From confined spinons to emergent fermions: Observation of elementary magnetic excitations in a transverse-field Ising chain

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We report on spectroscopy study of elementary magnetic excitations in an Ising-like antiferromagnetic chain compound SrCo2V2O8 as a function of temperature and applied transverse magnetic field up to 25 T. An optical as well as an acoustic branch of confined spinons, the elementary excitations at zero field, are identified in the antiferromagnetic phase below the Néel temperature of 5 K and described by a one-dimensional Schrödinger equation. The confinement can be suppressed by an applied transverse field and a quantum disordered phase is induced at 7 T. In this disordered paramagnetic phase, we observe three emergent fermionic excitations with different transverse-field dependencies. The nature of these modes is clarified by studying spin dynamic structure factor of a 1D transverse-field Heisenberg-Ising (XXZ) model using the method of infinite time evolving block decimation. Our work reveals emergent quantum phenomena and provides a concrete system for testing theoretical predications of one-dimension quantum spin models.

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I. INTRODUCTION

Emergent states of matter in quantum magnets are characterized by their elementary excitations that can be induced and tuned in an external magnetic field. Due to enhanced quantum fluctuations and reduced dimensionality, the low-dimensional quantum magnets host very often exotic quantum states of matter [1–13]. Especially for the one-dimensional (1D) spin systems, realization of novel quantum spin states is not only of lively experimental interest [7–10,12,13] but also a constant focus of theoretical study [1,2,14–17], since quantitative description of the elementary excitations can be provided by rigorous theoretical approaches and precise comparisons between theoretical predictions and experimental results can be made [7–10,12,13,15].

In the 1D spin-1/2 systems, the elementary magnetic excitations are spinons with fractional spin quantum number $S = 1/2$ [18]. Confinement of the fractional spinon excitations can be even realized [9,10,19–22] as an analogy with quark confinement in particle physics [23]. The concept of spinon confinement is illustrated in Figs. 1(a)–1(d)]. In a classical picture, the ground state of a spin-1/2 Ising antiferromagnetic chain corresponds to an antiparallel alignment of neighboring spins [Fig. 1(a)]. A single spin flip frustrates the intrachain exchange interactions $J$ and creates two spinons, each with spin-1/2 [Fig. 1(b)]. Subsequent spin flips can lead to propagation of the spinons along the chain [Fig. 1(c)], In the presence of weak interchain exchange interactions $J_{\perp} \ll J$ [Fig. 1(d)], a 3D Néel antiferromagnetic order can be stabilized at low temperature $T < T_N$ and the separation of spinons from each other is unfavorable due to the frustrated interchain couplings. Thus the two spinons feel a confining potential, increasing with the distance between them, and form quantized spinon bound states [Figs. 1(e) and 1(f)] [24,25]. Since the spinons are confined by the interchain couplings only in the antiferromagnetic phase, the magnetic phase transition can be thought as accompanied by a spinon confinement-deconfinement transition.

SrCo$_2$V$_2$O$_8$ is a representative realization of the paradigmatic 1D Heisenberg-Ising (XXZ) model. It crystallizes in the tetragonal structure with space group $I4_{1}cd$. Screw spin chains in SrCo$_2$V$_2$O$_8$ are constituted by edge-sharing CoO$_6$ octahedra and propagate along the crystallographic c axis [Fig. 1(g)]. Each screw period consists of four Co$^{2+}$ ions with effective spin $S = 1/2$, corresponding to the lattice constant $c$. The unit cell contains four screw chains with either left- or right-hand chirality [Fig. 1(b)]. Due to spin-orbit coupling and crystal field splitting, the Co$^{2+}$ ions have an Ising-like anisotropy along the c axis [26,27]. The dominant exchange interaction $J$ is antiferromagnetic and between the nearest-neighbor Co$^{2+}$ spins in the chain. The interchain coupling $J_{\perp}$ is much smaller and dominated by exchange between the nearest-neighbor Co$^{2+}$ ions in the ab plane that are from the neighboring chains with the same chirality [Fig. 1(h)].

