF) production cross section is treated
as a nuisance parameter in the fit and profiled, in the same
way as dealing with the normalisation factors of the differ-
rent background processes. This approach avoids making any
assumption about the presence or absence of the signal in
any of these production modes.

Figure 5 shows the 95% CL upper limits on $\sigma_H \times B(H \rightarrow W W)$ as a function of $m_H$ for a Higgs boson in the NWA sce-
nario in the mass range $200 \text{ GeV} \leq m_H \leq 4 \text{ (3) TeV}$ for the
ggF (VBF) production. Values above $6.4 \text{ pb} \ (1.3 \text{ pb})$ at $m_H =
200 \text{ GeV}$ and above $0.008 \text{ pb} \ (0.006 \text{ pb})$ at $4 \ (3) \text{ TeV}$ are
excluded at 95% CL by the quasi-inclusive ggF (VBF) NWA
analysis. The main systematic uncertainties affecting the lim-
its are the $p_T$ correction for the leading lepton in the top-quark
background, scale variations for the top-quark background,
the parton shower modelling of the $WW$ MC generator, and
the jet energy scale and resolution uncertainties. Limits are
consistent with those expected in the absence of a signal over
the investigated mass range. The fact that the observed lim-
its are more stringent than the expected ones for mass values
beyond 2 TeV is explained by the deficit in data at the high $m_T$
tail in Fig. 4. These limits are extracted using the asymptotic
approximation and their accuracy is verified to be consistent
within about 5% at 800 GeV and better than 20% at 2 TeV and
beyond using pseudo-experiments.

The analysis can be extended to a more general case where
the relative fraction of the ggF production cross section varies
over the total ggF and VBF production cross section. The
Corresponding 95% CL upper exclusion limits for a signal at
$800 \text{ GeV}$ are shown in Fig. 6. The dependence of the limits
on the ggF fraction for other masses is similar but becomes
slightly stronger (weaker) for lower (higher) mass values.

The VBF limits are tighter than the ggF limits since
the VBF $N_{\text{jet}} \geq 2$ signal region has the smallest background
contribution and thus is the most sensitive.

The NWA exclusion limit shown above can be further
translated to exclusion contours in the 2HDM for the phase
space where the narrow-width approximation is valid. The
95% CL exclusion contours for Type I and Type II in the
plane of $\tan \beta$ and $m_H$ for a Higgs boson mass $m_H$ is
excluded at 95% CL by the quas-inclusive ggF (VBF) NWA
analysis. The main systematic uncertainties affecting the lim-
its are the $p_T$ correction for the leading lepton in the top-quark
background, scale variations for the top-quark background,
the parton shower modelling of the $WW$ MC generator, and
the jet energy scale and resolution uncertainties. Limits are
consistent with those expected in the absence of a signal over
the investigated mass range. The fact that the observed lim-
its are more stringent than the expected ones for mass values
beyond 2 TeV is explained by the deficit in data at the high $m_T$
tail in Fig. 4. These limits are extracted using the asymptotic
approximation and their accuracy is verified to be consistent
within about 5% at 800 GeV and better than 20% at 2 TeV and
beyond using pseudo-experiments.

The coupling of the heaviest CP-even Higgs boson to vec-
tor bosons is proportional to $\cos(\beta - \alpha)$ and in the decou-
pling limit $\cos(\beta - \alpha) \to 0$, the light CP-even Higgs boson
is indistinguishable from a SM Higgs boson with the same
mass. The range of $\cos(\beta - \alpha)$ and $\tan \beta$ explored is limited
to the region where the assumption of a heavy narrow-width
Higgs boson with negligible interference is valid. When cal-
culating the limits at a given choice of $\cos(\beta - \alpha)$ and $\tan \beta$,
Fig. 9 Upper limits at 95% CL on the Higgs boson production cross section times branching fraction $\sigma_H \times B(H \rightarrow WW)$ in the $e\nu\mu\nu$ channel, for a signal with a width of 15% of the mass (top) and the comparison of three different widths (bottom) for the ggF (left) and VBF (right) production. The inner and outer bands show the $\pm 1\sigma$ and $\pm 2\sigma$ ranges around the expected limit.

