The following full text is a publisher's version.

For additional information about this publication click this link.
http://hdl.handle.net/2066/143915

Please be advised that this information was generated on 2019-12-23 and may be subject to change.
Centrality and rapidity dependence of inclusive jet production in $\sqrt{s_{NN}} = 5.02$ TeV proton–lead collisions with the ATLAS detector

ATLAS Collaboration*

**Abstract**

Measurements of the centrality and rapidity dependence of inclusive jet production in $\sqrt{s_{NN}} = 5.02$ TeV proton–lead ($p + Pb$) collisions and the jet cross-section in $\sqrt{s} = 2.76$ TeV proton–proton collisions are presented. These quantities are measured in datasets corresponding to an integrated luminosity of 278 nb$^{-1}$ and 4.0 pb$^{-1}$, respectively, recorded with the ATLAS detector at the Large Hadron Collider in 2013. The $p + Pb$ collision centrality was characterised using the total transverse energy measured in the pseudorapidity interval $-4.9 < \eta < -3.2$ in the direction of the lead beam. Results are presented for the double-differential per-collision yields as a function of jet rapidity and transverse momentum ($p_T$) for minimum-bias and centrality-selected $p + Pb$ collisions, and are compared to the jet rate from the geometric expectation. The total jet yield in minimum-bias events is slightly enhanced above the expectation in a $p_T$-dependent manner but is consistent with the expectation within uncertainties. The ratios of jet spectra from different centrality selections show a strong modification of jet production at all $p_T$ at forward rapidities and for large $p_T$ at mid-rapidity, which manifests as a suppression of the jet yield in central events and an enhancement in peripheral events. These effects imply that the factorisation between hard and soft processes is violated at an unexpected level in proton–nucleus collisions. Furthermore, the modifications at forward rapidities are found to be a function of the total energy only, implying that the violations may have a simple dependence on the hard parton–parton kinematics.

© 2015 CERN for the benefit of the ATLAS Collaboration. Published by Elsevier B.V. This is an open access article under the CC BY license (http://creativecommons.org/licenses/by/4.0/). Funded by SCOAP3.

1. Introduction

Proton–lead ($p + Pb$) collisions at the Large Hadron Collider (LHC) provide an excellent opportunity to study hard scattering processes involving a nuclear target [1]. Measurements of jet production in $p + Pb$ collisions provide a valuable benchmark for studies of jet quenching in lead–lead collisions by, for example, constraining the impact of nuclear parton distributions on inclusive jet yields. However, $p + Pb$ collisions also allow the study of possible violations of the QCD factorisation between hard and soft processes which may be enhanced in collisions involving nuclei.

Previous studies in deuteron–gold ($d + Au$) collisions at the Relativistic Heavy Ion Collider (RHIC) observed such violations, manifested in the suppressed production of very forward hadrons with transverse momenta up to 4 GeV [2–4]. Studies of forward di-hadron angular correlations at RHIC also showed a much weaker dijet signal in $d + Au$ collisions than in $p p$ collisions [4,5]. These effects have been attributed to the saturation of the parton distributions in the gold nucleus [6–8], to the modification of the nuclear parton distribution function [9], to the higher-twist contributions to the cross-section enhanced by the forward kinematics of the measurement [10], or to the presence of a large nucleus [11]. The extended kinematic reach of $p + Pb$ measurements at the LHC allows the study of hard scattering processes that produce forward hadrons or jets over a much wider rapidity and transverse momentum range. Such measurements can determine whether the factorisation violations observed at RHIC persist at higher energy and, if so, how the resulting modifications vary as a function of particle or jet momentum and rapidity. The results of such measurements could test the competing descriptions of the RHIC results and, more generally, provide new insight into the physics of hard scattering processes involving a nuclear target.

This paper reports the centrality dependence of inclusive jet production in $p + Pb$ collisions at a nucleon–nucleon centre-of-mass energy $\sqrt{s_{NN}} = 5.02$ TeV. The measurement was performed using a dataset corresponding to an integrated luminosity of 278 nb$^{-1}$ recorded in 2013. The $p + Pb$ jet yields were
compared to a nucleon–nucleon reference constructed from a measurement of jet production in pp collisions at a centre-of-mass energy $\sqrt{s} = 2.76$ TeV using a dataset corresponding to an integrated luminosity of 4.0 pb$^{-1}$ also recorded in 2013. Jets were reconstructed from energy deposits measured in the calorimeter using the anti-$k_t$ algorithm with radius parameter $R = 0.4$ [12].

The centrality of $p + Pb$ collisions was characterised using the total transverse energy measured in the pseudorapidity$^1$ interval $-4.9 < \eta < -3.2$ in the direction of the lead beam. Whereas in nucleus–nucleus collisions centrality reflects the degree of nuclear overlap between the colliding nuclei, centrality in $p + Pb$ collisions is sensitive to the multiple interactions between the proton and nucleons in the lead nucleus. Centrality has been successfully used at lower energies in $d + Au$ collisions at RHIC as an experimental handle on the collision geometry [2,13,14].