At $T_N = 5$ K a long-range Néel-type antiferromagnetic order along the chain is stabilized [Fig. 1(g)] [26–29].
Here we perform terahertz spectroscopy on the low-energy spin excitations in the 1D Ising-like antiferromagnetic-chain SrCo2V2O8 as a function of transverse magnetic field, supplemented by studying spin dynamics of the corresponding 1D transverse-field XXZ model using the rigorous method of infinite time evolving block decimation (iTEBD). In zero field and below TN, confinement of spinons of optical branch as well as of acoustic branch is observed and identified. Both the optical and the acoustic series of confined spinons can be described by a 1D Schrödinger equation with linear confining potential. In small transverse fields, where the system remains antiferromagnetically ordered, the confining effects are found to be reduced: the confining potential becomes shallower with increasing field and the higher energy spinon bound states tend to collapse on the lowest level. At the transverse-field-induced order-disorder transition, we observe collapse of the quantized confined-spinon levels. Above this transition, the low-energy spin excitations in SrCo2V2O8 are not deconfined spinons, but characterized by emergent fermions of the 1D transverse-field Ising class. We successfully identify the observed emergent fermions and their field dependencies by performing theoretical study of the 1D spin-1/2 XXZ model using the iTEBD method.

II. EXPERIMENTAL DETAILS AND THEORETICAL METHOD

High-quality single crystals of SrCo2V2O8 were grown using the floating-zone method and characterized by x-ray diffraction, neutron diffraction, and magnetization measurements [26]. Single crystals for the optical study were oriented using Laue diffraction and cut perpendicular to the tetragonal a axis with a typical surface of 4 × 4 mm² and a thickness of 1 mm. High-field optical measurements were performed in the High Field Magnet Laboratory in Nijmegen. A Michelson interferometer with mercury lamp was used to generate electromagnetic waves in the terahertz frequency range which were detected by a silicon bolometer cooled at 1.6 K. The sample transmission spectrum was measured in the spectral range of 1.2–10 meV (0.3–2.5 THz) at magnetic fields up to 25 T. Magnetic fields were applied perpendicular to the crystallographic ac plane and parallel to the propagation direction of the electromagnetic wave (Faraday configuration).

Standard infinite time evolving block decimation (iTEBD) method [33,34] was used to study spin dynamics of the 1D XXZ antiferromagnetic model in a transverse field. This method has been applied to other 1D systems [33,35] and turned out to be efficient to study 1D spin dynamics. Similar to the method of density matrix renormalization group, the iTEBD method assumes that the ground state (GS) of the spin system can be expressed as a matrix product state [33,34]. The two-point correlation function

\[
C(r,t) = \frac{\langle S_y(r,t)S_y(0,0) + S_y(r,t)S_y(0,0) \rangle_{GS}}{\langle S_y(r,t)S_y(0,0) \rangle_{GS}}
\]

is calculated to compare with experimental results. The spin components S_y and S_x are coupled to the THz magnetic field and the applied external magnetic field is set along the x direction. The spin dynamic structure factor (DSF) can be

Below TN confinement of spinons is characterized by a series of quantized spinon bound states [Fig. 1(f)] [22], which with spin S = +1 or −1 can be split in longitudinal magnetic field and exhibit linear field dependence [20,22]. By applying a transverse magnetic field, the Néel temperature can be reduced according to the magnetic susceptibility measurements [26,27,30–32]. At 2 K, a phase transition from the Néel antiferromagnetic (AFM) phase to a disordered paramagnetic (PM) phase is induced by a transverse field of 7 T [26,27].

The reduction of Néel temperature suggests that the applied transverse magnetic field competes with the weak interchain couplings, and thus reduces the confining effects. This motivates us to utilize a transverse field to tune the spinon confinement. By continuously changing the transverse field, the confinement is expected to be suppressed with deconfined spinons observed.
obtained via Fourier transformation
\[ S^{yz+\zeta}(q,\omega) = \int_{-\infty}^{\infty} dr \int_{-\infty}^{\infty} dt \, C(r,t) e^{i(qz-\omega t)} , \]
(2)
corresponding to the magnetic-dipole excitations observed in the optical measurements. In our calculation, the bond dimension is taken to be 40 and the discrete time step in the real time evolution is 0.02/\( J \). The maximal time difference \( |t - t'| \) is set to 32/\( J \), and we apply the linear prediction method [35] to further improve the resolution.

