Measurement of the Inelastic Proton-Proton Cross Section at $\sqrt{s} = 13$ TeV with the ATLAS Detector at the LHC

M. Aaboud et al.*
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This Letter presents a measurement of the inelastic proton-proton cross section using 60 $\mu$b$^{-1}$ of $pp$ collisions at a center-of-mass energy $\sqrt{s}$ of 13 TeV with the ATLAS detector at the LHC. Inelastic interactions are selected using rings of plastic scintillators in the forward region ($2.07 < |\eta| < 3.86$) of the detector. A cross section of $68.1 \pm 1.4$ mb is measured in the fiducial region $\xi = M_X^2/\sqrt{s} > 10^{-6}$, where $M_X$ is the larger invariant mass of the two hadronic systems separated by the largest rapidity gap in the event. In this $\xi$ range the scintillators are highly efficient. For diffractive events this corresponds to cases where at least one proton dissociates to a system with $M_X > 13$ GeV. The measured cross section is compared with a range of theoretical predictions. When extrapolated to the full phase space, a cross section of $78.1 \pm 2.9$ mb is measured, consistent with the inelastic cross section increasing with center-of-mass energy.

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The rise of the total proton-proton ($pp$) cross section with center-of-mass energy $\sqrt{s}$, predicted by Heisenberg [1] and observed at the CERN Intersecting Storage Rings [2], probes the nonperturbative regime of quantum chromodynamics (QCD). Arguments based on unitarity, analyticity, and factorization imply an upper bound on the chromodynamics (QCD). Arguments based on unitarity, analyticity, and factorization imply an upper bound on the chromodynamics (QCD). Arguments based on unitarity, analyticity, and factorization imply an upper bound on the chromodynamics (QCD). Arguments based on unitarity, analyticity, and factorization imply an upper bound on the chromodynamics (QCD).

Many experiments have measured $\sigma_{\text{inel}}$ and found an increase with $\sqrt{s}$ [6]. The TOTEM and ATLAS collaborations determined $\sigma_{\text{inel}}$ at $\sqrt{s} = 7$ and 8 TeV using the optical theorem and a measurement of the elastic cross section with Roman pot detectors [7–11]. Using a variety of alternative techniques, the ATLAS, CMS, ALICE, and LHCb experiments have made measurements of $\sigma_{\text{inel}}$ at $\sqrt{s} = 7$ TeV [12–15] and $\sqrt{s} = 2.76$ TeV (ALICE) [14]. The Pierre Auger Collaboration measured the inelastic $p$-air cross section at $\sqrt{s} = 57$ TeV and extracted $\sigma_{\text{inel}}$ using the Glauber model [16].

This Letter presents a measurement of the inelastic cross section $\sigma_{\text{inel}}$ using $pp$ collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector at the Large Hadron Collider (LHC). It is performed using two sets of scintillation counters in a data set corresponding to an integrated luminosity of $60.1 \pm 1.1$ $\mu$b$^{-1}$ collected in June 2015. In inelastic interactions, one or both protons dissociate as a result of colored (nondiffractive) or colorless (diffractive) exchange. The counters are insensitive to elastic $pp$ scattering and diffractive dissociation processes in which neither proton dissociates into a system, $X$, of mass $M_X > 13$ GeV, or equivalently, $\xi = M_X^2/\sqrt{s} > 10^{-6}$. The cross-section measurement is reported in this fiducial region, $\xi > 10^{-6}$, and after extrapolation to the total inelastic cross section using models of inelastic interactions.

The ATLAS detector is a cylindrical particle detector composed of several subdetector layers [17]. The inner tracking detector (ID) is immersed in a 2 T magnetic field provided by a superconducting solenoid. Around the tracker is a system of electromagnetic and hadronic calorimeters, which use liquid argon and lead, copper, or tungsten absorber for the electromagnetic and forward ($|\eta| > 1.7$) [18] hadronic components of the detector, and scintillator-tile active material and steel absorber for the central ($|\eta| < 1.7$) hadronic component.

