I. INTRODUCTION

The last few years it has become possible to obtain state-to-state cross sections for transitions between rotation–inversion \( (j^I\gamma) \) states of \( \text{NH}_3 \), induced by collisions with various perturbers. Advances in molecular beam techniques and in laser spectroscopy have made it feasible to discriminate between the symmetric \((+\)) and antisymmetric \((-\)) states of the inversion doublets. The close coupling (CC) method for the accurate quantum mechanical treatment of the problem has been well established theoretically for quite some time. However, computer systems that meet the computational demands have only recently become available.

In this paper we consider collisions of \( \text{NH}_3 \) with \( \text{Ar} \). This study was undertaken mainly for two reasons. First, comparison of theoretical results with experimental data enables us to determine the accuracy of an \textit{ab initio} intermolecular potential energy surface,\(^1\) in the region that is probed in scattering experiments. In addition, we used a slightly different potential in order to gain some understanding of the sensitivity of the cross sections to variations in the potential surface. This second potential contains a scaling parameter that was chosen to account for spectroscopic data regarding the bound states of \( \text{Ar}–\text{NH}_3 \). In particular, we wanted to see whether a variation that improved bound state results would also improve the outcome of the scattering calculations.

Second, we investigate how the description of the umbrella inversion of \( \text{NH}_3 \) influences the cross sections. Until now, this inversion has been included in scattering calculations on \( \text{NH}_3 \) only via a model in which the inversion-tunneling wave function is a linear combination of two delta functions centered at the equilibrium positions \( \text{Davis} \text{ and } \text{Boggs,}^2 \text{ Green}^3 \). Although the model has its justification in the fact that the period for inversion is much longer than the duration of a collision, it was not clear whether experimentally found deviations from predicted propensity rules could not be attributed to the neglect of the inversion motion in the description of the intermolecular potential.\(^4,5\) Here we take the inversion degree of freedom explicitly into account, in order to assess how severe an approximation is made in neglecting it.

Finally, we have performed some calculations using the much cheaper coupled states (CS) approximation, to find out how this approximation affects the calculated cross sections. In calculations on \( \text{He}–\text{NH}_3 \) using the CS method\(^6\) certain theoretical cross sections are found to vanish or almost vanish, whereas the experimental cross sections are significantly different from zero.\(^4\) By applying the CS approximation, together with the full CC method on the same \( \text{Ar}–\text{NH}_3 \) potential surface, we can establish to what extent deviations are caused by the theoretical scattering method.

II. THEORY

The coordinate system used in the CC method is the space-fixed frame.\(^7\) The vector \( \mathbf{R} \), with polar angles \((\beta, \alpha)\) in this frame, points from the \( \text{NH}_3 \) center of mass to the \( \text{Ar} \) nucleus. The orientation of \( \text{NH}_3 \) is given by the Euler angles \((\gamma, \delta, \varphi)\), where \( \gamma \) and \( \delta \) are the usual spherical polar angles of the symmetry axis of \( \text{NH}_3 \) with respect to the space fixed frame and \( \varphi \) is the third Euler angle describing a rotation of the symmetric top around its symmetry axis. In the geometry \( \gamma = \delta = \varphi = 0 \), the nitrogen is on the posi-
tive z-axis and one of the protons is in the xz plane with a positive x-component. The inversion coordinate \( \rho \) is defined as the angle between the \( C_3 \) axis and one of the N-H bonds.

The rotation–inversion scattering Hamiltonian can be written as

\[
\hat{H} = \hat{H}_{\text{umb}}(\rho) + \hat{H}_{\text{vdw}}(\gamma, \beta, \varphi, R, \alpha, \rho).
\]

The Hamiltonian for the umbrella motion of the \( \text{NH}_3 \) monomer, which depends only on the internal coordinate \( \rho \), is designated by \( \hat{H}_{\text{umb}} \). It describes both the fast umbrella vibration (\( \nu_2 \)) and the slow inversion tunneling. If the threefold symmetry is retained and the N–H distance is kept fixed at \( r_0 \), \( \hat{H}_{\text{umb}} \) is given by

\[
\hat{H}_{\text{umb}} = -\frac{1}{2} \frac{\partial^2}{\partial \rho^2} I_{\rho \rho}(\rho) + V_{\text{umb}}(\rho),
\]

where

\[
g(\rho) = I_{xx}(\rho) I_{yy}(\rho) I_{zz}(\rho) I_{pp}(\rho),
\]

\[
I_{xx}(\rho) = 3 m_H \rho_0^2 \left( \frac{1}{3} \sin^2 \rho + \frac{1}{2} \xi \cos^2 \rho \right),
\]

\[
I_{yy}(\rho) = I_{xx}(\rho),
\]

\[
I_{zz}(\rho) = 3 m_H \rho_0^2 \sin^2 \rho,
\]

\[
I_{pp}(\rho) = 3 m_H \rho_0^2 (\cos^2 \rho + \xi \sin^2 \rho),
\]

\[
\xi = m_N / (3 m_H + m_N).
\]

Here \( g(\rho) \) is the determinant of the metric tensor \( g = \text{det}(I_{xx} I_{yy} I_{zz} I_{pp}) \) in a curvilinear coordinate system; \( m_H \) and \( m_N \) are the masses of the hydrogen and nitrogen nuclei. The quantities \( I_{xx}, I_{yy} \) and \( I_{zz} \) are the moments of inertia of \( \text{NH}_3 \), which depend on the inversion coordinate \( \rho \). The generalized moment of inertia \( I_{pp} \) is associated with the umbrella motion and depends also on the inversion coordinate. The double well potential \( V_{\text{umb}}(\rho) \) is represented by a harmonic force field augmented by a Gaussian

\[
V_{\text{umb}}(\rho) = \frac{1}{2} k (\rho - \frac{1}{2} \pi)^2 + a \exp[-b (\rho - \frac{1}{2} \pi)^2].
\]

The van der Waals Hamiltonian can be written as

\[
\hat{H}_{\text{vdw}}(\gamma, \beta, \varphi, R, \alpha, \rho) = B(\rho) \hat{J}_z^2 + [C(\rho) - B(\rho)] \hat{J}_z^2
\]

\[
+ V_{\text{vdw}}(\gamma, \beta, \varphi, R, \alpha, \rho).
\]

The first two terms in \( \hat{H}_{\text{vdw}} \) represent the symmetric top Hamiltonian of \( \text{NH}_3 \). The rotational constants are related to the moments of inertia given in Eq. (3), \( B(\rho) = [2I_{xx}(\rho)]^{-1} \) and \( C(\rho) = [2I_{zz}(\rho)]^{-1} \). The third and fourth term give the kinetic energy of the “diatom,” with \( \hat{J}_z \) being the relative angular momentum. The intermolecular potential \( V_{\text{vdw}} \) is expanded in spherical harmonics \( Y_{\ell \mu} \).

\[
V_{\text{vdw}}(R, \Theta, \Phi, \rho) = \sum_{\ell \mu} v_{\ell \mu}(R, \rho) Y_{\ell \mu}(\Theta, \Phi),
\]

where \( \Theta \) and \( \Phi \) are the polar angles of the Ar projectile with respect to the principal axes frame of the \( \text{NH}_3 \) rotor. In the space fixed frame \( Y_{\ell \mu}(\Theta, \Phi) \) becomes

\[
Y_{\ell \mu}(\Theta, \Phi) = \sum_{\ell \mu} D_{\phi \ell \mu}(\gamma, \beta, \varphi) Y_{\ell \mu}(\beta, \alpha),
\]

where \( D_{\phi \ell \mu}(\gamma, \beta, \varphi) \) is the usual Wigner rotation matrix.\(^{11}\) The expansion coefficients \( v_{\ell \mu}(R, \rho) \) have been taken from Bulski et al.,\(^3\) who calculated the \textit{ab initio} potential for four different umbrella angles \( \rho \), and expanded it in tesseral harmonics. Due to the threefold symmetry of the ammonia...
only terms with \( m = 0,3,6,... \) are present. The first 15 terms with \( l \leq 7 \) have been included, which leads to an accuracy of \(-0.6\%\) in the convergence of the expansion. The expansion coefficients can be written as a sum of a short and a long range contribution,

\[
v_{\lambda \mu} (R, \rho) = v_{\lambda \mu}^{\text{SR}} (R, \rho) + v_{\lambda \mu}^{\text{LR}} (R, \rho),
\]

where

\[
v_{\lambda \mu}^{\text{SR}} (R, \rho) = F_{\lambda \mu} (\rho) \left[ 1 + \delta_{\lambda \mu} (\rho) R \right] \exp \left[ - \alpha_{\lambda \mu} (\rho) R \right] - \beta_{\lambda \mu} (\rho) R^2,
\]

\[
v_{\lambda \mu}^{\text{LR}} (R, \rho) = - \sum_{n=0}^{\infty} f_{\mu n} (R, \rho) C_{\mu n} (\rho) R^{-n}.
\]

The \( C_{\mu n} (\rho) \) are the induction and dispersion coefficients, the \( f_{\mu n} (R, \rho) \) are Tang and Toennies type damping functions. The values of all coefficients are given in Ref. 1. The rotational constants for a given tunneling state \( \nu \) are given by

\[
B_\nu = \langle \nu | B(\rho) | \nu \rangle,
\]

\[
C_\nu = \langle \nu | C(\rho) | \nu \rangle.
\]

Their values are listed in Table I.

In the CC method\(^7\) the angular basis functions are usually formed by Clebsch–Gordan coupling of the relative angular momentum functions \( Y_{lm} (\beta, \alpha) = | lm \rangle \) and the symmetric top functions \(| jkm \rangle\) to a total angular momentum \( J \) with space fixed \( z \)-component \( M \). In the case of permutation-inversion symmetry \( \Pi(D_{3h}) \), which we have here, it is convenient to take linear combinations \(| jkme \rangle\) of the symmetric top functions, defined by

\[
| jkme \rangle = \sum_{m' m''} [2(1 + \delta_{00})]^{-1/2} (| jkm \rangle + \epsilon | j-km \rangle),
\]

where \( k \geq 0 \) and \( \epsilon = \pm 1 \), except for \( k = 0 \) when obviously only \( \epsilon = +1 \) is allowed.

Since we take the umbrella motion explicitly into account, the basis has to be extended by taking the tensor product with the tunneling functions \( v(\rho) \),

\[
| jkIMe\nu \rangle = \sum_{m,m_1} | jkme \rangle | lm_1 \rangle | \nu \rangle | jmlm_1 \rangle | JM \rangle,
\]

where \( | jmlm_1 \rangle | JM \rangle \) is a Clebsch–Gordan coefficient.\(^11\)

From symmetry considerations it follows\(^3\) that the symmetric (antisymmetric) inversion function can combine to a state adapted to \( \Pi(D_{3h}) \) with only one of the two \(| jkme \rangle\) functions, so that \( \epsilon = \mp (1/\nu) \) for \( \nu = | \pm \rangle \). We can therefore omit the quantum number \( \nu \) and label the basis functions by \(| jkIMe\varepsilon \rangle\), instead of using the labeling given in Eq. (12). In the CC equations the noninteracting blocks for different \( J \) are separated into two parity blocks, each containing channels \(| jk\rangle\) having different values of \( \epsilon = (1/\nu) \) for \( \nu = | \pm \rangle \). States of the free \( \text{NH}_3 \) can be designated by \( f^e_\nu \), thereby uniquely specifying the inversion function.

The coupling between the channels that originates from the potential matrix elements is given by

\[
\langle \nu | v_{\lambda \mu} (R, \rho) | \nu' \rangle = \sum_{\lambda \mu} c_{\mu n} (R) \langle \nu | (\rho - \pi)^n \nu' \rangle,
\]

where the \( c_{\mu n} (R) \) are the polynomial expansion coefficients. The above relationship between the values of \( \nu \) and \( \nu' \) is used to insert the correct \( \langle \nu, \nu' \rangle \) combination into Eq. (13).

According to Ref. 12 the \( \text{ab initio} \) potential has to be scaled to give good agreement with spectroscopic data for the bound \( \text{Ar–NH}_3 \) complex. This scaling consists of multiplying the short range parameter \( F_3 \) in Eq. (9a) by a factor of 1.43. Here, calculations have been performed using both the original \( \text{ab initio} \) potential and a modified potential in which the same scaling was applied for all values of the inversion coordinate \( \rho \).

In addition to the calculation with the inversion averaged matrix elements in the way we have just described (henceforth referred to as the "exact" inversion method), we have used the model developed by Davis and Boggs\(^2\) and Green.\(^3\) In this model the inversion functions are taken to be delta functions, \( | \pm \rangle = [\delta (\rho - \rho_p) \pm \delta (\rho - \pi + \rho_p)] / \sqrt{2}, \)

where \( \rho_p \) is the value of the inversion coordinate in the equilibrium configuration. In this case the intermolecular potential needs to be known only for the equilibrium angle, since the expansion coefficients, averaged over the inversion functions, are now given by\(^2\)^{13}

\[
\langle \pm | v_{\lambda \mu} (R, \rho) | \pm \rangle = \begin{cases} | v_{\lambda \mu} (R, \rho_p) |, & \text{for } \lambda + \mu \text{ even} \\ 0, & \text{for } \lambda + \mu \text{ odd} \end{cases}
\]

\[
\langle \pm | v_{\lambda \mu} (R, \rho) | \mp \rangle = \begin{cases} | v_{\lambda \mu} (R, \rho_p) |, & \text{for } \lambda + \mu \text{ even} \\ 0, & \text{for } \lambda + \mu \text{ odd} \end{cases}
\]
Using this model for the inversion functions, together with the neglect of the inversion splitting, the scattering equations for para-NH₃ are invariant to a simultaneous change of parity in the incoming and outgoing channels, i.e.,

\[ \sigma(f_j \rightarrow j' k') = \sigma(j_k \rightarrow j'_k). \]  

(16)

In the CS method⁷ the scattering equations are expressed in a body-fixed coordinate system. The fourth term in the van der Waals Hamiltonian, Eq. (5), is approximated by putting \( \hat{J}^2 = \hat{J}^2 + 2 \hat{J} \cdot \hat{J} \) equal to \( J^2 \). This implies that the Coriolis interactions are neglected and that \( \Omega \), the projection of both \( J \) and \( \hat{J} \) on the vector \( R \), is a good quantum number, i.e., there is no coupling between channels with different \( \Omega \). The molecular symmetry group of the dimer is thus enlarged from \( \Pi(D_3) \) to the semidirect product of \( C_{w} \) with \( \Pi(D_{3h}) \).

### III. COMPUTATIONAL ASPECTS

The calculations were carried out with the HIBRIDON inelastic scattering code.¹³ The total collision energy \( E \), the maximum value of the total angular momentum \( J \) and the values of \( j \) and \( k \) at which the rotational basis set is truncated, are input parameters of the program. The values for \( l \) are then given by triangular inequalities [cf. Eq. (12)]. The program has the possibility of further reducing the size of the basis set as the overall rotation takes up more and more of the available energy. So, from a chosen value of \( J \) onwards, the program includes only open channels. To keep the calculations feasible even at higher energies, an interpolation scheme for the total cross sections as a function of \( J \) can be used, leading to a substantial reduction of the required CPU time.

The values of the total energies are determined by the two relative kinetic energies in the experiment, 280 and 485 cm⁻¹. The ortho-NH₃ with initial state \( j=k=0 \) has zero internal energy, so the total energies are equal to the relative kinetic energy. The initial \( j=k=1 \) state of para-NH₃, which is the ground state of this species, has an internal energy of 16.245 cm⁻¹. The total energies are consequently set equal to 296.245 and 501.245 cm⁻¹. The molecular levels in the basis set are retained up to \( j=9 \) inclusive, with all allowed values of \( k \). This means that for ortho-NH₃ 34 levels are included (with a maximum energy of 895 cm⁻¹), 11 of which are asymptotically accessible in the lower, and 19 in the higher energy case. Out of the 66 levels for para-NH₃ (with a maximum energy of 891 cm⁻¹), 24 and 40 levels are accessible, respectively. The \( J \) value at which we start to neglect closed channels is 78. The interpolation step size \( \Delta J \) is taken to be six, so that calculations are actually performed for \( J=0,6,12,\ldots,150 \). As explained in Sec. II, the NH₃ inversion is taken into account by calculation of the matrix elements according to the "exact" inversion method, given in Eq. (14), or according to the delta function model, given in Eq. (15).

Convergence with respect to relevant parameters in the propagator, such as the step size \( \Delta R \), was better than 1%. Table II shows the dependence of the cross sections on the magnitude of the rotational basis set. Going from a maximum \( J \) value of 9 to a maximum of 11 in the rotational basis set, induced changes in the cross sections of ~6%. The neglect of closed channels for \( J>78 \) did not affect the results. The effect of interpolation step size \( \Delta J \) on the cross sections is shown in Table III. The error due to the step size used here is found to be <5%.

Since the averaging of the rotational constants over the inversion wave functions has a small effect (cf. Table I), we have taken the same value for both inversion states,¹⁴ \( B=9.9402 \) cm⁻¹ and \( C=6.3044 \) cm⁻¹. The maximum number of channels used in the calculation was 219 per parity block for ortho-NH₃, taking ~24 cpu hours for a full calculation, and 441 per parity block for para-NH₃, taking ~241 cpu hours on an IBM RS/6000 model 320 workstation.

In the CS calculation we used only the "exact" inversion method. The value of \( \Omega \) ranged from 0 to 7 for ortho and from 0 to 8 for para; the maximum \( J \) value in the rotational basis set was 9 and \( J \) was varied from 0 to 150 at an energy of 485 and 501.245 cm⁻¹, respectively. The in-

<table>
<thead>
<tr>
<th>( \Delta J )</th>
<th>( \Delta J=1 )</th>
<th>( \Delta J=4 )</th>
<th>( \Delta J=6 )</th>
<th>( \Delta J=8 )</th>
</tr>
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<tbody>
<tr>
<td>( 1^2 )</td>
<td>10.83</td>
<td>10.93</td>
<td>11.35</td>
<td>9.06</td>
</tr>
<tr>
<td>( 2^2 )</td>
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<td>9.00</td>
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<td>4.32</td>
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<tr>
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<tr>
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<td>1.63</td>
<td>1.62</td>
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</tr>
<tr>
<td>( 6^2 )</td>
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<td>9.25</td>
<td>9.33</td>
<td>8.42</td>
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<tr>
<td>( 7^2 )</td>
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<td>10.93</td>
<td>11.35</td>
<td>9.06</td>
</tr>
<tr>
<td>( 8^2 )</td>
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<td>9.00</td>
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<th>( \Delta J )</th>
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<td>0.05</td>
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<tr>
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<td>0.22</td>
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<td>0.24</td>
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<tr>
<td>( 3^2 )</td>
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<td>0.05</td>
<td>0.05</td>
<td>0.05</td>
</tr>
<tr>
<td>( 4^2 )</td>
<td>0.06</td>
<td>0.05</td>
<td>0.05</td>
<td>0.05</td>
</tr>
<tr>
<td>( 5^2 )</td>
<td>0.06</td>
<td>0.05</td>
<td>0.05</td>
<td>0.05</td>
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</tbody>
</table>


* TABLE II. Effect of the maximum \( (j,k) \) values in the basis set on the cross sections \( \sigma(0^0 \rightarrow f_k) \) (in Å²) for ortho-NH₃-Ar at an energy of 485 cm⁻¹ (*exact* inversion, ab initio potential).
TABLE IV. State-to-state cross sections $\sigma(0_1^+ \rightarrow f_k^*)$ for ortho-NH$_3$–Ar in $\AA^2$ at an energy of 280 cm$^{-1}$. The cross sections given in parentheses are corrected for the incomplete initial state preparation in the measurement, as follows [cf. Eq. (17)]: $\sigma(0_1^+ \rightarrow f_k^*) = 0.92\sigma(0_1^+ \rightarrow f_k^*) - 0.08\Sigma_{f_k^* \neq f_k^*} \sigma(1_2^+ \rightarrow f_k^*),$ for $f_k^* \neq f_k^*$. A dash (—) in the last column indicates that the corresponding cross section has not been measured.

<table>
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<tr>
<th>$\tilde{A}$</th>
<th>$V_{11}^a$</th>
<th>$V_{11}^b$</th>
<th>$V_{11}^c$</th>
<th>$V_{11}^d$</th>
<th>Expt.</th>
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<tr>
<td>1$^o_1$</td>
<td>11.35(7.64)</td>
<td>4.03(1.09)</td>
<td>11.06(7.41)</td>
<td>3.82(0.87)</td>
<td>6.07</td>
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<td>2$^o_1$</td>
<td>8.74(8.69)</td>
<td>13.35(12.65)</td>
<td>8.45(8.42)</td>
<td>13.20(12.49)</td>
<td>6.73</td>
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<td>3$^o_1$</td>
<td>4.34(4.29)</td>
<td>2.56(2.98)</td>
<td>4.16(4.12)</td>
<td>2.51(2.94)</td>
<td>2.19</td>
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<td>4$^o_1$</td>
<td>0.39(0.53)</td>
<td>2.41(3.77)</td>
<td>0.35(0.49)</td>
<td>2.44(2.40)</td>
<td>0.48</td>
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<td>5$^o_1$</td>
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<td>0.75(1.53)</td>
<td>1.54(2.17)</td>
<td>0.70(1.53)</td>
<td>3.43</td>
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<tr>
<td>6$^o_1$</td>
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<td>10.9(9.47)</td>
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<td>10.65(9.96)</td>
<td>12.23</td>
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<tr>
<td>7$^o_1$</td>
<td>3.25(3.18)</td>
<td>1.94(1.97)</td>
<td>3.38(3.29)</td>
<td>1.87(1.91)</td>
<td>5.77</td>
</tr>
<tr>
<td>8$^o_1$</td>
<td>0.25(0.33)</td>
<td>0.45(0.51)</td>
<td>0.23(0.32)</td>
<td>0.44(0.50)</td>
<td>1.31</td>
</tr>
<tr>
<td>5$^1_1$</td>
<td>0.05(0.06)</td>
<td>0.39(0.40)</td>
<td>0.06(0.07)</td>
<td>0.41(0.42)</td>
<td>—</td>
</tr>
<tr>
<td>6$^1_1$</td>
<td>0.36(0.36)</td>
<td>0.28(0.28)</td>
<td>0.34(0.35)</td>
<td>0.31(0.31)</td>
<td>—</td>
</tr>
</tbody>
</table>

* $V_{11}^a$: delta function model inversion, $ab$ initio potential.
* $V_{11}^b$: delta function model inversion, scaled potential.
* $V_{11}^c$: “exact” inversion $ab$ initio potential.
* $V_{11}^d$: “exact” inversion, scaled potential.

Both for ortho and para-NH$_3$ the scaling in the potential has a large effect. In the case of ortho-NH$_3$, this effect is about the same for the various cross sections at both energies. Especially transitions to $1_2^+$, $2_2^+$, $3_2^+$, and $4_2^+$ are strongly affected in the lower energy case, and transitions to $1_3^+$, $1_4^+$, and $3_4^+$ in the higher energy case. Use of the delta function model for inversion does not affect the influence of the scaling. For para-NH$_3$, the scaling in the potential decreases some of the cross sections at the lower energy and increases them at the higher energy, for other cross sections it is vice versa. For the lower energy the scaling reduces the size of most of the para cross sections, except for the $2_1^+$, $4_1^+$, and $2_3^+$ states. The scaling induces large changes in the relative magnitudes for the $\pm$ inversion states for transitions to the $2_1^+$, $3_2^+$, and $4_2^+$ states.

Comparison with the experiment shows that the calculations using the original $ab$ initio potential give a better overall agreement than calculations using the scaled potential. Particularly, cross sections to the $1_2^+$, $2_2^+$, $3_2^+$, and $4_2^+$ states come out better. In a few cases, however, the cross sections from the scaled potential are closer to the experimental ones.

It has been debated whether it is necessary to include higher anisotropic terms, such as a $v_{33}$ term, in the description of the intermolecular potential, since the observed far-infrared Ar-NH$_3$ spectrum could also be explained with an effective angular potential that contains only terms up to $v_{20}$. When we look at the results of the scattering calculations, we observe, for example, that in the case of ortho-NH$_3$ the experimental cross sections $\sigma(0_1^+ \rightarrow 3_3^+)$ are reproduced fairly well. In the first Born approximation these transitions are solely due to the $v_{33}$ term in the potential. It seems unlikely that the agreement between experiment and theory could be maintained if this important first order contribution were zero, as it would be when $v_{33}$ would vanish.

Inspection of the influence of the scaling in $v_{33}$ on the cross sections of ortho-NH$_3$ shows that not only $\Delta f=\Delta k$ = 3 transitions are affected, but other transitions as well.
This would not be so in the first Born approximation. It means that the contributions from higher Born approximations are significant. The importance of higher order effects in the interaction between Ar and NH$_3$ is confirmed by calculations on bound states. Although the $j_{23}$ term does not contribute to the lower bound states in a first order perturbation theory, which has the isotropic Hamiltonian as its zeroth order Hamiltonian, it proves to be one of the dominant terms in determining the rovibrational energy levels of the Ar--NH$_3$ complex.

