Search for high-mass diphoton resonances in $pp$ collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector

The ATLAS Collaboration

Abstract

This article describes a search for high-mass resonances decaying to a pair of photons using a sample of 20.3 fb$^{-1}$ of $pp$ collisions at $\sqrt{s} = 8$ TeV recorded with the ATLAS detector at the Large Hadron Collider. The data are found to be in agreement with the Standard Model prediction, and limits are reported in the framework of the Randall-Sundrum model. This theory leads to the prediction of graviton states, the lightest of which could be observed at the Large Hadron Collider. A lower limit of 2.66 (1.41) TeV at 95% confidence level is set on the mass of the lightest graviton for couplings of $k/M_{Pl} = 0.1 (0.01)$.
1 Introduction

One of the primary goals of the experiments at the Large Hadron Collider (LHC) is the search for new phenomena which might manifest themselves in the high-energy regime made accessible for the first time by this machine. This article reports on a search for new high-mass resonances in the diphoton channel using the full data set of proton-proton (pp) collisions recorded with the ATLAS detector \[1\] at $\sqrt{s} = 8$ TeV in 2012.

The diphoton channel is important for searches for physics beyond the Standard Model (SM) since new high-mass states decaying to two photons are predicted in many extensions of the SM, and since searches in the diphoton channel benefit from a clean experimental signature: excellent mass resolution and modest backgrounds. One popular example of theories predicting a new high-mass diphoton resonance is the Randall-Sundrum (RS) model \[2\]. This model attempts to solve the so-called hierarchy problem of the SM, i.e. the fact that the electroweak scale is 16 orders of magnitude smaller than the Planck scale ($M_{Pl} \approx 10^{16}$ TeV). The hierarchy problem poses the question of the naturalness of the SM, and despite significant efforts over the past few decades, this problem has not been solved so far. In the RS model, the observed hierarchy of scales is accommodated assuming an extra spatial dimension. In this paradigm, the fundamental strength of gravity is comparable to the strength of the electroweak interaction, but gravity appears much weaker than the other interactions because it is diluted by the presence of the extra spatial dimension.

The RS model postulates a “warped” five-dimensional spacetime configuration. In this configuration there are two 3-branes that are the boundaries of the five-dimensional spacetime. The 3-branes can support $(3+1)$-dimensional field theories. The SM fields are located on one of the two 3-branes, the so-called TeV brane, while gravity originates from the other brane, called the Planck brane. Gravitons propagate in the five-dimensional bulk. The fifth dimension is compactified with length $r_c$. The solution to Einstein’s equations for this configuration is given by a spacetime metric that couples the standard four dimensions to the extra dimension with an exponential “warp factor”. In this configuration, a field with the fundamental mass parameter $m_0$ will appear to have the physical mass $m = e^{-kr_c} m_0$, where $k$ is the curvature scale of the extra dimension. Scales of the order of 1 TeV are therefore generated from fundamental scales of the order of $M_{Pl}$ if $kr_c \approx 12$ \[3\]. The compactification of the extra dimension gives rise to a Kaluza-Klein (KK) \[4\] tower of graviton excitations $G^*$ – a set of particles on the TeV brane with increasing mass. The phenomenology can be described in terms of the mass of the lightest KK graviton excitation ($m_{G^*}$) and the dimensionless coupling to the SM fields, $k/M_{Pl}$, where $M_{Pl} = M_{Pl}/\sqrt{8\pi}$ is the reduced Planck scale. The masses and couplings of the individual graviton states are determined by the scale $\Lambda_e = e^{-kr_c} M_{Pl} \approx O(\text{TeV})$ \[3\]. Values of $k/M_{Pl}$ in the range $[0.01, 0.1]$ are preferred from theoretical arguments \[3\]. The lightest graviton state is expected to be a fairly narrow resonance for $k/M_{Pl} < 0.3$ \[3\]. Its natural width varies like the square of $k/M_{Pl}$. For $k/M_{Pl} = 0.1$, the natural width increases from $\sim 8$ GeV at $m_{G^*} = 800$ GeV to $\sim 30$ GeV at $m_{G^*} = 2200$ GeV. In the first example, the natural width and the experimental mass resolution are similar, in the second example the natural width dominates the total width \(^1\) of the mass peak.

Searches for high-mass diphoton resonances at the LHC have been reported by both the ATLAS and CMS collaborations, and the latest results are summarized in Table 1. The most stringent experimental limits on RS gravitons to date have been obtained at the LHC, using a combination of the dielectron and dimuon channels. Limits from searches for narrow resonances in other channels are available in, for example, the

\(^1\) The experimental mass resolution at $m_{G^*} = 2200$ GeV is only marginally (6%) better than at $m_{G^*} = 800$ GeV.
the range of the pseudorapidity interval measurements of the muon track curvature in the region of absorbers. The muon spectrometer features three stations of precision chambers that allow for accurate material. In the region 1, the ECAL is used to correct for this leakage beyond the ECAL from a high-energy shower, and the measurements from the third layer are used to estimate the energy from photon showers. Significant energy deposited in the third layer is an indication for a candidate at low-mass scale. The first layer consists of highly granular strips segmented longitudinally in shower depth into three layers for secondary vertices, permitting the identification of photon conversions. The EM calorimeter serves as a thin presampler layer, installed in front of the ECAL and covering the pseudorapidity interval $|\eta| < 2.5$, is used to correct for energy loss before the ECAL. In the region $|\eta| < 1.8$, the hadronic sampling calorimeter uses iron absorbers and plastic scintillator tiles as the active material. In the region $1.5 < |\eta| < 4.9$, LAr is used as the active material, with copper or tungsten absorbers. The muon spectrometer features three stations of precision chambers that allow for accurate measurements of the muon track curvature in the region $|\eta| < 2.7$. Fast detectors enable muon triggering in the range $|\eta| < 2.4$. A three-level trigger system is used to select the events to be recorded for subsequent offline analysis.