A Glauber model [15] was used to determine the average number of nucleon–nucleon collisions, $(N_{\text{coll}})$, and the mean value of the overlap function, $T_{pA}(b) = \int_0^\infty \rho(b, z) dz$, where $\rho(b, z)$ is the nucleon density at impact parameter $b$ and longitudinal position $z$, in each centrality interval. Per-event jet yields, $(1/N_{\text{coll}})(d^2 N_{\text{jet}}/dp_T dy^\ast)$, were measured as a function of jet centre-of-mass rapidity, $y^\ast$, and transverse momentum, $p_T$, where $N_{\text{jet}}$ is the number of jets measured in $N_{\text{coll}}$ $p + Pb$ events analysed. The centrality dependence of the per-event jet yields was evaluated using the nuclear modification factor,

$$R_{pPb} \equiv \frac{1}{T_{pA}} \frac{(1/N_{\text{coll}}) d^2 N_{\text{jet}}/dp_T dy^\ast}{d^2 \sigma_{pp}/dp_T dy^\ast} \bigg|_{\text{cent}},$$

(1)

for a given centrality selection “cent”, where $d^2 \sigma_{pp}/dp_T dy^\ast$ is determined using the jet cross-section measured in pp collisions at $\sqrt{s} = 2.76$ TeV. The factor $R_{pPb}$ quantifies the absolute modification of the jet rate relative to the geometric expectation. In each centrality interval, the geometric expectation is the jet rate that would be produced by an incoherent superposition of the number of nucleon–nucleon collisions corresponding to the mean nuclear thickness in the given class of $p + Pb$ collisions.

Results are also presented for the central-to-peripheral ratio,

$$R_{CP} \equiv \frac{1}{T_{\text{coll}}} \frac{(1/N_{\text{coll}}) d^2 N_{\text{jet}}/dp_T dy^\ast}{d^2 \sigma_{pp}/dp_T dy^\ast} \bigg|_{\text{per}},$$

(2)

where $T_{\text{coll}}$ represents the ratio of $(N_{\text{coll}})$ in a given centrality interval to that in the most peripheral interval, $T_{\text{coll}} \equiv (N_{\text{cent}}/N_{\text{per}})$. The $R_{CP}$ ratio is sensitive to relative deviations in the jet rate from the geometric expectation between the $p + Pb$ event centralties. The $R_{pPb}$ and $R_{CP}$ measurements are presented as a function of inclusive jet $y^\ast$ and $p_T$.

For the 2013 $p + Pb$ run, the LHC was configured with a 4 TeV proton beam and a 1.57 TeV per-nucleon Pb beam that together produced collisions with $\sqrt{s_{NN}} = 5.02$ TeV and a rapidity shift of the centre-of-mass frame of 0.465 units relative to the ATLAS rest frame. The run was split into two periods, with the directions of the proton and lead beams being reversed at the end of the first period. The first period provided approximately 55% of the integrated luminosity with the Pb beam travelling to positive rapidity and the proton beam to negative rapidity, and the second period provided the remainder with the beams reversed. The analysis in this paper uses the events from both periods of data-taking and $y^\ast$ is defined so that $y^\ast > 0$ always refers to the downstream proton direction.

2. Experimental setup

The measurements presented in this paper were performed using the ATLAS inner detector (ID), calorimeters, minimum-bias trigger scintillator (MBTS), and trigger and data acquisition systems [16]. The ID measures charged particles within $|\eta| < 2.5$ using a combination of silicon pixel detectors, silicon microstrip detectors, and a straw-tube transition radiation tracker, all immersed in a 2 T axial magnetic field [17]. The calorimeter system consists of a liquid argon (LAr) electromagnetic (EM) calorimeter covering $|\eta| < 3.2$, a steel/scintillator sampling hadronic calorimeter covering $|\eta| < 1.7$, a LAr hadronic calorimeter covering $1.5 < |\eta| < 3.2$, and two LAr electromagnetic and hadronic forward calorimeters (FCal) covering $3.2 < |\eta| < 4.9$. The EM calorimeters use lead plates as the absorbers and are segmented longitudinally in shower depth into three compartments with an additional presampler layer in front for $|\eta| < 1.8$. The granularity of the EM calorimeter varies with layer and pseudorapidity. The middle sampling layer, which typically has the largest energy deposit in EM showers, has a $\Delta \eta \times \Delta \phi$ granularity of 0.025 $\times$ 0.025 within $|\eta| < 2.5$. The hadronic calorimeter uses steel as the absorber and has three segments longitudinal in shower depth with cell sizes $\Delta \eta \times \Delta \phi = 0.1 \times 0.1$ for $|\eta| < 2.5^1$ and 0.2 $\times$ 0.2 for $2.5 < |\eta| < 4.9$. The two FCal modules are composed of tungsten and copper absorbers with LAr as the active medium, which together provide ten interaction lengths of material. The MBTS detects charged particles over 2.1 $< |\eta| < 3.9$ using two hodoscopes of 16 counters each, positioned at $z = \pm 3.6$ m.

