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Evidence for light-by-light scattering in heavy-ion collisions with the ATLAS detector at the LHC

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Light-by-light scattering ($\gamma\gamma \rightarrow \gamma\gamma$) is a quantum-mechanical process that is forbidden in the classical theory of electrodynamics. This reaction is accessible at the Large Hadron Collider thanks to the large electromagnetic field strengths generated by ultra-relativistic colliding lead ions. Using 480 μb^{-1} of lead–lead collision data recorded at a centre-of-mass energy per nucleon pair of 5.02 TeV by the ATLAS detector, here we report evidence for light-by-light scattering. A total of 13 candidate events were observed with an expected background of 2.6 ± 0.7 events. After background subtraction and analysis corrections, the fiducial cross-section of the process Pb + Pb ($\gamma\gamma$) → $Pb^{(*)}$ + $Pb^{(*)}\gamma\gamma$, for photon transverse energy $E_T > 3$ GeV, photon absolute pseudorapidity $|\eta| < 2.4$, diphoton invariant mass greater than 6 GeV, diphoton transverse momentum lower than 2 GeV and diphoton acoplanarity below 0.01, is measured to be 70 ± 24 (stat.) ±17 (syst.) nb, which is in agreement with the standard model predictions.

One of the key features of Maxwell's equations is their linearity in both the sources and the fields, from which follows the superposition principle. This forbids effects such as light-by-light (LbyL) scattering, $\gamma\gamma \rightarrow \gamma\gamma$, which is a purely quantum-mechanical process. It was realized in the early history of quantum electrodynamics (QED) that LbyL scattering is related to the polarization of the vacuum[1]. In the standard model of particle physics, the virtual particles that mediate the LbyL coupling are electrically charged fermions or $W^{\pm}$ bosons. In QED, the $\gamma\gamma \rightarrow \gamma\gamma$ reaction proceeds at lowest order in the fine-structure constant (α_{em}) via virtual one-loop box diagrams involving fermions (Fig. 1a), which is an $\mathcal{O}(\alpha_{em}^4 \approx 3 \times 10^9)$ process, making it challenging to test experimentally. Indeed, the elastic LbyL scattering has remained unobserved: even the ultra-intense laser experiments are not yet powerful enough to probe this phenomenon[2].

LbyL scattering via an electron loop has been precisely, albeit indirectly, tested in measurements of the anomalous magnetic moment of the electron and muon[3,4] where it is predicted to contribute substantially, as one of the QED corrections[5]. The $\gamma\gamma \rightarrow \gamma\gamma$ reaction has been measured in photon scattering in the Coulomb field of a nucleus (Delbrück scattering) at fixed photon energies below 7 GeV (refs 6–9). The analogous process, where a photon splits into two photons by interaction with external fields (photon splitting), has been observed in the energy region of 0.1–0.5 GeV (ref. 10). A related process involving only real photons, in which several photons fuse to form an electron–positron pair (e^+e^-), has been measured in ref. 11. Similarly, the multiphoton Compton scattering, in which up to four laser photons interact with an electron, has been observed[12].

An alternative way by which LbyL interactions can be studied is by using relativistic heavy-ion collisions. In 'ultra-peripheral collision' (UPC) events, with impact parameters larger than twice the radius of the nuclei[13,14], the strong interaction does not play a role. The electromagnetic (EM) field strengths of relativistic ions scale with the proton number (Z). For example, for a lead (Pb) nucleus with $Z = 82$ the field can be up to 10^{25} V m^{-1} (ref. 15), much larger than the Schwinger limit[16] above which QED corrections become important. In the 1930s it was found that highly relativistic charged particles can be described by the equivalent photon approximation (EPA)[17–19], which is schematically shown in Fig. 1b. The EM fields produced by the colliding Pb nuclei can be treated as a beam of quasi-real photons with a small virtuality of $Q^2 < 1/R^2$, where R is the radius of the charge distribution and so $Q^2 < 10^{-3}$ GeV^2. Then, the cross-section for the reaction Pb + Pb ($\gamma\gamma$) → Pb + Pb $\gamma\gamma$ can be calculated by convolving the respective photon flux with the elementary cross-section for the process $\gamma\gamma \rightarrow \gamma\gamma$. Since the photon flux associated with each nucleus scales as Z^2, the cross-section is extremely enhanced as compared with proton–proton (pp) collisions.