III. EXPERIMENTAL RESULTS

A. Confinement of acoustic and optical spinons in zero field

A spin-flip excitation can be triggered by absorbing a photon, when the photon has the same energy as the two spinons created [Figs. 1(a) and 1(b)]. The resonance absorption of photons is manifested by a transmission minimum (or an absorption maximum) at the corresponding photon energy. After the creation, the two spinons can propagate along the chain [Fig. 1(c)]. In presence of an effective linear confining potential \( \lambda |z| \), the two-spinon bound state \( \psi(z) \) as a function of the distance \( z \) between them can be described by the one-dimensional Schrödinger equation
\[ -\frac{\hbar^2}{2\mu} \frac{d^2\psi}{dz^2} + \lambda |z| \psi = (E - 2E_0)\psi , \]
(3)
with \( \hbar \) the Planck constant, \( \mu \) an effective mass, and \( 2E_0 \) the energy threshold for creating two spinons. \( \lambda = 2J_\perp (S_z)^2/c \) is an effective description of the interchain exchange interaction \( J_\perp \), with the lattice constant \( c \). The solution of Eq. (3) is given by the linear dependence of \( E_i \) on \( \zeta_i \)
\[ E_i = 2E_0 + \zeta_i ^{\lambda^{2/3}} \left( \frac{\hbar^2}{\mu} \right)^{1/3} , \]
(4)
where \( i = 1, 2, 3, \ldots \) and \( \zeta_i \) are given by the negative zeros of the Airy function \( A(-\zeta_i) = 0 \) [36].

Figure 2(a) shows the ratio of two transmission spectra of \( \text{SrCo}_2\text{V}_2\text{O}_8 \) that are measured above and below the Néel temperature \( T_N = 5 \text{ K} \). Spin excitations due to the magnetic phase transition are manifested by the peaks. One can see three groups of excitations that are denoted by \( \alpha_i \), \( \beta_i \), and \( \gamma \). Eigenenergies \( E_{\alpha_i} \) of the \( \alpha_i \) series follow the linear dependence on \( \zeta_i \) as given in Eq. (4) with \( 2E_0 = 0.68 \text{ meV} \) and \( \lambda^{2/3} \left( \frac{\hbar^2}{\mu} \right)^{1/3} = 0.38 \text{ meV} \) as shown in Fig. 2(b). This indicates that the \( \alpha_i \) series is confined spinon excitations, which is consistent with previous study [22]. In the same representation [Fig. 2(b)], one can see that the \( \beta_i \) modes with higher eigenenergies \( E_{\beta_i} \) also follow a linear dependence and the linear fit is given by \( 2E_0 = 7.55 \text{ meV} \) and \( \lambda^{2/3} \left( \frac{\hbar^2}{\mu} \right)^{1/3} = 0.23 \text{ meV} \). Thus the \( \beta_i \) modes can be identified as another branch of confined spinons. Compared to the \( \alpha_i \) series, the \( \beta_i \) modes have higher energy threshold and larger effective mass, if the same confining potential is assumed for both branches.

It will be shown in the following that we are able to assign the \( \alpha_i \) modes as confined spinons of acoustic branch and the \( \beta_i \) modes as of optical branch, in analogy with the concept of optical and acoustic phonons. In contrast, the mode \( \gamma \) at 7.72 meV does not exhibit a series of excitations with higher energies and thus not fit into the same scheme as the \( \alpha_i \) and \( \beta_i \) modes. It is worth noting that the \( \beta_i \) modes are observed in the same energy range as the lower-lying optical phonon bands, making the higher energy levels difficult to resolve.