The $p + Pb$ and pp events used in this analysis were recorded using a combination of minimum-bias (MB) and jet triggers [18]. In $p + Pb$ data-taking, the MB trigger required hits in at least one counter in each side of the MBTS detector. In pp collisions the MB condition was the presence of hits in the pixel and microstrip detectors reconstructed as a track by the high-level trigger system. Jets were selected using high-level jet triggers implemented with a reconstruction algorithm similar to the procedure applied in the offline analysis. In particular, it used the anti-$k_t$ algorithm with $R = 0.4$, a background subtraction procedure, and a calibration of the jet energy to the full hadronic scale. The high-level jet triggers were seeded from a combination of low-level MB and jet hardware-based triggers. Six jet triggers with transverse energy thresholds ranging from 20 GeV to 75 GeV were used to select jets within $|\eta| < 3.2$ and a separate trigger with a threshold of 15 GeV was used to select jets with $3.2 < |\eta| < 4.9$. The triggers were prestcaled in a fashion which varied with time to accommodate the evolution of the luminosity within an LHC fill.

3. Data selection

In the offline analysis, charged-particle tracks were reconstructed in the ID with the same algorithm used in pp collisions [19]. The $p + Pb$ events used for this analysis were required to have

---

1 ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z-axis along the beam pipe. The x-axis points from the IP to the centre of the LHC ring, and the y-axis points upward. Cylindrical coordinates (r, \phi) are used in the transverse plane, \phi being the azimuthal angle around the beam pipe. The pseudorapidity is defined in laboratory coordinates in terms of the polar angle \theta as \eta = -\ln \tan(\theta/2). During 2013 \ p + Pb data-taking, the beam directions were reversed approximately half-way through the running period, but in presenting results the direction of the proton beam is always chosen to point to positive \eta.

2 The jet rapidity $y^\ast$ is defined as $y^\ast = 0.5 \ln \frac{E + p_z}{E - p_z}$ where $E$ and $p_z$ are the energy and the component of the momentum along the proton beam direction in the nucleon–nucleon centre-of-mass frame.

3 An exception is the third (outermost) sampling layer, which has a segmentation of $0.2 \times 0.1$ up to $|\eta| = 1.7$. 
a reconstructed vertex containing at least two associated tracks with $p_T > 0.1$ GeV, at least one hit in each of the two MBTS hodoscopes, and a difference between times measured on the two MBTS sides of less than 10 ns. Events containing multiple $p + Pb$ collisions (pileup) were suppressed by rejecting events having two or more reconstructed vertices, each associated with reconstructed tracks with a total transverse momentum scalar sum of at least 5 GeV. The fraction of events with one $p + Pb$ interaction rejected by this requirement was less than 0.1%. Events with a pseudorapidity gap (defined by the absence of clusters in the calorimeter with more than 0.2 GeV of transverse energy) of greater than two units on the Pb-going side of the detector were also removed from the analysis. Such events arise primarily from electromagnetic or diffractive excitation of the proton. After accounting for event selection, the number of $p + Pb$ events sampled by the highest-luminosity jet trigger (which was unprescaled) was 53 billion. The event selection criteria described here were designed to select a sample of $p + Pb$ events to which a centrality analysis can be applied and for which meaningful geometric parameters can be determined.

The $pp$ events used in this analysis were required to have a reconstructed vertex, with the same definition as the vertices in $p + Pb$ events above. No other requirements were applied.

4. Centrality determination

The centrality of the $p + Pb$ events selected for analysis was characterised by the total transverse energy $\Sigma E^p_T$ in the FCal module on the Pb-going side. The $\Sigma E^p_T$ distribution for minimum-bias $p + Pb$ collisions passing the event selection described in Section 3 is presented in Fig. 1. Following standard techniques [20], centrality intervals were defined in terms of percentiles of the $\Sigma E^p_T$ distribution after accounting for an estimated inefficiency of $(2 \pm 2)\%$ for inelastic $p + Pb$ collisions to pass the applied event selection. The following centrality intervals were used in this analysis, in order from the most central to the most peripheral: 0–10%, 10–20%, 20–30%, 30–40%, 40–60%, and 60–90%, with the 60–90% interval serving as the reference in the RCP ratio. Events with a centrality beyond 90% were not used in the analysis, since the uncertainties on the composition of the event sample and in the determination of the geometric quantities are large for these events.

A Glauber Monte Carlo (MC) [15] analysis was used to calculate $R_{coll}$ and $T_{PA}$ for each centrality interval. First, a Glauber MC program [21] was used to simulate the geometry of inelastic $p + Pb$ collisions and calculate the probability distribution of the number of nucleon participants $N_{part}$, $P(N_{part})$. The simulations used a Woods–Saxon nucleon density distribution and an inelastic nucleon–nucleon cross-section, $\sigma_{NN}$, of 70 ± 5 mb. Separately, PYTHIA 8 [22,23] simulations of 4 TeV on 1.57 TeV pp collisions provided a detector-level $\Sigma E^p_T$ distribution for nucleon–nucleon collisions, to be used as input to the Glauber model. This distribution was fit to a gamma distribution.