2 ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point in the center of the detector and the z-axis along the beam pipe. The x-axis points from the interaction point to the center of the LHC ring, and the y-axis points upward. Cylindrical coordinates (r, \(\phi\)) are used in the transverse plane, \(\phi\) being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle \(\theta\) as \(\eta = -\ln \tan(\theta/2)\).

### Table 1: Overview of previous limits on the mass of the lightest RS graviton state, expressed at 95% confidence level. The notation \(\ell^+\ell^-\) refers to a combination of the dielectron and dimuon channels. The CDF and D0 experiments operated at the Tevatron pp collider.

<table>
<thead>
<tr>
<th>Experiment</th>
<th>(\sqrt{s}) [TeV]</th>
<th>Luminosity [fb(^{-1})]</th>
<th>Channel</th>
<th>Limit</th>
<th>Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>ATLAS</td>
<td>7</td>
<td>4.9</td>
<td>(\gamma\gamma)</td>
<td>2.06 (1.00) TeV for (k/M_{Pl} = 0.1) (0.01)</td>
<td>[5]</td>
</tr>
<tr>
<td>CMS</td>
<td>7</td>
<td>2.2</td>
<td>(\gamma\gamma)</td>
<td>1.84 (0.86) TeV for (k/M_{Pl} = 0.1) (0.01)</td>
<td>[6]</td>
</tr>
<tr>
<td>ATLAS</td>
<td>8</td>
<td>20.5</td>
<td>(\ell^+\ell^-)</td>
<td>2.68 (1.25) TeV for (k/M_{Pl} = 0.1) (0.01)</td>
<td>[7]</td>
</tr>
<tr>
<td>CMS</td>
<td>7 and 8</td>
<td>5.3 and 4.1</td>
<td>(\ell^+\ell^-)</td>
<td>2.39 TeV for (k/M_{Pl} = 0.1)</td>
<td>[8]</td>
</tr>
<tr>
<td>ATLAS</td>
<td>7</td>
<td>4.7</td>
<td>jet-jet</td>
<td>1.45 TeV for (k/M_{Pl} = 0.1)</td>
<td>[9]</td>
</tr>
<tr>
<td>CMS</td>
<td>8</td>
<td>7</td>
<td>(W^+W^-)</td>
<td>1.23 TeV for (k/M_{Pl} = 0.1)</td>
<td>[10]</td>
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<tr>
<td>CDF</td>
<td>1.96</td>
<td>5.7 and 5.4</td>
<td>(e^+e^-) and (\gamma\gamma)</td>
<td>1.06 TeV for (k/M_{Pl} = 0.1)</td>
<td>[11]</td>
</tr>
<tr>
<td>D0</td>
<td>1.96</td>
<td>5.4</td>
<td>(e^+e^-) and (\gamma\gamma)</td>
<td>1.05 (0.56) TeV for (k/M_{Pl} = 0.1) (0.01)</td>
<td>[12]</td>
</tr>
</tbody>
</table>
3 Event reconstruction and selection

The events used in this analysis were recorded using a diphoton trigger. At the highest, software-based level of the trigger system, this trigger requires at least one photon with transverse energy $E_T^\gamma > 35$ GeV and at least one other photon with $E_T^\gamma > 25$ GeV. Both photons need to satisfy requirements on the shape of the energy deposit in the calorimeter. This trigger is nearly fully efficient for diphoton pairs that pass the offline event selection.

In the offline analysis, only events that pass the standard ATLAS data quality requirements are considered. Photons are reconstructed from fixed-size clusters of cells in the EM calorimeter. To ensure good photon mass resolution, photons are required to be located inside the precision region of the EM calorimeter ($|\eta| < 1.37$ or $1.52 < |\eta| < 2.37$). They are further required to satisfy $E_T^\gamma > 50$ GeV as well as the loose identification requirements from Ref. [13] updated for 2012 run conditions. The two highest-$E_T^\gamma$ photon candidates that pass these requirements are retained for further analysis. Events are retained for the main analysis if both candidates also satisfy the tight identification requirements [13] updated for the 2012 run conditions. These requirements include upper limits on the energy leakage into the hadronic calorimeter and on the width of the shower in the first and second layers of the EM calorimeter. Events where at least one candidate fails the tight photon identification are retained for the study of backgrounds with misidentified (fake) photons in data control samples.

The invariant mass of the diphoton system is evaluated using the photon energies measured in the calorimeter as well as the direction of the photon momenta determined from the positions of the photons in the calorimeter and the position of the primary vertex of the hard interaction. Primary vertices are reconstructed from the tracks in the ID and are required to be formed from at least three tracks. The primary vertex of the hard interaction is identified using information on the directions of flight of the photons as determined using the longitudinal segmentation of the ECAL (calorimeter pointing), the parameters of the beam spot and the properties of the tracks associated with each reconstructed vertex [14]. This procedure ensures good diphoton mass resolution despite the presence of pileup. Pileup refers to the additional $pp$ interactions that occur in the same bunch crossing as the $pp \rightarrow \gamma\gamma + X$ interaction or in the adjacent crossings.