At $z = \pm 3.6$ m, thin plastic scintillation counters, the minimum-bias trigger scintillators (MBTS), are installed on the front face of each endcap calorimeter. These detectors cover the region $2.07 < |\eta| < 3.86$. They are similar to those described in Ref. [17] but were rebuilt during 2014, when the coverage was slightly extended from $2.08 < |\eta| < 3.75$ after the $\sqrt{s} = 7$ TeV run. The MBTS are divided into inner (4 counters in $149 < r < 445$ mm) and outer (8 counters in $444.5 < r < 895$ mm) octagonal rings.

The ATLAS experiment uses a multistage trigger to select events at about 1 kHz for offline analysis. Three trigger configurations were used to collect data for this analysis. The primary triggers use the MBTS detector and constant-fraction discriminators to select events when two proton bunches collide in the detector. To facilitate background studies, data were also collected with the same selection when no proton bunch ("empty") or a single proton bunch from only one of the two beams ("single
beams”) was passing through the center of ATLAS. All of these triggers require at least one MBTS hit above threshold. Two additional triggers were used to collect data to determine the MBTS trigger efficiency, requiring either hits in a forward (5.6 < |η| < 5.9) Cherenkov detector (LUCID) or a far forward (|η| > 8.4) tungsten-scintillator calorimeter detector (LHCf [19]) located at z = ±17 m and ±140 m, respectively. The LHCf detector is an independent detector, but for the runs considered in this analysis, its trigger signals were incorporated into the ATLAS readout.

Monte Carlo (MC) simulation samples were produced to correct the fiducial measurement and to compare to the data. The detector response is modeled using a simulation based on Geant4 [20–22]. The data and MC simulated events are passed through the same reconstruction and analysis software.

The primary MC samples are based on the Pythia8 generator [23,24] either with the A2 [25] set of tuned underlying-event parameters and the MSTW 2008 LO PDF set [26] or with the Monash [27] set of tuned parameters and the NNPDF 2.3 LO PDF set [28]. The samples are divided into four components: single-dissociation (SD, pp → pX), double-dissociation (DD, pp → XY), central-dissociation (CD, pp → pXp), all involving colorless exchange, and nondiffractive dissociation (ND) wherein color flow is present between the two colliding protons. For all dissociation event types, the Monash tune is used.

Pythia8 uses a pomeron-based diffraction model [29] to describe colorless exchange with a default pomeron flux model by Schuler and Sjöstrand (SS) [30,31]. Alternative MC samples are generated with the pomeron flux model of Donnachie and Landshoff (DL) [32] and with the minimum-bias Rockefeller (MBR) model [33]. In the DL model, the pomeron Regge trajectory is given by α(t) = 1 + e + α′t, where e and α′ are free parameters. In most samples used for this analysis, the value of α′ is 0.25, the Pythia8 default. The e parameter is varied from 0.06 to 0.10 (the Pythia8 default is 0.085). An additional sample produced with α′ = 0.35 is found to be statistically consistent with the α′ = 0.25 default samples in each aspect of this analysis. The ranges of e and α′ considered are motivated by previous total, inelastic, elastic, and diffractive cross-section measurements, including measurements of low-mass diffraction by the ATLAS and CMS collaborations [34,35]. For the DL and SS models the CD component is neglected. The MBR model is tuned to data as described in Ref. [33] and includes a small CD component.

The Epos LHC and QGSJet-II event generators are also used to simulate pp collisions. Epos LHC [36] uses a “cut pomeron” model for diffraction and differs significantly from Pythia8 in its modeling of hadronization and the underlying event. QGSJet-II [37,38] uses Reggeon field theory to describe pomeron-pomeron interactions. Both Epos LHC and QGSJet-II have been developed primarily to model cosmic-ray showering in the atmosphere.

The fiducial region of the measurement is determined using MC simulation. In each generated event, the largest rapidity gap between any two final-state hadrons is used to define the boundary between two collections of hadrons. These definitions define the dissociation systems in an event-generator-independent manner. The invariant mass of each collection is calculated, and the larger of the two masses, denoted M_X, is used to define ξ = M_X^2/s. The variable ξ is constrained to be >6 × 10^{-6} by the elastic limit of m_p^2/s where m_p is the proton mass. This measurement is restricted to ξ > 10^{-6}, the region in which the event selection efficiency exceeds 50%.