Fig. 10 Upper limits at 95% CL on the resonance production cross section times branching fraction $\sigma_X \times B(X \rightarrow WW)$ (left) and on $\sin \theta_H$ (right) in the $e\nu\mu\nu$ channel, for a GM signal. The inner and outer bands show the $\pm 1\sigma$ and $\pm 2\sigma$ ranges around the expected limit. The full curves without dots correspond to the predicted theoretical cross section and the model parameter used in the benchmark model, respectively.
Fig. 11 Upper limits at 95% CL on the resonance production cross section times branching fraction $\sigma_X \times B(X \rightarrow WW)$ in the $\nu\mu\nu$ channel, for HVT qqA (left) and VBF (right) signals. The inner and outer bands show the $\pm 1\sigma$ and $\pm 2\sigma$ ranges around the expected limit. The full curves without dots correspond to the predicted theoretical cross sections.

Fig. 12 Upper limits at 95% CL on the resonance production cross section times branching fraction $\sigma_X \times B(X \rightarrow WW)$ in the $\nu\mu\nu$ channel, for a graviton signal with two different couplings of $k/M_{Pl} = 1$ (left) and $k/M_{Pl} = 0.5$ (right), and for an ELM spin-2 VBF signal (bottom). The inner and outer bands show the $\pm 1\sigma$ and $\pm 2\sigma$ ranges around the expected limit. The full curves without dots correspond to the predicted theoretical cross sections.
the relative rate of ggF and VBF production in the fit is set to the prediction of the 2HDM for that parameter choice. The white regions in the exclusion plots indicate regions of parameter space which are not excluded by the present analysis.

For the LWA scenario, the interference effects among the heavy boson, the light Higgs boson at 125 GeV and the SM WW continuum background were studied and found to have negligible impact on the exclusion limits. The 95% CL upper limits are shown in Fig. 9. The limits for signal widths of 5, 10 and 15% are comparable with those from the NWA scenario for the VBF signals while for the ggF signals, the limits weaken slightly at high masses as the width increases. For the LWA 15% case, the upper exclusion limit ranges between 5.2 pb (1.3 pb) at $m_H = 200$ GeV and 0.02 pb (0.006 pb) at 4 (3) TeV for the ggF (VBF) signals.

Figure 10 shows the limits on the resonance production cross section times branching fraction $\sigma \times B (X \rightarrow WW)$ and $\sin \theta_H$ for a scalar GM signal with masses between 200 GeV and 1 TeV. At the observed limit, the width is narrower than the experimental resolution [46]. The current sensitivity is not sufficient to exclude the benchmark model with $\sin \theta_H = 0.4$.

Limits are derived in the mass range from 250 GeV to 5 TeV and from 300 GeV to 1 TeV for a qqA and VBF HVT signal, respectively, as shown in Fig. 11. For the qqA production, signals below about 1.3 TeV are excluded at 95% CL upper limits are shown in Fig. 9. The limits for signal masses below 1.3, 1.1 TeV and 750 GeV are excluded, respectively.

Figure 12 shows the limits on a $G_{KK} \rightarrow WW$ signal for two different couplings: $k/\bar{M}_{Pl} = 1$ and $k/\bar{M}_{Pl} = 0.5$, for masses between 200 GeV and 5 TeV, and for an ELM spin-2 VBF signal for masses between 200 GeV and 1 TeV. The observed limits exclude a KK graviton signal lighter than 1.1 TeV (750 GeV) with the higher (lower) coupling, while the current sensitivity is not sufficient to exclude the ELM spin-2 VBF signal.