B. Tuning the spinon confinement in low transverse magnetic field

Figures 3(a) and 3(b) show transmission spectra of the acoustic \( \alpha_i \) and optical \( \beta_i \) spinons of \( \text{SrCo}_2\text{V}_2\text{O}_8 \) as a function of applied transverse magnetic field along the crystallographic \( a \) axis (\( H \parallel a \)) up to 9 T. At zero field the series of

FIG. 2. (a) Ratio of transmission spectra that are measured at 7.4 K (above \( T_N \)) and at 1.5 K (below \( T_N \)). Two series of excitations \( \alpha_i \) and \( \beta_i \) are observed and indicated by the arrows, corresponding to confined spinons of acoustic and optical branch, respectively. Another mode \( \gamma \) at 7.72 meV is marked by the star. (b) Eigenenergies \( E_{\alpha_i} \) and \( E_{\beta_i} \) of the modes \( \alpha_i \) and \( \beta_i \), respectively, as a function of \( \zeta_i \) the negative zeros of the Airy function [36]. Linear fits according to the 1D Schrödinger equation are shown by the solid lines [see Eqs. (3) and (4)].

FIG. 3. Transmission spectra of (a) the acoustic \( \alpha_i \) and (b) the optical \( \beta_i \) spinons in applied transverse magnetic fields along the crystallographic \( a \) axis (\( H \parallel a \)) up to 9 T. Series of absorption modes \( \alpha_i \) can be observed at each field in (a). In (b) the zero-field mode \( \beta_1 \) splits into two modes in magnetic fields due to Zeeman interaction, as marked by the plus and minus signs. Below 7 T, shifting of the mode \( \beta_2 \) to higher energy with increasing fields is observed and indicated by the arrows. The \( \gamma \) mode at 7.72 meV, marked by the star, has almost field-independent eigenenergy below 7 T and is not resolvable at 7 T or above.
Interestingly, the dependence of the resonance energy on field and the confining-potential profile is axis, where the threshold energy exhibits Zeeman-type linear fit determines the parameters of the magnetic field. At 7 T, the higher levels collapse to the lowest-level mode, for the acoustic and optical spinons: the γ mode is suppressed; and only three modes, α, β−, and β+, remain. This indicates a field-induced phase transition from the antiferromagnetically (AFM) ordered to a paramagnetic (PM) quantum-disordered phase [26,27] at $H_c = 7$ T.

Transmission minima, corresponding to confined acoustic and optical spinons, can be identified in agreement with the temperature-dependent measurements [Fig. 2(a)]. For the acoustic branch [Fig. 3(a)], a series of transmission minima for 1 T is observed at almost the same energies as for 0 T, while at 3 T the resolved modes are shifted to higher energies [see also Fig. 4(e)]. Interestingly, the dependence of the resonance energies on field still follows the characteristic dependence at finite transverse fields, as shown in Figs. 4(a) and 4(b). This indicates that the spinon confinement is still a valid description of the spin dynamics in low transverse fields. The linear fit determines the parameters $\lambda = 2/3(\frac{2E_0}{\mu_0 T})^{1/3}$ = 0.36 meV and $2E_0$ = 0.87 meV for 3 T [Fig. 4(b)], where the former decreased while the latter significantly increased compared to the results at zero field. Thus, by applying a transverse magnetic field, one not only changes the energy threshold of the confined spinons, but also tunes the profile of the confining potential.

As shown in Figs. 4(c) and 4(d), the confining potential and threshold energy show a nonlinear decrease and increase, respectively, with increasing transverse field. This is different from the situation of small longitudinal fields along the Ising axis, where the threshold energy exhibits Zeeman-type linear dependence on field and the confining-potential profile is almost unchanged [20,22], since the $U(1)$ symmetry of the spin Hamiltonian is broken in the transverse field. The results in Figs. 4(c) and 4(d) suggest that the confining potential becomes shallower in the transverse field. In a shallower confining potential, as illustrated in Figs. 1(e) and 1(f), the difference between two neighboring spinon bound states reduces. When the difference becomes very small, the higher level cannot be distinguished experimentally. With increasing transverse field, the highest resolvable mode shifts to lower energy levels and, finally, when the second level ($\alpha_2$) collapses onto the first one ($\alpha_1$), the spinon hierarchy breaks down and the spinon confinement is completely suppressed.