Two samples of data events passing the MBTS trigger requirements are selected: an inclusive sample and a single-sided sample. The inclusive selection requires at least two MBTS counters with a charge above 0.15 pC (n_{MBTS} ≥ 2). This threshold is chosen to be well above the electronic noise level of the counters. Requiring two hits rather than one substantially reduces background due to collision-induced radiation and activation. To constrain the diffractive component of the cross section and reduce the uncertainty in extrapolation to σ_{inel}, an additional single-sided selection is defined, requiring hits in at least two counters on one side of the detector and no hits on the other. In the data, 4,159,074 events pass the inclusive selection and 442,192 events pass the single-sided selection.

The fiducial cross section is determined by

\[ \sigma_{\text{fid}}^{el}(\xi > 10^{-6}) = \frac{N - N_{\text{BG}}}{\epsilon_{\text{trig}} \times \mathcal{L}} \times \frac{1 - f_{\xi < 10^{-6}}}{\epsilon_{\text{sel}}} , \]

where N is the number of observed events passing the inclusive selection, N_{BG} is the number of background events, ε_{trig} and ε_{sel} are factors accounting for the trigger and event selection efficiencies, 1 - f_{\xi < 10^{-6}} accounts for the migration of events with \xi < 10^{-6} into the fiducial region, and \mathcal{L} is the integrated luminosity of the sample.

Sources of background include interactions between the beam and residual gas in the beam pipe; interactions between the beam and collimators upstream of the detector, which can send charged particles through the detector parallel to the beam; collision-induced radiation; and activation backgrounds. Backgrounds from cosmic rays and instrumental noise are negligible. The mean number of pp collisions in the same LHC bunch crossing was 2.3 × 10^{-3} for the recorded data set. Thus, the contribution from multiple collisions is also negligible. The beam-related background components are extracted from single-beam events and dominate the total background. They are normalized by scaling the number of selected single-beam events by a factor of 37/4 × 2, accounting for the 37 colliding pairs of bunches and 4 bunches producing the single-beam data in this run. The factor of 2 accounts for the presence of two colliding bunches. The number of protons per bunch producing these single-beam events
agrees with that in the colliding bunches to within 10%. The radiation and activation-induced backgrounds are implicitly part of this background estimate. Double-counting of these components is removed using estimates from empty events. The total background contributions to the inclusive and single-sided data samples are determined to be 1.2% and 5.8%, respectively. The classification of single-sided events as double-sided due to noise or other backgrounds is estimated to be below 0.1%. A systematic uncertainty of 50% is assigned to the background based on studies of the background composition and the relative contributions of the background components. This uncertainty is treated as fully correlated between the single-sided and inclusive selections.

The trigger efficiency for events passing the inclusive selection, $\epsilon_{\text{trig}}$, is measured with respect to events selected with the LUCID detector after subtracting the background. A trigger efficiency of 99.7% (97.4%) is measured for the inclusive (single-sided) event sample. In both cases the statistical uncertainty is below 0.1%. The efficiency is also measured with events selected by the LHCf detector and agrees within ±0.3% with the LUCID determination. This difference is taken as a systematic uncertainty.

The ratio of the number of events passing the single-sided event selection to the number passing the inclusive selection ($R_{SS}$) is used to adjust, for each model, the fractional contribution of the single- and double-diffractive dissociative cross section ($\sigma_{SD} + \sigma_{DD}$) to the inelastic cross section, $f_D = (\sigma_{SD} + \sigma_{DD})/\sigma_{\text{inel}}$ [12]. The measured value is $R_{SS} = 10.4\%$ with a total uncertainty of ±0.4%. The dominant systematic uncertainty arises from the background subtraction in the single-sided sample. For each MC model, $f_D$ is varied until it matches the observed $R_{SS}$ value in data. The data uncertainty is used to set the error in the constrained $f_D$ for each model. An additional uncertainty in the ratio of single- to double-diffractive events is determined by taking the diffractive events to be entirely SD or to be evenly divided between SD and DD.