A calorimetric isolation observable is defined for photon candidates. This observable provides further discrimination between genuine photons and fake candidates from jets ($j$). It is used for two purposes in this analysis: to quantify the composition of the data sample in terms of true and fake photon candidates, and to suppress the reducible background by requiring both photons to be isolated. The isolation energy $E_T^{\text{iso}}$ is defined as the sum of the transverse energies of cells inside clusters of variable size (also called “topo-clusters” [15]) inside a cone of $\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2} = 0.4$ around the photon candidate, excluding a $5 \times 7$ grid of EM cells in the center of the cone.

The measurement of $E_T^{\text{iso}}$ is corrected for the following two effects: i) soft energy contributions from the underlying event and pileup that fall inside the isolation cone and ii) leakage of photon energy outside the central core, causing the isolation energy to grow as a function of photon $E_T^\gamma$. The expected soft contributions are subtracted from $E_T^{\text{iso}}$ on an event-by-event basis; the correction is based on a measurement of the properties of low-$p_T$ jets [16] in the event. The average expected contributions to $E_T^{\text{iso}}$ from leakage are subtracted based on simulations. The fluctuations (from one photon shower to another) around the average expected leakage contribution to $E_T^{\text{iso}}$ become more important as $E_T^\gamma$ increases. In particular, the tail at large $E_T^{\text{iso}}$ becomes more pronounced and would cause significant inefficiencies at large $E_T^\gamma$ if a fixed upper limit on $E_T^{\text{iso}}$ were imposed in the event selection. To avoid inefficiencies, a negative $E_T^\gamma$-dependent offset is added to $E_T^{\text{iso}}$, and the requirement $E_T^{\text{iso}} < 8$ GeV is imposed on both photon candidates after the addition of the offset. In the following, the symbol $E_T^{\text{iso}}$ refers to the isolation energy with the offset...
applied. The $E_T^{\gamma}$-dependence of the offset is chosen such that the efficiency of the $E_T^{\text{iso}}$ requirement is independent of $E_T^{\gamma}$ in simulated samples, and the size of the offset varies from zero at low $E_T^{\gamma}$ to $-5\text{ GeV}$ at $E_T^{\gamma} = 1\text{ TeV}$. The corresponding efficiency of the isolation requirement is $96\%$ per photon. For the purpose of the determination of the background composition, a less stringent requirement of $E_T^{\text{iso}} < 14\text{ GeV}$ is used. Examples of the $E_T^{\text{iso}}$ distributions will be shown after the discussion of the backgrounds.

### 4 Background estimate

An important contribution to the background arises from prompt $\gamma\gamma$ production via Standard Model processes. This contribution is irreducible and, as quantified below, represents the dominant source of background. Another significant contribution is due to events in which one or both of the photon candidates arise from other objects, such as misidentified jets or electrons. This background component is dominated by $\gamma + j$ and $j + j$ events with one or two jets reconstructed as photons. Backgrounds with electrons misidentified as photons – e.g. electron-positron events from Drell-Yan production, in $W/Z + \gamma$ or in $t\bar{t}$ events – were verified to be negligible after the event selection.

The irreducible background is estimated using Monte Carlo (MC) simulations normalized to the data in a control region defined below, and the reducible background is estimated using control samples in data. A combination of the leading-order (LO) event generator Pythia [17] and the next-to-leading-order (NLO) parton-level simulation Diphox [18] is used to simulate the irreducible background. Monte Carlo samples are generated using Pythia 8.163 with the CTEx6L1 [19] parton distribution functions (PDFs) and the AU2 Pythia parameter tune [20]. The detector response is simulated using Geant4 [21, 22], including pileup conditions similar to those observed in data. The simulated events are reweighted such that the number of $pp$ interactions per bunch crossing has the same distribution as in data. The NLO correction and the contributions of the fragmentation process are evaluated using Diphox 1.3.2 with the CTEx6.6M [23] PDF set and are used to scale the Pythia prediction by a factor that depends on the generated mass of the diphoton system. The so-called box contribution $gg \rightarrow \gamma\gamma$ through a quark loop is included in both the Pythia and Diphox predictions. From the point of view of power counting, this diagram is a next-to-next-to-leading-order (NNLO), i.e. $O(\alpha^2_{\text{EM}}\alpha_s^2)$, contribution. But given the large gluon luminosity at the LHC compared to the quark-antiquark one, the contribution of the box diagram is comparable to that of the $q\bar{q} \rightarrow \gamma\gamma$ process, which corresponds to a LO, i.e. $O(\alpha^2_{\text{EM}})$, diagram.

To reduce the impact of the theory uncertainties in the absolute normalization of the production cross-section of the irreducible background, the background estimates are scaled to the data at low mass. In the mass region $m_{\gamma\gamma} < 600\text{ GeV}$, where $m_{\gamma\gamma}$ denotes the measured invariant mass of the diphoton system, the most sensitive searches for narrow resonances in the diphoton channel rely on background estimates that make use of low-mass and high-mass sidebands [24]. The search described in the present article focuses on higher diphoton masses where the sidebands, especially at high mass, provide insufficient event yields. This search therefore uses a different approach to obtain the background prediction at high $m_{\gamma\gamma}$.

