PDF hosted at the Radboud Repository of the Radboud University Nijmegen

The following full text is a publisher’s version.

For additional information about this publication click this link.
http://hdl.handle.net/2066/168636

Please be advised that this information was generated on 2019-07-24 and may be subject to change.
Search for new phenomena in events containing a same-flavour opposite-sign dilepton pair, jets, and large missing transverse momentum in $\sqrt{s} = 13$ TeV $pp$ collisions with the ATLAS detector

ATLAS Collaboration
CERN, 1211 Geneva 23, Switzerland

Received: 18 November 2016 / Accepted: 16 February 2017
© CERN for the benefit of the ATLAS collaboration 2017. This article is published with open access at Springerlink.com

Abstract Two searches for new phenomena in final states containing a same-flavour opposite-sign lepton (electron or muon) pair, jets, and large missing transverse momentum are presented. These searches make use of proton–proton collision data, collected during 2015 and 2016 at a centre-of-mass energy $\sqrt{s} = 13$ TeV by the ATLAS detector at the large hadron collider, which correspond to an integrated luminosity of 14.7 fb$^{-1}$. Both searches target the pair production of supersymmetric particles, squarks or gluinos, which decay to final states containing a same-flavour opposite-sign lepton pair via one of two mechanisms: a leptonically decaying $Z$ boson in the final state, leading to a peak in the dilepton invariant-mass distribution around the $Z$ boson mass; and decays of neutralinos (e.g. $\tilde{\chi}^0_2 \rightarrow \ell^+ \ell^- \tilde{\chi}^0_1$), yielding a kinematic endpoint in the dilepton invariant-mass spectrum. The data are found to be consistent with the Standard Model expectation. Results are interpreted in simplified models of gluino-pair (squark-pair) production, and provide sensitivity to gluinos (squarks) with masses as large as 1.70 TeV (980 GeV).

Contents

1 Introduction ...........................................
2 ATLAS detector ........................................
3 SUSY signal models .................................
4 Data and Monte Carlo samples ......................
5 Analysis object identification and selection ....
6 Event selection ........................................
7 Background estimation ................................
   7.1 Flavour-symmetric backgrounds ............
   7.2 $Z/\gamma^* + \text{jets}$ background ..............
   7.3 Fake-lepton background ........................
   7.4 Diboson and rare top processes ..............
   7.5 Results in validation regions ............... 
8 Systematic uncertainties .........................
9 Results ..........................................
   9.1 Results in SRZ ................................
   9.2 Results in the edge SRs ....................
10 Interpretation ....................................
11 Conclusion ......................................
References ...........................................

1 Introduction

Supersymmetry (SUSY) [1–7] is an extension of the Standard Model (SM) that introduces partner particles (called sparticles) that differ by half a unit of spin from their SM counterparts. The squarks ($\tilde{q}$) and sleptons ($\tilde{\ell}$) are the scalar partners of the quarks and leptons, respectively, and the gluinos ($\tilde{g}$) are the fermionic partners of the gluons. The charginos ($\tilde{\chi}^\pm_i$) and neutralinos ($\tilde{\chi}^0_i$) are the mass eigenstates (where the index $i$ is ordered from the lightest to the heaviest) formed from the linear superpositions of the SUSY partners of the Higgs bosons (higgsinos) and electroweak gauge bosons.

If the masses of the gluino, higgsinos, and top squarks are close to the TeV scale, SUSY may offer a solution to the SM hierarchy problem [8–11]. In this case, strongly interacting sparticles should be produced at a high enough rate to be detected by the experiments at the large hadron collider (LHC). For models with R-parity conservation [12], such sparticles would be pair-produced and are expected to decay into jets, perhaps leptons, and the lightest stable SUSY particle (LSP). The LSP is assumed to be only weakly interacting and therefore escapes the detector, resulting in events with potentially large missing transverse momentum ($\not{p}_T$, with magnitude $E_T^{\text{miss}}$). In such a scenario the LSP could be a dark-matter candidate [13,14].
Final states containing pairs of leptons may arise from the cascade decays of squarks and gluinos via several mechanisms. In this paper, two search channels are considered that target scenarios with same-flavour (SF) opposite-sign (OS) lepton (electron or muon) pairs. The first channel requires a lepton pair with an invariant mass \( m_{\ell\ell} \) that is consistent with the Z boson mass \( m_Z \) (“on-shell Z” channel), while the second channel considers all SFOS lepton pairs (”edge” channel). The presence of two leptons in the final state suppresses large SM backgrounds from, e.g., QCD multijet and \( W + \text{jets} \) production, providing a clean environment in which to search for new physics. As discussed further below, in such events the distribution of dilepton mass \( m_{\ell\ell} \) may be used to characterise the nature of the SUSY particle decay and constrain mass differences between SUSY particles.

The SFOS lepton pairs may be produced in the decay \( \tilde{\chi}_2^0 \rightarrow \ell^+\ell^-\tilde{s}_1^0 \) (or, in models of generalised gauge mediation with a gravitino LSP [15–17], via \( \tilde{\chi}_1^0 \rightarrow \ell^+\ell^-\tilde{G} \)). The properties of the \( \tilde{\chi}_2^0 \) decay depend on the mass difference \( \Delta m_{\chi} \equiv m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0} \), the mixing of the charginos and neutralinos, and on whether there are additional sparticles with masses less than \( m_{\tilde{\chi}_2^0} \) that may be produced in the decay of the \( \tilde{\chi}_2^0 \) particle. For \( \Delta m_{\chi} > m_Z \), SFOS lepton pairs may be produced in the decay \( \tilde{\chi}_2^0 \rightarrow Z\tilde{\chi}_1^0 \rightarrow \ell^+\ell^-\tilde{\chi}_1^0 \), leading to a peak in the invariant-mass distribution near \( m_{\ell\ell} \approx m_Z \). Such models are the target of the on-shell Z search. For \( \Delta m_{\chi} < m_Z \), the decay \( \tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \rightarrow \ell^+\ell^-\tilde{\chi}_1^0 \) leads to a rising \( m_{\ell\ell} \) distribution that is truncated at a kinematic endpoint, whose position is given by \( m_{\ell\ell}^{\text{max}} = \Delta m_{\chi} - m_Z \), below the Z boson mass peak. If there are sleptons with masses less than \( m_{\tilde{\chi}_2^0} \), the \( \tilde{\chi}_2^0 \) particle may decay as \( \tilde{\chi}_2^0 \rightarrow \ell^+\ell^\mp \rightarrow \ell^+\ell^-\tilde{\chi}_1^0 \), also leading to a kinematic endpoint but with a different shape and \( m_{\ell\ell} \) endpoint position, given by \( m_{\ell\ell}^{\text{max}} = \sqrt{(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\chi}_1^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\chi}_1^0}^2)/m_{\ell\ell}^2} \), which may occur below, on, or above the Z boson mass peak. The latter two scenarios are targeted by the “edge” search channel, which considers the full \( m_{\ell\ell} \) range.

This paper reports on a search for SUSY in the same-flavour dilepton final state with 14.7 fb\(^{-1}\) of \( pp \) collision data at \( \sqrt{s} = 13 \text{ TeV} \) recorded in 2015 and 2016 by the ATLAS detector at the LHC. Searches for SUSY in the \( Z + \text{jets} + E_T^{\text{miss}} \) final state have previously been performed at \( \sqrt{s} = 8 \text{ TeV} \) by the CMS \([18,19]\) and ATLAS \([20]\) collaborations using Run-1 LHC data. In the ATLAS analysis performed with 20.3 fb\(^{-1}\) of \( \sqrt{s} = 8 \text{ TeV} \) data reported in Ref. \([20]\), an excess of events above the SM background with a significance of 3.0 standard deviations was observed. The event selection criteria for the on-shell Z search in this paper are almost identical, differing only in the details of the analysis object definitions and missing transverse momentum. CMS performed a search with \( \sqrt{s} = 13 \text{ TeV} \) data in a similar kinematic region but did not observe evidence to corroborate this excess \([21]\).

Searches for an edge in the \( m_{\ell\ell} \) distribution in events with \( 2\ell + \text{jets} + E_T^{\text{miss}} \) have been performed by the CMS \([19,22]\) and ATLAS \([20]\) collaborations. In Ref. \([19]\), CMS reported an excess above the SM prediction with a significance of 2.6 standard deviations. In a similar search region, however, the Run-1 ATLAS analysis \([20]\) and Run-2 CMS analysis \([21]\) observed results consistent with the SM prediction.

### 2 ATLAS detector

The ATLAS detector \([23]\) is a general-purpose detector with almost 4\pi coverage in solid angle.\(^1\) The detector comprises an inner tracking detector, a system of calorimeters, and a muon spectrometer.

The inner tracking detector (ID) is immersed in a 2 T magnetic field provided by a superconducting solenoid and allows charged-particle tracking out to \( |\eta| = 2.5 \). It includes silicon-pixel and silicon-strip tracking detectors inside a straw-tube tracking detector. In 2015 the detector received a new innermost layer of silicon pixels, which improves the track impact parameter resolution by almost a factor of two in both the transverse and longitudinal directions \([24]\).

High-granularity electromagnetic and hadronic calorimeters cover the region \( |\eta| < 4.9 \). All the electromagnetic calorimeters, as well as the endcap and forward hadronic calorimeters, are sampling calorimeters with liquid argon as the active medium and lead, copper, or tungsten as the absorber. The central hadronic calorimeter is a sampling calorimeter with scintillator tiles as the active medium and steel as the absorber.

The muon spectrometer uses several detector technologies to provide precision tracking out to \( |\eta| = 2.7 \) and triggering in \( |\eta| < 2.4 \), making use of a system of three toroidal magnets.

The ATLAS detector incorporates a two-level trigger system, with the first level implemented in custom hardware and the second level implemented in software. This trigger system selects events of interest at an output rate of about 1 kHz.

---

\(^1\) ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z-axis along the beam pipe. The x-axis points from the IP to the centre of the LHC ring, and the y-axis points upward. Cylindrical coordinates \((r, \phi)\) are used in the transverse plane, \(\phi\) being the azimuthal angle around the z-axis. The pseudorapidity is defined in terms of the polar angle \(\eta = -\ln \tan(\theta/2)\) and the rapidity is defined as \(y = 1/2 \cdot \ln((E + p_z)/(E - p_z)))\), where \(E\) is the energy and \(p_z\) the longitudinal momentum of the object of interest. The opening angle between two analysis objects in the detector is defined as \(\Delta R = \sqrt{(\Delta y)^2 + (\Delta \phi)^2}\).
3 SUSY signal models

SUSY-inspired simplified models are considered as signal scenarios for these analyses. In all of these models, squarks or gluinos are directly pair-produced, decaying via an intermediate neutralino, $\tilde{\chi}_1^0$, into the LSP ($\tilde{\chi}_1^0$). All particles not directly involved in the decay chains considered are effectively decoupled. Two example decay topologies are shown in Fig. 1. For all models with gluino-pair production, a three-body decay for $\tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_2^0$ is used. Signal models are generated in a grid over a two-dimensional space, varying the gluino or squark mass and the mass of either the $\tilde{\chi}_2^0$ or the $\tilde{\chi}_1^0$.

Three models, one with squark-pair production and two with gluino-pair production, which result exclusively in events with two on-shell $Z$ bosons in the final state are considered for the on-shell search. For two of these models, signal mass points are generated across the $\tilde{g}−\tilde{\chi}_2^0$ (or $\tilde{\tilde{q}}−\tilde{\chi}_2^0$) plane. These models are produced following the decays $\tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_2^0$ or $\tilde{\tilde{q}} \rightarrow q\tilde{q}\tilde{\chi}_2^0$, with the $\tilde{\chi}_1^0$ (LSP) mass set to 1 GeV, inspired by SUSY scenarios with a low-mass LSP (e.g. generalised gauge mediation). These two models are referred to here as the $\tilde{g}−\tilde{\chi}_2^0$ on-shell and $\tilde{\tilde{q}}−\tilde{\chi}_2^0$ on-shell grids, respectively, and are summarised in Table 1. The third model is based on MSSM-inspired topologies [25–27] with potentially higher mass LSPs. Signal points are generated across the $\tilde{g}−\tilde{\chi}_1^0$ plane, and this model is thus referred to as the $\tilde{g}−\tilde{\chi}_1^0$ on-shell grid. In this case the $\tilde{\chi}_2^0$ mass is set to be 100 GeV above the $\tilde{\chi}_1^0$ mass, which in many models maximises the branching fraction of the $\tilde{\chi}_2^0$ decay to $Z$ bosons. For the two models with gluino-pair production, since the gluino coupling to $q\tilde{q}$ is flavour independent and the corresponding flavours of squarks are assumed to be mass degenerate, the branching fractions for $q = u, d, c, s$ are each 25%. Other ATLAS searches are dedicated to final states with two leptons and heavy flavour jets [28,29]. For the model involving squark-pair production, the superpartners of the $u, d, c$ and $s$ quarks have the same mass, with the superpartners of the $b$ and $t$ quarks being decoupled.

The edge search considers two scenarios, both of which involve the direct pair production of gluinos and differ by the decay mode of the $\tilde{\chi}_2^0$. These signal models are also summarised in Table 1. In the $Z^{(*)}$ model the $\tilde{\chi}_2^0$ decays as $\tilde{\chi}_2^0 \rightarrow Z^{(*)}\tilde{\chi}_1^0$. For $\Delta m_{\chi} = m(\tilde{\chi}_2^0) - m(\tilde{\chi}_1^0) > m_Z$, the $Z$ boson is on-shell, leading to a peak in the $m_{\ell\ell}$ distribution at $m_Z$, while for $\Delta m_{\chi} < m_Z$, the $Z$ boson is off-shell, leading to an edge in the dilepton mass distribution with a position below $m_Z$. The slepton model assumes that the sleptons are lighter than the $\tilde{\chi}_2^0$, which decays as $\tilde{\chi}_2^0 \rightarrow \tilde{\ell}\tilde{\ell}\ell\bar{\ell}$ with $\ell = e, \mu, \tau$ and $\tilde{\ell} = \tilde{\nu}_e, \tilde{\nu}_\mu, \tilde{\nu}_\tau$. The endpoint position can occur at any mass, highlighting the need to search over the full dilepton mass distribution. The gluino decays as $\tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_2^0$, and both models have equal branching fractions for $q = u, d, c, s, b$. The $\tilde{\chi}_2^0$ mass is set to the average of the gluino and $\tilde{\chi}_1^0$ masses. For the slepton model, the masses of the superpartners of the left-handed sleptons are set as the average of the $\tilde{\chi}_2^0$ and $\tilde{\chi}_1^0$ masses, while the superpartners of the right-handed sleptons are decoupled. The three slepton flavours are mass-degenerate. In both these models the $\tilde{g}$ and $\tilde{\chi}_1^0$ masses are free parameters that are varied to produce the two-dimensional signal grid. The mass splittings are chosen to maximise the differences between these simplified models and other models with only one intermediate particle between the gluino and the LSP [30].