In this article, a measurement of LbyL scattering in Pb + Pb collisions at the Large Hadron Collider (LHC) is reported, following the approach recently proposed in ref. 20. The final-state signature of interest is the exclusive production of two photons, Pb + Pb ($\gamma\gamma$) → $Pb^{(*)}$ + $Pb^{(*)}\gamma\gamma$, where a possible EM excitation of the outgoing ions[21] is denoted by (*). Hence, the expected signature is two photons and no further activity in the central detector, since the $Pb^{(*)}$ ions escape into the LHC beam pipe. Moreover, it is predicted that the background is relatively low in heavy-ion collisions and is dominated by exclusive dielectron ($\gamma\gamma \rightarrow e^+e^-$) production[20,22]. The misidentification of electrons as photons can occur when the electron track is not reconstructed or the electron emits a hard-bremsstrahlung photon. The fiducial cross-section of the process $\gamma\gamma \rightarrow \gamma\gamma$ in Pb + Pb collisions is measured, using a data set recorded at a nucleon–nucleon centre-of-mass energy ($\sqrt{s_{NN}}$) of 5.02 TeV. This data set was recorded with the ATLAS detector at the LHC in 2015 and corresponds to an integrated luminosity of 480 ± 30 μb^{-1}. In addition to the measured fiducial cross-section, the significance of the observed number of signal candidate events is given, assuming the background-only hypothesis.

Experimental set-up

ATLAS is a cylindrical particle detector composed of several sub-detectors[23]. ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point in the centre of the detector and the z axis along the beam pipe. The x axis points from the interaction point to the centre of the LHC ring, and the y

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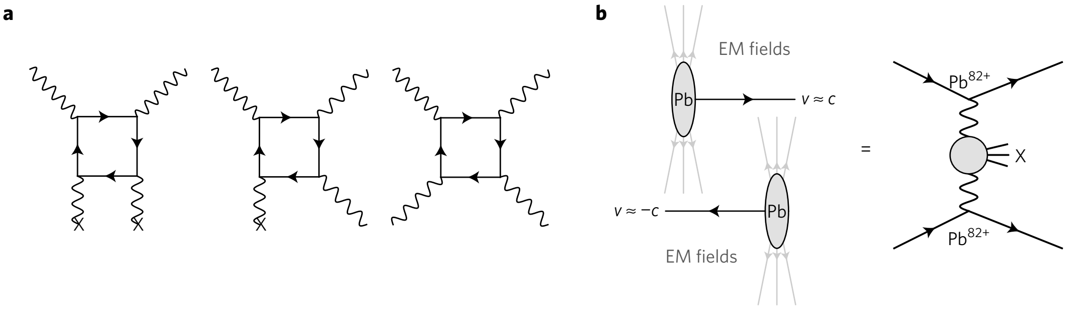

Figure 1 | Diagrams illustrating the QED LbyL interaction processes and the equivalent photon approximation. a, Diagrams for Delbrück scattering (left), photon splitting (middle) and elastic LbyL scattering (right). Each cross denotes external field legs, for example, an atomic Coulomb field or a strong background magnetic field. **b**, Illustration of an ultra-peripheral collision of two lead ions. Electromagnetic interaction between the ions can be described as an exchange of photons that can couple to form a given final state X. The flux of photons is determined from the Fourier transform of the electromagnetic field of the ion, taking into account the nuclear electromagnetic form factors.

axis points upwards. Cylindrical coordinates (r, ϕ) are used in the transverse plane, with ϕ being the azimuthal angle around the z axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$.

Angular distance is measured in units of $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$. The photon or electron transverse energy is $E_T = E \sin(\theta)$, where E is its energy. The inner tracking detector (ITD) consists of a silicon pixel system, a silicon microstrip detector and a straw-tube tracker immersed in a 2T magnetic field provided by a superconducting solenoid. The ITD track reconstruction efficiency is estimated in ref. 24 for minimum-bias pp events that, like UPC Pb + Pb events, have a low average track multiplicity. For charged hadrons in the transverse momentum range $100 < p_T < 200$ MeV the efficiency is about 50% and grows to 80% for $p_T > 200$ MeV. Around the tracker there is a system of EM and hadronic calorimeters, which use liquid argon and lead, copper or tungsten absorbers for the EM and forward ($|\eta| > 1.7$) hadronic components of the detector, and scintillator-tile active material and steel absorbers for the central ($|\eta| < 1.7$) hadronic component. The muon spectrometer consists of separate trigger and high-precision tracking chambers measuring the trajectory of muons in a magnetic field generated by superconducting air-core toroids. The ATLAS minimum-bias trigger scintillators (MBTSs) consist of scintillator slabs positioned between the ITD and the endcap calorimeters with each side having an outer ring of four slabs segmented in azimuthal angle, covering $2.07 < |\eta| < 2.76$ and an inner ring of eight slabs, covering $2.76 < |\eta| < 3.86$. The ATLAS zero-degree calorimeters (ZDCs), located along the beam axis at 140 m from the interaction point on both sides, detect neutral particles (including neutrons emitted from the nucleus). The ATLAS trigger system[25] consists of a Level-1 trigger implemented using a combination of dedicated electronics and programmable logic, and a software-based high-level trigger.