Specifically, the simulated irreducible background is normalized to the data in a low-mass control region after subtraction of the reducible backgrounds. This control region is defined to coincide with the first 22 bins in Fig. 3, i.e. as $179 < m_{\gamma\gamma} < 409\text{ GeV}$. To determine the composition of the data sample in the low-mass control region in terms of irreducible and reducible backgrounds, a template fit to the $E_T^{\text{iso}}$ distributions of the leading and subleading photon candidates is performed. A similar technique was used in earlier analyses of low-mass diphoton events [25, 26]. The $E_T^{\text{iso}}$ distributions, along with the result of the template fit, are shown in Fig. 1 (top and middle). The full event selection described in Sec. 3, with the relaxed requirement $E_T^{\text{iso}} < 14\text{ GeV}$, has been applied. The fit is performed simultaneously to the
two distributions, and the normalizations for four background components are allowed to float in the fit, namely the irreducible $\gamma\gamma$ component plus three reducible components: $\gamma + j$ (subleading photon candidate is fake), $j + \gamma$ and $j + j$. Templates for the $E_T^{\gamma\gamma}$ distribution of true photons and of fake photons are determined from data. The template for fake photons is obtained using a sample of photon candidates that pass the non-tight selection: photon candidates must satisfy the loose identification criteria and fail to meet a subset of the criteria in the tight selection. This subset has been chosen to include criteria that have been found to be only weakly correlated with $E_T^{\gamma\gamma}$. The template for true photons is obtained from the sample of tight photon candidates after subtraction of the fake-photon template normalized to the tight data sample at $E_T^{\gamma\gamma} > 10$ GeV where fake photons dominate.

The fit is performed separately in five subsamples of the data, because the relative contributions of the four background components vary significantly from one subsample to another. The five subsamples correspond to different combinations of pseudorapidities of the leading and subleading photon candidates: both photons in the “central” region ($|\eta| < 1.37$), leading photon central and subleading photon in the “endcap” region ($1.52 < |\eta| < 2.37$), leading photon in endcap and subleading central, both photons in same endcap and both photons in endcaps of opposite-sign pseudorapidity. In Fig. 1, the data from the five subsamples, as well as the corresponding fit results are combined, and only the sum of the photon and jet components over the four background components is shown. To illustrate the broadening of the $E_T^{\gamma\gamma}$ distribution at large $E_T^{\gamma}$ described above, Fig. 1 also includes the $E_T^{\gamma\gamma}$ distribution for photons with $E_T^{\gamma} > 400$ GeV.

The backgrounds in the search region, i.e. at masses higher than the low-mass control region, are predicted using estimates of the $m_{\gamma\gamma}$ shape of each background component and the normalizations obtained in the low-mass control region described above. As discussed above, the shape of the irreducible component is obtained using simulation. The shapes of the reducible components are obtained from data control samples. The control sample used to extract the shape of the $\gamma + j$ background is selected in the same way as the signal sample, except that the subleading photon candidate has to satisfy the loose identification criteria and fail to meet the tight criteria. The control samples for the $j + \gamma$ and $j + j$ backgrounds are defined using the same approach. Since these data control samples contain relatively few events at high $m_{\gamma\gamma}$, a fit to a smooth function of the form $f(m_{\gamma\gamma}) = p_1 \times (m_{\gamma\gamma})^{p_2} + p_3 \log(m_{\gamma\gamma})$, where the $p_i$ are free parameters, is used to extrapolate the shapes of the reducible backgrounds to high $m_{\gamma\gamma}$. This functional form describes well the shapes of the data in the control samples of this analysis. It was also successful in describing the shapes of the background control samples in earlier searches in the dijet [27, 28], dilepton [29], photon-jet [30] and diphoton [5] channels.

The shapes of the distribution of the isolation energy for genuine photons extracted from data (in bins of $E_T^}\gamma$) are compared to the shapes predicted using simulation. The impact of the small observed differences in the isolation distributions on the predicted $m_{\gamma\gamma}$ shape of the irreducible background and on the signal efficiency is quantified and taken into account in the systematic uncertainties.

The uncertainties in the prediction of the total background are shown in Fig. 2 as a function of $m_{\gamma\gamma}$. The uncertainties in the shape of the $m_{\gamma\gamma}$ distribution of the irreducible background are dominated by the uncertainties in the PDF set used for the Dithox simulation. The set of PDF variations provided with the CTEQ6.6M PDF set are used to propagate the PDF uncertainty to the shape of the irreducible background. The comparatively small difference between the irreducible background shapes obtained with the CTEQ6.6M and the MSTW2008 [31] PDF sets is added in quadrature to the uncertainty. An additional uncertainty arises from higher-order contributions that are not included in the Dithox generator. This uncertainty is evaluated by varying in a correlated and anticorrelated way the renormalization scale and the factorization scales in the hard scattering process and in the photon fragmentation function [18] by a
Figure 1: (Top) Distribution of the isolation energy $E_T^{iso}$ for the leading photon candidate in events in the low-mass control region. (Middle) Corresponding distribution for the subleading photon. The black points represent the data. The blue line shows the result of the template fit to the data. The red and green lines show the templates for photon and jet background respectively, normalized according to the fit result. The red and green bands represent the uncertainties in the photon and jet contributions due to the uncertainties in the template shapes. (Bottom) Distribution of $E_T^{iso}$ for photon candidates with $E_T^\gamma > 400$ GeV at any diphoton mass. The $E_T^{iso}$ values always include the $E_T^\gamma$-dependent offset discussed in the text.
Figure 2: Uncertainty on the total background estimate as a function of $m_{\gamma\gamma}$. The individual contributions due to the uncertainties in the irreducible and reducible components, as well as the contribution from the uncertainties in the normalization of the different background components (labeled “background composition”), are also shown.

factor two around their nominal value, which is set to the mass of the diphoton system. A small uncertainty arises from residual imperfections in the simulation of the isolation energy $E_{\text{iso}}$. The uncertainties in the shape of the reducible background arise from the finite size of the corresponding control samples, and from the extrapolation of the background shapes from the control samples to the signal region. The latter are assessed by varying the definition of the loose selection requirement that is used to define the control samples. The uncertainties in the normalizations of the individual background components are dominated by the uncertainties in the template shapes that are used in the fit to the isolation energy. They are assessed by varying the definition of the non-tight selection requirements that are used to obtain the template shapes.