As shown in Fig. 3(b), the zero-field optical mode $\beta_1$ splits into two in the transverse fields due to Zeeman interaction, one shifting up while the other shifting down, as marked by the plus and minus signs. The higher levels of confined optical spinons should be also split in the transverse field, since they belong to the same series. Due to strong phonon absorption at the same spectral range, fewer up-shifting confined modes of the optical branch can be resolved by the field-dependent measurements [marked by the arrows in Fig. 3(b), while higher levels of the down-shifting modes are covered by the strong $\beta_1$ absorptions. Nevertheless, one can clearly observe some similar tuning behavior of the confined optical spinons as of the acoustic modes: the higher-level modes move towards the first level with increasing transverse field, although they all shift to higher energies with increasing fields; above 3 T the $\beta_3$ mode cannot be distinguished from the $\beta_2$ mode, and at 7 T only the main modes ($\beta_+ and \beta_-$) can be observed indicating the suppression of the spinon confinement.

In addition, the γ mode observed at 7.72 meV in zero field [Fig. 2(a)] shifts only slightly to higher energy and becomes broader with the increasing transverse fields. At 7 T, the γ mode is already suppressed. Although the origin of this mode remains unclear, it is natural to relate the γ mode to the 3D magnetic order because it is observed only below the magnetic phase transition.

The eigenenergies of the confined acoustic and optical spinons, and of the γ mode, are summarized in Fig. 4(e) as a function of the applied transverse magnetic field. Transverse-field tuning of the spinon confinement is obvious: starting from zero field, both the acoustic and optical spinons are confined due to the interchain couplings and exhibit clear hierarchy of the spinon bound states. With increasing field, the interchain couplings are effectively reduced and the higher-level bound states are shifted downwards until the second level collapses onto the first level at 7 T. The breakdown of the hierarchy, suggesting the suppression of the interchain couplings, is an evident signature of the field-induced phase transition [26,27].

### C. Emergent fermions of 1D transverse-field Ising class

The above section documents the deconfining behavior of the confined spinons in a transverse magnetic field below the phase transition at $H_c = 7$ T. Above $H_c$ in the quantum disordered phase, one would expect emergent elementary excitations that characterize spin state of the disordered phase. The characteristic excitations can be revealed by performing magneto-optic experiment in high magnetic fields.

A typical absorption-coefficient spectrum is shown in Fig. 5(a) for a transverse field of 11 T above the field-induced phase transition. One can observe three absorption maxima at 3.1, 6.4, and 8.8 meV, corresponding to the $\alpha, \beta_-, and \beta_+$ modes, respectively. Their evolution with the transverse field is displayed in Fig. 5(c) by a contour plot of the absorption coefficient as a function of photon energy and transverse field, together with colored symbols highlighting their eigenenergies. The three modes $\alpha, \beta_-, and \beta_+$ are the
dominant excitations that can be unambiguously followed in higher fields.

The $\alpha$ mode, evolving from the lowest-lying confined spinon excitation $\alpha_1$ [Figs. 3(a) and 4(e)], increases in energy with increasing fields. The $\beta_-$ and $\beta_+$ modes correspond to the degenerate $\beta_1$ mode at zero field which is split in finite fields. At 22 T the $\alpha$ and $\beta_-$ modes do not cross but form a small hybridization gap of 0.68 meV [Fig. 5(c)]. Above 22 T, the mode $\alpha$ shifts slightly downwards while the $\beta_-$ mode shifts upwards.

IV. DISCUSSION

In order to understand origin of the observed modes in high fields and clarify the two branches of confined spinons, we use the standard iTEBD method [33,34] to study the spin dynamics of the spin-1/2 1D Heisenberg-Ising XXZ model, which describes the magnetic interactions of the cobalt spin chain. In a transverse magnetic field, the spin Hamiltonian is given by

$$ J \sum_{\langle i,j \rangle} (S_{x,i}S_{x,j} + S_{y,i}S_{y,j} + \Delta S_{z,i}S_{z,j}) - g_s \mu_B H_s \sum_i S_{z,i}, $$

where $\langle i,j \rangle$ denotes the nearest neighbors, $J$ is the nearest-neighbor intrachain antiferromagnetic exchange interaction, and $\Delta$ refers to an Ising-like anisotropy [37,38]. The last term is the Zeeman interaction term with the transverse magnetic field $H_s$ along the $a$ axis, the Landé $g$ factor, and the Bohr magneton $\mu_B$.