Monte Carlo simulation and theoretical predictions

Several Monte Carlo (MC) samples are produced to estimate background contributions and corrections to the fiducial measurement. The detector response is modelled using a simulation based on GEANT4 (refs 26,27). The data and MC simulated events are passed through the same reconstruction and analysis procedures.

LbyL signal events are generated taking into account box diagrams with charged leptons and quarks in the loops, as detailed in ref. 28. The contributions from W-boson loops are omitted in the calculations since they are mostly important for diphoton masses $m_{\gamma\gamma} > 2m_W$ (ref. 29). The calculations are then convolved with the Pb + Pb EPA spectrum from the STARlight 1.1 MC generator[30]. Next, various diphoton kinematic distributions are cross-checked with predictions from ref. 20 and good agreement is found. The theoretical uncertainty on the cross-section is mainly due to limited knowledge of the nuclear electromagnetic form factors and the related initial photon fluxes. This is studied in ref. 20 and the relevant uncertainty is conservatively estimated to be 20%. Higher-order corrections (not included in the calculations) are also part of the theoretical uncertainty and are of the order of a few per cent for diphoton invariant masses below 100 GeV (refs 31,32).

The sources of background considered in this analysis are: $\gamma\gamma \to e^+e^-$, central exclusive production (CEP) of photon pairs, exclusive production of quark–antiquark pairs ($\gamma\gamma \to q\bar{q}$) and other backgrounds that could mimic the diphoton event signatures. The $\gamma\gamma \to e^+e^-$ background is modelled with STARlight 1.1 (ref. 30), in which the cross-section is computed by combining the Pb + Pb EPA with the leading-order formula for $\gamma\gamma \to e^+e^-$. This process has been recently measured by the ALICE Collaboration, and a good agreement with STARlight is found[33]. The exclusive diphoton final state can be also produced via the strong interaction through a quark loop in the exchange of two gluons in a colour-singlet state (see Supplementary Fig. 2). This CEP process, $gg \to \gamma\gamma$, is modelled using SuperChic 2.03 (ref. 34), in which the pp cross-section has been scaled by $A^2 R_g^4$ as suggested in ref. 20, where $A = 208$ and $R_g \approx 0.7$ is a gluon shadowing correction[35]. This process has a large theoretical uncertainty, of $\mathcal{O}$ (100%), mostly related to incomplete knowledge of gluon densities[36]. The $\gamma\gamma \to q\bar{q}$ contribution is estimated using Herwig++ 2.7.1 (ref. 37) where the EPA formalism in pp collisions is implemented. The $\gamma\gamma \to q\bar{q}$ sample is then normalized to the corresponding cross-section in Pb + Pb collisions[30].

Event selection

Candidate diphoton events were recorded in the Pb + Pb run in 2015 using a dedicated trigger for events with moderate activity in the calorimeter but little additional activity in the entire detector. At Level-1 the total E_T registered in the calorimeter after noise suppression was required to be between 5 and 200 GeV. Then at the high-level trigger, events were rejected if more than one hit was found in the inner ring of the MBTS (MBTS veto) or if more than ten hits were found in the pixel detector.

The efficiency of the Level-1 trigger is estimated with $\gamma\gamma \to e^+e^-$ events passing an independent supporting trigger. This trigger is designed to select events with mutual dissociation of Pb nuclei and small activity in the ITD. It is based on a coincidence of signals in both ZDC sides and a requirement on the total E_T in the calorimeter below 50 GeV. Event candidates are required to have only two reconstructed tracks and two EM energy clusters. Furthermore, to reduce possible backgrounds, each pair of clusters (cl1, cl2) is required to have a small acoplanarity ($1 - \Delta\phi_{cl1,cl2}/\pi < 0.2$). The extracted Level-1 trigger efficiency is provided as a function of the