5 Signal response

Simulated event samples of RS graviton signals at different values of $m_G$ are used to study the response of the detector for signal events. The samples are generated using Pythia 8.163 with the CTEQ6L1 PDF set, the AU2 tune and the same detector and pileup simulation as for the SM $\gamma\gamma$ samples described above. The sources of uncertainty in the expected signal yield are summarized in Table 2. The systematic uncertainty is dominated by the uncertainty in the efficiency due to residual imperfections in the simulation of the observables that are used in the tight photon identification criteria and in the simulation of the isolation energy. Smaller contributions arise due to the limited size of the simulated samples and the determination of the trigger efficiency. The uncertainty in the integrated luminosity is also propagated to the expected yield. Uncertainties in the RS resonance shape due to our current knowledge of the photon energy scale and resolution as well as imperfections in the simulation of the pileup conditions were verified to have a negligible effect on the result.
Table 2: Summary of systematic uncertainties on the expected signal yield, excluding those on the production cross section and the acceptance. The total systematic uncertainty is obtained by summing the individual contributions in quadrature.

<table>
<thead>
<tr>
<th>Source</th>
<th>Uncertainty in signal yield [%]</th>
</tr>
</thead>
<tbody>
<tr>
<td>Integrated luminosity</td>
<td>2.8</td>
</tr>
<tr>
<td>MC statistics</td>
<td>1.0</td>
</tr>
<tr>
<td>Trigger efficiency</td>
<td>1.0</td>
</tr>
<tr>
<td>Photon ID efficiency</td>
<td>3.0</td>
</tr>
<tr>
<td>Photon isolation efficiency</td>
<td>0.3–2.1</td>
</tr>
<tr>
<td>Total</td>
<td>≈ 5</td>
</tr>
</tbody>
</table>

Figure 3: Observed invariant mass distribution of the selected diphoton events (black dots; the vertical and horizontal axes on logarithmic scales). To compensate for the rapid decrease of the spectrum, the bins are chosen to have constant logarithmic width. Specifically, the ratio of the upper to lower bin boundary is equal to 1.038 for all bins, and the first bin starts at 179 GeV. Superimposed are the SM background prediction including irreducible and reducible components and two examples of signal predictions. The low-mass control region is indicated by an arrow. (Bottom) Bin-by-bin significance of the difference between data and background expectation.
Table 3: Number of events expected from the reducible and irreducible background components as well as the total background and observed number of events in different mass bins. Each bin boundary in the table corresponds to a bin boundary in Fig. 3, but the binning of the table is more coarse than in the figure. Since the combined statistical and systematic uncertainties are strongly anticorrelated between the two components, the uncertainty in the total background prediction is smaller than the sum in quadrature of the uncertainties of the individual components.

<table>
<thead>
<tr>
<th>Mass window [GeV]</th>
<th>Background expectation (number of events)</th>
<th>Observed events</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Irreducible</td>
<td>Reducible</td>
</tr>
<tr>
<td>[179, 409]</td>
<td>23800 ± 2400</td>
<td>9100 ± 2400</td>
</tr>
<tr>
<td>Control region</td>
<td>1465</td>
<td>524</td>
</tr>
<tr>
<td>[409, 513]</td>
<td>1070 ± 110</td>
<td>400 ± 100</td>
</tr>
<tr>
<td>[513, 596]</td>
<td>369 ± 37</td>
<td>129 ± 34</td>
</tr>
<tr>
<td>[596, 719]</td>
<td>240 ± 24</td>
<td>74 ± 20</td>
</tr>
<tr>
<td>[719, 805]</td>
<td>75.8 ± 7.7</td>
<td>20.6 ± 5.5</td>
</tr>
<tr>
<td>[805, 901]</td>
<td>46.6 ± 4.8</td>
<td>11.5 ± 3.2</td>
</tr>
<tr>
<td>[901, 1009]</td>
<td>28.2 ± 3.0</td>
<td>6.3 ± 1.8</td>
</tr>
<tr>
<td>[1009, 1129]</td>
<td>16.8 ± 1.9</td>
<td>3.4 ± 1.0</td>
</tr>
<tr>
<td>[1129, 1217]</td>
<td>6.92 ± 0.89</td>
<td>1.35 ± 0.46</td>
</tr>
<tr>
<td>[1217, 1312]</td>
<td>4.85 ± 0.73</td>
<td>0.88 ± 0.39</td>
</tr>
<tr>
<td>[1312, 1415]</td>
<td>3.11 ± 0.54</td>
<td>0.58 ± 0.28</td>
</tr>
<tr>
<td>[1415, 1644]</td>
<td>3.39 ± 0.59</td>
<td>0.61 ± 0.29</td>
</tr>
<tr>
<td>[1644, 3000]</td>
<td>2.12 ± 0.61</td>
<td>0.41 ± 0.22</td>
</tr>
</tbody>
</table>