Using this method, the fitted $f_D$ in the PYTHIA8 samples is between 25% and 31%, depending on the model (the default value is 28%). For the QGSJET-II (EPOS LHC) model the fitted $f_D$ is 35% (37%), differing significantly from the default value of 21% (28%). The observed $R_{SS}$ and the MC predictions of its dependence on $f_D$ are shown in Fig. 1. The fitted $f_D$ is used when determining the acceptance corrections $\epsilon_{\text{sel}}$ and $f_{<10^{-6}}$ for each model.

In Fig. 2 the $n_{\text{MBTS}}$ distributions in data are compared to the ones from MC simulated samples utilizing the fitted $f_D$ values for both the inclusive and single-sided selections. The estimated background is subtracted from the measured distribution, and the trigger efficiency measured in data is applied to the simulation. The data distributions and MC simulation are peaked at high multiplicity values. In the single-sided case, $n_{\text{MBTS}} = 12$ corresponds to hits in all counters on one side of the detector. The data agree best with the DL models, particularly in the low-$n_{\text{MBTS}}$ range. The MBR-based distribution provides a slightly worse description of the data. The PYTHIA8 sample using the SS model does not describe data well in the low-multiplicity region. EPOS LHC and QGSJET-II also do not describe the data well, particularly in the single-sided hit multiplicity distribution. Therefore, the PYTHIA8 DL model with $\epsilon = 0.085$ is chosen as the nominal MC model for the $\epsilon_{\text{sel}}$ and $f_{<10^{-6}}$ corrections, and only the DL and MBR models are considered for systematic uncertainties related to the MC corrections.

The event selection efficiency, $\epsilon_{\text{sel}}$, depends upon the MBTS counter sensitivity. This sensitivity is tested using isolated charged particles, reconstructed as ID tracks in the region $2.07 < |y| < 2.5$ where the coverages of the MBTS and ID overlap. Over the full coverage of the MBTS counters, the calorimeter is used to measure the counter efficiency with respect to particles that deposit sufficient energy in the calorimeter to seed a topological energy cluster [39]. Differences between the efficiencies in data and MC simulation are accounted for by adjusting the MBTS charge threshold in MC simulation until the simulated efficiencies match those observed in the data. The residual uncertainty in the counter efficiency after these corrections is ±0.5% for the outer and ±1.0% for the inner counters. Additionally, an uncertainty arising from the knowledge of the material in front of the MBTS detector is estimated using MC samples with an increased amount of material in front of the MBTS. Based on the MC samples, the uncertainty in the efficiency measurement due to modeling of hadronization and the underlying event is estimated to be negligible.

After adjusting the counter charge threshold, $\epsilon_{\text{sel}}$ is determined from the nominal PYTHIA8 DL MC simulations, using the fitted $f_D$ corresponding to this model, to be 99.34% with a statistical uncertainty of ±0.03%. The uncertainty in the MBTS counter efficiencies results in...
only a ±0.1% uncertainty in the overall event selection efficiency, because many counters are hit in typical events. In addition, an uncertainty of ±0.2% in \( \epsilon_{\text{sel}} \) arises from the knowledge of the material in front of the MBTS.

The fraction of events passing the inclusive selection with \( \xi < 10^{-6} \) represents an additional background component in the fiducial cross-section measurement. It is determined using the same PYTHIA8 DL MC to be \( f_{\xi<10^{-6}} = (1.37 ± 0.05)\% \), where the uncertainty is statistical.