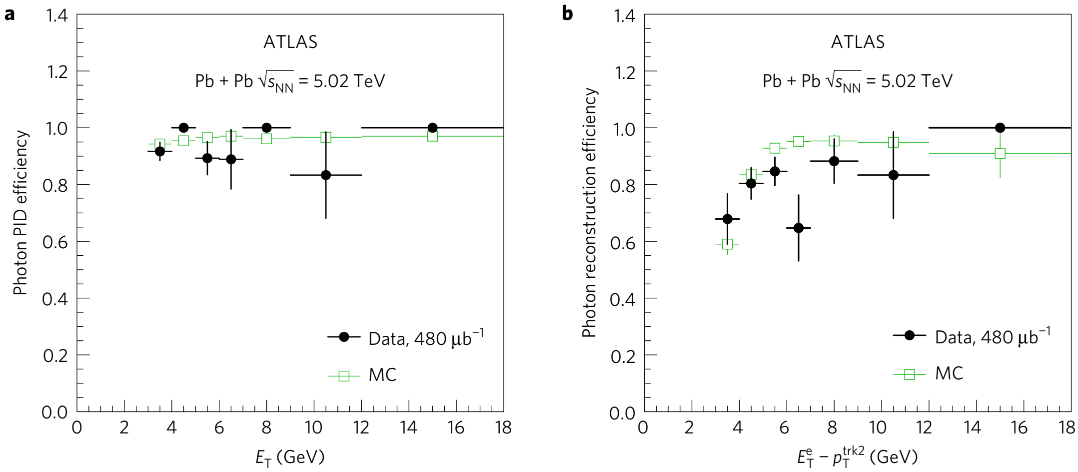

Figure 2 | Photon identification and reconstruction efficiencies. a, Photon PID efficiency as a function of photon E_T extracted from FSR event candidates. **b**, Photon reconstruction efficiency as a function of photon E_T (approximated with $E_T^e - p_T^{trk2}$) extracted from $\gamma\gamma \to e^+e^-$ events with a hard-bremsstrahlung photon. Data (filled markers) are compared with MC simulations (open markers). The statistical uncertainties arising from the finite size of the data and simulation samples are indicated by vertical bars.

sum of cluster transverse energies ($E_T^{cl1} + E_T^{cl2}$). The efficiency grows from about 70% at ($E_T^{cl1}+E_T^{cl2}$) = 6 GeV to 100% at ($E_T^{cl1}+E_T^{cl2}$) > 9 GeV. The efficiency is parameterized using an error function fit, which is then used to reweight the simulation. Due to the extremely low noise, very high hit reconstruction efficiency and low conversion probability of signal photons in the pixel detector (around 10%), the uncertainty due to the requirement for minimal activity in the ITD is negligible. The MBTS veto efficiency was studied using $\gamma\gamma \to \ell^+\ell^-$ events ($\ell = e, \mu$) passing a supporting trigger and it is estimated to be (98 ± 2)%.

Photons are reconstructed from EM clusters in the calorimeter and tracking information provided by the ITD, which allows the identification of photon conversions. Selection requirements are applied to remove EM clusters with a large amount of energy from poorly functioning calorimeter cells, and a timing requirement is made to reject out-of-time candidates. An energy calibration specifically optimized for photons[38] is applied to the candidates to account for upstream energy loss and both lateral and longitudinal shower leakage. A dedicated correction[39] is applied for photons in MC samples to correct for potential mismodelling of quantities that describe the properties ('shapes') of the associated EM showers.

The photon particle identification (PID) in this analysis is based on three shower-shape variables: the lateral width of the shower in the middle layer of the EM calorimeter, the ratio of the energy difference associated with the largest and second largest energy deposits to the sum of these energies in the first layer, and the fraction of energy reconstructed in the first layer relative to the total energy of the cluster. Only photons with $E_T > 3$ GeV and $|\eta| < 2.37$, excluding the calorimeter transition region $1.37 < |\eta| < 1.52$, are considered. The pseudorapidity requirement ensures that the photon candidates pass through regions of the EM calorimeter where the first layer is segmented into narrow strips, allowing for good separation between genuine prompt photons and photons coming from the decay of neutral hadrons. A constant photon PID efficiency of 95% as a function of η with respect to reconstructed photon candidates is maintained. This is optimized using multivariate analysis techniques[40], such that EM energy clusters induced by cosmic-ray muons are rejected with 95% efficiency.

Preselected events are required to have exactly two photons satisfying the above selection criteria, with a diphoton invariant mass greater than 6 GeV. To reduce the dielectron background, a veto on the presence of any charged-particle tracks (with $p_T > 100$ MeV, $|\eta| < 2.5$ and at least one hit in the pixel detector) is imposed. This requirement further reduces the fake-photon background from the dielectron final state by a factor of 25, according to simulation. It has almost no impact on $\gamma\gamma \to \gamma\gamma$ signal events, since the probability of photon conversion in the pixel detector is relatively small and converted photons are suppressed at low E_T (3–6 GeV) by the photon selection requirements. According to MC studies, the photon selection requirements remove about 10% of low-E_T photons. To reduce other fake-photon backgrounds (for example, cosmic-ray muons), the transverse momentum of the diphoton system ($p_T^{\gamma\gamma}$) is required to be below 2 GeV. To reduce background from CEP $gg \to \gamma\gamma$ reactions, an additional requirement on diphoton acoplanarity, $\mathrm{Aco} = 1 - \Delta\phi_{\gamma\gamma}/\pi < 0.01$, is imposed. This requirement is optimized to retain a high signal efficiency and reduce the CEP background significantly, since the transverse momentum transferred by the photon exchange is usually much smaller than that due to the colour-singlet-state gluons[41].