6 Results

The observed diphoton mass spectrum is shown in Fig. 3, together with the background expectation and the predicted signal for two examples of values of the RS model parameters. The acceptance (i.e. the fraction of simulated events in which both photons satisfy the requirements on $|\eta|$ and $E_T^\gamma$ described above) for the RS scenario with $m_{G^*} = 2.0$ TeV and $k/M_{Pl} = 0.1$ is 89.3% and the selection efficiency for events within acceptance is (55.9 ± 2.6)%. The bottom inset in Fig. 3 shows, bin-by-bin, the statistical significance of the difference between the data and the expected background. This significance is calculated using the prescription in Ref. [32] and is plotted as positive (negative) for bins with an excess (deficit) of data. The observed event yields and expected backgrounds are also summarized in Table 3. No significant deviation of the data from the expected background is observed. To further quantify the level of agreement, the BUMP HUNTER [33] algorithm is used to search for any “bump” on top of the expected background, i.e. any excess of events localized in a relatively narrow mass region. Specifically, the BUMP HUNTER algorithm scans the mass distribution in windows of progressively increasing width in order to identify the window with the most significant excess of observed events over the background anywhere in a predefined search region. In this analysis, the binned $m_{\gamma\gamma}$ spectrum from Fig. 3 is used, the search region is defined as 409 < $m_{\gamma\gamma}$ < 3000 GeV, and the window size is allowed to vary from one bin to half the number of bins in the search region. The most significant excess is found in the region 494 < $m_{\gamma\gamma}$ < 745 GeV. The probability to observe anywhere in the search region an excess that is at least as significant as the one observed in data and that arises from fluctuations in the background is 58%. This confirms the absence of any significant excess in the data.
Table 4: Expected and observed lower limits at 95% CL on $m_{G^*}$ for different values of $k/M_{Pl}$.

<table>
<thead>
<tr>
<th>$k/M_{Pl}$</th>
<th>$-2\sigma$</th>
<th>$-1\sigma$</th>
<th>central</th>
<th>$+1\sigma$</th>
<th>$+2\sigma$</th>
<th>Observed limit [TeV]</th>
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</thead>
<tbody>
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<td>0.010</td>
<td>1.45</td>
<td>1.38</td>
<td>1.30</td>
<td>1.20</td>
<td>1.11</td>
<td>1.41</td>
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<tr>
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<td>1.80</td>
<td>1.78</td>
<td>1.70</td>
<td>1.60</td>
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<td>1.62</td>
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<tr>
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<tr>
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As no evidence for a signal is found, upper limits at 95% confidence level (CL) are set on the production cross section times branching ratio, $\sigma \times BR(G^* \rightarrow \gamma\gamma)$, for the lightest RS graviton. A Bayesian approach is used to calculate the limits, and the corresponding computer program is implemented using the toolkit from Refs. [34, 35]. A likelihood function is used to describe the probability to obtain the $m_{\gamma\gamma}$ distribution observed in data, given the background and signal predictions. The likelihood is defined as the product of the per-bin probabilities over all bins in the search region ($m_{\gamma\gamma} > 409$ GeV) in Fig. 3. In each bin, the probability is computed from a Poisson distribution given the observed number of events and the model prediction, where the model prediction is the sum of the expected numbers of background and signal events. The expected signal yield depends on $m_{G^*}$ and $\sigma \times BR(G^* \rightarrow \gamma\gamma)$. Using Bayes’ theorem, the likelihood is converted into a posterior probability. A uniform positive prior in $\sigma \times BR(G^* \rightarrow \gamma\gamma)$ is assumed, and systematic uncertainties are incorporated using nuisance parameters with Gaussian priors. For a given value of $m_{G^*}$, the posterior probability is reduced to one parameter of interest, $\sigma \times BR(G^* \rightarrow \gamma\gamma)$, and the 95% CL limit on $\sigma \times BR(G^* \rightarrow \gamma\gamma)$ is obtained via integration of the posterior probability. The resulting limits on $\sigma \times BR(G^* \rightarrow \gamma\gamma)$, as a function of the mass hypothesis (value of $m_{G^*}$), are shown in Fig. 4. The phenomenology of the RS model can be described in terms of the parameters $m_{G^*}$ and $k/M_{Pl}$, i.e. $\sigma \times BR(G^* \rightarrow \gamma\gamma)$ can be calculated [36, 37] for a given pair of values for these parameters. The results of this calculation are represented by the lines for given values of $k/M_{Pl}$ in Fig. 4. They are obtained using the Pythia generator, which implements the calculations from Ref. [37]. A $K$-factor is applied to these predictions to account for QCD corrections at NLO. The value of the $K$-factor varies between 1.6 and 1.9, depending on $m_{G^*}$ and $k/M_{Pl}$. These results for the $K$-factor were provided by the authors of Refs. [38, 39], using updated calculations at $\sqrt{s} = 8$ TeV. Given these theory predictions for $\sigma \times BR(G^* \rightarrow \gamma\gamma)$, the upper limits on $\sigma \times BR(G^* \rightarrow \gamma\gamma)$ from Fig. 4 can be translated into upper limits on $k/M_{Pl}$. The results are shown, as a function of $m_{G^*}$, in Fig. 5. Alternatively, the upper limits on $\sigma \times BR(G^* \rightarrow \gamma\gamma)$ can be translated into lower limits on the graviton mass. The results are shown in Table 4, for different values of $k/M_{Pl}$. The limits presented in Fig. 5 and in Table 4 do not take into account any theory uncertainties in the prediction of $\sigma \times BR(G^* \rightarrow \gamma\gamma)$ for the RS signal. The size of one of the dominant theory uncertainties (the PDF uncertainty) is included in Fig. 4; uncertainties due to missing higher-order contributions in the calculation of the $K$-factor are estimated to be of comparable size.
7 Conclusions

In summary, a search for high-mass diphoton resonances has been performed using the full 2012 dataset collected by the ATLAS detector at the LHC (20.3 fb$^{-1}$ of pp collisions at $\sqrt{s} = 8$ TeV). No significant excess over the expected background is observed, and upper limits on the production cross section times branching fraction $\sigma \times \text{BR}(G^* \rightarrow \gamma\gamma)$ for narrow resonances are reported as a function of the resonance mass. High-mass diphoton resonances are predicted e.g. in models that postulate the existence of extra spatial dimensions in order to address the hierarchy problem, and could be observed at the LHC. Limits on the mass of the lightest graviton state in the framework of the Randall-Sundrum model of extra dimensions are derived from the limits on $\sigma \times \text{BR}(G^* \rightarrow \gamma\gamma)$. The only free parameters of this model are the mass of the lightest graviton ($m_{G^*}$) and the coupling to the SM fields ($k/M_{Pl}$), i.e. $\sigma \times \text{BR}(G^* \rightarrow \gamma\gamma)$ can be calculated once the values of these two parameters are specified. The limits can therefore be expressed in terms of $m_{G^*}$ for a given value of $k/M_{Pl}$. A lower limit of 2.66 (1.41) TeV at 95% confidence level is obtained on the mass of the lightest graviton for a coupling $k/M_{Pl} = 0.1$ (0.01). The results reported in this article constitute a significant improvement over the results in Ref. [5] obtained at $\sqrt{s} = 7$ TeV. The upper limits on $\sigma \times \text{BR}(G^* \rightarrow \gamma\gamma)$ are reduced by a factor of 2.1 (3.6) for a resonance with a mass of 0.5 TeV (2.75 TeV). Furthermore, the new limits are expressed in terms of the production cross section at $\sqrt{s} = 8$ TeV, compared to $\sqrt{s} = 7$ TeV for the old limits, and the production cross section for any new heavy particle is expected to be larger at $\sqrt{s} = 8$ TeV. The lower limits on the mass of the lightest
Figure 5: Expected and observed upper limits on $k/\mathcal{M}_{Pl}$ expressed at 95% CL, as a function of the graviton mass.

RS graviton are also significantly improved compared to the previous analysis.

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References

The ATLAS Collaboration

Physics Department, National Technical University of Athens, Zografou, Greece
Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
Institut de Física d’Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain
Institute of Physics, University of Belgrade, Belgrade, Serbia
Department for Physics and Technology, University of Bergen, Bergen, Norway
Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley CA, United States of America
Department of Physics, Humboldt University, Berlin, Germany
Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
Department of Physics, Bogazici University, Istanbul; Department of Physics, Dogus University, Istanbul; Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey
INFN Sezione di Bologna; Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
Department of Physics, Boston University, Boston MA, United States of America
Department of Physics, Brandeis University, Waltham MA, United States of America
Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; Electrical Circuits Department, Federal University of Juiz de Fora (UFJF), Juiz de Fora; Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei; Instituto de Física, Universidade de Sao Paulo, Sao Paulo, Brazil
Physics Department, Brookhaven National Laboratory, Upton NY, United States of America
National Institute of Physics and Nuclear Engineering, Bucharest; National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca; University Politehnica Bucharest, Bucharest; West University in Timisoara, Timisoara, Romania
Department of Physics, Universidad de Buenos Aires, Buenos Aires, Argentina
Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
Department of Physics, Carleton University, Ottawa ON, Canada
CERN, Geneva, Switzerland
Enrico Fermi Institute, University of Chicago, Chicago IL, United States of America
Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
Department of Modern Physics, University of Science and Technology of China, Anhui; Department of Physics, Nanjing University, Jiangsu; School of Physics, Shandong University, Shandong; Department of Physics and Astronomy, Shanghai Key Laboratory for Particle Physics and Cosmology, Shanghai Jiao Tong University, Shanghai; Physics Department, Tsinghua University, Beijing 100084, China
Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France
Nevis Laboratory, Columbia University, Irvington NY, United States of America
Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark
INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati; Dipartimento di Fisica, Università della Calabria, Rende, Italy
AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland
Louisiana Tech University, Ruston LA, United States of America

Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France

Fysiska institutionen, Lunds universitet, Lund, Sweden

Departamento de Física Teórica C-15, Universidad Autonoma de Madrid, Madrid, Spain

Institut für Physik, Universität Mainz, Mainz, Germany

School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom

CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France

Department of Physics, University of Massachusetts, Amherst MA, United States of America

Department of Physics, McGill University, Montreal QC, Canada

School of Physics, University of Melbourne, Victoria, Australia

Department of Physics, The University of Michigan, Ann Arbor MI, United States of America

Department of Physics and Astronomy, Michigan State University, East Lansing MI, United States of America

(a) INFN Sezione di Milano; (b) Dipartimento di Fisica, Università di Milano, Milano, Italy

B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus

National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus

Department of Physics, Massachusetts Institute of Technology, Cambridge MA, United States of America