Because the efficiency and migration corrections are correlated, they are combined in a single correction factor, \( C_{\text{MC}} = (1 - f_{\xi<10^{-6}})/\epsilon_{\text{sel}} \), for which systematic uncertainties are assessed. The systematic uncertainties include the counter efficiency variations, the impact of the material uncertainty, the uncertainty in the fitted value of \( f_D \), and the variation in \( C_{\text{MC}} \) found by comparing the PYTHIA8 DL and MBR models. Of these sources of uncertainty, the last is most important at ±0.5%. The value of \( C_{\text{MC}} \) is (99.3 ± 0.5)%. The uncertainty also implicitly contains an uncertainty due to the CD contribution, since this is included in only some of the models.

The uncertainty in the integrated luminosity is ±1.9%. It is derived, following a methodology similar to that detailed in Refs. [40,41], from a calibration of the luminosity scale using \( x\)-\( y \) beam-separation scans performed in August 2015. This calibration uncertainty is slightly smaller than what has been reported in Ref. [42] because the low-luminosity data set used in this Letter is not affected by the uncertainties related to high-luminosity runs.

The components of the fiducial cross-section calculation [Eq. (1)] are shown in Table 1 with their systematic uncertainties. The statistical uncertainties are negligible. The measured fiducial cross section is determined to be

\[
\sigma_{\text{inel}}^{\text{fid}} = 68.1 ± 0.6(\text{exp}) ± 1.3(\text{lum}) \text{ mb},
\]

where the first uncertainty refers to all experimental uncertainties apart from the luminosity and the second refers to the luminosity only.

The PYTHIA8 DL model predicts values of 71.0 mb, 69.1 mb, and 68.1 mb for \( \xi = 0.06, 0.085, \) and 0.10, respectively, all of which are compatible with the measurement. The PYTHIA8 MBR model predicts 70.1 mb, also in agreement with the measurement. The Epos LHC (71.2 mb) and QGSJET-II (72.7 mb) predictions exceed the data by 2–3\( \sigma \). The PYTHIA8 SS model predicts 74.4 mb, and thus exceeds the measured value by \( \sim 4\sigma \).

The extrapolation to \( \sigma_{\text{inel}}^{\text{fid}} \) uses constraints from previous ATLAS measurements to minimize the model dependence of the component that falls outside the fiducial region. \( \sigma_{\text{inel}} \) can be written as

\[
\sigma_{\text{inel}} = \sigma_{\text{inel}}^{\text{fid}} + \sigma_{\text{inel}}^{\text{7 TeV}}(\xi < 5 \times 10^{-6}) \\
\times \frac{C_{\text{MC}}(\xi < 10^{-6})}{\sigma_{\text{inel}}^{\text{7 TeV,MC}}(\xi < 5 \times 10^{-6})}.
\]

The term \( \sigma_{\text{inel}}^{\text{7 TeV}}(\xi < 5 \times 10^{-6}) = \sigma_{\text{inel}}^{\text{7 TeV}}(\xi > 5 \times 10^{-6}) = 9.9 ± 2.4 \text{ mb} \) is the difference between \( \sigma_{\text{inel}}^{\text{7 TeV}}(\xi < 5 \times 10^{-6}) \) and the background-subtracted distribution of the number of MBTS counters (\( n_{\text{MBTS}} \)) above threshold in data and MC simulation for (top) the inclusive selection and (bottom) the single-sided selection. The ratio of the MC models to the data is also shown. The experimental uncertainty is shown as a shaded band around the data points. The models shown here use the \( f_D \) value determined from the \( R_{\text{SS}} \) measurement.

**TABLE 1.** Inputs to the calculation of the measured cross section and their systematic uncertainties.

<table>
<thead>
<tr>
<th>Factor</th>
<th>Value</th>
<th>Relative ( \text{uncertainty} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>Number of events passing the</td>
<td>4,159,074</td>
<td>⋯</td>
</tr>
<tr>
<td>inclusive selection (N)</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Number of background events (( N_{\text{BG}} ))</td>
<td>51,187 ±50%</td>
<td></td>
</tr>
<tr>
<td>Integrated luminosity [( \text{pb}^{-1} )] (( L ))</td>
<td>60.1 ±1.9%</td>
<td></td>
</tr>
<tr>
<td>Trigger efficiency (( \epsilon_{\text{trg}} ))</td>
<td>99.7% ±0.3%</td>
<td></td>
</tr>
<tr>
<td>MC correction factor (( C_{\text{MC}} ))</td>
<td>99.3% ±0.5%</td>
<td></td>
</tr>
</tbody>
</table>
FIG. 3. The inelastic proton-proton cross section versus √s. Measurements from other hadron collider experiments [6,7,9,14,15] and the Pierre Auger experiment [16] are also shown. Some LHC data points have been slightly shifted in the horizontal position for display purposes. The data are compared to the PYTHIA8, EPOS LHC and QGSJet-II MC generator predictions. The uncertainty in the ATLAS ALFA measurement is smaller than the marker size.