Performance and validation of photon reconstruction

Since the analysis requires the presence of low-energy photons, which are not typically used in ATLAS analyses, detailed studies of photon reconstruction and calibration are performed.

High-p_T $\gamma\gamma \to \ell^+\ell^-$ production with a final-state radiation (FSR) photon is used for the measurement of the photon PID efficiency. Events with a photon and two tracks corresponding to oppositely charged particles with $p_T > 1$ GeV are required to pass the same trigger as in the diphoton selection or the supporting trigger. The ΔR between a photon candidate and a track is required to be greater than 0.2 to avoid leakage of the electron clusters from the $\gamma\gamma \to e^+e^-$ process to the photon cluster. The FSR event candidates are identified using a $p_T^{tt\gamma} < 1$ GeV requirement, where $p_T^{tt\gamma}$ is the transverse momentum of the three-body system consisting of two charged-particle tracks and a photon. The FSR photons are then used to extract the photon PID efficiency, which is defined as the probability for a reconstructed photon to satisfy the identification criteria. Figure 2a shows the photon PID efficiencies in data and simulation as a function of reconstructed photon E_T. Within their statistical precision the two results agree.

The photon reconstruction efficiency is extracted from data using $\gamma\gamma \to e^+e^-$ events where one of the electrons emits a hard-bremsstrahlung photon due to interaction with the material of the detector. Events with exactly one identified electron, two reconstructed charged-particle tracks and exactly one photon are studied. The electron E_T is required to be above 5 GeV and the p_T

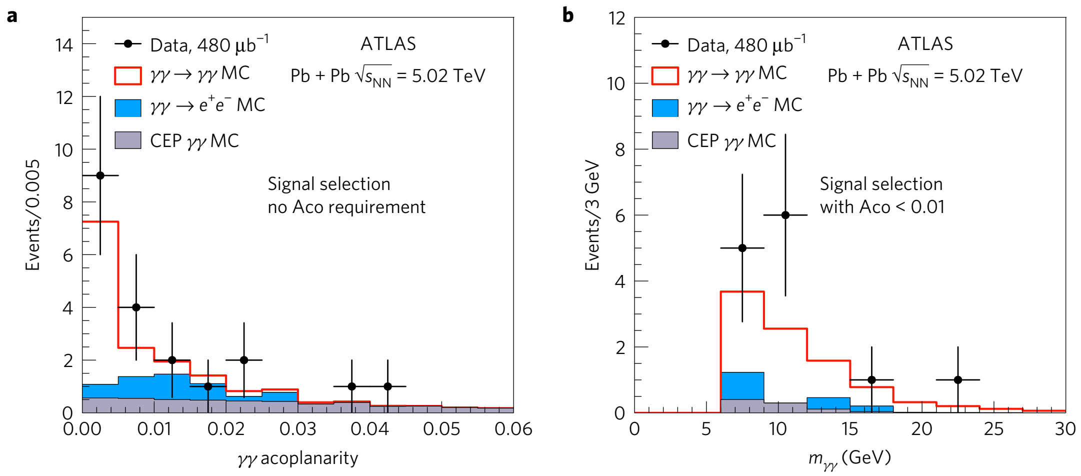

Figure 3 | Kinematic distributions for $\gamma\gamma \rightarrow \gamma\gamma$ event candidates. a, Diphoton acoplanarity before applying the Aco $<$ 0.01 requirement. **b**, Diphoton invariant mass after applying the Aco $<$ 0.01 requirement. Data (points) are compared to MC predictions (histograms). The statistical uncertainties on the data are shown as vertical bars.

of the track that is unmatched with the electron (trk2) is required to be below 2 GeV. The additional hard-bremsstrahlung photon is expected to have $E_T^{\gamma} \approx (E_T^{e} - p_T^{trk2})$. The $p_T^{trk2} < 2$ GeV requirement ensures a sufficient ΔR separation between the expected photon and the second electron, extrapolated to the first layer of the EM calorimeter. The data sample contains 247 $\gamma\gamma \rightarrow e^+e^-$ events that are used to extract the photon reconstruction efficiency, which is presented in Fig. 2b. Good agreement between data and $\gamma\gamma \rightarrow e^+e^-$ MC simulation is observed and the photon reconstruction efficiency is measured with a 5–10% relative uncertainty at low E_T (3–6 GeV).

