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I. INTRODUCTION

The last few years it has become possible to obtain state-to-state cross sections for transitions between rotation–inversion (\(j_f^e\)) states of NH\(_3\), induced by collisions with various perturbers. Advances in molecular beam techniques and in laser spectroscopy have made it feasible to discriminate between the symmetric (\(+\)) and antisymmetric (\(-\)) states of the inversion doublets. The close coupling (CC) method for the accurate quantum mechanical treatment of the problem has been well established theoretically for quite some time. However, computer systems that meet the computational demands have only recently become available.

In this paper we consider collisions of NH\(_3\) with Ar. This study was undertaken mainly for two reasons. First, comparison of theoretical results with experimental data enables us to determine the accuracy of an \(ab\) \textit{initio} intermolecular potential energy surface,\(^1\) in the region that is probed in scattering experiments. In addition, we used a slightly different potential in order to gain some understanding of the sensitivity of the cross sections to variations in the potential surface. This second potential contains a scaling parameter that was chosen to account for spectroscopic data regarding the bound states of Ar–NH\(_3\). In particular, we wanted to see whether a variation that improved bound state results would also improve the outcome of the scattering calculations.

Second, we investigate how the description of the umbrella inversion of NH\(_3\) influences the cross sections. Until now, this inversion has been included in scattering calculations on NH\(_3\) only via a model in which the inversion-tunneling wave function is a linear combination of two delta functions centered at the equilibrium positions (Davis and Boggs,\(^2\) Green\(^3\)). Although the model has its justification in the fact that the period for inversion is much longer than the duration of a collision, it was not clear whether experimentally found deviations from predicted propensity rules could not be attributed to the neglect of the inversion motion in the description of the intermolecular potential.\(^4,5\) Here we take the inversion degree of freedom explicitly into account, in order to assess how severe an approximation is made in neglecting it.

Finally, we have performed some calculations using the much cheaper coupled states (CS) approximation, to find out how this approximation affects the calculated cross sections. In calculations on He–NH\(_3\) using the CS method\(^6\) certain theoretical cross sections are found to vanish or almost vanish, whereas the experimental cross sections are significantly different from zero.\(^4\) By applying the CS approximation, together with the full CC method on the same Ar–NH\(_3\) potential surface, we can establish to what extent deviations are caused by the theoretical scattering method.

II. THEORY

The coordinate system used in the CC method is the space-fixed frame.\(^7\) The vector \(\mathbf{R}\), with polar angles \((\beta,\alpha)\) in this frame, points from the NH\(_3\) center of mass to the Ar nucleus. The orientation of NH\(_3\) is given by the Euler angles \((\gamma,\delta,\varphi)\), where \(\gamma\) and \(\delta\) are the usual spherical polar angles of the symmetry axis of NH\(_3\) with respect to the space fixed frame and \(\varphi\) is the third Euler angle describing a rotation of the symmetric top around its symmetry axis. In the geometry \(\gamma = \delta = \varphi = 0\), the nitrogen is on the posi-
tive z-axis and one of the protons is in the xz plane with a positive x-component. The inversion coordinate ρ is defined as the angle between the C axis and one of the N–H bonds.

The rotation–inversion scattering Hamiltonian can be written as

$$\hat{H} = \hat{H}_{\text{umb}}(\rho) + \hat{H}_{\text{vdw}}(\gamma, \delta, \varphi, R, \beta, \alpha, \rho).$$

(1)

The Hamiltonian for the umbrella motion of the NH₃ monomer, which depends only on the internal coordinate ρ, is designated by \(\hat{H}_{\text{umb}}\). It describes both the fast umbrella vibration \(v_2\) and the slow inversion tunneling. If the threefold symmetry is retained and the N–H distance is kept fixed at \(r_0\), \(\hat{H}_{\text{umb}}\) is given by\(^9\)

$$\hat{H}_{\text{umb}} = -\frac{1}{2}\frac{\partial^2}{\partial \rho^2} I_{xx}(\rho) \sin^2 \rho + \frac{1}{2}\frac{\partial^2}{\partial \rho^2} I_{zz}(\rho) \cos^2 \rho + V_{\text{umb}}(\rho),$$

(2)

where

$$g(\rho) = I_{xx}(\rho) I_{yy}(\rho) I_{zz}(\rho) I_{pp}(\rho),$$

(3a)

$$I_{xx}(\rho) = 3m_H \rho_0^2 \left[ \frac{1}{2} \sin^2 \rho + \xi \cos^2 \rho \right],$$

(3b)

$$I_{zz}(\rho) = 3m_N \rho_0^2 \sin^2 \rho,$$

(3c)

$$I_{pp}(\rho) = 3\bar{m} \rho_0^2 \left[ \cos^2 \rho + \xi \sin^2 \rho \right],$$

(3d)

$$\xi = m_N / (3m_H + m_N).$$

(3e)

Here \(g(\rho)\) is the determinant of the metric tensor \(g = \text{diag}(I_{xx}, I_{yy}, I_{zz}, I_{pp})\) in a curvilinear coordinate system; \(m_H\) and \(m_N\) are the masses of the hydrogen and nitrogen nuclei. The quantities \(I_{xx}, I_{yy}, I_{zz}\) and \(I_{pp}\) are the moments of inertia of NH₃, which depend on the inversion coordinate \(\rho\). The generalized moment of inertia \(I_{pp}\) is associated with the umbrella motion and depends also on the inversion coordinate. The double well potential \(V_{\text{umb}}(\rho)\) is represented by a harmonic force field augmented by a Gaussian

$$V_{\text{umb}}(\rho) = \frac{1}{2}k(\rho - \frac{1}{2} \pi)^2 + a \exp \left[ -b(\rho - \frac{1}{2} \pi)^2 \right].$$

(4)

The parameters \(k, a, b\) are chosen such that the measured inversion tunneling splitting in the \(v_2\) ground state and both transitions to the \(v_2\) first excited state are reproduced to an accuracy better than 0.1%. The form of the resulting potential is shown in Fig. 1.

The associated eigenvalue problem is solved in Ref. 10 with the use of a basis of functions sin \(m\varphi\) (\(m = 1, \ldots, 100\)). Here we consider only the lowest two eigenfunctions \(\mid \psi\rangle\), also shown in Fig. 1. They describe the lower and upper inversion states that are separated by 0.8 cm⁻¹. The lowest of the two, which is designated by \(\psi = +\), is symmetric with respect to \(\rho = \pi - \rho\). The upper level, designated by \(\psi = -\), is antisymmetric with respect to this operation.

The van der Waals Hamiltonian can be written as

$$\hat{H}_{\text{vdw}}(\gamma, \delta, \varphi, R, \beta, \alpha, \rho) = B(\rho) \xi_j \frac{\partial^2}{\partial \rho^2} \xi_j^2 + \frac{1}{2} \frac{\partial^2}{\partial \rho^2} \xi_j^2 + V_{\text{vdw}}(\rho, \varphi, R, \beta, \alpha, \rho).$$

(5)

The first two terms in \(\hat{H}_{\text{vdw}}\) represent the symmetric top Hamiltonian of NH₃. The rotational constants are related to the moments of inertia given in Eq. (3), \(B(\rho) = [2I_{xx}(\rho)]^{-1}\) and \(C(\rho) = [2I_{zz}(\rho)]^{-1}\). The third and fourth term give the kinetic energy of the “diatom,” with \(\xi_j\) being the relative angular momentum. The intermolecular potential \(V_{\text{vdw}}\) is expanded in spherical harmonics \(Y_{\ell m}\),

$$V_{\text{vdw}}(R, \Theta, \Phi, \rho) = \sum_{\ell m} v_{\ell m}(R, \rho) Y_{\ell m}(\Theta, \Phi),$$

(6)

where \(\Theta\) and \(\Phi\) are the polar angles of the Ar projectile with respect to the principal axes frame of the NH₃ rotor. In the space fixed frame \(Y_{\ell m}(\Theta, \Phi)\) becomes

$$Y_{\ell m}(\Theta, \Phi) = \sum_{\ell} D_{\ell m}^{(4)}(\gamma, \delta, \varphi) Y_{4m}(\beta, \alpha),$$

(7)

where \(D_{\ell m}^{(4)}(\gamma, \delta, \varphi)\) is the usual Wigner rotation matrix.\(^{11}\) The expansion coefficients \(v_{\ell m}(R, \rho)\) have been taken from Bulski et al.,\(^1\) who calculated the \(ab\) \(initio\) potential for four different umbrella angles \(\rho\), and expanded it in tesseral harmonics. Due to the threefold symmetry of the ammonia
TABLE I. Rotational constants for the lowest two inversion states (in cm\(^{-1}\)).

<table>
<thead>
<tr>
<th>(v)</th>
<th>(B)</th>
<th>(C)</th>
<th>(B)</th>
<th>(C)</th>
</tr>
</thead>
<tbody>
<tr>
<td>+</td>
<td>10.000</td>
<td>6.337</td>
<td>9.9402</td>
<td>6.3044</td>
</tr>
</tbody>
</table>

\(^a\)Reference 10.  
\(^b\)Reference 14.

only terms with \(m=0,3,6,...\) are present. The first 15 terms with \(l\leq 7\) have been included, which leads to an accuracy of \(-0.6\%\) in the convergence of the expansion. The expansion coefficients are written as a sum of a short and a long range contribution,

\[
v_{\lambda \mu}(R,p) = v_{\lambda \mu}^{SR}(R,p) + v_{\lambda \mu}^{LR}(R,p),
\]

where

\[
v_{\lambda \mu}^{SR}(R,p) = F_{\lambda \mu}(p) [1 + \delta_{\lambda \mu}(p) R] \exp[-\alpha_{\lambda \mu}(p) R - \beta_{\lambda \mu}(p) R^2],
\]

and

\[
v_{\lambda \mu}^{LR}(R,p) = - \sum_{n=6}^{10} f_{n}^{\mu}(R,p) C_{n}^{\lambda \mu}(p) R^{-n}.
\]

The \(C_{n}^{\lambda \mu}(p)\) are the induction and dispersion coefficients, the \(f_{n}^{\mu}(R,p)\) are Tang and Toennies type damping functions. The values of all coefficients are given in Ref. 1. The rotational constants for a given tunneling state \(v\) are given by

\[
B_{\lambda} = \langle v | B(p) | v \rangle
\]

and

\[
C_{\lambda} = \langle v | C(p) | v \rangle.
\]

Their values are listed in Table I.

In the CC method\(^7\) the angular basis functions are usually formed by Clebsch–Gordan coupling of the relative angular momentum functions \(Y_{lm\ell}(\beta,\alpha) = |lm\ell\rangle\) and the symmetric top functions \(|jkme\rangle\) to a total angular momentum \(J\) with space fixed z-component \(M\). In the case of permutation-inversion symmetry \(\Pi(D_{3h})\), which we have here, it is convenient to take linear combinations \(|jkme\rangle\) of the symmetric top functions, defined by

\[
|jkme\rangle = \sum_{m,m_{1}} (jkm + \epsilon |j-km\rangle),
\]

where \(k \geq 0\) and \(\epsilon = \pm 1\), except for \(k=0\) when obviously only \(\epsilon = +1\) is allowed.