measured at 7 TeV using the ALFA detector [8], σinel 7 TeV, and σinel measured at 7 TeV for ξ > 5 × 10^-6 using the MBTS [12] (The 7 TeV result is corrected upward by 1.9% following an improved luminosity calibration [40]). The uncertainties of the two measurements are uncorrelated.

The PYTHIA8 DL and PYTHIA8 MBR MC samples are used to assess the systematic uncertainty in the MC-derived ratio of cross sections in Eq. (2), which is determined to be 1.015 ± 0.081. (The value of the ratio arises from an approximately 20% increased cross section from increasing √s which is largely compensated by a 15% decrease due to the change in the ξ distribution.) These models also agree with the measurement of σinel 7 TeV (ξ < 5 × 10^-6) to within 2σ.

The measured value for σinel is

σinel = 78.1 ± 0.6(exp) ± 1.3(lum) ± 2.6(extrap) mb.

This and other inelastic cross-section measurements are compared to several Monte Carlo models in Fig. 3. Additional predictions range between 76.6 and 81.6 mb [43–47]. Compared to the measurement with the ALFA detector at √s = 7 TeV the cross section is higher by (9 ± 4)%.

In summary, a measurement of the inelastic cross section in 60 μb^-1 of proton-proton collision data at √s = 13 TeV collected with the ATLAS detector at the LHC is presented. The measurement is performed in a fiducial region ξ > 10^-6, and the result is extrapolated to the inelastic cross section using measurements at √s = 7 TeV. The measured cross section agrees well with a variety of theoretical predictions and is consistent with the inelastic cross section increasing with center-of-mass energy, as observed at lower energies.

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[18] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the z axis along the beam pipe. The x axis points from the IP to the center of the LHC ring, and the y axis points upward. Cylindrical coordinates \((r, \phi)\) are used in the transverse plane, \(\phi\) being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle \(\eta = -\ln \tan(\theta/2)\).


Fysiska institutionen, Lunds universitet, Lund, Sweden

Departamento de Física Teorica C-15, Universidad Autonoma de Madrid, Madrid, Spain

Institut für Physik, Universität Mainz, Mainz, Germany

School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom

CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France

Department of Physics, University of Massachusetts, Amherst, Massachusetts, USA

School of Physics, University of Melbourne, Victoria, Australia

Department of Physics, The University of Michigan, Ann Arbor, Michigan, USA

Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA

INFN Sezione di Milano, Italy

Dipartimento di Fisica, Università di Milano, Milano, Italy

B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus

National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus

Group of Particle Physics, University of Montreal, Montreal, QC, Canada

P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia

Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia

National Research Nuclear University MEPhI, Moscow, Russia

D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia

Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany

Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany

Nagasaki Institute of Applied Science, Nagasaki, Japan

Graduate School of Science and Kobe yashi-Maskawa Institute, Nagoya University, Nagoya, Japan