In addition, a cross-check is performed on $Z \rightarrow \mu^+\mu^-\gamma$ events identified in pp collision data from 2015 corresponding to an integrated luminosity of 1.6 fb^{-1}. The results support (in a similar way to ref. 42) the choice to use the three shower-shape variables in this photon PID selection in an independent sample of low-E_T photons.

The photon cluster energy resolution is extracted from data using $\gamma\gamma \rightarrow e^+e^-$ events. The electrons from the $\gamma\gamma \rightarrow e^+e^-$ reaction (see Supplementary Information) are well balanced in their transverse momenta, with very small standard deviation, $\sigma_{p_T^{e+} - p_T^{e-}} <$ 30 MeV, much smaller than the expected EM calorimeter energy resolution. Therefore, by measuring $(E_T^{cl1} - E_T^{cl2})$ distributions in $\gamma\gamma \rightarrow e^+e^-$ events, one can extract the cluster energy resolution, $\sigma_{E_T^{cl}}$. For electrons with $E_T < 10$ GeV, the $\sigma_{E_T^{cl}}/E_T^{cl}$ is observed to be approximately 8% both in data and simulation. An uncertainty of $\delta\sigma_{E_T^{\gamma}}/\sigma_{E_T^{\gamma}} = 15\%$ is assigned to the simulated photon energy resolution and takes into account differences between $\sigma_{E_T^{cl}}$ in data and $\sigma_{E_T^{\gamma}}$ in simulation.

Similarly, the EM cluster energy scale can be studied using the $(E_T^{cl1} + E_T^{cl2})$ distribution. It is observed that the simulation provides a good description of this distribution, within the relative uncertainty of 5% that is assigned to the EM cluster energy-scale modelling.

Background estimation

Due to its relatively high rate, the exclusive production of electron pairs ($\gamma\gamma \rightarrow e^+e^-$) can be a source of fake diphoton events. The contribution from the dielectron background is estimated using $\gamma\gamma \rightarrow e^+e^-$ MC simulation (which gives 1.3 events) and is verified using the following data-driven technique. Two control regions are defined that are expected to be dominated by $\gamma\gamma \rightarrow e^+e^-$ backgrounds. The first control region is defined by requiring events with exactly one reconstructed charged-particle track and two identified photons that satisfy the same preselection criteria as for the signal definition. The second control region is defined similarly to the first one, except exactly two tracks are required ($N_{trk} = 2$). Good agreement is observed between data and MC simulation in both control regions, but the precision is limited by the number of events in data. A conservative uncertainty of 25% is therefore assigned to the $\gamma\gamma \rightarrow e^+e^-$ background estimation, which reflects the statistical uncertainty of data in the $N_{trk} = 1$ control region. The contribution from a related QED process, $\gamma\gamma \rightarrow e^+e^-\gamma\gamma$, is evaluated using the MadGraph5_aMC@NLO MC generator[43] and is found to be negligible.

The Aco $<$ 0.01 requirement significantly reduces the CEP $gg \rightarrow \gamma\gamma$ background. However, the MC prediction for this process has a large theoretical uncertainty; hence, an additional data-driven normalization is performed in the region Aco $> b$, where b is a value greater than 0.01 which can be varied. Three values of b (0.01, 0.02, 0.03) are used, where the central value $b = 0.02$ is chosen to derive the nominal background prediction and the values $b = 0.01$ and $b = 0.03$ to define the systematic uncertainty. The normalization is performed using the condition: $f_{gg\rightarrow\gamma\gamma}^{norm,b} = (N_{data}\ (\text{Aco} > b) - N_{sig}\ (\text{Aco} > b) - N_{\gamma\gamma\rightarrow e^+e^-}(\text{Aco} > b))/N_{gg\rightarrow\gamma\gamma}\ (\text{Aco} > b)$, for each value of b, where N_{data} is the number of observed events, N_{sig} is the expected number of signal events, $N_{\gamma\gamma\rightarrow e^+e^-}$ is the expected background from $\gamma\gamma \rightarrow e^+e^-$ events and $N_{gg\rightarrow\gamma\gamma}$ is the MC estimate of the expected background from CEP $gg \rightarrow \gamma\gamma$ events. The normalization factor is found to be $f_{gg\rightarrow\gamma\gamma}^{norm} = 0.5 \pm 0.3$ and the background due to CEP $gg \rightarrow \gamma\gamma$ is estimated to be $f_{gg\rightarrow\gamma\gamma}^{norm} \times N_{gg\rightarrow\gamma\gamma}$ (Aco $<$ 0.01) $= 0.9 \pm 0.5$ events. To verify the CEP $gg \rightarrow \gamma\gamma$ background estimation method, energy deposits in the ZDC are studied for events before the Aco selection. It is expected that the outgoing ions in CEP events predominantly dissociate, which results in the emission of neutrons detectable in the ZDC[20]. Good agreement between the normalized CEP $gg \rightarrow \gamma\gamma$ MC expectation and the observed events with a ZDC signal corresponding to at least 1 neutron is observed in the full Aco range (see Supplementary Information for details).

Low-p_T dijet events can produce multiple π^0 mesons, which could potentially mimic diphoton events. The event selection requirements are efficient in rejecting such events, and based on studies performed with a supporting trigger, the background from hadronic processes is estimated to be 0.3 ± 0.3 events. MC studies show that the background from $\gamma\gamma \rightarrow q\bar{q}$ processes is negligible.

Exclusive neutral two-meson production can be a potential source of background for LbyL events, mainly due to their back-to-back topology being similar to that of the CEP $gg \rightarrow \gamma\gamma$ process. The cross-section for this process is calculated to be below 10% of the CEP $gg \rightarrow \gamma\gamma$ cross-section[44,45] and it is therefore considered to

Table 1 | The number of events accepted by the sequential selection requirements for data, compared with the number of background and signal events expected from the simulation.

Selection	$\gamma\gamma \to e^+e^-$	CEP $gg \to \gamma\gamma$	Hadronic fakes	Other fakes	Total background	Signal	Data
Preselection	74	4.7	6	19	104	9.1	105
$N_{\mathrm{trk}} = 0$	4.0	4.5	6	19	33	8.7	39
$p_{\mathrm{T}}^{\gamma\gamma} < 2\,\mathrm{GeV}$	3.5	4.4	3	1.3	12.2	8.5	21
Aco < 0.01	1.3	0.9	0.3	0.1	2.6	7.3	13
Uncertainty	0.3	0.5	0.3	0.1	0.7	1.5	

The signal simulation is based on calculations from ref. 28. In addition, the uncertainties on the expected number of events passing all selection requirements are given.

give a negligible contribution to the signal region. The contribution from bottomonia production (for example, $\gamma\gamma \to \eta_b \to \gamma\gamma$ or $\gamma\mathrm{Pb} \to \Upsilon \to \gamma\eta_b \to 3\gamma$) is calculated using parameters from refs 46, 47 and is found to be negligible.

The contribution from other fake diphoton events (for example those induced by cosmic-ray muons) is estimated using photons that fail to satisfy the longitudinal shower-shape requirement. The total background due to other fake photons is found to be 0.1 ± 0.1 events. As a further cross-check, additional activity in the muon spectrometer is studied. It is observed that out of 18 events satisfying the inverted $p_{\mathrm{T}}^{\gamma\gamma}$ requirement, 13 have at least one additional reconstructed muon. In the region $p_{\mathrm{T}}^{\gamma\gamma} < 2\,\mathrm{GeV}$, no events with muon activity are found, which is compatible with the above-mentioned estimate of 0.1 ± 0.1.

The contribution from UPC events where both nuclei emit a bremsstrahlung photon is estimated using calculations from ref. 13 and is found to be negligible for photons with $|\eta| < 2.4$ and $E_{\mathrm{T}} > 3\,\mathrm{GeV}$.

Results

Photon kinematic distributions for events satisfying the selection criteria are shown in Fig. 3. The shape of the diphoton acoplanarity distribution for $\gamma\gamma \to e^+e^-$ events in Fig. 3a reflects the trajectories of the electron and positron in the detector magnetic field, before they emit hard photons in their collisions with the ITD material. In total, 13 events are observed in data whereas 7.3 signal events and 2.6 background events are expected. In general, good agreement between data and MC simulation is observed. The effect of sequential selection requirements on the number of events selected is shown in Table 1, for each of the data, signal and background samples.

To quantify an excess of events over the background expectation, a test statistic based on the profile likelihood ratio[48] is used. The p value for the background-only hypothesis, defined as the probability for the background to fluctuate and give an excess of events as large or larger than that observed in the data, is found to be 5×10^{-6}. The p value can be expressed in terms of Gaussian tail probabilities, which, given in units of standard deviation (σ), corresponds to a significance of 4.4σ. The expected p value and significance (obtained before the fit of the signal-plus-background hypothesis to the data and using standard model predictions from ref. 28) are 8×10^{-5} and 3.8σ, respectively.