Since we take the umbrella motion explicitly into account, the basis has to be extended by taking the tensor product with the tunneling functions \(v(p)\),

\[
|jkJM\rangle = \sum_{m,m_{1}} |jkme\rangle |lm\rangle |\langle jmlm_{1}\rangle |JM\rangle,
\]

where \(|jm\rangle\) is a Clebsch–Gordan coefficient.\(^1\) From symmetry considerations it follows\(^3\) that the symmetric (antisymmetric) inversion function can combine to a state adapted to \(\Pi(D_{3h})\) with only one of the two \(|jkme\rangle\) functions, so that \(\epsilon = \pm (-1)^{l}\) for \(v = |\pm\rangle\). We can therefore omit the quantum number \(v\) and label the basis functions by \(|jkJM\rangle\), instead of using the labeling given in Eq. (12). In the CC equations the noninteracting blocks for different \(J\) are separated into two parity blocks, each containing channels \(|jk\rangle\) having different values of \(\epsilon\) \((-1)^{l} / k^{+} v^{\prime}\). States of the free NH\(_3\) can be designated by \(j_{k}^{\prime}\), thereby uniquely specifying the inversion function.

The coupling between the channels that originates from the potential matrix elements is given by

\[
\langle jkJM\rangle = \sum_{\lambda \mu} c_{\lambda \mu}^{\lambda \mu}(R,p) Y_{\lambda \mu}(\Theta,\Phi) |jk'JM\rangle.
\]

The integration over all the relevant angles \(\alpha,\beta,\gamma,\delta,\rho\) is performed, after substitution of Eq. (7), by angular momentum techniques, as is usual in the CC method.\(^7\) In order to calculate the matrix elements between the two tunneling states \(v(p)\) we obtained an analytical representation of the \(\rho\)-dependence of the expansion coefficients \(v_{\lambda \mu}(R,p)\) by fitting, for each value of \(R\), a fifth order polynomial in \(\rho - \frac{1}{2} \pi\) through the \(ab\) initio values of the coefficients for different \(\rho\). The fit contains only even or only odd powers depending on whether \(\lambda + \mu\) is even or odd. The matrix elements can then be evaluated as follows:

\[
\begin{align*}
\langle v | v_{\lambda \mu}(R,p) | v' \rangle &= \sum_{n=0}^{5} c_{n}^{\lambda \mu}(R) \langle v | (\rho - \frac{1}{2} \pi)^{n} | v' \rangle,
\end{align*}
\]

where the \(c_{n}^{\lambda \mu}(R)\) are the polynomial expansion coefficients. The above relationship between the values of \(\epsilon\) and \(v\) is used to insert the correct \((v,v')\) combination into Eq. (13).

According to Ref. 12 the \(ab\) initio potential has to be scaled to give good agreement with spectroscopic data for the bound Ar–NH\(_3\) complex. This scaling consists of multiplying the short range parameter \(F_{33}\) in Eq. (9a) by a factor of 1.43. Here, calculations have been performed using both the original \(ab\) initio potential and a modified potential in which the same scaling was applied for all values of the inversion coordinate \(\rho\).

In addition to the calculation with the inversion averaged matrix elements in the way we have just described (henceforth referred to as the "exact" inversion method), we have used the model developed by Davis and Boggs\(^2\) and Green.\(^3\) In this model the inversion functions are taken to be delta functions, \(|\pm\rangle = [\delta(\rho - \rho_{e}) \pm \delta(\rho - \pi - \rho_{e})] / \sqrt{2}\), where \(\rho_{e}\) is the value of the inversion coordinate in the equilibrium configuration. In this case the intermolecular potential needs to be known only for the equilibrium angle, since the expansion coefficients, averaged over the inversion functions, are now given by\(^2,3\)

\[
\begin{align*}
\langle \pm | v_{\lambda \mu}(R,p) | \pm \rangle &= v_{\lambda \mu}(R,\rho_{e}), & \text{for } \lambda + \mu & \text{even}, \\
&= 0, & \text{for } \lambda + \mu & \text{odd}, \\
\langle \pm | v_{\lambda \mu}(R,p) | \mp \rangle &= 0, & \text{for } \lambda + \mu & \text{even}, \\
&= v_{\lambda \mu}(R,\rho_{e}), & \text{for } \lambda + \mu & \text{odd}.
\end{align*}
\]
Using this model for the inversion functions, together with
the neglect of the inversion splitting, the scattering equa-
tions for para-\(\text{NH}_3\) are invariant to a simultaneous change
of parity in the incoming and outgoing channels, i.e.,
\[
\sigma(j_k^e \rightarrow j_k'^e) = \sigma(j_k^e \rightarrow j_k'^e).
\]

In the CS method\(^7\) the scattering equations are ex-
pressed in a body-fixed coordinate system. The fourth term
in the van der Waals Hamiltonian, Eq. (5), is approxi-
mated by putting \(\vec{J} = J_1 + J_2 - 2 j_0 \vec{J}\) equal to \(\vec{J}\). This
implies that the Coriolis interactions are neglected and that
\(\Omega\), the projection of both \(J\) and \(\vec{J}\) on the vector \(R\), is a good
quantum number, i.e., there is no coupling between chan-
nels with different \(\Omega\). The molecular symmetry group of
the dimer is thus enlarged from \(\Pi(D_{3h})\) to the semidirect
product of \(C_\infty\) with \(\Pi(D_{3h})\).

III. COMPUTATIONAL ASPECTS

The calculations were carried out with the HIBRIDON
inelastic scattering code.\(^1\) The total collision energy \(E\), the
maximum value of the total angular momentum \(J\) and the
values of \(j\) and \(k\) at which the rotational basis set is trun-
cated, are input parameters of the program. The values for
\(l\) are then given by triangular inequalities [cf. Eq. (12)].
The program has the possibility of further reducing the size
of the basis set as the overall rotation takes up more and
more of the available energy. So, from a chosen value of \(J\)
onwards, the program includes only open channels. To
keep the calculations feasible even at higher energies, an
interpolation scheme for the total cross sections as a func-
tion of \(J\) can be used, leading to a substantial reduction of
the required cpu time.

The values of the total energies are determined by the
two relative kinetic energies in the experiment, 280 and 485 cm\(^{-1}\). The ortho-\(\text{NH}_3\) with initial state \(J = k = 0\)
has zero internal energy, so the total energies are equal to
the relative kinetic energy. The initial \(j = k = 1\) state of
para-\(\text{NH}_3\), which is the ground state of this species, has an
internal energy of 16.245 cm\(^{-1}\). The total energies are con-
sequently set equal to 296.245 and 501.245 cm\(^{-1}\). The mo-
lecular levels in the basis set are retained up to \(j = 9\) inclu-
sive, with all allowed values of \(k\). This means that for
ortho-\(\text{NH}_3\), 34 levels are included (with a maximum energy of
895 cm\(^{-1}\)), 11 of which are asymptotically accessible in
the lower, and 19 in the higher energy case. Out of the 66
levels for para-\(\text{NH}_3\) (with a maximum energy of 891
\(\text{cm}^{-1}\)), 24 and 40 levels are accessible, respectively. The \(J\)
value at which we start to neglect closed channels is 78.
The interpolation step size \(\Delta J\) is taken to be six, so that
calculations are actually performed for \(J = 0, 6, 12, ..., 150\).
As explained in Sec. II, the \(\text{NH}_3\) inversion is taken into
account by calculation of the matrix elements according to
the "exact" inversion method, given in Eq. (14), or ac-
cording to the delta function model, given in Eq. (15).

Convergence with respect to relevant parameters in the
propagator, such as the step size \(\Delta R\), was better than 1%
Table II shows the dependence of the cross sections on
the magnitude of the rotational basis set. Going from a maxi-

\begin{table}
\centering
\begin{tabular}{|c|c|c|c|c|}
\hline
\(J\) & \(\Delta J = 1\) & \(\Delta J = 4\) & \(\Delta J = 6\) & \(\Delta J = 8\) \\
\hline
1\(^2\) & 10.83 & 10.93 & 11.35 & 9.06 \\
2\(^2\) & 9.02 & 9.00 & 8.74 & 9.63 \\
3\(^2\) & 4.43 & 4.32 & 4.34 & 4.29 \\
4\(^2\) & 0.42 & 0.40 & 0.39 & 0.38 \\
5\(^2\) & 1.64 & 1.63 & 1.62 & 1.59 \\
6\(^2\) & 9.26 & 9.25 & 9.33 & 8.42 \\
7\(^2\) & 3.21 & 3.33 & 3.25 & 3.34 \\
8\(^2\) & 0.25 & 0.22 & 0.25 & 0.24 \\
9\(^2\) & 0.06 & 0.05 & 0.05 & 0.05 \\
10\(^2\) & 0.34 & 0.36 & 0.36 & 0.34 \\
\hline
\end{tabular}
\caption{Effect of the interpolation step size \(\Delta J\) on the cross sections \(\sigma(0^+ \rightarrow j_0)\) (in \(\text{Å}^2\)) for ortho-\(\text{NH}_3\)-Ar at an energy of 280 cm\(^{-1}\) (delta function model inversion, \textit{ab initio} potential).
}
\end{table}

mum \(J\) value of 9 to a maximum value of 11 in the rota-
tional basis set, induced changes in the cross sections of
\(\sim 6\%\). The neglect of closed channels for \(J > 78\) did not
affect the results. The effect of interpolation step size \(\Delta J\)
on the cross sections is shown in Table III. The error due to
the step size used here is found to be \(< 5\%\).

Since the averaging of the rotational constants over the
inversion wave functions has a small effect (cf. Table I),
we have taken the same value for both inversion states,\(^1\)
\(B = 9.9402\) cm\(^{-1}\) and \(C = 6.3044\) cm\(^{-1}\). The maximum
number of channels used in the calculation was 219 per
parity block for ortho-\(\text{NH}_3\), taking \(\sim 24\) cpu hours for a
full calculation, and 441 per parity block for para-\(\text{NH}_3\),
taking \(\sim 241\) cpu hours on an IBM RS/6000 model 320
workstation.

In the CS calculation we used only the "exact" invo-
version method. The value of \(\Omega\) ranged from 0 to 7 for ortho
and from 0 to 8 for para; the maximum \(J\) value in the
rotational basis set was 9 and \(J\) was varied from 0 to 150 at
an energy of 485 and 501.245 cm\(^{-1}\), respectively. The in-

\begin{table}
\centering
\begin{tabular}{|c|c|c|c|}
\hline
\(J\) & \(\Delta J = 1\) & \(\Delta J = 4\) & \(\Delta J = 6\) \\
\hline
1\(^2\) & 9.67 & 10.10 & 10.50 \\
2\(^2\) & 3.94 & 3.84 & 3.87 \\
3\(^2\) & 0.46 & 0.52 & 0.54 \\
4\(^2\) & 0.13 & 0.13 & 0.13 \\
5\(^2\) & 0.07 & 0.07 & 0.07 \\
6\(^2\) & 1.32 & 1.25 & 1.25 \\
7\(^2\) & 11.70 & 11.96 & 11.98 \\
8\(^2\) & 2.52 & 2.27 & 2.26 \\
9\(^2\) & 0.65 & 0.72 & 0.72 \\
10\(^2\) & 0.10 & 0.10 & 0.10 \\
11\(^2\) & 0.63 & 0.60 & 0.61 \\
12\(^2\) & 0.06 & 0.07 & 0.07 \\
13\(^2\) & 0.01 & 0.02 & 0.02 \\
14\(^2\) & 0.04 & 0.04 & 0.04 \\
15\(^2\) & 0.41 & 0.42 & 0.42 \\
16\(^2\) & 0.11 & 0.13 & 0.13 \\
17\(^2\) & 0.01 & 0.01 & 0.01 \\
\hline
\end{tabular}
\caption{Effect of the interpolation step size \(\Delta J\) on the cross sections \(\sigma(0^+ \rightarrow f_0)\) (in \(\text{Å}^2\)) for ortho-\(\text{NH}_3\)-Ar at an energy of 485
\(\text{cm}^{-1}\) ("exact" inversion, \textit{ab initio} potential).
}\end{table}
TABLE IV. State-to-state cross sections \( \sigma(0_{0_0}^+ \rightarrow f_j^i) \) for ortho-NH\(_3\)–Ar in \( \AA^2 \) at an energy of 280 cm\(^{-1}\). The cross sections given in parentheses are corrected for the incomplete initial state preparation in the measurement, as follows [cf. Eq. (17)]: 
\[
\sigma(-\rightarrow f_j^i) = 0.92\sigma(0_{0_0}^+ \rightarrow f_j^i) + 0.08\sigma(1_{0_1}^{-} \rightarrow f_j^i), \quad \text{for } f_j^i \neq 1_{0_1}^{-}.
\]
A dash (—) in the last column indicates that the corresponding cross section has not been measured.

<table>
<thead>
<tr>
<th>( \Lambda )</th>
<th>( V_{i_0}^b )</th>
<th>( V_{i_0}^b )</th>
<th>( V_{ii}^b )</th>
<th>( V_{iv}^b )</th>
<th>Expt</th>
</tr>
</thead>
<tbody>
<tr>
<td>1^+</td>
<td>11.35 (7.64)</td>
<td>4.03 (1.09)</td>
<td>11.06 (7.41)</td>
<td>3.82 (0.87)</td>
<td>6.07</td>
</tr>
<tr>
<td>2^+</td>
<td>8.74 (8.69)</td>
<td>13.35 (12.65)</td>
<td>8.45 (8.42)</td>
<td>13.20 (12.49)</td>
<td>6.73</td>
</tr>
<tr>
<td>3^+</td>
<td>4.34 (4.29)</td>
<td>2.56 (2.98)</td>
<td>4.16 (4.12)</td>
<td>2.51 (2.94)</td>
<td>2.19</td>
</tr>
<tr>
<td>4^+</td>
<td>0.28 (0.07)</td>
<td>0.35 (0.35)</td>
<td>0.31 (0.31)</td>
<td>0.41 (0.42)</td>
<td>3.43</td>
</tr>
<tr>
<td>5^+</td>
<td>3.32 (3.19)</td>
<td>1.94 (1.97)</td>
<td>3.38 (3.29)</td>
<td>1.87 (1.91)</td>
<td>5.77</td>
</tr>
<tr>
<td>6^+</td>
<td>0.05 (0.06)</td>
<td>0.39 (0.40)</td>
<td>0.06 (0.07)</td>
<td>0.41 (0.42)</td>
<td>—</td>
</tr>
<tr>
<td>7^+</td>
<td>0.36 (0.36)</td>
<td>0.28 (0.28)</td>
<td>0.34 (0.35)</td>
<td>0.31 (0.31)</td>
<td>—</td>
</tr>
</tbody>
</table>

\( V_{i_0}^b \), delta function model inversion, \( ab\ initio \) potential.
\( V_{ii}^b \), delta function model inversion, scaled potential.
\( V_{iv}^b \), “exact” inversion \( ab\ initio \) potential.
\( V_{iv}^b \), “exact” inversion, scaled potential.

Both for ortho and para-NH\(_3\), the scaling in the potential has a large effect. In the case of ortho-NH\(_3\), this effect is about the same for the various cross sections at both energies. Especially transitions to 1\(_{0_1}^+\), 2\(_{0_1}^+\), 3\(_{0_1}^+\), and 4\(_{0_1}^+\) are strongly affected in the lower energy case, and transitions to 1\(_{0_1}^+\), 2\(_{0_1}^+\), and 3\(_{0_1}^+\) in the higher energy case. Use of the delta function model for inversion does not affect the influence of the scaling. For para-NH\(_3\), the scaling in the potential decreases some of the cross sections at the lower energy and increases them at the higher energy, for other cross sections it is vice versa. For the lower energy the scaling reduces the size of most of the para cross sections, except for the 2\(_{1_0}^+\), 3\(_{1_0}^+\), and 4\(_{1_0}^+\) states. The scaling induces large changes in the relative magnitudes for the \( \pm \) inversion states for transitions to the 2\(_{1_0}^+\), 3\(_{1_0}^+\), and 4\(_{1_0}^+\) states.

Comparison with the experiment shows that the calculations using the original \( ab\ initio \) potential give a better overall agreement than calculations using the scaled potential. Particularly, cross sections to the 1\(_{0_1}^+\), 2\(_{0_1}^+\), and 2\(_{1_0}^+\) states come out better. In a few cases, however, the cross sections from the scaled potential are closer to the experimental ones.

It has been debated whether it is necessary to include higher anisotropic terms, such as a \( v_{33} \) term,\(^1\) in the description of the intermolecular potential, since the observed far-infrared Ar–NH\(_3\) spectrum could also be explained with an effective angular potential that contains only terms up to \( v_{20} \).\(^1\) When we look at the results of the scattering calculations, we observe, for example, that in the case of ortho-NH\(_3\) the experimental cross sections \( \sigma(0_{0_0}^+ \rightarrow 3\_3^+) \) are reproduced fairly well. In the first Born approximation these transitions are solely due to the \( v_{33} \) term in the potential. It seems unlikely that the agreement between experiment and theory could be maintained if this important first order contribution were zero, as it would be when \( v_{33} \) would vanish.

Inspection of the influence of the scaling in \( v_{33} \) on the cross sections of ortho-NH\(_3\) shows that not only \( \Delta j = \Delta k = 3 \) transitions are affected, but other transitions as well.
TABLE V. State-to-state cross sections $\sigma(0\bar{4}^+ \rightarrow f\bar{v})$ for ortho-NH$_3$-Ar in Å$^2$ at an energy of 485 cm$^{-1}$. The corrected cross sections given in parentheses are obtained as indicated in Table IV. A dash (—) in the last column indicates that the corresponding cross section has not been measured.

<table>
<thead>
<tr>
<th>$f\bar{v}$</th>
<th>$V_{III}^a$</th>
<th>$V_{IV}^b$</th>
<th>$V_{III}^c$</th>
<th>$V_{IV}^d$</th>
<th>Expt.</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\Delta^2$</td>
<td>7.90(4.30)</td>
<td>5.94(2.40)</td>
<td>7.92(4.33)</td>
<td>5.53(2.04)</td>
<td>4.33</td>
</tr>
<tr>
<td>$\Delta^1$</td>
<td>10.20(9.85)</td>
<td>10.73(10.17)</td>
<td>10.10(9.76)</td>
<td>11.02(10.42)</td>
<td>5.72</td>
</tr>
<tr>
<td>$\Delta^3$</td>
<td>4.07(4.07)</td>
<td>1.84(2.18)</td>
<td>3.94(3.94)</td>
<td>1.77(2.14)</td>
<td>2.38</td>
</tr>
<tr>
<td>$\Delta^4$</td>
<td>0.50(0.63)</td>
<td>3.00(2.87)</td>
<td>0.46(0.59)</td>
<td>3.04(2.91)</td>
<td>1.73</td>
</tr>
<tr>
<td>$\Delta^5$</td>
<td>0.13(0.19)</td>
<td>1.07(1.06)</td>
<td>0.13(0.18)</td>
<td>1.06(1.04)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^6$</td>
<td>0.07(0.07)</td>
<td>0.13(0.14)</td>
<td>0.07(0.07)</td>
<td>0.12(0.13)</td>
<td>1.01</td>
</tr>
<tr>
<td>$\Delta^7$</td>
<td>1.39(2.11)</td>
<td>0.88(1.96)</td>
<td>1.32(2.07)</td>
<td>0.83(1.92)</td>
<td>1.68</td>
</tr>
<tr>
<td>$\Delta^8$</td>
<td>12.14(11.35)</td>
<td>12.83(11.96)</td>
<td>11.70(10.93)</td>
<td>13.12(12.21)</td>
<td>14.23</td>
</tr>
<tr>
<td>$\Delta^9$</td>
<td>2.47(2.56)</td>
<td>2.14(2.21)</td>
<td>2.52(2.59)</td>
<td>2.07(2.15)</td>
<td>4.90</td>
</tr>
<tr>
<td>$\Delta^{10}$</td>
<td>0.66(0.76)</td>
<td>0.87(0.92)</td>
<td>0.65(0.75)</td>
<td>0.82(0.87)</td>
<td>2.70</td>
</tr>
<tr>
<td>$\Delta^{11}$</td>
<td>0.11(0.13)</td>
<td>0.25(0.28)</td>
<td>0.10(0.12)</td>
<td>0.22(0.25)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^{12}$</td>
<td>0.66(0.71)</td>
<td>1.18(1.16)</td>
<td>0.63(0.69)</td>
<td>1.17(1.15)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^{13}$</td>
<td>0.06(0.09)</td>
<td>0.27(0.26)</td>
<td>0.06(0.09)</td>
<td>0.26(0.26)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^{14}$</td>
<td>0.01(0.02)</td>
<td>0.04(0.05)</td>
<td>0.01(0.02)</td>
<td>0.04(0.05)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^{15}$</td>
<td>0.04(0.07)</td>
<td>0.06(0.11)</td>
<td>0.04(0.06)</td>
<td>0.06(0.11)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^{16}$</td>
<td>0.38(0.39)</td>
<td>0.91(0.87)</td>
<td>0.41(0.43)</td>
<td>0.09(0.11)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^{17}$</td>
<td>0.12(0.12)</td>
<td>0.12(0.11)</td>
<td>0.11(0.12)</td>
<td>0.11(0.11)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^{18}$</td>
<td>0.01(0.01)</td>
<td>0.02(0.02)</td>
<td>0.01(0.01)</td>
<td>0.02(0.02)</td>
<td>—</td>
</tr>
</tbody>
</table>

$^a$ $V_{III}$, delta function model inversion, ab initio potential.
$^b$ $V_{IV}$, delta function model inversion, scaled potential.
$^c$ $V_{III}$, "exact" inversion, ab initio potential.
$^d$ $V_{IV}$, "exact" inversion, scaled potential.

This would not be so in the first Born approximation. It means that the contributions from higher Born approximations are significant. The importance of higher order effects in the interaction between Ar and NH$_3$ is confirmed by calculations on bound states. Although the $\Delta^3$ term does not contribute to the lower bound states in a first order perturbation theory, which has the isotropic Hamiltonian as its zeroth order Hamiltonian, it proves to be one of the dominant terms in determining the rovibrational energy levels of the Ar-NH$_3$ complex.$^{12}$

At the energies used here, application of the delta function model for the inversion of NH$_3$ has only a small effect on the cross sections, as compared to the "exact" calculations, in the order of 3%. The parity propensities, which

TABLE VI. State-to-state cross sections $\sigma(1\bar{1}^0 \rightarrow f\bar{v})$ for para-NH$_3$-Ar in Å$^2$. The relative kinetic energies are as indicated. The cross sections given in parentheses are corrected for the incomplete initial state preparation in the measurement, as follows [cf. Eq. (17)]: $\sigma(\bar{f}\bar{v}) = 0.95\sigma(1\bar{1}^0) + 0.05\sigma(1\bar{1}^0 - \bar{f}\bar{v})$. A dash (—) in the last column indicates that the corresponding cross section has not been measured.

<table>
<thead>
<tr>
<th>$f\bar{v}$</th>
<th>$V_{III}^a$</th>
<th>$V_{IV}^b$</th>
<th>Expt.</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\Delta^2$</td>
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<td>2.82(3.21)</td>
<td>4.70</td>
</tr>
<tr>
<td>$\Delta^3$</td>
<td>4.40(4.39)</td>
<td>10.73(10.33)</td>
<td>6.32</td>
</tr>
<tr>
<td>$\Delta^4$</td>
<td>3.00(2.06)</td>
<td>1.57(1.66)</td>
<td>0.89</td>
</tr>
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<td>$\Delta^5$</td>
<td>3.35(3.28)</td>
<td>3.27(3.18)</td>
<td>1.35</td>
</tr>
<tr>
<td>$\Delta^6$</td>
<td>0.57(0.56)</td>
<td>1.42(1.37)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^7$</td>
<td>0.33(0.34)</td>
<td>0.50(0.55)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^8$</td>
<td>0.99(1.61)</td>
<td>0.47(1.11)</td>
<td>1.19</td>
</tr>
<tr>
<td>$\Delta^9$</td>
<td>13.48(12.85)</td>
<td>13.30(12.65)</td>
<td>12.94</td>
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</table>

<table>
<thead>
<tr>
<th>$f\bar{v}$</th>
<th>$V_{III}^a$</th>
<th>$V_{IV}^b$</th>
<th>Expt.</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\Delta^2$</td>
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<td>3.06(2.96)</td>
<td>3.24</td>
</tr>
<tr>
<td>$\Delta^3$</td>
<td>3.90(3.84)</td>
<td>1.05(1.15)</td>
<td>2.56</td>
</tr>
<tr>
<td>$\Delta^4$</td>
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<td>0.97(0.96)</td>
<td>0.97</td>
</tr>
<tr>
<td>$\Delta^5$</td>
<td>0.59(0.58)</td>
<td>0.75(0.76)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^6$</td>
<td>0.12(0.12)</td>
<td>0.22(0.22)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^7$</td>
<td>0.03(0.04)</td>
<td>0.09(0.10)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^8$</td>
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<td>0.80</td>
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<td>$\Delta^9$</td>
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<td>2.15(2.07)</td>
<td>2.43</td>
</tr>
<tr>
<td>$\Delta^{10}$</td>
<td>0.39(0.39)</td>
<td>0.39(0.38)</td>
<td>0.44</td>
</tr>
<tr>
<td>$\Delta^{11}$</td>
<td>0.42(0.42)</td>
<td>0.22(0.23)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^{12}$</td>
<td>0.22(0.24)</td>
<td>0.25(0.27)</td>
<td>—</td>
</tr>
<tr>
<td>$\Delta^{13}$</td>
<td>0.54(0.52)</td>
<td>0.67(0.65)</td>
<td>—</td>
</tr>
</tbody>
</table>

$^a$ $V_{III}$ "exact" inversion, ab initio potential.
$^b$ $V_{IV}$ "exact" inversion, scaled potential.

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are defined as the ratios between the difference and the sum of the cross sections to the ± inversion states, are hardly affected. The suggestion that deviations of theoretical propensities from experimental ones in He-NH₃ can be accounted for by taking the inversion motion explicitly into account is hereby contradicted to all likelihood.

Our calculations show further that the invariance to a simultaneous change of parity in the incoming and outgoing channels [cf. Eq. (16)] for para-NH₃ holds to an accuracy of ~0.2% when the inversion is included in the potential matrix elements, but the inversion splitting of 0.8 cm⁻¹ is neglected. When this splitting is included also, the deviations are of the order of 3%. Billing has found in He-NH₃ calculations at an energy of 65 meV (525 cm⁻¹), using a semiclassical approach, that the inversion is obeyed to within 5% when the splitting of the inversion doublet is taken into account. The effect of including the tunneling splitting on the invariance is therefore larger than the effect of the “exact” calculation of potential matrix elements over the inversion wave functions.

The limited influence of the delta function model for inversion seems surprising at first sight, because the wave functions in Fig. 1 do not resemble delta functions. It means, however, as is noted by Davis and Boggs, that the are functions of that vary slowly enough for the delta function approach to be valid.

Application of the CS approximation gives a reasonable agreement for para-NH₃ (Table VII). For ortho-NH₃, however, there are strong deviations, both from the CC method and from experiment. The CS calculation gives zero cross sections to the 3f and 4f states, whereas both experimental and CC cross sections to these states are different from zero. This means that these transitions are caused by the Coriolis terms in the Hamiltonian. Even though these terms are small, their long range apparently gives rise to significant transition probabilities.

In calculations on He-NH₃ scattering at an energy of 98 meV (792 cm⁻¹) using the CS method Meyer et al. have found the cross sections to the 3f and 4f states to be exactly zero and the cross section to the 2f state to be

FIG. 2. Experimental and theoretical cross sections for ortho-NH₃ at a relative kinetic energy of 280 cm⁻¹. The theoretical values are given in parentheses in Table IV, for the 1f state we have used the expression given in that table.

FIG. 3. Experimental and theoretical cross sections for ortho-NH₃ at a relative kinetic energy of 485 cm⁻¹. The theoretical values are given in parentheses in Table V, for the 1f state we have used the expression given in Table IV.
V. CONCLUSIONS

We have calculated close coupling state-to-state cross sections for the inelastic scattering of \( \text{NH}_3 \) with \( \text{Ar} \) at two different collision energies and compared them with experimentally derived cross sections. The inversion motion of \( \text{NH}_3 \) has been taken into account explicitly. Comparison with calculations that use a delta function model description of the inversion motion, shows that this model leads to errors of 3% only, at the energies used here. Previously found deviations from experimentally determined parity propensities for \( \text{NH}_3-\text{He} \), cannot be attributed to use of the delta function model. It is more likely that these discrepancies arise from shortcomings of the intermolecular potential used for that system.

In calculations on bound states of \( \text{Ar}-\text{NH}_3 \), the use of a potential in which a single term in the angular expansion of the \textit{ab initio} potential of Ref. 1 was scaled by a factor of 1.43, gave better agreement with spectroscopic data than the use of the original \textit{ab initio} potential. In the present scattering calculations the opposite is true. This can be seen as a manifestation of the fact that scattering and bound states probe different regions of the intermolecular potential surface. The applied scaling is too crude to obtain a fully realistic potential surface. For ortho-\( \text{NH}_3 \) the calculated \textit{ab initio} cross sections reproduce the experimental ones fairly well. For para-\( \text{NH}_3 \) the overall agreement is good too, but differences remain between theoretical and experimental parity propensities, indicating that the \textit{ab initio} potential needs further improvement.

Comparison of the results for the two potentials shows further that the cross sections are very sensitive to variations in the potential surface. The changes in the cross sections for transitions to the various rotation–inversion states induced by the scaling, show also that higher order effects play a role in the scattering process, just as they do in bound state interactions.

Application of the coupled states approximation leads...
### Table VII. Results of CC and CS calculations ("exact" inversion, \textit{ab initio} potential, without correction for initial state preparation, at a relative kinetic energy of 485 cm\(^{-1}\)) and experimental results. The cross sections are in \(\text{Å}^2\). A dash (—) in the last column indicates that the corresponding cross section has not been measured.

<table>
<thead>
<tr>
<th>(J)</th>
<th>(\alpha)</th>
<th>CC</th>
<th>CS</th>
<th>Expt.</th>
<th>(\beta)</th>
<th>CC</th>
<th>CS</th>
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</thead>
<tbody>
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<td>1(^{0}_{\sigma})</td>
<td>7.92</td>
<td>9.24</td>
<td>4.33</td>
<td></td>
<td>2(^{1}_{\sigma})</td>
<td>3.62</td>
<td>3.30</td>
<td>5.51</td>
</tr>
<tr>
<td>2(^{0}_{\sigma})</td>
<td>10.10</td>
<td>8.17</td>
<td>5.72</td>
<td></td>
<td>2(^{1}_{\sigma})</td>
<td>4.39</td>
<td>3.39</td>
<td>3.76</td>
</tr>
<tr>
<td>3(^{0}_{\sigma})</td>
<td>3.94</td>
<td>5.17</td>
<td>2.38</td>
<td></td>
<td>3(^{1}_{\sigma})</td>
<td>1.65</td>
<td>1.31</td>
<td>1.35</td>
</tr>
<tr>
<td>4(^{0}_{\sigma})</td>
<td>0.46</td>
<td>0.71</td>
<td>1.73</td>
<td></td>
<td>3(^{1}_{\sigma})</td>
<td>3.24</td>
<td>3.35</td>
<td>1.37</td>
</tr>
<tr>
<td>5(^{0}_{\sigma})</td>
<td>0.13</td>
<td>0.13</td>
<td>—</td>
<td></td>
<td>4(^{1}_{\sigma})</td>
<td>1.10</td>
<td>0.91</td>
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<tr>
<td>6(^{0}_{\sigma})</td>
<td>0.07</td>
<td>0.08</td>
<td>1.01</td>
<td></td>
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<td>0.27</td>
<td>0.28</td>
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<td>3(^{1}_{\sigma})</td>
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<td>0</td>
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<td>4.26</td>
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<tr>
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<td>0</td>
<td>2.70</td>
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<td>5(^{4}_{\sigma})</td>
<td>3.94</td>
<td>4.93</td>
<td>1.95</td>
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<tr>
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<td>—</td>
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<td>0.06</td>
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<td>0</td>
<td>—</td>
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<td>1.25</td>
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<tr>
<td>6(^{3}_{\sigma})</td>
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<td>—</td>
<td></td>
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<td>0.49</td>
<td>0.58</td>
<td>0.75</td>
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<tr>
<td>7(^{0}_{\sigma})</td>
<td>0.11</td>
<td>0.12</td>
<td>—</td>
<td></td>
<td>5(^{9}_{\sigma})</td>
<td>0.40</td>
<td>0.23</td>
<td>—</td>
</tr>
<tr>
<td>7(^{1}_{\sigma})</td>
<td>0.01</td>
<td>0</td>
<td>—</td>
<td></td>
<td>5(^{10}_{\sigma})</td>
<td>0.94</td>
<td>0.99</td>
<td>—</td>
</tr>
</tbody>
</table>

...to artificial selection rules for the cross sections from the \(0^{0}_{\sigma}\) state. Some of the cross sections, which should vanish according to these rules, are not even small, in the CC calculations, nor in the experiment. Consequently, the overall agreement with close coupling results is poor for ortho-NH\(_3\), reasonable for para-NH\(_3\).

### Acknowledgments

We are grateful to Professor Millard Alexander and Dr. Pierre Valiron for making available the HIBRIDON computer program and to Dr. Pierre Valiron also for valuable discussions. This work is part of the research program of the "Stichting voor Fundamenteel Onderzoek der Materie (FOM)," which is financially supported by the "Netherlands Organization for Scientific Research (NWO)."

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