INFN Sezione di Napoli, Italy

Dipartimento di Fisica, Università di Napoli, Napoli, Italy

Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA

Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands

Department of Physics, Northern Illinois University, DeKalb, Illinois, USA

Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia

Department of Physics, New York University, New York, New York, USA

Ohio State University, Columbus, Ohio, USA

Faculty of Science, Okayama University, Okayama, Japan

Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, Oklahoma, USA

Department of Physics, Oklahoma State University, Stillwater, Oklahoma, USA

Palacký University, RCPTM, Olomouc, Czech Republic

Center for High Energy Physics, University of Oregon, Eugene, Oregon, USA

LAL, Univ. Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France

Graduate School of Science, Osaka University, Osaka, Japan

Department of Physics, University of Oslo, Oslo, Norway

Department of Physics, Oxford University, Oxford, United Kingdom

INFN Sezione di Pavia, Italy

Dipartimento di Fisica, Università di Pavia, Pavia, Italy

Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania, USA

National Research Centre “Kurchatov Institute” B.P. Konstantinov Petersburg Nuclear Physics Institute, St. Petersburg, Russia

INFN Sezione di Pisa, Italy

Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy

Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, Pennsylvania, USA

Laboratório de Instrumentação e Física Experimental de Partículas—LIP, Lisboa, Portugal

Faculdade de Ciências, Universidade de Lisboa, Lisboa, Portugal

Department of Physics, University of Coimbra, Coimbra, Portugal

Centro de Física Nuclear da Universidade de Lisboa, Lisboa, Portugal

Departamento de Física, Universidade do Minho, Braga, Portugal

Departamento de Física Teorica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain

Dep Física and CEFITEC of Faculdade de Ciencias e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal

Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic

Czech Technical University in Prague, Praha, Czech Republic

Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic

State Research Center Institute for High Energy Physics (Protvino), NRC KI, Russia

Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
1124 INFN Sezione di Roma, Italy
1125 Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy
1126 INFN Sezione di Roma Tor Vergata, Italy
1127 Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
1128 INFN Sezione di Roma Tre, Italy
1129 Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
1130 Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies—Université Hassan II, Casablanca, Morocco
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1132 Faculté des Sciences Semlalia, Université Cadi Ayyad, LPTP, Marrakech, Morocco
1133 Faculté des Sciences, Université Mohamed Premier and LPTP, Oujda, Morocco
1134 Faculté des sciences, Université Mohammed V, Rabat, Morocco
1135 DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France
1136 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, California, USA
1137 Department of Physics, University of Washington, Seattle, Washington, USA
1138 Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
1139 Department of Physics, Shinshu University, Nagano, Japan
1140 Fachbereich Physik, Universität Siegen, Siegen, Germany
1141 Department of Physics, Simon Fraser University, Burnaby, BC, Canada
1142 SLAC National Accelerator Laboratory, Stanford, California, USA
1143 Department of Mathematics, Physics & Informatics, Comenius University, Bratislava, Slovak Republic
1144 Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
1145 Department of Physics, University of Cape Town, Cape Town, South Africa
1146 School of Physics, University of the Witwatersrand, Johannesburg, South Africa
1147 Department of Physics, Stockholm University, Sweden
1148 The Oskar Klein Centre, Stockholm, Sweden
1149 Physics Department, Royal Institute of Technology, Stockholm, Sweden
1150 Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook, New York, USA
1151 Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
1152 School of Physics, University of Sydney, Sydney, Australia
1153 Institute of Physics, Academia Sinica, Taipei, Taiwan
1154 Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel
1155 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
1156 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
1157 International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan
1158 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
1159 Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
1160 Department of Physics, University of Toronto, ON, Canada
1161 TRIUMF, Vancouver, BC, Canada
1162 Department of Physics and Astronomy, York University, Toronto, ON, Canada
1163 Faculty of Pure and Applied Sciences, and Center for Integrated Research in Fundamental Science and Engineering, University of Tsukuba, Tsukuba, Japan
1164 Department of Physics and Astronomy, Tufts University, Medford, Massachusetts, USA
1165 INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy
1166 ICTP, Trieste, Italy
1167 Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy
1168 Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
1169 Department of Physics, University of Illinois, Urbana, Illinois, USA
1170 Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain
1171 Department of Physics, University of British Columbia, Vancouver, BC, Canada
1172 Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada
1173 Department of Physics, University of Warwick, Coventry, United Kingdom
1174 Waseda University, Tokyo, Japan
1175 Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel
1176 Department of Physics, University of Wisconsin, Madison, Wisconsin, USA
1177 Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany