$B^0$ lifetime measurement in the CP-odd decay channel $B^0_s \rightarrow J/\psi f_0(980)$

The lifetime of the $B_s^0$ meson is measured in the decay channel $B_s^0 \rightarrow J/\psi \pi^+\pi^-$ with $880 \leq M_{\pi^+\pi^-} \leq 1080$ MeV/$c^2$, which is mainly a CP-odd state and dominated by the $f_0(980)$ resonance. In 10.4 fb$^{-1}$ of data collected with the D0 detector in Run II of the Tevatron, the lifetime of the $B_s^0$ meson is measured to be $\tau(B_s^0) = 1.70 \pm 0.14 \text{ (stat)} \pm 0.05 \text{ (syst)}$ ps. Neglecting CP violation in $B_s^0/B_s^0$ mixing, the measurement can be translated into the width of the heavy mass eigenstate of the $B_s^0$, $\Gamma_H = 0.59 \pm 0.05 \text{ (stat)} \pm 0.02 \text{ (syst)}$ ps$^{-1}$.

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The $B_s^0$ and $B^0$ mesons are produced as flavor eigenstates at hadron colliders, but the particles propagate as mass eigenstates. There are two mass eigenstates, the so-called heavy and light states, which are linear combinations of the flavor eigenstates. In the absence of CP-violation in mixing, the mass eigenstates are also CP eigenstates, with the heavier state expected to be the CP-odd state. The lifetimes of the two mass eigenstates can be different from each other and different from the average $B_s^0$ lifetime. A measurement of the $B_s^0$ lifetime in either a pure CP-odd state or pure CP-even state would give important additional information about the $B_s^0$ system.

The $B_s^0 \rightarrow J/\psi f_0(980)$ decay channel corresponds to a pure CP-odd eigenstate decay due to angular momen-


dation conservation, since the parent $B_s^0$ is spin 0, the $f_0(980)$ has $J^{PC} = 0^{++}$, and the $J/\psi$ has $J^{PC} = 1^{--}$. Throughout this Letter, the appearance of a specific charge state also implies its charge conjugate. This decay channel was first observed by the LHCb collaboration [1], and later confirmed by the Belle [2], CDF [3] and D0 [4] collaborations. A measurement of the $B_s^0$ lifetime in this channel gives access to the lifetime of the heavy mass eigenstate. The lifetime measurement can be transformed into a measurement of the parameter $\Gamma_H$, the decay width of the heavy $B^0$ mass eigenstate. CDF [3] and LHCb [5] have measured this lifetime, reporting $\tau(B_s^0) = (1.70 \pm 0.12 \pm 0.03)$ ps and $\tau(B_s^0) = (1.70 \pm 0.04 \pm 0.026)$ ps respectively, which are in good agreement with each other and somewhat longer than the mean lifetime $\tau(B_s^0) = (1.52 \pm 0.007)$ ps [6].

In this analysis, we report the lifetime of the $B_s^0$ meson measured in the decay channel $B_s^0 \rightarrow J/\psi(\rightarrow \mu^+\mu^-)\pi^+\pi^-$ with $880 \leq M_{\pi^+\pi^-} \leq 1080$ MeV/$c^2$, which is dominated by the $f_0(980)$ resonance and which is CP-odd at the 99% level [7][8]. The data used in this analysis were collected with the D0 detector during Run II of the Tevatron collider at a center-of-mass energy of 1.96 TeV, and correspond to an integrated luminosity of 10.4 fb$^{-1}$.

The D0 detector is described in detail elsewhere [9]. The detector components most relevant to this analysis are the central tracking and the muon systems. The former consists of a silicon microstrip tracker (SMT) and a central scintillating fiber tracker (CFT) surrounded by a 2 T superconducting solenoidal magnet. The SMT has a design optimized for tracking and vertexing for pseudo-

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rapidity of $|\eta| < 3$ [10]. For charged particles, the resolution on the distance of closest approach as provided by the tracking system is approximately 50 $\mu$m for tracks with $p_T \approx 1$ GeV/c, where $p_T$ is the component of the momentum perpendicular to the beam axis. It improves asymptotically to 15 $\mu$m for tracks with $p_T > 10$ GeV/c. Preshower detectors and electromagnetic and hadronic calorimeters surround the tracker. The muon system is located outside the calorimeter, and consists of multilayer drift chambers and scintillation counters inside 1.8 T iron toroidal magnets, and two similar layers outside the toroids. Muon identification and tracking for 1.8 T iron toroidal magnets, and two similar layers out-tilayer drift chambers and scintillation counters inside calorimeters surround the tracker. The muon system Preshower detectors and electromagnetic and hadronic decay vertex, and $p_T$ is the transverse momentum vector of the $B_s^0$ candidate. The event-by-event value of the proper transverse decay length, $\lambda$, for the $B_s^0$ candidate is given by:

$$\lambda = L_{xy} \frac{cM_B}{p_T},$$

(1)

where $M_B$ is the world average mass value of the $B_s^0$ meson [13]. In order to remove background, $B_s^0$ candidates are required to have $\lambda > 0.02$ cm and uncertainties on $\lambda$ of less than 0.01 cm.

A simultaneous unbinned maximum likelihood fit to the mass and proper decay length distributions is performed to measure the lifetime. The likelihood function $L$ is defined by:

$$L = \prod_{j=1}^{N} \left[ N_{\text{sig}} F_{\text{sig}}^j + N_{\text{comb}} F_{\text{comb}}^j + N_{\text{sf}} F_{\text{sf}}^j + N_{B^+} F_{B^+}^j \right],$$

(2)

where $N$ is the total number of events and $N_{\text{sig}}$, $N_{\text{comb}}$, $N_{\text{sf}}$ and $N_{B^+}$ are the expected number of signal, combinatorial background, cross-feed contamination and $B^\pm \rightarrow J/\psi K^\pm$ events in the sample, respectively. All these parameters are determined in the fit. The different background contributions are discussed below.