10 Conclusion

A search for heavy neutral resonances decaying into a WW boson pair in the $e\nu\mu\nu$ channel performed by the ATLAS Collaboration at the LHC is presented. The search uses proton–proton collision data collected at a centre-of-mass energy of 13 TeV corresponding to an integrated luminosity of 36.1 fb$^{-1}$. No significant excess of events beyond the Standard Model background prediction is found in the mass range between 200 GeV and up to 5 TeV. Upper limits are set on the product of the production cross section and the $X \rightarrow WW$ branching fraction in several scenarios: a high-mass Higgs boson with a narrow width or with intermediate widths (of 5, 10 and 15% of the heavy Higgs boson mass), as well as other spin-0, spin-1, and spin-2 signals. For the narrow-width heavy Higgs boson signals, values above 6.4 pb at $m_H = 200$ GeV and above 0.008 pb at 4 TeV are excluded at 95% confidence level for the gluon–gluon fusion production mode. The corresponding values for the vector-boson fusion production modes are 1.3 pb and 0.006 pb at 200 GeV and 3 TeV, respectively. For the signals of the heavy vector triplet model $A$ produced by quark–antiquark annihilation and of the Randall–Sundrum graviton model with $k/\bar{M}_{Pl} = 1$ and 0.5, mass values below 1.3, 1.1 TeV and 750 GeV are excluded, respectively.

Acknowledgements

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently. We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWFW and FWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPEP and FAPES, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSM CR, MPO CR and VSC CR, Czech Republic; DNRF and DSNRC, Denmark; IN2P3-CNRS, CEA-DRF/IRFU, France; SRNSF, Georgia; BMBF, HGF, and MPG, Germany; GSRT, Greece; RGC, Hong Kong SAR, China; ISF, I-CORE and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRSST, Morocco; NWO, Netherlands; RCN, Norway; MNiSW and NCN, Poland; FCT, Portugal; MNE/IFA, Romania; MES of Russia and NRC KI, Russian Federation; JINR; MESTD, Serbia; MSSR, Slovakia; ARRS and MIZS, Slovenia; DST/NRF, South Africa; MINECO, Spain; SRC and Wallenberg Foundation, Sweden; SERI, SNSF and Cantons of Bern and Geneva, Switzerland; MOST, Taiwan; TAEK, Turkey; STFC, United Kingdom; DOE and NSF, United States of America. In addition, individual groups and members have received support from BCKDF, the Canada Council, CANARIE, CRC, Compute Canada, QFRNT, and the Ontario Innovation Trust, Canada; EPLANET, ERC, ERDF, FP7, Horizon 2020 and Marie Sklodowska-Curie Actions, European Union; Investissements d’Avenir Labex and Idex, ANR, Région Auvergne and Fondation Partager le Savoir, France; DFG and AvH Foundation, Germany; Herakleitos, Thales and Aristeia programmes co-financed by EU-ESF and the Greek NSRF; BSF, GIF and Minerva, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRSST, Morocco; NWO, Netherlands; RCN, Norway; MNiSW and NCN, Poland; FCT, Portugal; MNE/IFA, Romania; MES of Russia and NRC KI, Russian Federation; JINR; MESTD, Serbia; MSSR, Slovakia; ARRS and MIZS, Slovenia; DST/NRF, South Africa; MINECO, Spain; SRC and Wallenberg Foundation, Sweden; SERI, SNSF and Cantons of Bern and Geneva, Switzerland; MOST, Taiwan; TAEK, Turkey; STFC, United Kingdom; DOE and NSF, United States of America. In addition, individual groups and members have received support from BCKDF, the Canada Council, CANARIE, CRC, Compute Canada, QFRNT, and the Ontario Innovation Trust, Canada; EPLANET, ERC, ERDF, FP7, Horizon 2020 and Marie Sklodowska-Curie Actions, European Union; Investissements d’Avenir Labex and Idex, ANR, Région Auvergne and Fondation Partager le Savoir, France; DFG and AvH Foundation, Germany; Herakleitos, Thales and Aristeia programmes co-financed by EU-ESF and the Greek NSRF; BSF, GIF and Minerva, Israel; BRF, Norway; CERCA Programme Generalitat de Catalunya, Generalitat Valenciana, Spain; the Royal Society and Leverhulme Trust, United Kingdom. The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN, the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CERN, IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA), the Tier-2 facilities worldwide and large non-WLCG resource providers. Major contributors of computing resources are listed in Ref. [112].

Open Access

This article is distributed under the terms of the Creative Commons Attribution 4.0 International License (http://creativecommons.org/licenses/by/4.0/), which permits unrestricted use, distribution, and reproduction in any medium, provided you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made. Funded by SCOAP3.
References


32. CMS Collaboration, Search for massive resonances decaying into WW, WZ or ZZ bosons in proton-proton collisions at \(\sqrt{s} = 13\) TeV. JHEP 03, 162 (2017). arXiv:1612.09159 [hep-ex]


35. ATLAS Collaboration, Searches for heavy ZZ and WZ resonances in the \(\ell\ell qq\) and \(q\bar{q}qq\) final states in pp collisions at \(\sqrt{s} = 13\) TeV with the ATLAS detector (2017). arXiv:1708.09638 [hep-ex]


44. J.F. Gunion et al., Higgs bosons in left-right symmetric models.


46. M. Zaro, H. Logan, Recommendations for the interpretation of LHC searches for $H^0, H^\pm$, and $H^{\pm\pm}$ in vector boson fusion with decays to vector boson pairs, LHC Higgs Cross Section Working Group. LHCHXSWG-2015-001. https://cds.cern.ch/record/2002500

47. ATLAS Collaboration, The ATLAS experiment at the CERN large hadron collider. JINST 3, S08003 (2008)


5 LAPP, CNRS/IN2P3 and Université Savoie Mont Blanc, Annecy-le-Vieux, France
6 High Energy Physics Division, Argonne National Laboratory, Argonne, IL, USA
7 Department of Physics, University of Arizona, Tucson, AZ, USA
8 Department of Physics, The University of Texas at Arlington, Arlington, TX, USA
9 Physics Department, National and Kapodistrian University of Athens, Athens, Greece
10 Physics Department, National Technical University of Athens, Zografou, Greece
11 Department of Physics, The University of Texas at Austin, Austin, TX, USA
12 Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
13 Institut de Física d'Altes Energies (IFAE), The Barcelona Institute of Science and Technology, Barcelona, Spain
14 Institute of Physics, University of Belgrade, Belgrade, Serbia
15 Department for Physics and Technology, University of Bergen, Bergen, Norway
16 Physics Division, Lawrence Berkeley National Laboratory, University of California, Berkeley, CA, USA
17 Department of Physics, Humboldt University, Berlin, Germany
18 Albert Einstein Center for Fundamental Physics, Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
19 School of Physics and Astronomy, University of Birmingham, Birmingham, UK
20 (a) Department of Physics, Bogazici University, Istanbul, Turkey; (b) Institute of Physics, Gaziantep University, Gaziantep, Turkey; (c) Faculty of Engineering and Natural Sciences, Istanbul Bilgi University, Istanbul, Turkey; (d) Physics Department, Bahcesehir University, Istanbul, Turkey
21 Centro de Investigaciones, Universidad Antonio Narino, Bogotá, Colombia
22 (a) INFN Sezione di Bologna, Bologna, Italy; (b) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
23 Physikalisches Institut, University of Bonn, Bonn, Germany
24 Department of Physics, Boston University, Boston, MA, USA
25 Department of Physics, Brandeis University, Waltham, MA, USA
26 (a) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil; (b) Electrical Circuits Department, Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil; (c) Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil; (d) Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil
27 Physics Department, Brookhaven National Laboratory, Upton, NY, USA
28 (a) Transilvania University of Brasov, Brasov, Romania; (b) Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania; (c) Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi, Romania; (d) Physics Department, National Institute for Research and Development of Isotopic and Molecular Technologies, Cluj Napoca, Romania; (e) University Politehnica Bucharest, Bucharest, Romania; (f) West University in Timisoara, Timisoara, Romania
29 Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
30 Cavendish Laboratory, University of Cambridge, Cambridge, UK
31 Department of Physics, Carleton University, Ottawa, ON, Canada
32 CERN, Geneva, Switzerland
33 Enrico Fermi Institute, University of Chicago, Chicago, IL, USA
34 (a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile; (b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaiso, Chile