The obtained spin dynamic structural factor (DSF) according to Eq. (2) is shown for the whole Brillouin zone at various transverse magnetic fields in Fig. 6. At zero field the known dispersion relation of spinons [10,39] is well reproduced in our calculations [Fig. 6(a)]. The dispersion relation is symmetric with respect to $q = \pi$. Two minima of the lower-energy boundary are degenerate and located at $q = 0$ and $q = \pi$, respectively. In a finite transverse field, the energy at $q = 0$ starts to increase and the corresponding minimum shifts to larger $q$ values, while at $q = \pi$ the energy minimum is only slightly lowered. Moreover, a clear split of the dispersion relation occurs between $q = 0$ and $q = \pi$ and becomes more significant with the increasing transverse fields. These features of the dispersion relation characterize emergent fermions of the transverse-field Ising-like chain [40,41].

Terahertz spectroscopy is expected to probe the dispersion relation close to the $\Gamma$ point ($q = 0$) in reciprocal space, since the wavelength of terahertz electromagnetic waves is much larger than the lattice constants of SrCo$_2$V$_2$O$_8$. To compare with the experimental results, in Fig. 5(d) the obtained DSF at $q = 0$ is shown as a function of energy and transverse magnetic field, which resembles the field dependence of the experimentally observed $\alpha$ mode. This indicates that the $\alpha$ fermionic excitations feature the spin dynamics at the $\Gamma$ point and the corresponding confined spinons observed at zero field belong to an acoustic branch.

To reconcile the experimental results, the DSF at $q = \pi/2$ is also shown in Fig. 5(d) as a function of transverse field. An excitation mode close to 8 meV at zero field splits into two branches in the finite transverse fields, one increasing in energy while the other decreasing. Figure 5(b) shows the DSF...
very close to each other.

structure factor represents the dispersion relation of (a) spinons and (b)–(e) emergent fermions.

be testified in real materials.

method [33,34].

[Figs.5(c)]. In the extended Brillouin zone the crossing of the

the experimental observation of the hybridization gap at 22 T

acoustic

ferromagnetic chain SrCo2V2O8 in transverse fields up to

studied low-energy spin dynamics of the Ising-like anti-

boundary conditions gives the interaction matrix, and thus mixes

transverse field can lead to nonzero off-diagonal elements of

The observation of the $q = \pi/2$ mode by optical spec-

troscopy indicates that the Brillouin zone is folded by a

factor of four and translational invariance of the 1D $XXZ$

model is broken due to the four spin sites per unit cell in

SrCo2V2O8 [Fig. 1(g)]. The zone-folding effect also explains

the parameters

$J = \frac{\pi}{2} \sqrt{2} \frac{\Delta}{\pi} + \frac{\pi}{2} \frac{\Delta}{\pi}$

and the

$gx\mu B H_x$

factor smaller than the longitudinal $g_\|= 5.5$ and the $\Delta > 1$

value reveals an Ising-like anisotropy [22]. The exchange

interaction is quite larger than the value of

$J = 2.8$ meV

in BaCo2V2O8 [20], while the anisotropy of the two systems is

very close to each other.

These results indicate that the experimentally observed $\beta_-$

and $\beta_+$ modes correspond to the emergent-fermion excita-

tions at $q = \pi/2$ of the 1D transverse-field $XXZ$ model, which are

the low-energy spin excitations characteristic for the quantum

disordered phase. It is also clarified that the $\beta_i$ modes observed

at zero field [Fig. 1(a)] can be assigned as confined spinons of

optical branch, in analogy with the concept of optical phonons. At

zero field the DSF at $q = \pi/2$ is much larger than that at

$q = 0$. This explains the stronger absorption of the $\beta_1$

mode compared to the $\alpha_1$ mode, as observed in Fig. 2(a).

With further increasing transverse fields, our calculation

suggests that the lower boundary at $q = \pi$ will be first

gapless at 55 T, where a transverse-field Ising quantum critical

point [40,41] is expected to emerge. At the quantum critical point, predicted dynamical features [16,42] are interesting to

be testified in real materials.