Group of Particle Physics, University of Montreal, Montreal QC, Canada

P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia

Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia

National Research Nuclear University MEPhI, Moscow, Russia

D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia

Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany

Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany

Nagasaki Institute of Applied Science, Nagasaki, Japan

Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan

(a) INFN Sezione di Napoli; (b) Dipartimento di Fisica, Università di Napoli, Napoli, Italy

Department of Physics and Astronomy, University of New Mexico, Albuquerque NM, United States of America

Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands

Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands

Department of Physics, Northern Illinois University, DeKalb IL, United States of America

Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia

Department of Physics, New York University, New York NY, United States of America

Ohio State University, Columbus OH, United States of America

Faculty of Science, Okayama University, Okayama, Japan

Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK, United States of America

Department of Physics, Oklahoma State University, Stillwater OK, United States of America

Palacký University, RCPTM, Olomouc, Czech Republic
Center for High Energy Physics, University of Oregon, Eugene OR, United States of America
LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
Graduate School of Science, Osaka University, Osaka, Japan
Department of Physics, University of Oslo, Oslo, Norway
Department of Physics, Oxford University, Oxford, United Kingdom
(a) INFN Sezione di Pavia; (b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
Department of Physics, University of Pennsylvania, Philadelphia PA, United States of America
Petersburg Nuclear Physics Institute, Gatchina, Russia
(a) INFN Sezione di Pisa; (b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA, United States of America
(a) INFN Sezione di Roma; (b) Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy
(a) INFN Sezione di Roma Tor Vergata; (b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
(a) INFN Sezione di Roma Tre; (b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
(a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca; (b) Centre National de l’Energie des Sciences Techniques Nucleaires, Rabat; (c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech; (d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; (e) Faculté des sciences, Université Mohammed V-Agdal, Rabat, Morocco
DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France
Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA, United States of America
Department of Physics, University of Washington, Seattle WA, United States of America
Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
Department of Physics, Shinshu University, Nagano, Japan
Fachbereich Physik, Universität Siegen, Siegen, Germany
Department of Physics, Simon Fraser University, Burnaby BC, Canada
SLAC National Accelerator Laboratory, Stanford CA, United States of America
(a) Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
(a) Department of Physics, University of Cape Town, Cape Town; (b) Department of Physics,
University of Johannesburg, Johannesburg; (c) School of Physics, University of the Witwatersrand, Johannesburg, South Africa
147 (a) Department of Physics, Stockholm University; (b) The Oskar Klein Centre, Stockholm, Sweden
148 Physics Department, Royal Institute of Technology, Stockholm, Sweden
149 Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook NY, United States of America
150 Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
151 School of Physics, University of Sydney, Sydney, Australia
152 Institute of Physics, Academia Sinica, Taipei, Taiwan
153 Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel
154 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
155 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
156 International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan
157 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
158 Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
159 Department of Physics, University of Toronto, Toronto ON, Canada
160 (a) TRIUMF, Vancouver BC; (b) Department of Physics and Astronomy, York University, Toronto ON, Canada
161 Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan
162 Department of Physics and Astronomy, Tufts University, Medford MA, United States of America
163 Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia
164 Department of Physics and Astronomy, University of California Irvine, Irvine CA, United States of America
165 (a) INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; (b) ICTP, Trieste; (c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy
166 Department of Physics, University of Illinois, Urbana IL, United States of America
167 Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
168 Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain
169 Department of Physics, University of British Columbia, Vancouver BC, Canada
170 Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada
171 Department of Physics, University of Warwick, Coventry, United Kingdom
172 Waseda University, Tokyo, Japan
173 Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel
174 Department of Physics, University of Wisconsin, Madison WI, United States of America
175 Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany
176 Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany
177 Department of Physics, Yale University, New Haven CT, United States of America
178 Yerevan Physics Institute, Yerevan, Armenia
179 Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France
a Also at Department of Physics, King’s College London, London, United Kingdom
b Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
c Also at Novosibirsk State University, Novosibirsk, Russia
Also at TRIUMF, Vancouver BC, Canada

Also at Department of Physics, California State University, Fresno CA, United States of America

Also at Department of Physics, University of Fribourg, Fribourg, Switzerland

Also at Departamento de Física e Astronomia, Faculdade de Ciencias, Universidade do Porto, Portugal

Also at Tomsk State University, Tomsk, Russia

Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France

Also at Università di Napoli Parthenope, Napoli, Italy

Also at Institute of Particle Physics (IPP), Canada

Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom

Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia

Also at Louisiana Tech University, Ruston LA, United States of America

Also at Institut Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain

Also at Department of Physics, National Tsing Hua University, Taiwan

Also at Department of Physics, The University of Texas at Austin, Austin TX, United States of America

Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia

Also at Georgian Technical University (GTU), Tbilisi, Georgia

Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan

Also at Manhattan College, New York NY, United States of America

Also at Institute of Physics, Academia Sinica, Taipei, Taiwan

Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France

Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan

Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia

Also at Section de Physique, Université de Genève, Geneva, Switzerland

Also at International School for Advanced Studies (SISSA), Trieste, Italy

Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, United States of America

Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China

Also at Faculty of Physics, M.V.Lomonosov Moscow State University, Moscow, Russia

Also at National Research Nuclear University MEPhI, Moscow, Russia

Also at Department of Physics, Stanford University, Stanford CA, United States of America

Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary

Also at Department of Physics, Oxford University, Oxford, United Kingdom

Also at Department of Physics, The University of Michigan, Ann Arbor MI, United States of America

Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa

Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia

* Deceased