The cross-section for the $\mathrm{Pb} + \mathrm{Pb}\ (\gamma\gamma) \to \mathrm{Pb}^{(*)} + \mathrm{Pb}^{(*)}\gamma\gamma$ process is measured in a fiducial phase space defined by the photon transverse energy $E_{\mathrm{T}} > 3\,\mathrm{GeV}$, photon absolute pseudorapidity $|\eta| < 2.4$, diphoton invariant mass greater than 6 GeV, diphoton transverse momentum lower than 2 GeV and diphoton acoplanarity below 0.01. Experimentally, the fiducial cross-section is given by

$$\sigma_{\mathrm{fid}} = \frac{N_{\mathrm{data}} - N_{\mathrm{bkg}}}{C \times \int L\mathrm{d}t} \quad (1)$$

where N_{data} is the number of selected events in data, N_{bkg} is the expected number of background events and $\int L\mathrm{d}t$ is the integrated luminosity. The factor C is used to correct for the net effect of the trigger efficiency, the diphoton reconstruction and PID efficiencies, as well as the impact of photon energy and angular resolution. It is defined as the ratio of the number of generated signal events satisfying the selection criteria after particle reconstruction and detector simulation to the number of generated events satisfying the fiducial criteria before reconstruction. The value of C and its total uncertainty is determined to be 0.31 ± 0.07. The dominant systematic uncertainties come from the uncertainties on the photon reconstruction and identification efficiencies. Other minor sources of uncertainty are the photon energy scale and resolution uncertainties and trigger efficiency uncertainty. To check for a potential model dependence, calculations from ref. 28 are compared with predictions from ref. 20, and a negligible impact on the C-factor uncertainty is found. Table 2 lists the separate contributions to the systematic uncertainty. The uncertainty on the integrated luminosity is 6%. It is derived following a methodology similar to that detailed in refs 49,50, from a calibration of the luminosity scale using x–y beam-separation scans performed in December 2015.

Table 2 | Summary of systematic uncertainties.

Source of uncertainty	Relative uncertainty
Trigger	5%
Photon reco. efficiency	12%
Photon PID efficiency	16%
Photon energy scale	7%
Photon energy resolution	11%
Total	24%

The table shows the relative systematic uncertainty on detector correction factor C broken into its individual contributions. The total is obtained by adding them in quadrature.

The measured fiducial cross-section is $\sigma_{\mathrm{fid}} = 70 \pm 24$ (stat.) ± 17 (syst.) nb, which is in agreement with the predicted values of 45 ± 9 nb (ref. 20) and 49 ± 10 nb (ref. 28) within uncertainties.

Conclusion

In summary, this article presents evidence for the scattering of LbyL in quasi-real photon interactions from $480\,\mu\mathrm{b}^{-1}$ of ultra-peripheral $\mathrm{Pb} + \mathrm{Pb}$ collisions at $\sqrt{s_{\mathrm{NN}}} = 5.02$ TeV by the ATLAS experiment at the LHC. The statistical significance against the background-only hypothesis is found to be 4.4 standard deviations. After background subtraction and analysis corrections, the fiducial cross-section for the $\mathrm{Pb} + \mathrm{Pb}\ (\gamma\gamma) \to \mathrm{Pb}^{(*)} + \mathrm{Pb}^{(*)}\gamma\gamma$ process was measured and is compatible with standard model predictions.

The analysis is mostly limited by the amount of data available and the lower limit on transverse energy for reconstructed photons ($E_{\mathrm{T}} = 3\,\mathrm{GeV}$), below which more signal is expected. Advancements on these two points would also allow for reconstruction of low-mass mesons decaying into two photons, which in turn could be used to improve detector calibration. The heavy-ion data yield is expected to double at the end of 2018 (and again increase tenfold after

LHC Run 4, scheduled to start in 2026), which would significantly reduce the statistical uncertainty. Future upgrades of ATLAS, such as extended tracking acceptance from $|\eta| < 2.5$ to $|\eta| < 4.0$, will further improve this.

Data availability. The experimental data that support the findings of this study are available in HEPData with the identifier http://dx.doi.org/10.17182/hepdata.77761.

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Author contributions

All authors have contributed to the publication, being variously involved in the design and the construction of the detectors, in writing software, calibrating subsystems, operating the detectors and acquiring data, and finally analysing the processed data. The ATLAS Collaboration members discussed and approved the scientific results. The manuscript was prepared by a subgroup of authors appointed by the collaboration and subject to an internal collaboration-wide review process. All authors reviewed and approved the final version of the manuscript.

Additional information

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Competing financial interests

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