V. SUMMARY

Using high-resolution terahertz spectroscopy, we have studied low-energy spin dynamics of the Ising-like anti-

ferromagnetic chain SrCo2V2O8 in transverse fields up to

25 T. According to the terahertz electrodynamic response, the

transverse field-tuned spin dynamics is clarified and it can be

categorized into three regimes:

(i) At zero field, the spin dynamics is characterized by

confined spinons. We have identified confined spinons of an

optical branch in addition to an acoustic branch. Both branches

can be described by a 1D Schrödinger equation with linear

confining potential. Compared to the acoustic one, the optical

branch is found at higher energy with larger effective mass.

(ii) We have shown that the confinement of the optical

and of the acoustic spinons can be tuned down by a small

transverse field: the confining potential becomes shallower

and the higher confinement levels are shifted closer to the

lowest one. These properties can be viewed as deconfining

behavior of the confined spinons, since the spinon series can

be still described by the 1D Schrödinger equation and the small

transverse field can be viewed as a perturbation.

(iii) We have also provided experimental evidence of the

suppression of the spinon confinement by large transverse

field, which is manifested by the collapse of the spinon

hierarchy. This, however, does not result in deconfined spinons,

but leads to the emergent-fermion excitations of the 1D

transverse-field Ising class, because the transverse field is

too strong to be treated as perturbation and even a phase

transition from the antiferromagnetic to paramagnetic phase is

induced. The evolution of emergent fermions in high transverse

fields has been revealed by the terahertz spectroscopy and by

studying low-energy spin dynamics of the 1D Ising-like

$XXZ$ model using the method of infinite time evolving block
decimation. A quantum critical point of transverse-field Ising

class is predicted to occur at 55 T yet to be testified.

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FIG. 6. (a)–(e) Spin dynamic structure factor $S(q, \omega$) [Eq. (2)] in the reciprocal space $(0, 2\pi/c_0)$ calculated up to 4$J$ for the 1D Heisenberg-Ising $XXZ$ model Eq. (5) in various transverse magnetic fields $g_\mu B H_x$ of (a) 0, (b) 0.32, (c) 0.64, (d) 0.96, and (e) 1.28$J$ by the iTEBD method [33,34]. $c_0$ is the natural lattice constant of the 1D $XXZ$ model and $J$ is the nearest-neighbor exchange interaction. The dynamic structure factor represents the dispersion relation of (a) spinons and (b)–(e) emergent fermions.

spectra of $q = 0$ and $q = \pi/2$ which agrees very well with the

absorption-coefficient spectrum of 11 T in Fig. 5(a). The good

agreement between experiment and theory is achieved with

the parameters $J = 3.55$ meV, $g_\|= 2.79$, and $\Delta = 2.00$. The

$g_i$ factor smaller than the longitudinal $g_\|= 5.5$ and the $\Delta > 1$

value reveals an Ising-like anisotropy [22]. The exchange

interaction is quite larger than the value of $J = 2.8$ meV

in BaCo2V2O8 [20], while the anisotropy of the two systems is

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The observation of the $q = \pi/2$ mode by optical spec-

troscopy indicates that the Brillouin zone is folded by a

factor of four and translational invariance of the 1D $XXZ$

model is broken due to the four spin sites per unit cell in

SrCo2V2O8 [Fig. 1(g)]. The zone-folding effect also explains

the experimental observation of the hybridization gap at 22 T

[Figs. 5(c)]. In the extended Brillouin zone the crossing of the

acoustic $\omega$ and optical $\beta_-$ modes [Figs. 5(d)] is protected by the translational symmetry, since they are located at different

$q$ vectors. Due to the broken symmetry in the material, the

transverse field can lead to nonzero off-diagonal elements of the interaction matrix, and thus mixes $\omega$ and $\beta_-$ states and opens the hybridization gap.

With further increasing transverse fields, our calculation

suggests that the lower boundary at $q = \pi$ will be first

gapless at 55 T, where a transverse-field Ising quantum critical

point [40,41] is expected to emerge. At the quantum critical point, predicted dynamical features [16,42] are interesting to

be testified in real materials.
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