<table>
<thead>
<tr>
<th>Model</th>
<th>Production mode</th>
<th>Quark flavours</th>
<th>$m(\tilde{g})/m(\tilde{q})$</th>
<th>$m(\tilde{\chi}_2^0)$</th>
<th>$m(\tilde{\chi}_1^0)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\tilde{g}−\tilde{\chi}_2^0$ on-shell</td>
<td>$\tilde{g}\tilde{g}$</td>
<td>$u, d, c, s$</td>
<td>$x$</td>
<td>$y$</td>
<td>1 GeV</td>
</tr>
<tr>
<td>$\tilde{g}−\tilde{\chi}_1^0$ on-shell</td>
<td>$\tilde{g}\tilde{g}$</td>
<td>$u, d, c, s$</td>
<td>$x$</td>
<td>$m(\tilde{\chi}_1^0) + 100$ GeV</td>
<td>$y$</td>
</tr>
<tr>
<td>$\tilde{\tilde{q}}−\tilde{\chi}_2^0$ on-shell</td>
<td>$\tilde{\tilde{q}}\tilde{g}$</td>
<td>$u, d, c, s$</td>
<td>$x$</td>
<td>$y$</td>
<td>1 GeV</td>
</tr>
<tr>
<td>$Z^{(*)}$</td>
<td>$\tilde{g}\tilde{g}$</td>
<td>$u, d, c, s, b$</td>
<td>$x$</td>
<td>$(m(\tilde{g}) + m(\tilde{\chi}_1^0))/2$</td>
<td>$y$</td>
</tr>
<tr>
<td>slepton</td>
<td>$\tilde{g}\tilde{g}$</td>
<td>$u, d, c, s, b$</td>
<td>$x$</td>
<td>$(m(\tilde{g}) + m(\tilde{\chi}_1^0))/2$</td>
<td>$y$</td>
</tr>
</tbody>
</table>
4 Data and Monte Carlo samples

The data used in this analysis were collected by ATLAS during 2015 and 2016, with a mean number of additional \(pp\) interactions per bunch crossing (pile-up) of approximately 14 in 2015 and 21 in 2016, and a centre-of-mass collision energy of 13 TeV. Following requirements based on beam and detector conditions and data quality, the data collision energy of 13 TeV. Following requirements based approximately 14 in 2015 and 21 in 2016, and a centre-of-mass x–y beam-separation scans performed in August 2015 and May 2016.

Data events are collected using a combination of single-lepton and dilepton triggers [33], in order to maximise the signal acceptance. The dielectron, dimuon, and electron–muon triggers have leading-lepton \(p_T\) thresholds in the range 12–24 GeV. Additional single-electron (single-muon) triggers are also used, with trigger \(p_T\) thresholds of 60 (50) GeV, to increase the trigger efficiency for models with high-\(p_T\) leptons. Events are required to contain at least two selected leptons with \(p_T > 25\) GeV, making the selection fully efficient with respect to the trigger \(p_T\) thresholds.

An additional control sample of events containing photons is collected using a set of single-photon triggers with \(p_T\) thresholds in the range 20–140 GeV. All triggers except for the one with threshold \(p_T = 120\) GeV in 2015, or the one with \(p_T = 140\) GeV in 2016, are prescaled. Events are required to contain a selected photon with \(p_T > 37\) GeV, so that they are selected efficiently by the lowest available \(p_T\) trigger in 2015, which had a threshold of \(p_T^\gamma = 35\) GeV.

Simulated event samples are used to aid in the estimation of SM backgrounds, validate the analysis techniques, optimise the event selection, and provide predictions for SUSY signal processes. All SM background samples used are listed in Table 2, along with the parton distribution function (PDF) set, the configuration of underlying-event and hadronisation parameters (underlying-event tune) and the cross-section calculation order in \(\alpha_S\) used to normalise the event yields for these samples.

Samples simulated using MG5_AMC@NLO v2.2.2 [34], interfaced with PYTHIA 8.186 [35] with the A14 underlying-event tune [36] to simulate the parton shower and hadronisation, are generated at leading order in \(\alpha_S\) (LO) with the NNPDF23LO PDF set [37]. For samples generated using Powheg Box V2 [38–40], PYTHIA 6.428 [41] is used to simulate the parton shower, hadronisation, and the underlying event. The CTEQ6L1 PDF set is used with the corresponding PERUGIA2012 [42] tune. In the case of both the MG5_AMC@NLO and Powheg samples, theEvtGen v1.2.0 program [43] is used for properties of the bottom and charm hadron decays. SHERPA 2.1.1 [44] simulated samples use the CT10 PDF set with SHERPA’s own internal parton shower [45] and hadronisation methods, as well as the SHERPA default underlying-event tune. Diboson processes with four charged leptons, three charged leptons and a neutrino or two charged leptons and two neutrinos are simulated using the SHERPA 2.1.1 generator. Matrix elements contain all diagrams with four electroweak vertices. They are calculated for up to one (4\(\ell\), 2\(\ell + 2\nu\)) or zero (3\(\ell + 1\nu\)) partons at next-to-leading order in \(\alpha_S\) (NLO) and up to three partons at LO using the Comix [46] and OpenLoops [47] matrix element generators and merged with the SHERPA parton shower using the ME+PS@NLO prescription [48]. For the Z/\(\gamma^*\) + jets background, SHERPA 2.1.1 is used to generate a sample with up to two additional partons at NLO and up to four at LO. For Monte Carlo (MC) closure studies, \(\gamma + \) jets events are generated at LO with up to four additional partons using SHERPA 2.1.1. Additional MC simulation samples of events with a leptonically decaying vector boson and photon (\(V\gamma\), where \(V = W, Z\)) are generated at LO using SHERPA 2.1.1. Matrix elements including all diagrams with three electroweak couplings are calculated with up to three partons. These samples are used to estimate backgrounds with real \(E_T^{\text{miss}}\) in \(\gamma + \) jets event samples.

Table 2 Simulated background event samples used in this analysis with the corresponding matrix element and parton shower generators, cross-section order in \(\alpha_S\) used to normalise the event yield, underlying-event tune and PDF set

<table>
<thead>
<tr>
<th>Physics process</th>
<th>Generator</th>
<th>Parton shower</th>
<th>Cross section</th>
<th>Tune</th>
<th>PDF set</th>
</tr>
</thead>
<tbody>
<tr>
<td>(t\bar{t} + W) and (t\bar{t} + Z) [60,61]</td>
<td>MG5_AMC@NLO</td>
<td>PYTHIA 8.186</td>
<td>NLO [62,63]</td>
<td>A14</td>
<td>NNPDF23LO</td>
</tr>
<tr>
<td>(t\bar{t} + WW) [60]</td>
<td>MG5_AMC@NLO</td>
<td>PYTHIA 8.186</td>
<td>LO [34]</td>
<td>A14</td>
<td>NNPDF23LO</td>
</tr>
<tr>
<td>(t\bar{t}) [64]</td>
<td>Powheg Box V2 r3026</td>
<td>PYTHIA 6.428</td>
<td>NNLO+NNLL [65,66]</td>
<td>PERUGIA2012</td>
<td>NLO CT10</td>
</tr>
<tr>
<td>Single-top ((Wt)) [64]</td>
<td>Powheg Box V2 r2856</td>
<td>PYTHIA 6.428</td>
<td>Approx. NNLO [67]</td>
<td>PERUGIA2012</td>
<td>NLO CT10</td>
</tr>
<tr>
<td>(WW, WZ) [68]</td>
<td>SHERPA 2.1.1</td>
<td>SHERPA 2.1.1</td>
<td>NLO [69,70]</td>
<td>SHERPA default</td>
<td>NLO CT10</td>
</tr>
<tr>
<td>(Z/\gamma^*(\rightarrow \ell\ell) + ) jets [71]</td>
<td>SHERPA 2.1.1</td>
<td>SHERPA 2.1.1</td>
<td>NNLO [72,73]</td>
<td>SHERPA default</td>
<td>NLO CT10</td>
</tr>
<tr>
<td>(\gamma + ) jets</td>
<td>SHERPA 2.1.1</td>
<td>SHERPA 2.1.1</td>
<td>LO [44]</td>
<td>SHERPA default</td>
<td>NLO CT10</td>
</tr>
<tr>
<td>(V(=W, Z)\gamma)</td>
<td>SHERPA 2.1.1</td>
<td>SHERPA 2.1.1</td>
<td>LO [44]</td>
<td>SHERPA default</td>
<td>NLO CT10</td>
</tr>
</tbody>
</table>
The SUSY signal samples are produced at LO using MG5_AMC@NLO with the NNPDF2.3LO PDF set, interfaced with PYTHIA 8.186. The scale parameter for CKKW-L matching [49, 50] is set at a quarter of the mass of the gluino. Up to one additional parton is included in the matrix element calculation. The underlying event is modelled using the A14 tune for all signal samples, and EvtGen is adopted to describe the properties of bottom and charm hadron decays. Signal cross sections are calculated at NLO in αs. This includes the resummation of soft gluon emission at next-to-leading-logarithm accuracy (NLO+NLL) [51–55].

All of the SM background MC samples are subject to a full ATLAS detector simulation [56] using GEANT4 [57]. A fast simulation [56], which uses a combination of a parameterisation of the response of the ATLAS electromagnetic and hadronic calorimeters and GEANT4, is used in the case of signal MC samples. This fast simulation is validated by comparing a few chosen signal samples to some fully simulated points. Minimum-bias interactions are generated and overlaid on the hard-scattering process to simulate the effect of multiple pp interactions occurring during the same (in-time) or a nearby (out-of-time) bunch-crossing (pile-up). These are produced using PYTHIA 8.186 with the A2 tune [58] and MSTW 2008 PDF set [59]. The pile-up distribution in MC samples is simulated to match that in data during 2015 and 2016 pp data-taking.

5 Analysis object identification and selection

All analysis objects are categorised as either “baseline” or “signal” based on various quality and kinematic requirements. Baseline objects are used in the calculation of missing transverse momentum and to disambiguate between the analysis objects in the event, while the jets and leptons entering the final analysis selection must pass more stringent signal requirements. The selection criteria for both the baseline and signal objects differ from the requirements used in the Run-1 ATLAS Z+jets+E_{T}^{miss} search reported in Ref. [20], owing to the new silicon-pixel tracking layer and significant changes to the reconstruction software since 2012 data-taking. In particular, improvements in the lepton identification criteria have reduced the background due to hadrons misidentified as electrons. The primary vertex in each event is defined as the reconstructed vertex [74] with the highest \( \sum p_T^2 \), where the summation includes all particle tracks with \( p_T > 400 \text{ MeV} \) associated with the vertex.

Electron candidates are reconstructed from energy clusters in the electromagnetic calorimeter matched to ID tracks. Baseline electrons are required to have transverse energy \( E_T > 10 \text{ GeV} \), satisfy the “loose likelihood” criteria described in Ref. [75] and reside within the region \(|\eta| < 2.47\). Signal electrons are further required to have \( p_T > 25 \text{ GeV} \), satisfy the “medium likelihood” criteria of Ref. [75], and be consistent with originating from the primary vertex. The signal electrons must originate from within \(|z_0 \sin \theta| = 0.5 \text{ mm} \) of the primary vertex along the direction of the beamline.\(^2\) The transverse-plane distance of closest approach of the electron to the beamline, divided by the corresponding uncertainty, must be \(|d_0/\sigma_{d_0}| < 5\). These electrons must also be isolated with respect to other objects in the event, according to a \( p_T \)-dependent isolation requirement. The isolation uses calorimeter- and track-based information to obtain 95\% efficiency at \( p_T = 25 \text{ GeV} \), rising to 99\% efficiency at \( p_T = 60 \text{ GeV} \).

Baseline muons are reconstructed from either ID tracks matched to muon segments (collections of hits in a single muon spectrometer layer) or combined tracks formed from the ID and muon spectrometer [76]. They must satisfy the “medium” selection criteria described in Ref. [76], and to satisfy \( p_T > 10 \text{ GeV} \) and \(|\eta| < 2.5\). Signal muon candidates are further required to have \( p_T > 25 \text{ GeV} \), be isolated, and have \(|z_0 \sin \theta| < 0.5 \text{ mm} \) and \(|d_0/\sigma_{d_0}| < 3\). Calorimeter- and track-based isolation criteria are used to obtain 95\% efficiency at \( p_T = 25 \text{ GeV} \), rising to 99\% efficiency at \( p_T = 80 \text{ GeV} \) [76]. Further, the relative uncertainties in the \( q/p \) of each of the ID track alone and muon spectrometer track alone are required to be less than 80\% of the uncertainty in the \( q/p \) of the combined track. This reduces the already low rate of grossly mismeasured muons. The combined isolation and identification efficiency for single leptons, after the trigger requirements, is about 70\% (80\%) for electrons (muons) with \( p_T \sim 25 \text{ GeV} \), rising to about 90\% for \( p_T > 200 \text{ GeV} \).

Jets are reconstructed from topological clusters of energy [77] in the calorimeter using the anti-\( k_T \) algorithm [78, 79] with a radius parameter of 0.4. Calibration corrections are applied to the jets based on a comparison to jets made of stable particles (those with lifetimes \( \tau > 0.3 \times 10^{-10} \text{ s} \)) in the MC simulation. A residual correction is applied to jets in data, based on studies of \( p_T \) balance between jets and well-calibrated objects in the MC simulation and data [80, 81]. Baseline jet candidates are required to have \( p_T > 20 \text{ GeV} \) and reside within the region \(|\eta| < 4.5\). Signal jets are further required to satisfy \( p_T > 30 \text{ GeV} \) and reside within the region \(|\eta| < 2.5\). Jets with \( p_T < 60 \text{ GeV} \) and \(|\eta| < 2.4 \) must meet additional criteria designed to select jets from the hard-scatter interaction and reject those originating from pile-up. This is enforced by using the jet vertex tagger described in Ref. [82]. Finally, events containing a jet that does not pass specific jet quality requirements are vetoed from the analysis selection in order to remove events impacted by detector noise and

\(^2\) The distance of closest approach between a particle object and the primary vertex (beamline) in the longitudinal (transverse) plane is denoted by \( z_0 \) (\( d_0 \)).
non-collision backgrounds [83,84]. The MV2c10 boosted
decision tree algorithm [85,86] identifies jets with |\eta| < 2.5
containing b-hadrons (b-jets) based on quantities such as the
impact parameters of associated tracks and any reconstructed
secondary vertices. A selection that provides 77% efficiency
for tagging b-jets in simulated \(t\bar{t}\) events is used. These tagged
jets are called b-tagged jets.