35 (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China; (b) Department of Physics, Nanjing University, Nanjing, Jiangsu, China; (c) Physics Department, Tsinghua University, Beijing 100084, China; (d) University of Chinese Academy of Science (UCAS), Beijing, China
36 (a) Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Anhui, China; (b) School of Physics, Shandong University, Shandong, China; (c) School of Physics and Astronomy, Key Laboratory for Particle Physics, Astrophysics and Cosmology, Ministry of Education, Shanghai Key Laboratory for Particle Physics and Cosmology, Tsung-Dao Lee Institute, Shanghai Jiao Tong University, Shanghai, China
37 Université Clermont Auvergne, CNRS/IN2P3, LPC, Clermont-Ferrand, France
38 Nevis Laboratory, Columbia University, Irvington, NY, USA
39 Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
131 Faculty of Mathematics and Physics, Charles University, Prague, Czech Republic
132 State Research Center Institute for High Energy Physics (Protvino), NRC KI, Protvino, Russia
133 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, UK
134 (a)INFN Sezione di Roma, Rome, Italy; (b)Dipartimento di Fisica, Sapienza Università di Roma, Rome, Italy
135 (a)INFN Sezione di Roma Tor Vergata, Rome, Italy; (b)Dipartimento di Fisica, Università di Roma Tor Vergata, Rome, Italy
136 (a)INFN Sezione di Roma Tre, Rome, Italy; (b)Dipartimento di Matematica e Fisica, Università Roma Tre, Rome, Italy
137 (a)Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies-Université Hassan II, Casablanca, Morocco; (b)Centre National de l’Energie des Sciences Techniques Nucleaires, Rabat, Morocco; (c)Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Marrakech, Morocco; (d)Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda, Morocco; (e)Faculté des Sciences, Université Mohammed V, Rabat, Morocco
138 DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France
139 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, USA
140 Department of Physics, University of Washington, Seattle, WA, USA
141 Department of Physics and Astronomy, University of Sheffield, Sheffield, UK
142 Department of Physics, Shinshu University, Nagano, Japan
143 Department Physik, Universität Siegen, Siegen, Germany
144 Department of Physics, Simon Fraser University, Burnaby, BC, Canada
145 SLAC National Accelerator Laboratory, Stanford, CA, USA
146 (a)Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovak Republic; (b)Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
147 (a)Department of Physics, University of Cape Town, Cape Town, South Africa; (b)Department of Physics, University of Johannesburg, Johannesburg, South Africa; (c)School of Physics, University of the Witwatersrand, Johannesburg, South Africa
148 (a)Department of Physics, Stockholm University, Stockholm, Sweden; (b)The Oskar Klein Centre, Stockholm, Sweden
149 Physics Department, Royal Institute of Technology, Stockholm, Sweden
150 Departments of Physics and Astronomy and Chemistry, Stony Brook University, Stony Brook, NY, USA
151 Department of Physics and Astronomy, University of Sussex, Brighton, UK
152 School of Physics, University of Sydney, Sydney, Australia
153 Institute of Physics, Academia Sinica, Taipei, Taiwan
154 Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel
155 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
156 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
157 International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan
158 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
159 Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
160 Tomsk State University, Tomsk, Russia
161 Department of Physics, University of Toronto, Toronto, ON, Canada
162 (a)INFN-TIFPA, Trento, Italy; (b)University of Trento, Trento, Italy
163 (a)TRIUMF, Vancouver, BC, Canada; (b)Department of Physics and Astronomy, York University, Toronto, ON, Canada
164 Faculty of Pure and Applied Sciences, and Center for Integrated Research in Fundamental Science and Engineering, University of Tsukuba, Tsukuba, Japan
165 Department of Physics and Astronomy, Tufts University, Medford, MA, USA
166 Department of Physics and Astronomy, University of California Irvine, Irvine, CA, USA
167 (a)INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy; (b)ICTP, Trieste, Italy; (c)Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy
168 Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
169 Department of Physics, University of Illinois, Urbana, IL, USA
170 Instituto de Fisica Corpuscular (IFIC), Centro Mixto Universidad de Valencia-CSIC, Valencia, Spain
171 Department of Physics, University of British Columbia, Vancouver, BC, Canada
172 Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada
Department of Physics, University of Warwick, Coventry, UK
Waseda University, Tokyo, Japan
Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel
Department of Physics, University of Wisconsin, Madison, WI, USA
Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany
Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany
Department of Physics, Yale University, New Haven, CT, USA
Yerevan Physics Institute, Yerevan, Armenia
Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France
Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan
Also at Department of Physics, King’s College London, London, UK
Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
Also at Novosibirsk State University, Novosibirsk, Russia
Also at TRIUMF, Vancouver, BC, Canada
Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY, USA
Also at Physics Department, An-Najah National University, Nablus, Palestine
Also at Department of Physics, California State University, Fresno, CA, USA
Also at Department of Physics, University of Fribourg, Fribourg, Switzerland
Also at II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain
Also at Departamento de Física e Astronomía, Faculdade de Ciencias, Universidade do Porto, Porto, Portugal
Also at Tomsk State University, Tomsk, and Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia
Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing, China
Also at Universita di Napoli Parthenope, Naples, Italy
Also at Institute of Particle Physics (IPP), Canada
Also at Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania
Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia
Also at Borough of Manhattan Community College, City University of New York, New York, USA
Also at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece
Also at Centre for High Performance Computing, CSIR Campus, Rosebank, Cape Town, South Africa
Also at Louisiana Tech University, Ruston, LA, USA
Also at Institucio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain
Also at Department of Physics, The University of Michigan, Ann Arbor MI, USA
Also at LAL, Univ. Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France
Also at Graduate School of Science, Osaka University, Osaka, Japan
Also at Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, The Netherlands
Also at Department of Physics, The University of Texas at Austin, Austin, TX, USA
Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia
Also at CERN, Geneva, Switzerland
Also at Georgian Technical University (GTU), Tbilisi, Georgia
Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan
Also at Manhattan College, New York, NY, USA
Also at The City College of New York, New York NY, USA
Also at Departamento de Física Teorica y del Cosmos, Universidad de Granada, Granada, Portugal
Also at Department of Physics, California State University, Sacramento, CA, USA
Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia
Also at Departement de Physique Nucleaire et Corpusculaire, Université de Genève, Geneva, Switzerland
am Also at Institut de Física d’Altes Energies (IFAE), The Barcelona Institute of Science and Technology, Barcelona, Spain
an Also at School of Physics, Sun Yat-sen University, Guangzhou, China
ao Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia, Bulgaria
ap Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
aq Also at National Research Nuclear University MEPhI, Moscow, Russia
ar Also at Department of Physics, Stanford University, Stanford, CA, USA
as Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary
at Also at Faculty of Engineering, Giresun University, Giresun, Turkey
au Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
av Also at Department of Physics, Nanjing University, Jiangsu, China
aw Also at Institute of Physics, Academia Sinica, Taipei, Taiwan
ax Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia
*Deceased