Then, an extension of the wounded-nucleon (WN) [24] model that included a non-linear dependence of $\Sigma E^p_T$ on $N_{part}$ was used to define $N_{part}$-dependent gamma distributions for $\Sigma E^p_T$, with the constraint that the distributions reduce to the PYTHIA distribution for $N_{part} = 2$. The non-linear term accounted for the possible variation of the effective FCal acceptance resulting from an $N_{part}$-dependent backward rapidity shift of the produced soft particles with respect to the nucleon–nucleon frame [25]. The gamma distributions were summed over $N_{part}$ with a $P(N_{part})$ weighting to produce a hypothetical $\Sigma E^p_T$ distribution. That distribution was fit to the measured $\Sigma E^p_T$ distribution shown in Fig. 1 with the parameters of the extended WN model allowed to vary freely. The best fit, which contained a significant non-linear term, successfully described the $\Sigma E^p_T$ distribution in data over several orders of magnitude. From the results of the fit, the distribution of $N_{part}$ values and the corresponding $[N_{part}]$ were calculated for each centrality interval. The resulting $R_{coll}$ and $T_{PA}$ values and correspondng systematic uncertainties, which are described in Section 8, are shown in Table 1.

5. Monte Carlo simulation

The performance of the jet reconstruction procedure was evaluated using a sample of 36 million events in which simulated $\sqrt{s} = 5.02$ TeV $pp$ hard-scattering events were overlaid with minimum-bias $p + Pb$ events recorded during the 2013 run. Thus the sample contains an underlying event contribution that is identical in all respects to the data. The simulated events were generated using PYTHIA 8 [22] (version 6.425, AUET2B tune [26], CTEQ6L1 parton distribution functions [27]) and the detector effects were fully simulated using GEANT4 [28,29]. These events were produced for different $p_T$ intervals of the generator-level (“truth”) $R = 0.4$ jets. In total, the generator-level spectrum spans $10 < p_T < 10^3$ GeV. Separate sets of 18 million events each were generated for the two different beam directions to take into account any $z$-axis asymmetries in the detector. For each beam direction, the four-momenta of the generated particles were longitudinally boosted by a rapidity of $\pm 0.465$ to match the corresponding beam conditions. The events were simulated using detector conditions appropriate to the two periods of the 2013 $p + Pb$ run and reconstructed using the same algorithms as were applied to the experimental data. A sep-
arate 9-million-event sample of fully simulated 2.76 TeV PYTHIA pp hard scattering events (with the same version, tune and parton distribution function set) was used to evaluate the jet performance in $\sqrt{s}=2.76$ TeV pp collisions during 2013 data-taking.

6. Jet reconstruction and performance

The jet reconstruction and underlying event subtraction procedures were adapted from those used by ATLAS in Pb+Pb collisions, which are described in detail in Refs. [30,31], and are summarised here along with any substantial differences from the referenced analyses.

An iterative procedure was used to obtain an event-by-event estimate of the underlying event energy density while excluding contributions from jets to that estimate. The modulation of the underlying event energy density to account for potential elliptic flow was not included in this analysis. Jets were reconstructed from the anti-$k_T$ algorithm with $R=0.4$ applied to calorimeter cells grouped into $\Delta y \times \Delta \phi = 0.1 \times 0.1$ towers, with the final jet kinematics calculated from the background-subtracted energy in the cells contained in the jet. The rate of jets reconstructed from the underlying event fluctuations of soft particles was negligible in the kinematic range studied and therefore no attempt to reject them was made. The mean subtracted transverse energy in $p+p$ collisions was 2.4 GeV (1.4 GeV) for jets with $|y^+|<1$ ($y^+>3$). In pp collisions, this procedure simply subtracts the underlying event pedestal deposited in the calorimeter which can arise, in part, from the presence of additional pp interactions in the same crossing (in-time pileup).

Following the above jet reconstruction, a small correction, typically a few percent, was applied to the transverse momentum of those jets which did not overlap with a region excluded from the background determination and thus were erroneously included in the initial estimate of the underlying event background. Then, the jet energies were corrected to account for the calorimeter energy response using an $y$- and $p_T$-dependent multiplicative factor that was derived from the simulations [32]. Following this calibration, a final multiplicative in situ calibration was applied to account for differences between the simulated detector response and data. The measured $p_T$ of jets recoiling against objects with an independently calibrated energy scale — such as Z bosons, photons, or jets in a different region of the detector — was investigated. The in situ calibration, which typically differed from unity by a few percent, was derived by comparing this $p_T$ balance in pp data with that in simulations in a fashion similar to that used previously within ATLAS [33].

The jet reconstruction performance was evaluated in the simulated samples by applying the same subtraction and reconstruction procedure as was applied to data. The resulting reconstructed jets with transverse momentum $p_T^{\text{jet}}$ were compared with their corresponding generator jets, which were produced by applying the anti-$k_T$ algorithm to the final-state particles produced by PYTHIA, excluding muons and neutrinos. Each generator jet was matched to a reconstructed jet, and the $p_T$ difference between the two jets was studied as a function of the generator jet transverse momentum, $p_T^{\text{gen}}$, and generator jet rapidity $y^*$, and in the six $p+Pb$ event centrality intervals.

The reconstruction efficiency for jets having $p_T^{\text{gen}}$ > 25 GeV was found to be greater than 99%. The performance was quantified by the means and standard deviations of the $\Delta p_T/p_T$ ($= p_T^{\text{rec}}/p_T^{\text{gen}} - 1$) distributions, referred to as the jet energy scale closure and jet energy resolution respectively. The closure in $p+Pb$ events was less than 2% for $p_T^{\text{gen}} > 25$ GeV jets and was better than 1% for $p_T^{\text{gen}} > 100$ GeV jets. At low $p_T^{\text{gen}}$, the energy scale closure and resolution exhibited a weak $p+Pb$ centrality dependence, with differences in the closure of up to 1% and differences in the resolution of up to 2% in the most central 0–10% events relative to the 60–90% peripheral events. At high jet $p_T$, the response was centrality independent within sensitivity. In pp events, the closure was less than 1% in the entire kinematic range studied.

In order to quantify the degree of $p_T$-bin migration introduced by the detector response and reconstruction procedure, response matrices were populated by recording the $p_T$ values of each generator–reconstructed jet pair. Separate matrices were constructed for each $y^*$ interval and $p+Pb$ centrality interval used in the analysis. The $p_T$ bins used were chosen to increase with $p_T$ such that the width of each bin was $\approx 0.25$ of the bin low edge. Using this binning, the proportion of jets with reconstructed $p_T$ in the same bin as their truth $p_T$ monotonically increased with truth $p_T$ and was 50–70%.

7. Data analysis

A combination of minimum-bias and jet triggered $p+Pb$ events were selected for analysis as described in Section 2. The sample luminosity (defined as the luminosity divided by the mean luminosity-weighted prescale) of the jet triggers increased with increasing $p_T$ threshold. Offline jets were selected for the analysis by requiring a match to an online jet trigger. The efficiency of the various triggers was determined with respect to the minimum-bias trigger and to lower threshold jet triggers. For simplicity, each $p_T$ bin used jets selected by only one jet trigger. In a given $p_T$ bin, jets were selected by the highest-threshold jet trigger for which the efficiency was determined to be greater than 99% in the bin. No additional corrections for the trigger efficiency were applied.

The double-differential per-event jet yields in $p+Pb$ collisions were constructed via

$$
\frac{1}{N_{\text{ext}}} \frac{d^2N_{\text{jet}}}{dp_T dy^*} = \frac{1}{N_{\text{ext}}} \frac{N_{\text{jet}}}{\Delta p_T \Delta y^*},
$$

(3)

where $N_{\text{ext}}$ is the total (unprescaled) number of MB $p+Pb$ events sampled, $N_{\text{jet}}$ is the yield of jets corrected for all detector effects and the instantaneous trigger prescale during data-taking, and $\Delta p_T$ and $\Delta y^*$ are the widths of the $p_T$ and $y^*$ bins. The centrality-dependent yields were constructed by restricting $N_{\text{ext}}$ and $N_{\text{jet}}$ to come from $p+Pb$ events within a given centrality interval. The double-differential cross-section in pp collisions was constructed via

$$
\frac{d^2\sigma}{dp_T dy^*} = \frac{1}{L_{\text{int}}} \frac{N_{\text{jet}}}{\Delta p_T \Delta y^*},
$$

(4)

where $L_{\text{int}}$ is the total integrated luminosity of the jet trigger used in the given $p_T$ bin. The $p_T$ binning in the pp cross-section was chosen such that the $x_T = 2p_T/\sqrt{s}$ binning between the $p+Pb$ and $pp$ datasets is the same.

Both the per-event yields in $p+Pb$ collisions and the cross-section in pp collisions were restricted to the $p_T$ range where the MC studies described in Section 6 show that the efficiency for a truth jet to remain in the same $p_T$ bin is $\geq 50\%$. This $p_T$ range was rapidity dependent, with the lowest $p_T$ bin edge used ranging from 50 GeV in the most backward rapidity intervals studied to 25 GeV in the most forward intervals.

The measured $p+Pb$ and pp yields were corrected for jet energy resolution and residual distortions of the jet energy scale which result in $p_T$-bin migration. For each rapidity interval, the yield was corrected by the use of $p_T$-dependent (and, in the $p+Pb$ case, centrality-dependent) bin-by-bin correction factors $C(p_T, y^*)$ obtained from the ratio of the reconstructed to the truth jet $p_T$ distributions for jets originating in a true $y^*$ bin, according to
\[ C(p_T, y^*) = \frac{N_{\text{jet}}^{\text{true}}(p_T, y^*)}{N_{\text{jet}}^{\text{reco}}(p_T, y^*)}, \]

where \( N_{\text{jet}}^{\text{true}} \) \( (N_{\text{jet}}^{\text{reco}}) \) is the number of truth jets in the given \( p_T^{\text{true}} \) \( (p_T^{\text{reco}}) \) bin in the corresponding MC samples.

Since the determination of the correction factors \( C(p_T, y^*) \) is sensitive to the shape of the jet spectrum in the MC sample, the response matrices used to generate them were reweighted to provide a better match between the reconstructed distributions in data and simulated events. The spectrum of generator jets was weighted jet-by-jet by the ratio of the reconstructed spectrum in data to that in simulation. This ratio was found to be approximately linear in the logarithm of reconstructed \( p_T \). A separate reweighting was performed for the \( p + \text{Pb} \) jet yield in each centrality interval, resulting in changes of \( \lesssim 10\% \) from the original correction factors before reweighting. The resulting corrections to the \( p + \text{Pb} \) and \( pp \) yields were at most \( 30\% \), and were typically \( \leq 10\% \) for jets with \( p_T > 100 \text{ GeV} \). These corrections were applied to the detector-level yield \( N_{\text{jet}}^{\text{reco}} \) to give the particle-level yield via

\[ N_{\text{jet}} = C(p_T, y^*) N_{\text{jet}}^{\text{reco}}. \]

A \( \sqrt{s} = 5.02 \text{ TeV} \) \( pp \) reference jet cross-section was constructed through the use of the corrected 2.76 TeV \( pp \) cross-section and a previous ATLAS measurement of the \( x_T \)-scaling between the inclusive jet cross-sections at \( \sqrt{s} = 2.76 \text{ TeV} \) (measured using 0.20 pb\(^{-1} \) of data collected in 2011) and 7 TeV (measured using 37 pb\(^{-1} \) of data collected in 2010) \[34]\). In this previous analysis, the \( \sqrt{s} \)-scaled ratio \( \rho \) of the 2.76 TeV cross-section to that at 7 TeV was evaluated at fixed \( x_T \),

\[ \rho(x_T; y^*) = \left( \frac{2.76 \text{ TeV}}{7 \text{ TeV}} \right)^3 \frac{d^2\sigma^{2.76 \text{ TeV}} / dp_T dy^*}{d^2\sigma^{7 \text{ TeV}} / dp_T dy^*}. \]

where \( d^2\sigma / dp_T dy^* \) is the \( pp \) jet cross-section at the given centre-of-mass energy \( \sqrt{s} \), and the numerator and denominator are evaluated at the same \( x_T \) (but different \( p_T = x_T \sqrt{s}/2 \)). Equation (7) can be rearranged to define the cross-section at \( \sqrt{s} = 7 \text{ TeV} \) in terms of that at 2.76 TeV times a multiplicative factor and divided by \( \rho \).

The \( \sqrt{s} = 5.02 \text{ TeV} \) \( pp \) cross-section at each \( p_T \) and \( y^* \) value was constructed by scaling the corrected \( \sqrt{s} = 2.76 \text{ TeV} \) \( pp \) cross-section measured at the equivalent \( x_T \) according to

\[ \frac{d^2\sigma^{5.02 \text{ TeV}}}{dp_T dy^*} = \rho(x_T; y^*)^{-0.643} \left( \frac{2.76 \text{ TeV}}{5.02 \text{ TeV}} \right)^3 \frac{d^2\sigma^{2.76 \text{ TeV}}}{dp_T dy^*}. \]

where the power \(-\ln(2.76/5.02)/\ln(2.76/7) \approx -0.643\) interpolates between 2.76 TeV and 7 TeV to 5.02 TeV using a power-law collision energy dependence at each \( p_T \) and \( y^* \). Since the jet energy scale and \( x_T \)-interpolation uncertainties are large for the \( pp \) data at large rapidities (\( |y^*| > 2.8 \)), a \( \sqrt{s} = 5.02 \text{ TeV} \) reference is not constructed in that rapidity region.

The \( pp \) jet cross-section at \( \sqrt{s} = 2.76 \text{ TeV} \) measured with the 2013 data was found to agree with the previous ATLAS measurement of the same quantity \[34\] within the systematic uncertainties.

8. Systematic uncertainties

The \( R_{CP} \) and \( R_{PbPb} \) measurements are subject to systematic uncertainties arising from a number of sources: the jet energy scale and resolution, differences in the spectral shape between data and simulation affecting the bin-by-bin correction factors, residual inefficiency in the trigger selection, and the estimates of the geometric quantities \( R_{\text{coll}} \) (in \( R_{CP} \)) and \( T_{\text{Pb}} \) (in \( R_{PbPb} \)). In addition to these sources of uncertainty, which are common to the \( R_{CP} \) and \( R_{PbPb} \) measurements, \( R_{PbPb} \) is also subject to uncertainties from the \( x_T \)-interpolation of the \( \sqrt{s} = 2.76 \text{ TeV} \) \( pp \) cross-section to the \( \sqrt{s} = 5.02 \text{ TeV} \) centre-of-mass energy and from the integrated luminosity of the \( pp \) dataset.

Uncertainties in the jet energy scale and resolution influence the correction of the \( p + \text{Pb} \) and \( pp \) jet spectra. The uncertainty in the scale was taken from studies of the \textit{in situ} calorimeter response and systematic variations of the jet response in simulation \[32\], as well as studies of the relative energy scale difference between the jet reconstruction procedure in heavy-ion collisions and the procedure used by ATLAS for inclusive jet measurements in 2.76 TeV and 7 TeV \( pp \) collisions \[34,35\]. The total energy scale uncertainty in the measured \( p_T \) range was \( \lesssim 4\% \) for jets in \( |y^*| < 2.8 \), and \( \lesssim 7\% \) for jets in \( |y^*| > 2.8 \). The sensitivity of the results to the uncertainty in the energy scale was evaluated separately for ten distinct sources of uncertainty. Each source was treated as fully uncorrelated with any other source, but fully correlated with itself in \( p_T \), \( \eta \), and \( \sqrt{s} \). The uncertainty in the resolution was taken from \textit{in situ} studies of the dijet energy balance \[36\]. The resolution uncertainty was generally \( < 10\% \), except for low-\( p_T \) jets where it was \( < 20\% \). The effects on the \( R_{CP} \) and \( R_{PbPb} \) measurements were evaluated through an additional smearing of the energy of reconstructed jets in the simulation such that the resolution uncertainty was added to the original resolution in quadrature.

The resulting systematic uncertainties on \( R_{CP} \) (\( \delta R_{CP} \)) and \( R_{PbPb} \) (\( \delta R_{PbPb} \)) were evaluated by producing new response matrices in accordance with each source of the energy scale uncertainty and the resolution uncertainty, generating new correction factors, and calculating the new \( R_{CP} \) and \( R_{PbPb} \) results. Each energy scale and resolution variation was applied to all rapidity bins and to both the \( p + \text{Pb} \) and \( pp \) response matrices simultaneously. The uncertainty on \( R_{CP} \) and \( R_{PbPb} \) from the total energy scale uncertainty was determined by adding the effects of the ten energy scale uncertainty sources in quadrature. Since the correction factors for the \( p + \text{Pb} \) spectra in different centrality intervals were affected to a similar degree by variations in the energy scale and resolution, the effects tended to cancel in the \( R_{CP} \) ratio, and the resulting \( \delta R_{CP} \) were small. The resulting \( \delta R_{PbPb} \) values were somewhat larger than the \( \delta R_{CP} \) values due to the relative centre-of-mass shift between the \( p + \text{Pb} \) and \( pp \) collision systems. The centrality dependence of the energy scale and resolution uncertainties in \( p + \text{Pb} \) events was negligible.

To achieve better correspondence with the data, the reconstructed jet spectrum was reweighted to match the spectral shape in data before deriving the bin-by-bin correction factors as described above. To determine the sensitivity of the results to this reweighting procedure, the slope of the fit to the ratio of the detector-level spectrum in data to that in simulation was varied by the fit uncertainty, and the correction factors were recomputed with this alternative weighting. The resulting \( \delta R_{PbPb} \) and \( \delta R_{CP} \) from the nominal values were included in the total systematic uncertainty.

As the jet triggers used for the data selection were evaluated to have greater than 99% efficiency in the \( p_T \) regions where they are used to select jets, an uncertainty of 1% was chosen for the centrality selected \( p + \text{Pb} \) yields and the \( pp \) cross-section in the range \( 20 < p_T < 125 \text{ GeV} \). This uncertainty was taken to be uncorrelated between the centrality-selected \( p + \text{Pb} \) yields and the \( pp \) cross-section, resulting in a 1.4% uncertainty on the \( R_{CP} \) and \( R_{PbPb} \) measurements.

The geometric quantities \( R_{\text{coll}} \) and \( T_{\text{Pb}} \) and their uncertainties are listed in Table 1. These uncertainties arise from uncertainties in the geometric modelling of \( p + \text{Pb} \) collisions and in modelling the \( N_{\text{part}} \) dependence of the forward particle production measured by \( \Sigma E_T^{\text{Pb}} \). In general, the uncertainties were asymmetric.
Uncertainties in $R_{\text{coll}}$ were largest for the ratio of the most central to the most peripheral interval (0–10%/60–90%), where they were $+17/–6$%, and smallest in the 40–60%/60–90% ratio, where they were $+4/–3$%. Uncertainties in $T_{\text{jet}}$ were largest in the most central (0–10%) and most peripheral (60–90%) centrality intervals, where the upper or lower uncertainty was as high as 10%, and smaller for intervals in the middle of the $p + \text{Pb}$ centrality range, where they reached a minimum of $+3/–2$% for the 20–30% interval.

The $x_T$-interpolation of the $\sqrt{s} = 2.76$ TeV $pp$ jet cross-section to 5.02 TeV is sensitive to uncertainties in $\rho(x_T, y^*)$, the $\sqrt{s}$-scaled ratio of jet spectra at 2.76 and 7 TeV. Following Eq. (8), the uncertainty in the interpolated $pp$ cross-section ($\sigma^{5.02 \text{ TeV}}$) at fixed $x_T$ is related to the uncertainty in $\rho (\delta \rho)$ via $\delta \rho / \rho = 0.643 (\delta \rho / \rho)$, where $\delta \rho$ was taken from Ref. [34]. The values of $\delta \rho$ ranged from 5% to 23% in the region of the measurement and were generally larger at lower $x_T$ and at larger rapidities.