Photon candidates must satisfy “tight” selection criteria
described in Ref. [87], have \(p_T > 25\) GeV and reside within
the region |\eta| < 2.37, excluding the transition region 1.37 <
|\eta| < 1.6 where there is a discontinuity in the calorimeter.
Signal photons are further required to have \(p_T > 37\) GeV
and to be isolated from other objects in the event, using \(p_T\)
-dependent requirements on both track- and calorimeter-based
isolation.

To avoid the duplication of analysis objects in more
than one baseline selection, an overlap removal procedure
is applied. Any baseline jet within \(\Delta R = 0.2\) of a baseline
electron is removed, unless the jet is b-tagged, in
which case the electron is identified as originating from
a heavy-flavour decay and is removed instead. Remaining
electrons residing within \(\Delta R = 0.4\) of a baseline jet are
then removed from the event. Subsequently, any baseline
muon residing within \(\Delta R = 0.2\) of a remaining baseline b-
tagged jet is discarded. If such a jet is not b-tagged then the
jet is removed instead. Any remaining muon found within
\(\min(0.04 + (10\) GeV)/\(p_T\), 0.4\) of a jet is also discarded.
This stage of the overlap removal procedure differs from
that used in Ref. [20]. It was improved to retain muons near
jet candidates mostly containing calorimeter energy from
final-state radiation from muons, while still rejecting muons
from heavy-flavour decays. Finally, to remove electron can-
didates originating from muon bremsstrahlung, any baseline
electron within \(\Delta R = 0.01\) of any remaining baseline muon
is removed from the event. Photons are removed if they
reside within \(\Delta R = 0.4\) of a baseline electron, and
any jet within \(\Delta R = 0.4\) of any remaining photon is dis-
carded.

The \(E_T^{\text{miss}}\) is defined as the magnitude of the negative vector
sum, \(p_T^{\text{miss}}\), of the transverse momenta of all baseline
electrons, muons, jets, and photons [88,89]. Low-momentum
contributions from particle tracks from the primary vertex
that are not associated with reconstructed analysis objects
are included in the calculation of \(E_T^{\text{miss}}\). This contribution to
the \(E_T^{\text{miss}}\) is referred to as the “soft term”.

Models with large hadronic activity are targeted by placing
additional requirements on the quantity \(H_T\), defined as the
scalar sum of the \(p_T\) values of all signal jets, or on \(H_T^{\text{incl}}\),
the scalar sum of the \(p_T\) values of all signal jets and the two
leptons with largest \(p_T\).

All MC samples have correction factors applied to take
into account small differences between data and MC simu-
lation in identification, reconstruction and trigger efficiencies
for leptons. The \(p_T\) values of leptons in MC samples are
additionally smeared to match the momentum resolution in data.

6 Event selection

For each search channel, signal regions (SRs) are designed
to target events from specific SUSY signal models. Con-
trol regions (CRs) are defined to be depleted in SUSY signal
events and enriched in specific SM backgrounds, and they are
used to assist in estimating these backgrounds in the SRs. To
validate the background estimation procedures, various vali-
dation regions (VRs) are defined to be analogous to the CRs
and SRs, but with less stringent requirements than the SRs on
\(E_T^{\text{miss}}, H_T^{\text{incl}}\) or \(H_T\). Other VRs with additional requirements
on the number of leptons are used to validate the modelling
of backgrounds in which more than two leptons are expected.

Events in SRs are required to contain at least two signal
leptons (electrons or muons). If more than two signal leptons
are present in a given event, the selection process continues
based on the two leptons with the highest \(p_T\) values in the
event.

The selected events must pass at least one of the lepton-
ic triggers. If an event is selected by a dilepton trigger, the two
leading, highest \(p_T\), leptons must be matched to one of the
objects that triggered the event. These leptons must also have
\(p_T\) higher than the threshold of the trigger in question. For
events selected by a single-lepton trigger, at least one of the
two leading leptons must be matched to the trigger object in
the same way. The leading two leptons in the event must have
\(p_T > 25\) GeV, and form an SFOS pair.

As at least two jets are expected in all signal models studied,
selected events are further required to contain at least
two signal jets. Furthermore, events in which the top quark,
\(t\bar{t}\), is raised to 50 GeV to increase the sensitivity to signal mod-
els with final-state \(Z\) bosons. This is an increased leading-
lepton \(p_T\) threshold relative to Ref. [20] and is found to better
reject fake-lepton candidates from misidentified jets, photon
conversions and b-hadron decays, while retaining high effi-
ciency for signal events, which tend to produce boosted \(Z\)
Table 3  Overview of all signal (SR), control (CR) and validation regions (VR) used in the on-shell \(Z\) search. The flavour combination of the dilepton pair is denoted as either “SF” for same-flavour or “DF” for different-flavour. All regions require at least two leptons, unless otherwise indicated. In the case of CR\(\gamma\), VR-WZ, VR-ZZ, and VR-3L the number of leptons, rather than a specific flavour configuration, is indicated. More details are given in the text. The main requirements that distinguish the control and validation regions from the signal region are indicated in bold. The kinematic quantities used to define these regions are discussed in the text. The quantity \(m_T(\ell_3, E_T^{\text{miss}})\) indicates the transverse mass formed by the \(E_T^{\text{miss}}\) and the lepton which is not assigned to either of the \(Z\)-decay leptons.

<table>
<thead>
<tr>
<th>On-shell (Z) regions</th>
<th>(E_T^{\text{miss}}) (GeV)</th>
<th>(H_T^{\text{incl}}) (GeV)</th>
<th>(n_{\text{jets}})</th>
<th>(m_{\ell\ell}) (GeV)</th>
<th>SF/DF</th>
<th>(\Delta \phi(\text{jet}_{12}, p_T^{\text{miss}}))</th>
<th>(m_T(\ell_3, E_T^{\text{miss}})) (GeV)</th>
<th>(n_{\text{jets}})</th>
</tr>
</thead>
<tbody>
<tr>
<td>Signal region</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>SRZ</td>
<td>&gt; 225</td>
<td>&gt; 600</td>
<td>(\geq 2)</td>
<td>(81 &lt; m_{\ell\ell} &lt; 101)</td>
<td>SF</td>
<td>&gt; 0.4</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Control regions</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>CRZ</td>
<td>&lt; 60</td>
<td>&gt; 600</td>
<td>(\geq 2)</td>
<td>(81 &lt; m_{\ell\ell} &lt; 101)</td>
<td>SF</td>
<td>&gt; 0.4</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>CR-FS</td>
<td>&gt; 225</td>
<td>&gt; 600</td>
<td>(\geq 2)</td>
<td>(61 &lt; m_{\ell\ell} &lt; 121)</td>
<td>DF</td>
<td>&gt; 0.4</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>CRT</td>
<td>&gt; 225</td>
<td>&gt; 600</td>
<td>(\geq 2)</td>
<td>(&gt; 45, m_{\ell\ell} \notin [81, 101])</td>
<td>SF</td>
<td>&gt; 0.4</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>CR(\gamma)</td>
<td>–</td>
<td>&gt; 600</td>
<td>(\geq 2)</td>
<td>–</td>
<td>0(\ell, 1\gamma)</td>
<td>–</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Validation regions</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>VRZ</td>
<td>&lt; 225</td>
<td>&gt; 600</td>
<td>(\geq 2)</td>
<td>(81 &lt; m_{\ell\ell} &lt; 101)</td>
<td>SF</td>
<td>&gt; 0.4</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>VRT</td>
<td>100–200</td>
<td>&gt; 600</td>
<td>(\geq 2)</td>
<td>(&gt; 45, m_{\ell\ell} \notin [81, 101])</td>
<td>SF</td>
<td>&gt; 0.4</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>VR-S</td>
<td>100–200</td>
<td>&gt; 600</td>
<td>(\geq 2)</td>
<td>(81 &lt; m_{\ell\ell} &lt; 101)</td>
<td>SF</td>
<td>&gt; 0.4</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>VR-FS</td>
<td>100–200</td>
<td>&gt; 600</td>
<td>(\geq 2)</td>
<td>(61 &lt; m_{\ell\ell} &lt; 121)</td>
<td>DF</td>
<td>&gt; 0.4</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>VR-WZ</td>
<td>100–200</td>
<td>–</td>
<td>(\geq 2)</td>
<td>–</td>
<td>3(\ell)</td>
<td>–</td>
<td>&lt; 100</td>
<td>0</td>
</tr>
<tr>
<td>VR-ZZ</td>
<td>&lt; 100</td>
<td>–</td>
<td>(\geq 2)</td>
<td>–</td>
<td>4(\ell)</td>
<td>–</td>
<td>–</td>
<td>0</td>
</tr>
<tr>
<td>VR-3L</td>
<td>60–100</td>
<td>&gt; 200</td>
<td>(\geq 2)</td>
<td>(81 &lt; m_{\ell\ell} &lt; 101)</td>
<td>3(\ell)</td>
<td>&gt; 0.4</td>
<td>–</td>
<td>–</td>
</tr>
</tbody>
</table>
Table 4 Overview of all signal (SR), control (CR) and validation regions (VR) used in the edge search. The flavour combination of the dilepton pair is denoted as either “SF” for same-flavour or “DF” for different-flavour. The charge combination of the leading lepton pairs are given as “SS” for same-sign or “OS” for opposite-sign. All regions require at least two leptons, with the exception of CR-real, which requires exactly two leptons, and the three \( \gamma \) CRs, which require no leptons and one photon. More details are given in the text. The main requirements that distinguish the control and validation regions from the signal regions are indicated in bold. The kinematic quantities used to define these regions are discussed in the text.

<table>
<thead>
<tr>
<th>Edge regions</th>
<th>( E_T^{\text{miss}} ) (GeV)</th>
<th>( H_T ) (GeV)</th>
<th>( n_{\text{jets}} )</th>
<th>( m_{\ell\ell} ) (GeV)</th>
<th>SF/DF</th>
<th>OS/SS</th>
<th>( \Delta \phi (\text{jet}_1, \text{jet}_2, \eta^{\text{miss}}) )</th>
<th>( m_{\ell\ell} ) ranges</th>
</tr>
</thead>
<tbody>
<tr>
<td><strong>Signal regions</strong></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>SR-low</td>
<td>&gt; 200</td>
<td>–</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>SF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>9</td>
</tr>
<tr>
<td>SR-medium</td>
<td>&gt; 200</td>
<td>&gt; 400</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>SF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>8</td>
</tr>
<tr>
<td>SR-high</td>
<td>&gt; 200</td>
<td>&gt; 700</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>SF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>7</td>
</tr>
<tr>
<td><strong>Control regions</strong></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>CRZ-low</td>
<td>&lt; 60</td>
<td>–</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>SF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>–</td>
</tr>
<tr>
<td>CRZ-medium</td>
<td>&lt; 60</td>
<td>&gt; 400</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>SF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>–</td>
</tr>
<tr>
<td>CRZ-high</td>
<td>&lt; 60</td>
<td>&gt; 700</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>SF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>–</td>
</tr>
<tr>
<td>CR-FS-low</td>
<td>&gt; 200</td>
<td>–</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>DF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>–</td>
</tr>
<tr>
<td>CR-FS-medium</td>
<td>&gt; 200</td>
<td>&gt; 400</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>DF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>–</td>
</tr>
<tr>
<td>CR-FS-high</td>
<td>&gt; 200</td>
<td>&gt; 700</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>DF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>–</td>
</tr>
<tr>
<td>CR( \gamma )-low</td>
<td>–</td>
<td>–</td>
<td>≥ 2</td>
<td>–</td>
<td>SF, 1( \gamma )</td>
<td>–</td>
<td>–</td>
<td></td>
</tr>
<tr>
<td>CR( \gamma )-medium</td>
<td>–</td>
<td>&gt; 400</td>
<td>≥ 2</td>
<td>–</td>
<td>SF, 1( \gamma )</td>
<td>–</td>
<td>–</td>
<td></td>
</tr>
<tr>
<td>CR( \gamma )-high</td>
<td>–</td>
<td>&gt; 700</td>
<td>≥ 2</td>
<td>–</td>
<td>SF, 1( \gamma )</td>
<td>–</td>
<td>–</td>
<td></td>
</tr>
<tr>
<td>CR-real</td>
<td>–</td>
<td>&gt; 200</td>
<td>≥ 2</td>
<td>81–101</td>
<td>2( \ell ) SF</td>
<td>OS</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>CR-fake</td>
<td>&lt; 125</td>
<td>–</td>
<td>–</td>
<td>( \epsilon [12, \infty], \not \in [81, 101] ) (SF)</td>
<td>2( \ell ) SF/DF</td>
<td>SS</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td><strong>Validation regions</strong></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>VR-low</td>
<td>100–200</td>
<td>–</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>SF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>–</td>
</tr>
<tr>
<td>VR-medium</td>
<td>100–200</td>
<td>&gt; 400</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>SF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>–</td>
</tr>
<tr>
<td>VR-high</td>
<td>100–200</td>
<td>&gt; 700</td>
<td>≥ 2</td>
<td>&gt; 12</td>
<td>SF</td>
<td>OS</td>
<td>&gt; 0.4</td>
<td>–</td>
</tr>
<tr>
<td>VR-fake</td>
<td>&gt; 50</td>
<td>–</td>
<td>≥ 2</td>
<td>( \epsilon [12, \infty], \not \in [81, 101] ) (SF)</td>
<td>SF/DF</td>
<td>SS</td>
<td>–</td>
<td>–</td>
</tr>
</tbody>
</table>
To select events containing a leptonically decaying Z boson, the invariant mass of the dilepton system is required to be $81 < m_{\ell\ell} < 101$ GeV. In the CRs and VRs that use the Z mass sidebands, only events with $m_{\ell\ell} > 45$ GeV are used to reject the lower $m_{\ell\ell}$ region dominated by Drell–Yan (DY) production. In Ref. [20] an “on-Z” SR, denoted SRZ, is defined requiring $E_T^{\text{miss}} > 225$ GeV and $H_T^{\text{incl}} > 600$ GeV. The region is motivated by SUSY signals with high gluino or squark mass and high jet activity. Since $b$-jets are not always expected in the simplified models considered here, no requirement is placed on $b$-tagged jet multiplicity ($n_{b-jets}$) so as to be as inclusive as possible and to be consistent with Ref. [20]. Dedicated CRs are defined, with selection criteria similar to those of SRZ, to estimate the contribution from the dominant SM backgrounds in SRZ. These CRs are discussed in more detail in Sect. 7.

The edge selection requires at least two leptons with $p_T > 25$ GeV. The search is performed across the full $m_{\ell\ell}$ spectrum, with the exception of the region with $m_{\ell\ell} < 12$ GeV, which is vetoed to reject low-mass DY events and the $J/\psi$ and $\Upsilon$ resonances. Three regions are defined to target signal models with low, medium and high values of $\Delta m_{\tilde{\chi}} = m(\tilde{g}) - m(\tilde{\chi}^0_1)$, denoted SR-low, SR-medium, and SR-high, respectively. All these regions require $E_T^{\text{miss}} > 200$ GeV. SR-medium and SR-high also include the requirements $H_T > 400$ GeV and $H_T > 700$ GeV, respectively, to further isolate high-$\Delta m_{\tilde{\chi}}$ events. Here the leptons are not included in the $H_T$ definition to avoid introducing any bias in the $m_{\ell\ell}$ distribution. Events selected in SR-low, SR-medium and SR-high are further grouped into non-orthogonal $m_{\ell\ell}$ windows, which represent the search regions used in the edge analysis. The dilepton mass ranges of these are chosen to maximise sensitivity to the targeted signal models, with the window boundaries being motivated by the dilepton mass endpoints of generated signal points. In total, 24 $m_{\ell\ell}$ windows are defined by selecting ranges with the best expected sensitivity to signal models. Of these windows, nine are in SR-low, eight are in SR-medium and seven are in SR-high. Details of the $m_{\ell\ell}$ definitions in these regions are given along with the results in Sect. 9. Models without light sleptons are targeted by windows with $m_{\tilde{\ell}} < 60$ GeV or $m_{\tilde{\ell}} < 80$ GeV for $\Delta m_{\chi} < m_Z$ leading to off-shell Z bosons, and by the window with $81 < m_{\tilde{\ell}} < 101$ GeV for $\Delta m_{\chi} > m_Z$ leading to on-shell Z bosons. Models with light sleptons are targeted by the remaining $m_{\ell\ell}$ windows, which cover the full $m_{\ell\ell}$ range. The edge selection and on-shell Z selection are not orthogonal. In particular, SR-medium in the range $81 < m_{\ell\ell} < 101$ GeV overlaps significantly with SRZ.

For the combined $ee + \mu\mu$ channels, the typical signal acceptance times efficiency values for the signal models considered in SRZ are 2–8%. They are 8–40%, 3–35%, and 1–35%, inclusively in $m_{\ell\ell}$, for SR-low, SR-medium and SR-high, respectively. The on-shell Z and edge analyses are each optimised for different signal models. There are models in which signal contamination in CRs or VRs can become significant. For example, CRT in Table 3 is used to normalise the $t\bar{t}$ MC sample to data as a cross-check in the on-shell Z search, but it is a region where the signal contamination from signal models targeted by the edge search can be up to 80% relative to the expected background. In addition, the contamination from on-shell Z signals in the region used to validate the $Z/\gamma^* +$ jets and flavour-symmetric estimates, VR-S, is up to 60% for models with $m(\tilde{g}) < 1$ TeV. The signal contamination from the slepton models in the DF regions used to estimate the flavour-symmetric backgrounds in the edge search, CR-FS-low/medium/high in Table 4, is less than 20% for models with $m(\tilde{g}) > 600$ GeV. It is only the contamination in these $e\mu$ CRs that is relevant in terms of the model-dependent interpretation of the results, and its impact is further discussed in Sect. 10. In general, for models giving substantial
contamination in the CRs, the signal-to-background ratio in the SRs is found to be large enough for this contamination to have negligible impact on the sensitivity of the search.

7 Background estimation

The dominant background processes in the SRs are “flavour-symmetric” (FS) backgrounds, where the ratio of $ee$, $\mu\mu$ and $e\mu$ dileptonic branching fractions is 1:1:2 because the two leptons originate from independent $W \to \ell\nu$ decays. This background is dominated by $t\bar{t}$ (50–70%) and also includes $WW$, $Wt$, and $Z \to \tau\tau$ processes. The FS background constitutes 60–90% of the expected SM background in the SRs, and is estimated using control samples of $e\mu$ events.

As all the SRs have a high-$E_T^{miss}$ requirement, $Z/\gamma^* + \text{jets}$ events only enter the SRs when there is large $E_T^{miss}$ originating from instrumental effects or from neutrinos in jet fragments. This background is generally small, but it is difficult to model with MC simulation and can mimic signal, particularly for the on-shell $Z$ search. This background is estimated using a control sample of $\gamma + \text{jets}$ events in data, which are kinematically similar to $Z/\gamma^* + \text{jets}$ and have similar sources of $E_T^{miss}$.

The production of $WZ/ZZ$ dibosons contributes approximately 30% of the SM background in SRZ and up to 20% of the background in the edge SR $m_{\ell\ell}$ windows. These backgrounds are estimated from MC simulation, after validation in dedicated 3$\ell$ ($WZ$) and 4$\ell$ ($ZZ$) VRs. Rare top backgrounds, which include $t\bar{t}W$, $t\bar{t}Z$ and $t\bar{t}WW$ processes, constitute <5% of the expected SM background in all SRs, and are estimated from MC simulation. The contribution from events with fake or misidentified leptons is at most 15% (in one of the edge $m_{\ell\ell}$ ranges in SR-low), but is generally <5% of the expected SM background in most SRs.

7.1 Flavour-symmetric backgrounds

The flavour-symmetric background is dominant in all SRs. To estimate the contribution of this background to each SR, the so-called “flavour-symmetry” method, detailed in Ref. [20], is used. In this method, data events from a DF control sample, which is defined with the same kinematic requirements as the SR, are used to determine the expected event yields in the SF channels. In the on-shell $Z$ analysis, the method is used to predict the background yield in the $Z$ mass window, defined as $81 < m_{\ell\ell} < 101$ GeV. In the edge analysis, the method is extended to predict the full dilepton mass shape, such that a prediction can be extracted in any of the predefined $m_{\ell\ell}$ windows.

For the edge search, the flavour-symmetric contribution to each $m_{\ell\ell}$ bin of the signal regions is predicted using data from the corresponding bin in an $e\mu$ control region. All edge CR-FS regions (definitions can be seen in Table 4) are 88–97% pure in flavour-symmetric processes (this purity is calculated from MC simulation).

For the on-shell search, this method is complicated slightly by a widening of the $m_{\ell\ell}$ window used in CR-FS, the $e\mu$ control region (defined in Table 3). The window is enlarged to $61 < m_{\ell\ell} < 121$ GeV to approximately triple the amount of data in the control region and thus increase the statistical precision of the method. This results in a region that is ~95% pure in flavour-symmetric processes (the expected composition of this 95% is ~80% $t\bar{t}$, ~10% $Wt$, ~10% $WW$ and <1% $Z \to \tau\tau$).

Apart from the $m_{\ell\ell}$ widening in CR-FS, the method used is identical for the on-shell and edge regions. Events in the control regions are subject to lepton $p_T^{\ell}$ and $\eta$-dependent correction factors measured in data and MC simulation. Because the triggers used are not identical in 2015 and 2016, these factors are measured separately for each year and account for the different identification and reconstruction efficiencies for electrons and muons, as well as the different trigger efficiencies for the dilepton, dimuon and electron–muon selections. The estimated numbers of events in the SF channels, $N_{ee/\mu\mu}^{\text{est}}$, are given by:

$$N_{ee/\mu\mu}^{\text{est}} = \frac{1}{2} \cdot f_{FS} \cdot f_{Z-mass} \cdot \sum_k k_e(p_T^{(3)\ell}, \eta^{(3)\ell}) \cdot \alpha(p_T^{(3)\ell}, \eta^{(3)\ell}),$$ (1)

$$N_{\mu\mu}^{\text{est}} = \frac{1}{2} \cdot f_{FS} \cdot f_{Z-mass} \cdot \sum_k k_\mu(p_T^{(2)\ell}, \eta^{(2)\ell}) \cdot \alpha(p_T^{(2)\ell}, \eta^{(2)\ell}),$$ (2)

where $N_{ee/\mu\mu}^{\text{data}}$ is the number of data events observed in a given control region, $\alpha(p_T^{(3)\ell}, \eta^{(3)\ell})$ accounts for the different trigger efficiencies for SF and DF events, and $k_e(p_T^{(3)\ell}, \eta^{(3)\ell})$ and $k_\mu(p_T^{(2)\ell}, \eta^{(2)\ell})$ are electron and muon selection efficiency factors for the kinematics of the lepton being replaced, in event $i$. The trigger and selection efficiency correction factors are derived from the events in an inclusive on-$Z$ selection ($81 < m_{\ell\ell} < 101$ GeV, $\geq 2$ jets), according to:

$$k_e(p_T, \eta) = \sqrt{\frac{N_{ee}^{\text{meas}(p_T, \eta)}}{N_{ee}^{\text{meas}(p_T, \eta)}}},$$ (3)

$$k_\mu(p_T, \eta) = \sqrt{\frac{N_{\mu\mu}^{\text{meas}(p_T, \eta)}}{N_{ee}^{\text{meas}(p_T, \eta)}}},$$ (4)

$$\alpha(p_T, \eta) = \sqrt{\frac{\epsilon_{ee}^{\text{trig}}(p_T^{(3)\ell}, \eta^{(3)\ell}) \times \epsilon_{\mu\mu}^{\text{trig}}(p_T^{(2)\ell}, \eta^{(2)\ell})}{\epsilon_{e\mu}^{\text{trig}}(p_T^{(3)\ell}, \eta^{(3)\ell})}},$$ (5)
where $\epsilon_{ee/\mu\mu}^{\text{trig}}$ is the trigger efficiency and $N_{ee/\mu\mu}^{\text{meas}}$ is the number of $ee/\mu\mu$ events in the inclusive on-Z region outlined above. Here $k_{e}(p_{T}, \eta)$ and $k_{\mu}(p_{T}, \eta)$ are calculated separately for leading and sub-leading leptons, while $\alpha$ is calculated for the leading lepton, $\ell_{1}$. The correction factors are typically within 10% of unity, except in the region $|\eta| < 0.1$ where, because of the lack of coverage by the muon spectrometer, they are up to 50% from unity. For all background estimates based on the flavour-symmetry method, results are computed separately for $ee$ and $\mu\mu$ and then summed to obtain the combined predictions. The resulting estimates from the DF channels are scaled according to the fraction of flavour-symmetric backgrounds in each $e\mu$ control sample, $f_{\text{FS}}$ (95% in CR-FS), which is determined by subtracting non-flavour-symmetric backgrounds taken from MC simulation from the data observed in the corresponding $e\mu$ region. In the on-shell case, the result is also scaled by the fraction of events in CR-FS expected to be contained within $81 < m_{\ell\ell} < 101$ GeV, $f_{Z,\text{mass}}$ (38%), which is otherwise set to 100% for the edge regions. The validity of extrapolating in $m_{\ell\ell}$ between CR-FS and SRZ was checked by comparing the $m_{\ell\ell}$ shape in data and MC simulation in a region similar to VR-S, but with the $m_{\ell\ell}$ requirement relaxed and $H_{\ell\ell}^{\text{incl}} > 300$ GeV to obtain a sample with a large number of events. The resulting on-Z fractions in MC simulation were found to agree with data within statistical uncertainties, which are summed in quadrature to assign a systematic uncertainty. In the case of the edge search the full $m_{\ell\ell}$ distribution is validated by applying a flavour-symmetry method to $t\bar{t}$ MC events in VR-low, VR-medium and VR-high. This procedure results in good closure, which is further discussed in Sect. 7.5. The difference between the prediction and the observed distribution is used to assign an MC non-closure uncertainty to the estimate.

The flavour-symmetry method in SRZ is further cross-checked by performing a profile likelihood fit [90] of MC yields to data in the $Z$-mass sidebands ($m_{\ell\ell} \not\in [81, 101]$ GeV), the region denoted CRT in Table 3, which is dominated by $t\bar{t}$ (with a purity of >75%) and contains 273 events in data. The other flavour-symmetric processes in this region contribute $\sim 12\%$ ($Wt$), $10\%$ ($WW$) and <1% ($Z \rightarrow \tau\tau$). All SM background processes are taken directly from MC simulation in this cross-check, including backgrounds also estimated using the flavour-symmetry method. The normalisation of the dominant $t\bar{t}$ background is a free parameter and is the only parameter affected by the fit. For this cross-check, the contamination from Beyond Standard Model processes in the $Z$-mass sidebands is assumed to be negligible. The fit results in a scale factor of 0.64 for the $t\bar{t}$ yield predicted by simulation. This result is extrapolated from the $Z$-mass sidebands to SRZ and gives a prediction of $29 \pm 7$ events, which is consistent with the nominal flavour-symmetry background estimate of $33 \pm 4$ in this region.

The sideband fit is repeated at lower $E_{T}^{\text{miss}}$ in VRT, with the results being propagated to VR-S, so as to test the $m_{\ell\ell}$ extrapolation used in the sideband fit method. The normalisation to data in this region, which is at lower $E_{T}^{\text{miss}}$ relative to CRT, results in a scale factor of 0.80 for the $t\bar{t}$ yield predicted by simulation. The number of FS events predicted in VR-S using the sideband fit in VRT is compatible with the number estimated by applying the FS method to data in VR-FS. The results of the background estimate in both VR-S and SRZ obtained from the flavour-symmetry method are compared with the values obtained by the sideband fit cross-check in Table 5. The methods result in consistent estimates in both regions. Further results in the edge VRs are discussed in Sect. 7.5.

A potential cause of the low scale factors obtained from the sideband fit at large $H_{T}$ and $E_{T}^{\text{miss}}$ is mismodelling of the top-quark $p_{T}$ distribution, where measurements of $t\bar{t}$ differential cross sections by the ATLAS and CMS experiments indicate that the top-quark $p_{T}$ distribution predicted by most generators is harder than that observed in data [91,92]. Corrections to the MC predictions according to NNLO calculations provided in Ref. [93] indicate an improvement in the top-quark pair modelling at high $H_{T}$, which should lead to scale factors closer to unity. With the data-driven method used to estimate $t\bar{t}$ contributions in this analysis, the results do not depend on these corrections. They are therefore not applied to the $t\bar{t}$ MC sample for the sideband-fit cross-check.

### 7.2 $Z/\gamma^{\ast} + \text{jets}$ background

The $Z/\gamma^{\ast} + \text{jets}$ background estimate is based on a data-driven method that uses $\gamma + \text{jets}$ events in data to model the $E_{T}^{\text{miss}}$ distribution of $Z/\gamma^{\ast} + \text{jets}$. The $\gamma + \text{jets}$ and $Z/\gamma^{\ast} + \text{jets}$ processes have similar event topologies, with a well-measured object recoiling against a hadronic system, and both tend to have $E_{T}^{\text{miss}}$ that stems from jet mismeasurements and neutrinos in hadronic decays. In this method, which has been used by CMS in a search in this final state [18], a sample of data events containing at least one photon and no leptons is constructed using the same kinematic selection as each of

<table>
<thead>
<tr>
<th>Region</th>
<th>Flavour-symmetry</th>
<th>Sideband fit</th>
</tr>
</thead>
<tbody>
<tr>
<td>SRZ</td>
<td>33 ± 4</td>
<td>29 ± 7</td>
</tr>
<tr>
<td>VR-S</td>
<td>99 ± 8</td>
<td>92 ± 25</td>
</tr>
</tbody>
</table>
the SRs, without the $E_{T}^{\text{miss}}$ and $\Delta \phi (\text{jet}_{12}, p_{T}^{\text{miss}})$ requirements (the CRγγ regions defined in Tables 3, 4).

The requirement $\Delta \phi (\text{jet}_{12}, p_{T}^{\text{miss}}) > 0.4$ applied in the SRs suppresses $E_{T}^{\text{miss}}$ from jet mismeasurements and increases the relative contributions to $E_{T}^{\text{miss}}$ from the photon, electrons, and muons. The difference in resolution between photons, electrons, and muons can be significant at high $p_{T}$. Therefore, before the $\Delta \phi (\text{jet}_{12}, p_{T}^{\text{miss}}) > 0.4$ requirement is applied, the photon $p_{T}$ is smeared according to a $Z \rightarrow ee$ or $Z \rightarrow \mu\mu$ resolution function. The smearing function is derived by comparing the $E_{T}^{\text{miss}}$-projection along the boson momentum in $Z/\gamma^{*} + \text{jets}$ and $\gamma + \text{jets}$ MC events in a 1-jet control region with no other event-level kinematic requirements. A deconvolution is applied to avoid including the photon resolution in the $Z$ resolution. For each event, a photon $p_{T}$ smearing $\Delta p_{T}$ is obtained by sampling the smearing function. The photon $p_{T}$ is shifted by $\Delta p_{T}$, with the parallel component of the $E_{T}^{\text{miss}}$ being correspondingly adjusted by $-\Delta p_{T}$.

The smeared $\gamma + \text{jets}$ events are then reweighted to match the boson $p_{T}$ distribution of the $Z/\gamma^{*} + \text{jets}$ events. This reweighting is applied separately in each region and accounts for small differences between the $\gamma + \text{jets}$ events and $Z/\gamma^{*} + \text{jets}$ events, which arise mainly from the mass of the $Z$ boson. The reweighting is done using $Z/\gamma^{*} + \text{jets}$ events in data, and is checked using $Z/\gamma^{*} + \text{jets}$ MC simulation in an MC closure test, as described further below. Following this smearing and reweighting procedure, the $E_{T}^{\text{miss}}$ of each $\gamma + \text{jets}$ event is recalculated, and the final $E_{T}^{\text{miss}}$ distribution is obtained after applying the $\Delta \phi (\text{jet}_{12}, p_{T}^{\text{miss}}) > 0.4$ requirement. For each SR, the resulting $E_{T}^{\text{miss}}$ distribution is normalised to data in a CRZ with the same requirements except that the SR $E_{T}^{\text{miss}}$ requirement is replaced by $E_{T}^{\text{miss}} < 60 \text{ GeV}$.

The shape of the $Z/\gamma^{*} + \text{jets}$ $m_{\ell\ell}$ distribution is extracted from MC simulation and validated by comparing to data in events with lower $E_{T}^{\text{miss}}$ requirements and a veto on $b$-tagged jets, to suppress the background from $t\bar{t}$. The $m_{\ell\ell}$ distribution is modelled by parameterising the $m_{\ell\ell}$ in $Z/\gamma^{*} + \text{jets}$ events as a function of the difference between reconstructed and true $Z$ boson $p_{T}$ in MC simulation. This parameterization ensures that the correlation between lepton momentum mismeasurement and observed $m_{\ell\ell}$ values far from the $Z$ boson mass is preserved. Each photon event is assigned an $m_{\ell\ell}$ via a random sampling of the corresponding distribution, equating photon $\Delta p_{T}$ and the difference between true and reconstructed $Z$ boson $p_{T}$. The resulting $m_{\ell\ell}$ distribution in $\gamma + \text{jets}$ MC simulation is compared to that extracted from $Z/\gamma^{*} + \text{jets}$ MC simulation and the difference is assessed as a systematic uncertainty in the background prediction for each $m_{\ell\ell}$ bin.

The full smearing, reweighting, and $m_{\ell\ell}$ assignment procedure is applied to the $V\gamma$ MC sample in parallel with the $\gamma + \text{jets}$ data sample. After applying all corrections to both samples, the $V\gamma$ contribution to the $\gamma + \text{jets}$ data sample is subtracted to remove contamination from backgrounds with real $E_{T}^{\text{miss}}$. Contamination by events with fake photons in these $\gamma + \text{jets}$ data samples is small, and this contribution is therefore neglected.

In the $H_{T}$-inclusive region corresponding to VR-low, there is a non-negligible contribution expected from $Z/\gamma^{*} + \text{jets}$ events with $p_{T}^{Z} < 37 \text{ GeV}$. Given the photon trigger strategy discussed in Sect. 4, no photons with $p_{T} < 37 \text{ GeV}$ are included in the event selection. To account for this photon $p_{T}$ threshold, a boson-$p_{T}$ correction of up to 50% is applied as a function of $E_{T}^{\text{miss}}$ in VR-low. This correction uses the fraction of $Z/\gamma^{*} + \text{jets}$ events in a given $E_{T}^{\text{miss}}$ bin expected to have $p_{T}^{Z} < 37 \text{ GeV}$, according to MC simulation. The $\gamma + \text{jets}$ data are then scaled according to this fraction, as a function of $E_{T}^{\text{miss}}$, to correct for the missing $p_{T}^{Z} < 37 \text{ GeV}$ contribution. The correction is found to be negligible in all signal regions.

The distribution of $E_{T}^{\text{miss}}$ obtained in SHERPA $Z/\gamma^{*} + \text{jets}$ MC simulation is compared to that obtained by applying this background estimation technique to SHERPA $\gamma + \text{jets}$ MC samples. In this check the $\gamma + \text{jets}$ MC simulation is reweighted according to the $p_{T}$ distribution given by the $Z/\gamma^{*} + \text{jets}$ MC simulation. The result of this MC closure check is shown in Fig. 3a for events in VRZ (without an upper $E_{T}^{\text{miss}}$ cut), where good agreement between $Z/\gamma^{*} + \text{jets}$ and corrected $\gamma + \text{jets}$ MC simulation can be seen across the entire $E_{T}^{\text{miss}}$ spectrum. A comparison between the full $E_{T}^{\text{miss}}$ spectrum in data and the $Z/\gamma^{*} + \text{jets}$ background estimated via the $\gamma + \text{jets}$ method is also shown in Fig. 3b for events in VRZ. The systematic uncertainties associated with this method are described in Sect. 8.

7.3 Fake-lepton background

Semileptonic $t\bar{t}$, $W \rightarrow \ell\nu$ and single top ($s$- and $t$-channel) events enter the dilepton channels via “fake” leptons. These can include misidentified hadrons, converted photons or non-prompt leptons from $b$-hadron decays. The extent of this background is estimated using the matrix method, detailed in Ref. [94]. Its contribution in regions with high lepton $p_{T}$ and dilepton invariant mass is negligible, but in the edge search, where lower-$p_{T}$ leptons are selected and events can have low $m_{\ell\ell}$, the fake-lepton background can make up to 15% of the total background. In this method a control sample is constructed using baseline leptons, thereby enhancing the probability of selecting a fake lepton due to the loosier lepton selection and identification criteria relative to the signal lepton selection. For each relevant CR, VR or SR, the region-specific kinematic requirements are placed upon this sample of baseline leptons. The number of events in this sample in which the selected leptons subsequently pass ($N_{\text{pass}}$) or fail ($N_{\text{fail}}$) the signal lepton requirements in Sect. 5 are
then counted. In the case of a one-lepton selection, the number of fake-lepton events in a given region is then estimated according to:

$$N_{\text{pass}}^{\text{fake}} = \frac{N_{\text{fail}} - (1/\epsilon_{\text{real}} - 1) \times N_{\text{pass}}^{\text{real}}}{1/\epsilon_{\text{fake}} - 1/\epsilon_{\text{real}}}.$$  \hfill (6)

Here $\epsilon_{\text{real}}$ is the relative identification efficiency (from baseline to signal) for genuine, prompt (“real”) leptons and $\epsilon_{\text{fake}}$ is the relative identification efficiency (again from baseline to signal) with which non-prompt leptons or jets might be misidentified as prompt leptons. This principle is then expanded to a dilepton selection by using a four-by-four matrix to account for the various possible real–fake combinations for the two leading leptons in an event.

The real-lepton efficiency, $\epsilon_{\text{real}}$, is measured in $Z \to \ell\ell$ data events using a tag-and-probe method in CR-real, defined in Table 4. In this region the $p_T$ of the leading lepton is required to be $>40$ GeV, and only events with exactly two SFOS leptons are selected. The fake-lepton efficiency, $\epsilon_{\text{fake}}$, is measured in CR-fake, a region enriched with fake leptons by requiring same-sign lepton pairs. The lepton $p_T$ requirements are the same as those in CR-real, with the leading lepton being tagged as the “real” lepton and the fake efficiency being evaluated based on the sub-leading lepton in the event. An $E_T^{\text{miss}}$ requirement of $<125$ GeV is used to reduce possible contamination from Beyond Standard Model processes. In this region the background due to prompt-lepton production, estimated from MC simulation, is subtracted from the total data contribution. Prompt-lepton production makes up 7% (11%) of the baseline electron (muon) sample and 10% (61%) of the signal electron (muon) sample in CR-fake. From the resulting data sample the fraction of events in which the baseline leptons pass a signal-like selection yields the fake efficiency. Both the real- and fake-lepton efficiencies are binned as a function of lepton $p_T$ and calculated separately for the 2015 and 2016 data sets.

This method is validated by checking the closure in MC simulation and data–background agreement in VR-fake.

7.4 Diboson and rare top processes

The remaining SM background contribution in the SRs is due to $WZ/ZZ$ diboson production and rare top processes ($t\bar{t}Z$, ...
Rare top events compose <5% of the expected SM background in the SRs and are taken directly from MC simulation.

Production of \( WZ/ZZ \) dibosons constitutes about 30% of the expected background in SRZ and up to 20% in some edge SR \( m_{\ell\ell} \) windows. In SRZ, this background is composed of roughly 70% \( WZ \), about 40% of which is \( WZ \rightarrow \ell \ell \tau \nu \). This is the largest background contribution that is estimated from MC simulation, and must be carefully validated, especially because these backgrounds contain \( Z \) bosons and can thus mimic a signal by producing a peak at \( m_{\ell\ell} \approx m_Z \). To validate the MC modelling of these backgrounds, VRs with three leptons (VR-WZ) and four leptons (VR-ZZ) are defined (selection shown in Table 3). In VR-WZ, from the three selected leptons in an event, the SFOS pair with \( m_{\ell\ell} \) most consistent with the \( Z \) mass is indentified as the \( Z \) candidate. The transverse mass of the remaining lepton and the \( \mathbb{E}_T^{\text{miss}} \), \( \mathbb{m}_T(\ell_3, \mathbb{E}_T^{\text{miss}}) \), is then required to be <100 GeV, forming the \( W \) candidate. In VR-ZZ an \( \mathbb{E}_T^{\text{miss}} < 100 \) GeV requirement is used to suppress \( WZ \) and top processes. The yields and kinematic distributions observed in these regions are well-modelled by MC simulation. In particular, the \( \mathbb{E}_T^{\text{miss}}, H_T, \) jet multiplicity, and boson \( p_T \) distributions show good agreement. An additional three-lepton VR (VR-3L) is defined to provide validation of the diboson background in a region of phase space closer to the SR; good agreement is observed in this region as well.

### Table 6

<table>
<thead>
<tr>
<th></th>
<th>VR-S</th>
<th>VR-WZ</th>
<th>VR-ZZ</th>
<th>VR-3L</th>
</tr>
</thead>
<tbody>
<tr>
<td>Observed events</td>
<td>236</td>
<td>698</td>
<td>132</td>
<td>32</td>
</tr>
<tr>
<td>Total expected background events</td>
<td>224 ± 41</td>
<td>622 ± 66</td>
<td>139 ± 25</td>
<td>35 ± 10</td>
</tr>
<tr>
<td>Flavour-symmetric (( t\bar{t}W/t\bar{t}WW ))</td>
<td>( 99 \pm 8 )</td>
<td>–</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>( WZ/ZZ ) events</td>
<td>( 27 \pm 13 )</td>
<td>( 573 \pm 66 )</td>
<td>( 139 \pm 25 )</td>
<td>( 25 \pm 10 )</td>
</tr>
<tr>
<td>Rare top events</td>
<td>( 11 \pm 3 )</td>
<td>( 14 \pm 3 )</td>
<td>( 0.44 \pm 0.11 )</td>
<td>( 9.1 \pm 2.3 )</td>
</tr>
<tr>
<td>( Z/\gamma^* + \text{jets} ) events</td>
<td>( 84 \pm 37 )</td>
<td>–</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Fake-lepton events</td>
<td>( 4 \pm 4 )</td>
<td>( 35 \pm 6 )</td>
<td>–</td>
<td>( 0.6 \pm 0.3 )</td>
</tr>
</tbody>
</table>

### Table 7

<table>
<thead>
<tr>
<th></th>
<th>VR-low</th>
<th>VR-medium</th>
<th>VR-high</th>
</tr>
</thead>
<tbody>
<tr>
<td>Observed events</td>
<td>16,253</td>
<td>1917</td>
<td>314</td>
</tr>
<tr>
<td>Total expected background events</td>
<td>16,500 ± 700</td>
<td>1990 ± 150</td>
<td>340 ± 60</td>
</tr>
<tr>
<td>Data-driven flavour-symmetry events</td>
<td>14,700 ± 600</td>
<td>1690 ± 120</td>
<td>250 ± 50</td>
</tr>
<tr>
<td>( WZ/ZZ ) events</td>
<td>250 ± 80</td>
<td>40 ± 19</td>
<td>9 ± 6</td>
</tr>
<tr>
<td>Data-driven ( Z/\gamma^* + \text{jets} (\gamma + \text{jets}) ) events</td>
<td>1100 ± 400</td>
<td>130 ± 70</td>
<td>50 ± 29</td>
</tr>
<tr>
<td>Rare top events</td>
<td>87 ± 23</td>
<td>27 ± 7</td>
<td>6.5 ± 1.8</td>
</tr>
<tr>
<td>Data-driven fake-lepton events</td>
<td>270 ± 100</td>
<td>98 ± 35</td>
<td>20 ± 11</td>
</tr>
</tbody>
</table>

### 7.5 Results in validation regions

The expected background yields in VR-S are shown in Table 6 and compared with the observed data yield. Agreement between the data and the expected Standard Model background is observed. The expected background yields in the three diboson VRs are also shown in Table 6. The data are consistent with the expected background. Similar information for the edge VRs is provided in Table 7. Data and background estimates are in agreement within uncertainties.

Figure 4 shows the observed and expected \( m_{\ell\ell} \) distributions in the same edge VRs. The same background estimation methods are applied to both MC simulation and data. In the MC studies, the flavour-symmetry method of Sect. 7.1 is applied to \( t\bar{t} \) MC simulation, and the observed SF \( m_{\ell\ell} \) distribution is compared to the prediction based on DF events. In the data studies, the observed SF \( m_{\ell\ell} \) distribution is compared to the sum of FS backgrounds from the extended flavour-symmetry method, the \( Z/\gamma^* + \text{jets} \) background from the
Fig. 4 Validation of the flavour-symmetry method for the edge search using MC events (left) and data (right), in the VR-low (top), VR-medium (middle), and VR-high (bottom) regions. In the MC plots the flavour-symmetry estimate from $e\mu\bar{t}$ MC samples is compared with the observed SF distribution from these MC samples, with the MC statistical uncertainty indicated by the hashed bands. In the data plots, all uncertainties in the background prediction are included in the hashed band. The rare top and data-driven fake-lepton backgrounds are grouped under “other” backgrounds. The bottom panel of each figure shows the ratio of the observation (left in MC simulation; right in data) to the prediction. In cases where the data point is not accommodated by the scale of this panel, a red arrow indicates the direction in which the point is out of range. The last bin contains the overflow.
y + jets method, and the WZ/ZZ diboson, rare top, and fake-lepton backgrounds.

The observed MC closure is good in all validation regions. The data agree with the expected background in the validation regions as well. No significant discrepancies or trends are apparent.

8 Systematic uncertainties

The data-driven background estimates are subject to uncertainties associated with the methods employed and the limited number of events used in their estimation. The dominant uncertainty (10%) for the flavour-symmetry-based background estimate in SRZ is due to the limited number of events in CR-FS. Other systematic uncertainties assigned to this background estimate include those due to MC closure (3%), the measurement of the efficiency correction factors (3%) and the extrapolation in m_{ll} (1%). In the case of the edge SRs the statistical uncertainty is also the dominant uncertainty in the flavour-symmetric background estimate in the case of SR-high, but for both SR-medium and SR-low the uncertainties from the MC non-closure and efficiency correction factors are comparable in size, or in some cases larger. These uncertainties can contribute up to 5% in SR-low and SR-medium and 10% in SR-high.

Several sources of systematic uncertainty are assessed for the Z/γ^* + jets background. The boson p_T reweighting procedure is assigned an uncertainty based on a comparison of the nominal results with those obtained by reweighting using three other kinematic variables, namely H_T, Z-boson E_T and jet multiplicity. For the smearing function, which is measured using MC events in a 1-jet control region, an uncertainty is derived by comparing the results obtained using the nominal smearing function with those obtained using a smearing function from a 2-jet sample of MC events, and also using a smearing function measured in a 1-jet data sample. An uncertainty of between 40–100% is assigned to account for different reweighting procedures and between 20–100% for the smearing procedure applied to γ + jets events. The smearing uncertainty dominates in SR-high, while the reweighting uncertainty dominates in SR-low and SR-medium, with both being around 60% in SRZ. The full reweighting and smearing procedure is carried out using γ + jets MC events such that an MC non-closure uncertainty can be derived by comparing the resulting γ + jets MC E_T^{miss} distribution to that in Z/γ^* + jets MC events. The resulting uncertainty of up to 35% is calculated in the VRs, so as to maximise the number of events that contribute. An uncertainty of 16% is assessed for the Vγ backgrounds, based on data-to-MC agreement in a Vγ-enriched control region. This uncertainty is propagated to the final Z/γ^* + jets estimate following the subtraction of the Vγ background. In VR-low, a correction is applied to the E_T^{miss} distribution in γ + jets events to account for the fraction of Z/γ^* + jets events in this H_T-inclusive region expected to have boson p_T less than 37 GeV. The full size of this correction (up to 50% for E_T^{miss} = 150 GeV) is applied as a systematic uncertainty. The m_{ll} distribution assigned to γ + jets MC events is compared to that of Z/γ^* + jets MC events, and the relative difference in a given m_{ll} bin is assigned as an uncertainty. Finally, the statistical precision of the estimate also enters as a systematic uncertainty of ~10% in the final background estimate. After applying the correction procedure, differences in the number of b-tagged jets between Z/γ^* + jets and γ + jets are found to be negligible, indicating good agreement in heavy-flavour content.

The uncertainties in the fake-lepton background stem from the number of events in the regions used to measure the real- and fake-lepton efficiencies, the limited size of the inclusive loose-lepton sample, and from varying the region used to measure the fake-lepton efficiency. The nominal fake-lepton efficiency is compared with those measured in a region with b-tagged jets and a region with a b-jet, as well as a region with the prompt-lepton subtraction varied by 20%. Varying the sample composition via b-jet tagging gives the largest uncertainty. The uncertainty for the edge SRs from the statistical component of the lepton efficiencies is 30–45%, and from varying the region for the fake-lepton efficiency it is 50–75%. The uncertainties in SRZ are generally larger due to the small number of events contributing to the estimate in this region.

Theoretical and experimental uncertainties are taken into account for the signal models, as well as background processes that rely on MC simulation. The estimated uncertainty in the luminosity measurement is 2.9% [31,32]. The jet energy scale is subject to uncertainties associated with the jet flavour composition, the pile-up and the jet and event kinematics [81]. Uncertainties in the jet energy resolution are included to account for differences between data and MC simulation [81]. An uncertainty in the E_T^{miss} soft-term resolution and scale is taken into account [88], and uncertainties due to the lepton energy scales and resolutions, as well as trigger, reconstruction, and identification efficiencies, are also considered.

The WZ/ZZ processes are assigned a cross-section uncertainty of 6% and an additional uncertainty based on comparisons between SHERPA and POWHEG MC samples, which is up to 50% in the SRs. Uncertainties due to the choice of factorisation and renormalisation scales are calculated by varying the nominal values up and down by a factor of two and can be up to 23%. For rare top processes, a 13% PDF and scale variation uncertainty is applied [34] in addition to a 22% cross-section uncertainty [61–63].

For signal models, the nominal cross section and the uncertainty are taken from an envelope of cross-section predictions using different PDF sets and factorisation and renormalisa-
tion scales, as described in Refs. [95, 96]. These are calculated at next-to-leading-logarithm accuracy (NLO + NLL) [51–55], and the resulting uncertainties range from 16 to 30%.

A breakdown of the dominant uncertainties in the background prediction in the SRs is provided in Table 8 for the on-shell Z and edge searches. Here these uncertainties are quoted relative to the total background. In the case of the edge regions a range is quoted, taking into account the relative contribution of the given uncertainty in each of the \(m_{\ell\ell}\) ranges in SR-low, SR-medium and SR-high. The largest uncertainties in the signal regions are due to the size of the \(e\mu\) data sample in CR-FS, used to provide the flavour-symmetric background estimate, the combined systematic uncertainty in the same background, the systematic uncertainty in \(\gamma +\) jets, or, in the case of SRZ, the \(WZ/ZZ\) generator uncertainty. The statistical component of the uncertainty from the flavour-symmetric estimate is largest for the edge analysis in SR-medium and SR-high in the highest \(m_{\ell\ell}\) regions. In the edge SRs the uncertainty in the \(WZ/ZZ\) background tends to be highest in the \(m_{\ell\ell}\) ranges that include the Z window. The uncertainty in the fake-lepton background is largest in SR-high, where fake leptons can compose a larger fraction of the background. Experimental uncertainties have a far lower impact on the systematic uncertainty of the total background (~2%).

### 9 Results

#### 9.1 Results in SRZ

For the on-shell Z search, the expected background and observed yields in the SR are shown in Table 9. A total of 60 events are observed in data with a predicted background of 53.5 ± 9.3 events. There are 35 events observed in the \(ee\) channel, and 25 events observed in the \(\mu\mu\) channel. The probability for the background to produce a fluctua-
The dilepton invariant-mass distribution for the $e\mu$ and $\mu\mu$ channels with the kinematic requirements of SRZ, but over the full $m_{\ell\ell}$ range, is shown in Fig. 6. Here the data are consistent with the expected background over the full $m_{\ell\ell}$ range. The dilepton invariant-mass, jet and $b$-tagged jet multiplicity, $E_{T}^{\text{miss}}$, $H_{T}^{\text{incl}}$ and $p_{T}^{\ell\ell}$ distributions in SRZ are shown in Fig. 7. The shapes of the background distributions in these figures are obtained from MC simulation, where the MC simulation is normalised according to the data-driven estimates in the SR. Here two representative examples of $\tilde{g} \rightarrow \gamma Z_2$ on-shell signal models, with $(m(\tilde{g}), m(\tilde{X}_2^0)) = (1095, 205)$ GeV and $(m(\tilde{g}), m(\tilde{X}_2^0)) = (1240, 960)$ GeV, are overlaid. To demonstrate the modelling of the $Z/\gamma^* + \text{jets}$ background in VR-S and SRZ, Fig. 8 shows the minimum $\Delta \phi(\text{jet}_{12}, p_{T}^{\text{miss}})$ distribution over the full range, where $\Delta \phi(\text{jet}_{12}, p_{T}^{\text{miss}}) > 0.4$ is required in VR-S and SRZ. Here the $Z/\gamma^* + \text{jets}$ distribution is modelled using the full data-driven prediction from $\gamma + \text{jets}$. Two of the events in the SR contain a third signal lepton.

9.2 Results in the edge SRs

The integrated yields in the edge signal regions are compared to the expected background in Table 10. To allow for the visualisation of a potential edge, the full $m_{\ell\ell}$ distributions in the three search regions are compared to the expected background in Fig. 9. In addition, the observed $m_{\ell\ell}$ distributions are compared to the predictions from MC simulation in Fig. 10, in which the $t\bar{t}$ background is scaled such that the total MC expected yield matches the data in the $e\mu$ CR. The $t\bar{t}$ normalisation factors are $\mu_{t\bar{t}} = 0.85 \pm 0.03, 0.75 \pm 0.04$, and $0.57 \pm 0.07$ in SR-low, SR-medium, and SR-high, respectively, where the uncertainty is the data statistical uncertainty. The data-driven flavour-symmetry prediction is used for the quantitative results of the analysis. This prediction does not rely on the $t\bar{t}$ normalisation scale factors discussed above.
The MC-based cross-check method is used to examine the $m_{ll}$ distribution in finer bins than can be achieved with the flavour-symmetry method, due to the limited statistical precision of the $e\mu$ CR.

As signal models may produce kinematic endpoints at any value of $m_{ll}$, any excess must be searched for across the $m_{ll}$ distribution. To do this a “sliding window” approach is used. The binning in the SRs, shown in Fig. 9, defines many possible dilepton mass windows. The 24 $m_{ll}$ ranges (9 for SR-low, 8 for SR-medium, and 7 for SR-high) are chosen because they are the most sensitive for at least one grid point in the signal model parameter space. Some of the ranges overlap. The results in these regions are summarised in Fig. 11, and the expected and observed yields in the combined $ee + \mu\mu$...
channel for all 24 $m_{\ell\ell}$ ranges are presented in Table 11. In SR-low and SR-medium, the data are consistent with the expected background across the full $m_{\ell\ell}$ range. In SR-high the data show a slight excess above the background at low $m_{\ell\ell}$. Of these 24 $m_{\ell\ell}$ ranges, the largest excess is observed in SR-high with $12 < m_{\ell\ell} < 101$ GeV. Here a total of 90 events are observed in data, compared to an expectation of $65 \pm 10$ events, corresponding to a local significance of $1.7\sigma$.

10 Interpretation

In this section, exclusion limits are shown for the SUSY models detailed in Sect. 3. The asymptotic $CL_S$ prescription [90,98], implemented in the HistFitter program [97], is used to determine cross-section upper limits at 95% confidence level (CL) for the on-$Z$ search. For the edge search, pseudo-experiments are used to evaluate the cross-section upper limits. A Gaussian model for nuisance parameters is used for all signal and background uncertainties. Exceptions are the statistical uncertainties of the flavour-symmetry method, $\gamma$ + jets method and MC-based backgrounds, all of which are treated as Poissonian nuisance parameters. The different experimental uncertainties are treated as correlated between signal and background events. The theoretical uncertainty of the signal cross section is not accounted for in the limit-setting procedure. Instead, following the initial limit determination, the impact of varying the signal cross section within its uncertainty is evaluated separately and indicated in the exclusion results. Limits are based on the combined $e\mu + \mu\mu$ results. Possible signal contamination in the CRs is neglected in the limit-setting procedure; the contamination is found to be negligible for signal points near the exclusion boundaries. Far from the exclusion boundary, although the signal contamination can be significant, the number of events appearing in the signal region is large enough that the points are still excluded, due to the relative branching fractions for the signal in the CR and SR. For example, for models with signal contamination of 50% in CR-FS the signal-to-background ratio in SRZ is $\sim 10$.

The results of the on-shell Z search are interpreted in a simplified model with gluino-pair production, where each gluino decays as $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_0^0$, $\tilde{\chi}_0^0 \rightarrow Z\tilde{\chi}_2^0$ and the $\tilde{\chi}_2^0$ mass is set to 1 GeV. The expected and observed exclusion contours for this $\tilde{g}-\tilde{\chi}_2^0$ on-shell grid are shown in the $m(\tilde{g})$--$m(\tilde{\chi}_2^0)$ plane in Fig. 12. The expected (observed) lower limit on the gluino mass is about 1.35 TeV (1.30 TeV) for a $\tilde{\chi}_2^0$ with a

---

**Table 10** Breakdown of the expected background and observed data yields for SR-low, SR-medium and SR-high, integrated over the $m_{\ell\ell}$ spectrum. The flavour-symmetric, $Z/\gamma + \text{jets}$ and fake-lepton components are all derived using data-driven estimates described in Sect. 7. All remaining backgrounds are taken from MC simulation. The quoted uncertainties include statistical and systematic contributions.

<table>
<thead>
<tr>
<th></th>
<th>SR-low</th>
<th>SR-medium</th>
<th>SR-high</th>
</tr>
</thead>
<tbody>
<tr>
<td>Observed events</td>
<td>1394</td>
<td>689</td>
<td>212</td>
</tr>
<tr>
<td>Total expected background events</td>
<td>1500 ± 100</td>
<td>700 ± 60</td>
<td>171 ± 18</td>
</tr>
<tr>
<td>Flavour-symmetric (t, W, W and Z → $\tau\tau$) events</td>
<td>1270 ± 70</td>
<td>584 ± 32</td>
<td>148 ± 14</td>
</tr>
<tr>
<td>$Z/\gamma + \text{jets}$ events</td>
<td>90 ± 50</td>
<td>50 ± 40</td>
<td>3.3 ± 3</td>
</tr>
<tr>
<td>WZ/ZZ events</td>
<td>63 ± 31</td>
<td>26 ± 11</td>
<td>7 ± 4</td>
</tr>
<tr>
<td>Rare top events</td>
<td>19 ± 5</td>
<td>11.3 ± 3.2</td>
<td>4.2 ± 1.4</td>
</tr>
<tr>
<td>Fake-lepton events</td>
<td>59 ± 34</td>
<td>32 ± 19</td>
<td>10 ± 8</td>
</tr>
</tbody>
</table>

---

**Fig. 8** The min. $\Delta\phi(jet_{12}, p_{T}^{\text{miss}})$ distribution in (left) VR-S and (right) SRZ, where the min. $\Delta\phi(jet_{12}, p_{T}^{\text{miss}}) > 0.4$ requirement has been lifted. The vertical dashed lines indicate the requirement in each region. The flavour-symmetric and $Z/\gamma + \text{jets}$ distributions are taken completely from the data-driven estimate. The rare top and data-driven fake-lepton backgrounds are grouped under “other” backgrounds. The hashed uncertainty bands include the statistical and systematic uncertainties in the background prediction.
Fig. 9 Expected and observed dilepton mass distributions, with the bin boundaries considered for the interpretation, in (top left) SR-low, (top-right) SR-medium, and (bottom) SR-high of the edge search. These bins, and sets of neighbouring bins, make up the mll windows used for the interpretation. The flavour-symmetric and $Z/\gamma^* + \text{jets}$ distributions are taken completely from the data-driven estimate. The rare top and data-driven fake-lepton backgrounds are grouped under “other” backgrounds. All statistical and systematic uncertainties are included in the hashed bands. The ratio of data to predicted background is shown in the bottom panels. In cases where the data point is not accommodated by the scale of this panel, a red arrow indicates the direction in which the point is out of range.

mass of 1.1 TeV in this model. The impact of the systematic uncertainties in the background and the experimental uncertainties in the signal, shown with a coloured band, is about 100 GeV on the gluino mass limit. The systematic uncertainty of the signal cross section, shown as dotted lines around the observed contour, has an impact of about 40 GeV. Figure 12 also shows the expected and observed exclusion limits for the $\tilde{q} - \tilde{\chi}_2^0$ on-shell model. This is a simplified model with squark-pair production, where each squark decays to a quark and a neutralino, with the neutralino subsequently decaying to a Z boson and an LSP with a mass of 1 GeV. In this model, exclusion is expected (observed) for squarks with masses below 1040 GeV (980 GeV) for a $\tilde{\chi}_2^0$ mass of 600 GeV.

Figure 13 shows the expected and observed exclusion contours for the $\tilde{g} - \tilde{\chi}_1^0$ on-shell model, in which the produced
Fig. 10 The dilepton mass distributions in the (top) SR-low (left) and CR-FS-low (right), (middle) SR-medium (left) and CR-FS-medium (right), and (bottom) SR-high (left) and CR-FS-high (right) regions of the edge search. The $t\bar{t}$ MC sample is normalised such that the total MC prediction matches data in the $e\mu$ channel for each region. The $m_{\ell\ell}$ shape and normalisation for the $Z/\gamma^* +$ jets background is taken from the $\gamma +$ jets method. The rare top and data-driven fake-lepton backgrounds are grouped under “other” backgrounds. Example signal benchmarks from the slepton and $Z(\ast)$ models are overlaid on the distributions. The first (second) number in parentheses is the gluino (LSP) mass. The overflow is included in the last bin.

gluinos follow the same decay chain as in the model above. In this case the mass difference $\Delta m = m(\tilde{\chi}^0_2) - m(\tilde{\chi}^0_1)$ is set to 100 GeV.

The results of the edge search are interpreted in two simplified models with gluino-pair production, in which each gluino decays as $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}^0_2$. For each point in the signal-
expected, the discovery number of events lower than significance is set to zero is truncated at 0.5 and the Gaussian significance (uncertainties, observed data, statistical and systematic backgrounds, with combined m ranges. The given for SR-low, SR-medium signal regions. The results are observed data yields in the edge expected background and Table 11 Breakdown of the expected background and observed data yields in the edge signal regions. The results are given for SR-low, SR-medium and SR-high in all 24 m ranges. The m range in units of GeV is indicated in the leftmost column of the table. Left to right: the total expected background, with combined statistical and systematic uncertainties, observed data, 95% CL upper limits on the visible cross section (⟨σ⟩obs) and on the number of signal events (Sobs). The sixth column (Sexp) shows the expected 95% CL upper limit on the number of signal events, given the expected number (and ±1σ excursions) of background events. The last two columns indicate the discovery p value (p(s = 0)) [97], and the Gaussian significance (Z(s = 0)). For an observed number of events lower than expected, the discovery p value is truncated at 0.5 and the significance is set to zero

<table>
<thead>
<tr>
<th>Signal region</th>
<th>Total Bkg.</th>
<th>Data</th>
<th>⟨σ⟩obs [fb]</th>
<th>Sobs [95]</th>
<th>Sexp [95]</th>
<th>p(s = 0)</th>
<th>Z(s = 0)</th>
</tr>
</thead>
<tbody>
<tr>
<td>SR-low</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>12–61</td>
<td>187 ± 18</td>
<td>175</td>
<td>2.68</td>
<td>39.4</td>
<td>48±23</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>12–81</td>
<td>330 ± 24</td>
<td>320</td>
<td>3.88</td>
<td>57.1</td>
<td>64±20</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>12–101</td>
<td>617 ± 63</td>
<td>534</td>
<td>4.64</td>
<td>68.2</td>
<td>96±26</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>81–101</td>
<td>287 ± 50</td>
<td>214</td>
<td>2.73</td>
<td>40.2</td>
<td>62±20</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>101–201</td>
<td>529 ± 34</td>
<td>540</td>
<td>6.80</td>
<td>99.9</td>
<td>91±29</td>
<td>0.40</td>
<td>0.26</td>
</tr>
<tr>
<td>101–301</td>
<td>741 ± 48</td>
<td>732</td>
<td>7.28</td>
<td>107</td>
<td>113±33</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>201–401</td>
<td>295 ± 30</td>
<td>262</td>
<td>3.43</td>
<td>50.5</td>
<td>70±27</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>301–501</td>
<td>113 ± 17</td>
<td>99</td>
<td>2.37</td>
<td>34.8</td>
<td>46±16</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>&gt;501</td>
<td>29 ± 10</td>
<td>29</td>
<td>1.88</td>
<td>27.7</td>
<td>27±16</td>
<td>0.50</td>
<td>0.01</td>
</tr>
<tr>
<td>SR-medium</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>12–61</td>
<td>119 ± 15</td>
<td>109</td>
<td>2.38</td>
<td>35.1</td>
<td>43±29</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>12–81</td>
<td>190 ± 18</td>
<td>191</td>
<td>3.57</td>
<td>52.5</td>
<td>51±15</td>
<td>0.48</td>
<td>0.06</td>
</tr>
<tr>
<td>12–101</td>
<td>315 ± 43</td>
<td>299</td>
<td>5.12</td>
<td>75.3</td>
<td>81±29</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>81–101</td>
<td>125 ± 35</td>
<td>108</td>
<td>3.18</td>
<td>46.7</td>
<td>51±12</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>101–201</td>
<td>235 ± 20</td>
<td>240</td>
<td>4.26</td>
<td>62.6</td>
<td>58±19</td>
<td>0.42</td>
<td>0.19</td>
</tr>
<tr>
<td>101–301</td>
<td>332 ± 25</td>
<td>336</td>
<td>4.92</td>
<td>72.3</td>
<td>69±22</td>
<td>0.45</td>
<td>0.14</td>
</tr>
<tr>
<td>201–401</td>
<td>126 ± 13</td>
<td>128</td>
<td>3.27</td>
<td>48.0</td>
<td>46±16</td>
<td>0.46</td>
<td>0.11</td>
</tr>
<tr>
<td>&gt;401</td>
<td>28 ± 8</td>
<td>22</td>
<td>1.09</td>
<td>16.1</td>
<td>21±9</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>SR-high</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>12–61</td>
<td>23 ± 5</td>
<td>27</td>
<td>1.84</td>
<td>27.0</td>
<td>20±8</td>
<td>0.27</td>
<td>0.62</td>
</tr>
<tr>
<td>12–81</td>
<td>39 ± 7</td>
<td>53</td>
<td>3.32</td>
<td>48.9</td>
<td>26±18</td>
<td>0.08</td>
<td>1.40</td>
</tr>
<tr>
<td>12–101</td>
<td>65 ± 10</td>
<td>90</td>
<td>4.00</td>
<td>58.8</td>
<td>31±17</td>
<td>0.04</td>
<td>1.73</td>
</tr>
<tr>
<td>81–101</td>
<td>26 ± 6</td>
<td>37</td>
<td>2.17</td>
<td>31.9</td>
<td>20±5</td>
<td>0.12</td>
<td>1.18</td>
</tr>
<tr>
<td>101–201</td>
<td>59 ± 9</td>
<td>75</td>
<td>3.68</td>
<td>54.1</td>
<td>31±11</td>
<td>0.10</td>
<td>1.27</td>
</tr>
<tr>
<td>201–401</td>
<td>39 ± 7</td>
<td>33</td>
<td>1.82</td>
<td>26.7</td>
<td>28±14</td>
<td>0.50</td>
<td>0.00</td>
</tr>
<tr>
<td>&gt;401</td>
<td>10 ± 5</td>
<td>14</td>
<td>2.04</td>
<td>30.0</td>
<td>21±9</td>
<td>0.27</td>
<td>0.62</td>
</tr>
</tbody>
</table>
model parameter space, limits on the signal strength are calculated using the $m_{\ell\ell}$ window with the best expected sensitivity. Details of the windows are described in Sect. 9.

The excluded regions in the $m(\tilde{g})$--$m(\tilde{\chi}^0_1)$ plane are presented in Fig. 14 for the slepton model. In this model, pair-produced gluinos each decay as $\tilde{g} \rightarrow q\tilde{q}_2^0, \tilde{\chi}_2^0 \rightarrow \ell^+\ell^-, \tilde{\tau} \rightarrow \ell^+\tilde{\chi}_1^0$. Here, the results exclude gluinos with masses as large as 1.7 TeV, with an expected limit of 1.75 TeV for small $m(\tilde{\chi}_1^0)$. The results probe kinematic endpoints as small as $m_{\text{min}}^{\text{max}} = m(\tilde{\chi}_1^0) - m(\tilde{\chi}^0_1) = 1/2(m(\tilde{g}) - m(\tilde{\chi}^0_1)) = 50$ GeV.

The $Z^{(*)}$ exclusion limits from the results in the edge SRs are compared with the same limits derived using the results in SRZ in Fig. 15. In this model, pair-produced gluinos each decay as $\tilde{g} \rightarrow q\tilde{q}_2^0, \tilde{\chi}_2^0 \rightarrow Z^{(*)}\tilde{\chi}_1^0$, and the mass splitting between the $\tilde{\chi}^0_2$ and the $\tilde{\chi}^0_1$ determines whether the $Z$ boson is produced on-shell. Here the edge limits extend into the more compressed region, whereas the expected SRZ exclusion probes higher $\tilde{\chi}^0_1$ masses in the on-shell regime. At high gluino masses, the edge SRs provide stronger limits. For the $Z^{(*)}$ model, the expected and observed gluino mass limits are 1.4 TeV and 1.34 TeV (1.35 and 1.3 TeV for the on-Z signal region), respectively, for $\tilde{\chi}^0_1$ masses below 400 GeV. The sensitivity in the $Z^{(*)}$ model is smaller than that of the slepton model because the leptonic branching fraction of the $Z$ boson suppresses the signal production rate.
11 Conclusion

This paper presents two searches for new phenomena in final states containing a same-flavour opposite-sign lepton (electron or muon) pair, jets, and large missing transverse momentum using 14.7 fb\(^{-1}\) of ATLAS data collected during 2015 and 2016 at the LHC at \(\sqrt{s} = 13\) TeV. The first search (on-shell Z search) targets lepton pairs consistent with Z boson decay, while the second search (edge search) targets a kinematic endpoint feature in the dilepton mass distribution. For the edge search, a set of 24 mass ranges are considered, with different requirements on \(E_T^{\text{miss}}\) and \(H_T\), and different kinematic endpoint values in the dilepton invariant-mass distribution. The data in both searches are found to be consistent with the Standard Model prediction. The results are interpreted in simplified models of gluino-pair production and squark-pair production, and exclude gluinos (squarks) with masses as large as 1.7 TeV (980 GeV).

Acknowledgements We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently. We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWF and FWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DNRF and DNSRC, Denmark; IN2P3-CNRS, CEA-DSM/IRFU, France; GNSF, Georgia; BMBF, HGF, and MPG, Germany; GSRT, Greece; RGC, Hong Kong SAR, China; ISF, I-CORE and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRST, Morocco; FOM and NWO, Netherlands; RCN, Norway; MNiSW and NCN, Poland; FCT, Portugal; MNE/IFA, Romania; MES of Russia and NRC KI, Russian Federation; JINR; MESTD, Serbia; MSSR, Slovakia; ARRS and MIZŠ, Slovenia; DST/NRF, South Africa; MINECO, Spain; SRC and Wallenberg Foundation, Sweden; SERI, SNSF and Cantons of Bern and Geneva, Switzerland; MOST, Taiwan; TAEK, Turkey; STFC, United Kingdom; DOE and NSF, United States of America; and, in addition, individual groups and members have received support from BCKDF, the Canada Council, CONACyT, CRC, Compute Canada, FQRNT, and the Ontario Innovation Trust, Canada; EPLANET, ERC, ERDF, FP7, Horizon 2020 and Marie Sklodowska-Curie Actions, European Union; Investissements d’Avenir Label and Idex, ANR, Région Auvergne and Fondation Partager le Savoir, France; DFG and AvH Foundation, Germany; Herakleitos, Thales and Aristeia programmes co-financed by EU-ESF and the Greek NSRF; BSF, GIF and Minerva, Israel; BRF, Norway; CERCA Programme Generalitat de Catalunya, Generalitat Valenciana, Spain; the Royal Society and Leverhulme Trust, United Kingdom. The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN, the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA), the Tier-2 facilities worldwide and large non-WLCG resource providers. Major contributors of computing resources are listed in Ref. [99].

Open Access This article is distributed under the terms of the Creative Commons Attribution 4.0 International License (http://creativecommons.org/licenses/by/4.0/), which permits unrestricted use, distribution, and reproduction in any medium, provided you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made. Funded by SCOAP3.

References

1. Y.A. Gol'fand, E.P. Likhtman, Extension of the algebra of Poincare group generators and violation of p invariance. JETP Lett. 13, 323–326 (1971)

ATLAS Collaboration

12 Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
13 Institut de Física d’Altes Energies (IFAE), The Barcelona Institute of Science and Technology, Barcelona, Spain
14 Institute of Physics, University of Belgrade, Belgrade, Serbia
15 Department for Physics and Technology, University of Bergen, Bergen, Norway
16 Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, USA
17 Department of Physics, Humboldt University, Berlin, Germany
18 Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
19 School of Physics and Astronomy, University of Birmingham, Birmingham, UK
20 (a) Department of Physics, Bogazici University, Istanbul, Turkey; (b) Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey; (c) Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul, Turkey; (d) Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul, Turkey
21 Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia
22 (a) INFN Sezione di Bologna, Bologna, Italy; (b) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
23 Physikalisches Institut, University of Bonn, Bonn, Germany
24 Department of Physics, Boston University, Boston, MA, USA
25 Department of Physics, Brandeis University, Waltham, MA, USA
26 (a) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil; (b) Electrical Circuits Department, Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil; (c) Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil; (d) Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil
27 Physics Department, Brookhaven National Laboratory, Upton, NY, USA
28 (a) Transilvania University of Brasov, Brasov, Romania; (b) Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania; (c) Physics Department, National Institute for Research and Development of Isotopic and Molecular Technologies, Cluj Napoca, Romania; (d) West University in Timisoara, Timisoara, Romania
29 Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
30 Cavendish Laboratory, University of Cambridge, Cambridge, UK
31 Department of Physics, Carleton University, Ottawa, ON, Canada
32 CERN, Geneva, Switzerland
33 Enrico Fermi Institute, University of Chicago, Chicago, IL, USA
34 (a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile; (b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
35 (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China; (b) Department of Physics, Nanjing University, Jiangsu, China; (c) Physics Department, Tsinghua University, Beijing 100084, China
36 (a) Department of Modern Physics, University of Science and Technology of China, Anhui, China; (b) School of Physics, Shandong University, Shandong, China; (c) Department of Physics and Astronomy, Shanghai Key Laboratory for Particle Physics and Cosmology, Shanghai Jiao Tong University (also affiliated with PKU-CHEP), Shanghai, China
37 Laboratoire de Physique Corpusculaire, Université Clermont Auvergne univernité Blaise Pascal, CNRS/IN2P3, Clermont-Ferrand, France
38 Nevis Laboratory, Columbia University, Irvington, NY, USA
39 Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
40 (a) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Frascati, Italy; (b) Dipartimento di Fisica, Università della Calabria, Rende, Italy
41 (a) Faculty of Physics and Applied Computer Science, AGH University of Science and Technology, Krakow, Poland; (b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Kraków, Poland
42 Institute of Nuclear Physics, Polish Academy of Sciences, Kraków, Poland
43 Physics Department, Southern Methodist University, Dallas, TX, USA
44 Physics Department, University of Texas at Dallas, Richardson, TX, USA
45 DESY, Hamburg and Zeuthen, Germany
46 Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
47 Institut für Kern-und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
48 Department of Physics, Duke University, Durham, NC, USA
49 SUPA-School of Physics and Astronomy, University of Edinburgh, Edinburgh, UK
50 INFN Laboratori Nazionali di Frascati, Frascati, Italy
51 Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
52 Département de Physique Nucléaire et Corpusculaire, Université de Genève, Geneva, Switzerland
53 (a) INFN Sezione di Genova, Genoa, Italy; (b) Dipartimento di Fisica, Università di Genova, Genoa, Italy
54 (a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia; (b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
55 II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
56 SUPA-School of Physics and Astronomy, University of Glasgow, Glasgow, UK
57 II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
58 Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France
59 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, USA
60 (a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany; (b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany; (c) ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
61 Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
62 (a) Department of Physics, The Chinese University of Hong Kong, Shatin, NT, Hong Kong; (b) Department of Physics, The University of Hong Kong, Hong Kong, China; (c) Department of Physics and Institute for Advanced Study, The Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
63 Department of Physics, National Tsing Hua University, Hsinchu, Taiwan, Taiwan
64 Department of Physics, Indiana University, Bloomington, IN, USA
65 Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
66 University of Iowa, Iowa City, IA, USA
67 Department of Physics and Astronomy, Iowa State University, Ames, IA, USA
68 Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
69 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
70 Graduate School of Science, Kobe University, Kobe, Japan
71 Faculty of Science, Kyoto University, Kyoto, Japan
72 Kyoto University of Education, Kyoto, Japan
73 Department of Physics, Kyushu University, Fukuoka, Japan
74 Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
75 Physics Department, Lancaster University, Lancaster, UK
76 (a) INFN Sezione di Lecce, Lecce, Italy; (b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
77 Oliver Lodge Laboratory, University of Liverpool, Liverpool, UK
78 Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia
79 School of Physics and Astronomy, Queen Mary University of London, London, UK
80 Department of Physics, Royal Holloway University of London, Surrey, UK
81 Department of Physics and Astronomy, University College London, London, UK
82 Louisiana Tech University, Ruston, LA, USA
83 Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
84 Fysiska institutionen, Lunds universitet, Lund, Sweden
85 Departamento de Física Teorica C-15, Universidad Autonoma de Madrid, Madrid, Spain
86 Institut für Physik, Universität Mainz, Mainz, Germany
87 School of Physics and Astronomy, University of Manchester, Manchester, UK
88 CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
89 Department of Physics, University of Massachusetts, Amherst, MA, USA
90 Department of Physics, McGill University, Montreal, QC, Canada
91 School of Physics, University of Melbourne, Melbourne, VIC, Australia
92 Department of Physics, The University of Michigan, Ann Arbor, MI, USA
93 Department of Physics and Astronomy, Michigan State University, East Lansing, MI, USA
94 (a) INFN Sezione di Milano, Milan, Italy; (b) Dipartimento di Fisica, Università di Milano, Milan, Italy
95 B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus
96 Research Institute for Nuclear Problems of Byelorussian State University, Minsk, Republic of Belarus
97 Group of Particle Physics, University of Montreal, Montreal, QC, Canada
98 P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia
99 Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
100 National Research Nuclear University MEPhI, Moscow, Russia
101 D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
102 Fakultät für Physik, Ludwig-Maximilians-Universität München, Munich, Germany
103 Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), Munich, Germany
104 Nagasaki Institute of Applied Science, Nagasaki, Japan
105 Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
106 (a) INFN Sezione di Napoli, Naples, Italy; (b) Dipartimento di Fisica, Università di Napoli, Naples, Italy
107 Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, USA
108 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, The Netherlands
109 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, The Netherlands
110 Department of Physics, Northern Illinois University, DeKalb, IL, USA
111 Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
112 Department of Physics, New York University, New York, NY, USA
113 Ohio State University, Columbus, OH, USA
114 Faculty of Science, Okayama University, Okayama, Japan
115 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, USA
116 Department of Physics, Oklahoma State University, Stillwater, OK, USA
117 Palacký University, RCPTM, Olomouc, Czech Republic
118 Center for High Energy Physics, University of Oregon, Eugene, OR, USA
119 LAL, Univ. Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France
120 Graduate School of Science, Osaka University, Osaka, Japan
121 Department of Physics, University of Oslo, Oslo, Norway
122 Department of Physics, Oxford University, Oxford, UK
123 (a) INFN Sezione di Pavia, Pavia, Italy; (b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
124 Department of Physics, University of Pennsylvania, Philadelphia, PA, USA
125 National Research Centre “Kurchatov Institute” B.P.Konstantinov Petersburg Nuclear Physics Institute, St. Petersburg, Russia
126 (a) INFN Sezione di Pisa, Pisa, Italy; (b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
127 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, USA
128 (a) Laboratório de Instrumentação e Física Experimental de Partículas-LIP, Lisbon, Portugal; (b) Faculdade de Ciências, Universidade de Lisboa, Lisbon, Portugal; (c) Department of Physics, University of Coimbra, Coimbra, Portugal; (d) Centro de Física Nuclear da Universidade de Lisboa, Lisbon, Portugal; (e) Departamento de Física, Universidade do Minho, Braga, Portugal; (f) Departamento de Física Teorica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain
129 Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
130 Czech Technical University in Prague, Prague, Czech Republic
131 Faculty of Mathematics and Physics, Charles University in Prague, Prague, Czech Republic
132 State Research Center Institute for High Energy Physics (Protvino), NRC KI, Russia
133 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, UK
134 (a) INFN Sezione di Roma, Rome, Italy; (b) Dipartimento di Fisica, Sapienza Università di Roma, Rome, Italy
135 (a) INFN Sezione di Roma Tor Vergata, Rome, Italy; (b) Dipartimento di Fisica, Università di Roma Tor Vergata, Rome, Italy
136 (a) INFN Sezione di Roma Tre, Rome, Italy; (b) Dipartimento di Matematica e Fisica, Università Roma Tre, Rome, Italy
137 (a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies-Université Hassan II, Casablanca, Morocco; (b) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPEHA-Marrakech, Marrakech, Morocco; (c) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda, Morocco; (d) Faculté des Sciences, Université Mohammed V, Rabat, Morocco

© Springer
138 DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France
139 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, USA
140 Department of Physics, University of Washington, Seattle, WA, USA
141 Department of Physics and Astronomy, University of Sheffield, Sheffield, UK
142 Department of Physics, Shinshu University, Nagano, Japan
143 Fachbereich Physik, Universität Siegen, Siegen, Germany
144 Department of Physics, Simon Fraser University, Burnaby, BC, Canada
145 SLAC National Accelerator Laboratory, Stanford, CA, USA
146 (a)Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovak Republic; (b)Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
147 (a)Department of Physics, University of Cape Town, Cape Town, South Africa; (b)Department of Physics, University of Johannesburg, Johannesburg, South Africa; (c)School of Physics, University of the Witwatersrand, Johannesburg, South Africa
148 (a)Department of Physics, Stockholm University, Stockholm, Sweden; (b)The Oskar Klein Centre, Stockholm, Sweden
149 Physics Department, Royal Institute of Technology, Stockholm, Sweden
150 Departments of Physics and Astronomy and Chemistry, Stony Brook University, Stony Brook, NY, USA
151 Department of Physics and Astronomy, University of Sussex, Brighton, UK
152 School of Physics, University of Sydney, Sydney, NSW, Australia
153 Institute of Physics, Academia Sinica, Taipei, Taiwan
154 Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel
155 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
156 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
157 International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan
158 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
159 Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
160 Tomsk State University, Tomsk, Russia
161 Department of Physics, University of Toronto, Toronto, ON, Canada
162 (a)INFN-TIFPA, Povo, Italy; (b)University of Trento, Trento, Italy
163 (a)TRIUMF, Vancouver, BC, Canada; (b)Department of Physics and Astronomy, York University, Toronto, ON, Canada
164 Faculty of Pure and Applied Sciences, and Center for Integrated Research in Fundamental Science and Engineering, University of Tsukuba, Tsukuba, Japan
165 Department of Physics and Astronomy, Tufts University, Medford, MA, USA
166 Department of Physics and Astronomy, University of California Irvine, Irvine, CA, USA
167 (a)INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy; (b)ICTP, Trieste, Italy; (c)Dipartimento di Chimica Fisica e Ambiente, Università di Udine, Udine, Italy
168 Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
169 Department of Physics, University of Illinois, Urbana, IL, USA
170 Instituto de Fisica Corpuscular (IFIC) and Departamento de Fisica Atomica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNMT), University of Valencia and CSIC, Valencia, Spain
171 Department of Physics, University of British Columbia, Vancouver, BC, Canada
172 Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada
173 Department of Physics, University of Warwick, Coventry, UK
174 Waseda University, Tokyo, Japan
175 Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel
176 Department of Physics, University of Wisconsin, Madison, WI, USA
177 Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany
178 Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany
179 Department of Physics, Yale University, New Haven, CT, USA
180 Yerevan Physics Institute, Yerevan, Armenia
181 Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France