At the energies used here, application of the delta function model for the inversion of NH$_3$ has only a small effect on the cross sections, as compared to the "exact" calculations, in the order of 3%. The parity propensities, which

<table>
<thead>
<tr>
<th>$J_l$</th>
<th>$V_{i}^{a}$</th>
<th>$V_{ii}^{a}$</th>
<th>$V_{iii}^{a}$</th>
<th>$V_{iv}^{a}$</th>
<th>Expt.</th>
</tr>
</thead>
<tbody>
<tr>
<td>2$^+$</td>
<td>2.82(2.31)</td>
<td>4.22(4.23)</td>
<td>7.90(4.30)</td>
<td>5.54(2.02)</td>
<td>4.33</td>
</tr>
<tr>
<td>3$^+$</td>
<td>10.73(10.33)</td>
<td>7.94(4.30)</td>
<td>10.20(9.85)</td>
<td>7.53(2.04)</td>
<td>5.72</td>
</tr>
<tr>
<td>4$^+$</td>
<td>3.27(3.18)</td>
<td>2.82(2.31)</td>
<td>4.22(4.23)</td>
<td>7.90(4.30)</td>
<td>2.38</td>
</tr>
<tr>
<td>5$^+$</td>
<td>0.97(0.96)</td>
<td>0.97(0.96)</td>
<td>0.97(0.96)</td>
<td>0.97(0.96)</td>
<td>1.73</td>
</tr>
<tr>
<td>6$^+$</td>
<td>1.57(1.66)</td>
<td>1.57(1.66)</td>
<td>1.57(1.66)</td>
<td>1.57(1.66)</td>
<td>1.01</td>
</tr>
<tr>
<td>7$^+$</td>
<td>2.14(2.21)</td>
<td>2.14(2.21)</td>
<td>2.14(2.21)</td>
<td>2.14(2.21)</td>
<td>1.68</td>
</tr>
</tbody>
</table>

$^aV_{i}$: delta function model inversion, ab initio potential.
$^bV_{ii}$: delta function model inversion, scaled potential.
$^cV_{iii}$: "exact" inversion, ab initio potential.
$^dV_{iv}$: "exact" inversion, scaled potential.

TABLE VI. State-to-state cross sections $\sigma(1f \rightarrow 4f)$ for ortho-NH$_3$--Ar in Å$^2$. The corrected cross sections given in parentheses are obtained as indicated in Table IV. A dash (—) in the last column indicates that the corresponding cross section has not been measured.

<table>
<thead>
<tr>
<th>$J_l$</th>
<th>$V_{iii}^{a}$</th>
<th>$V_{iv}^{a}$</th>
<th>Expt.</th>
</tr>
</thead>
<tbody>
<tr>
<td>2$^+$</td>
<td>3.27(3.18)</td>
<td>4.22(4.23)</td>
<td>7.90(4.30)</td>
</tr>
<tr>
<td>3$^+$</td>
<td>1.57(1.66)</td>
<td>1.57(1.66)</td>
<td>1.57(1.66)</td>
</tr>
<tr>
<td>4$^+$</td>
<td>2.14(2.21)</td>
<td>2.14(2.21)</td>
<td>2.14(2.21)</td>
</tr>
<tr>
<td>5$^+$</td>
<td>2.82(2.31)</td>
<td>4.22(4.23)</td>
<td>7.90(4.30)</td>
</tr>
<tr>
<td>6$^+$</td>
<td>3.27(3.18)</td>
<td>4.22(4.23)</td>
<td>7.90(4.30)</td>
</tr>
</tbody>
</table>

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are defined as the ratios between the difference and the sum of the cross sections to the ± inversion states, are hardly affected. The suggestion that deviations of theoretical propensities from experimental ones in He-NH₃ can be accounted for by taking the inversion motion explicitly into account is hereby contradicted to all likelihood.

Our calculations show further that the invariance to a simultaneous change of parity in the incoming and outgoing channels [cf. Eq. (16)] for para-NH₃ holds to an accuracy of ~0.2% when the inversion is included in the potential matrix elements, but the inversion splitting of 0.8 cm⁻¹ is neglected. When this splitting is included also, the deviations are of the order of 3%. Billing has found in He–NH₃ calculations at an energy of 65 meV (525 cm⁻¹), using a semiclassical approach, that the ψ^j^p (R, ρ) are functions of ρ that vary slowly enough for the delta function approach to be valid.

Application of the CS approximation gives a reasonable agreement for para-NH₃ (Table VII). For ortho-NH₃, however, there are strong deviations, both from the CC method and from experiment. The CS calculation gives zero cross sections to the 3^± states, whereas both experimental and CC cross sections to these states are different from zero. This means that these transitions are caused by the Coriolis terms in the Hamiltonian. Even though these terms are small, their long range apparently gives rise to significant transition probabilities.

In calculations on He–NH₃ scattering at an energy of 98 meV (792 cm⁻¹) using the CS method Meyer et al. have found the cross sections to the 3^± and 4^± states to be exactly zero and the cross section to the 2^± state to be...
The fact that the $3^+_1$ and $4^+_1$ cross sections are zero is an artifact of the CS method. The smallness of the $2^+_1$ cross section, however, cannot be attributed to this method.

V. CONCLUSIONS

We have calculated close coupling state-to-state cross sections for the inelastic scattering of NH$_3$ with Ar at two different collision energies and compared them with experimentally derived cross sections. The inversion motion of NH$_3$ has been taken into account explicitly. Comparison with calculations that use a delta function model description of the inversion motion, shows that this model leads to errors of 3% only, at the energies used here. Previously found deviations from experimentally determined parity propensities for NH$_3$-He, cannot be attributed to use of the delta function model. It is more likely that these discrepancies arise from shortcomings of the intermolecular potential used for that system.

In calculations on bound states of Ar-NH$_3$, the use of a potential in which a single term in the angular expansion of the ab initio potential of Ref. 1 was scaled by a factor of 1.43, gave better agreement with spectroscopic data than the use of the original ab initio potential. In the present scattering calculations the opposite is true. This can be seen as a manifestation of the fact that scattering and bound states probe different regions of the intermolecular potential surface. The applied scaling is too crude to obtain a fully realistic potential surface. For ortho-NH$_3$ the calculated ab initio cross sections reproduce the experimental ones fairly well. For para-NH$_3$ the overall agreement is good too, but differences remain between theoretical and experimental parity propensities, indicating that the ab initio potential needs further improvement.

Comparison of the results for the two potentials shows further that the cross sections are very sensitive to variations in the potential surface. The changes in the cross sections for transitions to the various rotation–inversion states induced by the scaling, show also that higher order effects play a role in the scattering process, just as they do in bound state interactions.

Application of the coupled states approximation leads
TABLE VII. Results of CC and CS calculations ("exact" inversion, \textit{ab initio} potential, without correction for initial state preparation, at a relative kinetic energy of 485 cm\(^{-1}\)) and experimental results. The cross sections are in Å\(^2\). A dash (\textendash) in the last column indicates that the corresponding cross section has not been measured.

<table>
<thead>
<tr>
<th>(\delta)</th>
<th>ortho-NH(_3)–Ar</th>
<th></th>
<th>para-NH(_3)–Ar</th>
</tr>
</thead>
<tbody>
<tr>
<td>(J)</td>
<td>CC</td>
<td>CS</td>
<td>Expt.</td>
</tr>
<tr>
<td>1(^3)</td>
<td>7.92</td>
<td>9.24</td>
<td>4.33</td>
</tr>
<tr>
<td>2(^3)</td>
<td>10.10</td>
<td>8.17</td>
<td>5.72</td>
</tr>
<tr>
<td>3(^3)</td>
<td>3.94</td>
<td>5.17</td>
<td>2.38</td>
</tr>
<tr>
<td>4(^3)</td>
<td>0.46</td>
<td>0.71</td>
<td>1.73</td>
</tr>
<tr>
<td>5(^3)</td>
<td>0.13</td>
<td>0.13</td>
<td></td>
</tr>
<tr>
<td>6(^3)</td>
<td>0.07</td>
<td>0.08</td>
<td>1.01</td>
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<td>3(^1)</td>
<td>1.32</td>
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<td>4(^1)</td>
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<td>2.27</td>
<td>4.26</td>
</tr>
<tr>
<td>4(^3)</td>
<td>0.25</td>
<td>0.23</td>
<td>1.01</td>
</tr>
<tr>
<td>5(^1)</td>
<td>0.06</td>
<td>0.06</td>
<td></td>
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<td>7(^1)</td>
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<td></td>
</tr>
<tr>
<td>7(^3)</td>
<td>0.11</td>
<td>0.12</td>
<td></td>
</tr>
<tr>
<td>8(^1)</td>
<td>0.01</td>
<td>0.01</td>
<td></td>
</tr>
</tbody>
</table>

...to artificial selection rules for the cross sections from the 0\(^3\) state. Some of the cross sections, which should vanish according to these rules, are not even small, in the CC calculations, nor in the experiment. Consequently, the overall agreement with close coupling results is poor for ortho-NH\(_3\), reasonable for para-NH\(_3\).

ACKNOWLEDGMENTS

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