The functions $F$ are the product of three probability density functions that model distributions of the mass $m$, the proper transverse decay length $\lambda$, and the uncertainty on the proper decay length $\sigma_\lambda$ for the signal, combinatorial background, cross-feed contamination, and $B^\pm$ events

$$F_a = M_a(m_j) T_a(\lambda_j) \sigma_\lambda(m_j) E_a(\sigma_\lambda),$$

(3)

where $m_j$, $\lambda_j$, and $\sigma_\lambda$ represent the mass, the transverse proper decay length, and its uncertainty, respectively, for a given event $j$. The use of the probability density functions $T$ and $E$ follows the method of reference [15]. The specific models and parameters used in the fit are described below.

For the signal, the mass distribution is modeled by a Gaussian function, $M_{\text{sig}}(m_j) = G(m_j; \mu_m, \sigma_m)$, where

$$G(m_j; \mu_m, \sigma_m) = \frac{1}{\sqrt{2\pi}\sigma_m} e^{-\frac{(m_j-\mu_m)^2}{2\sigma_m^2}},$$

(4)

with $\mu_m$ and $\sigma_m$ the mean and the width of the Gaussian, determined from the fit.

The combinatorial background is primarily due to random combinations of $J/\psi$’s with additional tracks in the event, and its mass distribution is described by an exponential function

$$M_{\text{comb}}(m_j; a_0) = e^{a_0 m_j},$$

(5)
with \( a_0 \) determined from the likelihood fit.

The physics cross-feed contamination is mainly produced by the combination of \( J/\psi \) mesons from \( b \) hadron decays with other particles produced in the collision, including from the same \( b \) hadron. Other \( b \) hadron decays with final states such as \( B^0 \rightarrow J/\psi K\pi, B^0 \rightarrow J/\psi \pi\pi \) and \( B^+ \rightarrow J/\psi KK \) are reconstructed at mass below the signal of the \( B_s^+ \), either due to the lower mass of the \( B^0 \) or the incorrect mass assignment of the pion mass to a kaon track. Simulations of these decays show that the cross-feed contamination can be described by a single Gaussian component

\[
M_{xf}(m_j) = G(m_j; \mu_{xf}, \sigma_{xf}),
\]

where \( \mu_{xf} \) and \( \sigma_{xf} \) are the mean and the width of the Gaussian, determined from the likelihood fit.

The final contribution arises from \( B^\pm \rightarrow J/\psi K^\pm \) decays in which the kaon has been assigned a pion mass, and an additional track accidentally forms a vertex with the \( J/\psi K^\pm \). The candidate mass is reconstructed in the region of real \( B^0 \) events. If the higher \( PT \) non-\( \mu \)-track in \( B_s^0 \) candidates is assigned a kaon mass, a clear \( B^\pm \) signal emerges. Events in this \( B^\pm \) mass peak, when interpreted as \( J/\psi \pi\pi \), are used as a template to determine the mass distribution of the \( B^\pm \rightarrow J/\psi K^\pm \) contamination in the \( B_s^0 \) candidates.

The \( \lambda \) distribution for the signal is parameterized by an exponential decay convoluted with a resolution function

\[
T_{\text{sig}}(\lambda | \sigma_{\lambda}) = \frac{1}{\lambda_B} \int_0^\infty G(x; \lambda, \sigma_{\lambda}) \exp \left( -\frac{x}{\lambda_B} \right) dx,
\]

with \( \lambda_B = c \tau \) of the \( B_s^0 \) to be measured. The \( \lambda \) distribution for the background components is parameterized by the sum of two exponential decay functions modeling combinatorial background \( T_{\text{comb}}(\lambda_j) \), an exponential decay for the cross-feed contamination \( T_{\text{xf}}(\lambda_j) \), and an exponential decay function that describes \( T_{B^\pm}(\lambda_j) \) for \( B^\pm \) contamination.

The distribution of the \( \lambda \) uncertainty \( E_{\text{sig}}(\sigma_{\lambda}) \) is described by a phenomenological model, using an exponential with decay constant \( 1/\zeta \), convoluted with a Gaussian with mean \( \epsilon \) and width \( \delta \):

\[
E_{\text{sig}}(\sigma_{\lambda}; \zeta, \epsilon, \delta) = \frac{1}{\zeta} e^{-\sigma_{\lambda}/\zeta} \otimes G(\sigma_{\lambda}; \epsilon, \delta),
\]

where the parameters \( \zeta, \epsilon \) and width \( \delta \) are determined from the fit in the sample of events. The uncertainties in \( \lambda \) for the background components are treated in the same manner.

The fit yields \( c\tau(B_s^0) = 504 \pm 42 \) \( \mu \)m and the numbers of signal decays to be 494 \( \pm 85 \). Figure 1 shows the mass, \( \lambda \) and \( \lambda \) uncertainty distributions for data with the fit results superimposed.

Each of the different background components is indicated in the figure. The fit yields \( c\tau(B_s^0) = 504 \pm 42 \) \( \mu \)m.

**FIG. 1:** Distributions of (a) invariant mass, (b) proper transverse decay length, and (c) proper transverse decay length uncertainty for \( B_s^0 \) candidates, with the fit results superimposed.
FIG. 2: $M(\pi^+\pi^-)$ distribution for events with $M(\mu^+\mu^-\pi^+\pi^-)$ within ±1σ of the $B_s^0$ mass.
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