The integrated luminosity for the 2013 $pp$ dataset was determined by measuring the interaction rate with several ATLAS subdetectors. The absolute calibration was derived from three van der Meer scans [37] performed during the $pp$ data-taking in 2013 in a fashion similar to that used previously within ATLAS [38] for $pp$ data-taking at higher energies. The systematic uncertainty on the integrated luminosity was estimated to be 3.1%.

The uncertainties from the jet energy scale, jet energy resolution, reweighting and $x_T$-interpolation are $p_T$ and $y^*$ dependent, while the uncertainties from the trigger, luminosity, and geometric factors are not. The total systematic uncertainty on the $R_{pT,\text{ pb}}$ measurement ranges from 7% at mid-rapidity and high $p_T$ to 18% at forward rapidities and low $p_T$. In most $p_T$ and rapidity bins, the dominant systematic uncertainty on $R_{pT,\text{ pb}}$ is from the $x_T$-interpolation. The $p_T$- and $y^*$-dependent systematic uncertainties on $R_{CP}$ are small. Near mid-rapidity or at high $p_T$, they are 2%, rising to approximately 12% at low $p_T$ in forward rapidities. Thus, in most of the kinematic region studied, the dominant uncertainty on $R_{CP}$ is from the geometric factors $R_{\text{coll}}$.

9. Results

Fig. 2 presents the fully corrected per-event jet yield as a function of $p_T$ in 0–90% $p + \text{Pb}$ collisions, for each of the jet centre-of-mass rapidity ranges used in this analysis. At mid-rapidity, the yields span over eight orders of magnitude.

The jet nuclear modification factor $R_{pT,\text{ pb}}$ for 0–90% $p + \text{Pb}$ events is presented in Fig. 3 in the eight rapidity bins for which the $pp$ reference was constructed. At most rapidities studied, the $R_{pT,\text{ pb}}$ values show a slight (≈10%) enhancement above one, although many bins are consistent with unity within the systematic uncertainties. At mid-rapidity, the $R_{pT,\text{ pb}}$ values reach a maximum near 100 GeV. No large modification of the total yield of jets relative to the geometric expectation (under which $R_{pT,\text{ pb}} = 1$) is observed. The data in Fig. 3 are compared to a next-to-leading order perturbative QCD calculation of $R_{pT,\text{ pb}}$ with the EPS09 parameterisation of nuclear parton distribution functions [9], using CT10 [39] for the free proton parton distribution functions and following the procedure for calculating jet production rates in $p + \text{Pb}$ collisions described in Refs. [1,40]. The data are slightly higher than the calculation, but generally consistent with it within systematic uncertainties.

The central-to-peripheral ratio $R_{CP}$ for jets in $p + \text{Pb}$ collisions is summarised in Fig. 4, where the $R_{CP}$ values for three centrality intervals are shown in all rapidity ranges studied. The $R_{CP}$ ratio shows a strong variation with centrality relative to the geometric expectation, under which $R_{CP} = 1$. The jet $R_{CP}$ for 0–10%/60–90% events is smaller than one at all rapidities for jet $p_T > 100$ GeV and at all $p_T$ at sufficiently forward (proton-going, $y^* > 0$) rapidities. Near mid-rapidity, the 40–60%/60–90% $R_{CP}$ values are consistent with unity up to 100–200 GeV, but indicate a small suppression at higher $p_T$. In all rapidity intervals studied, $R_{CP}$ decreases with increasing $p_T$ and in increasingly more central collisions. Furthermore, at fixed $p_T$, $R_{CP}$ decreases systematically at more forward rapidities. At the highest $p_T$ in the most forward rapidity bin, the 0–10%/60–90% $R_{CP}$ value is ≈0.2. In the backward rapidity direction (lead-going, $y^* < 0$), $R_{CP}$ is found to be enhanced by 10–20% for low-$p_T$ jets.

Fig. 5 summarises the jet $R_{pT,\text{ pb}}$ in central, mid-central and peripheral events in all rapidity intervals studied. The patterns observed in the centrality-dependent $R_{pT,\text{ pb}}$ values are a consequence of the near-geometric scaling of the minimum-bias $R_{pT,\text{ pb}}$ values along with the strong modifications of the central-to-peripheral ratio $R_{CP}$. At sufficiently high $p_T$, $R_{CP}$ in central events is found to be suppressed ($R_{pT,\text{ pb}} < 1$) and in peripheral events to be enhanced ($R_{pT,\text{ pb}} > 1$). Generally, these respective deviations from the geometric expectation (under which $R_{pT,\text{ pb}} = 1$ for all centrality intervals) increase with $p_T$ and, at fixed $p_T$, increase as the rapidity becomes more forward. Thus, the large effects in $R_{CP}$ are consistent with a combination of modifications that have opposite sign in the centrality-dependent $R_{pT,\text{ pb}}$ values but have little effect on the centrality-inclusive (0–90%) $R_{pT,\text{ pb}}$ values. At backward-going rapidities ($y^* < 0$) the $R_{pT,\text{ pb}}$ value for low-$p_T$ jets in all centrality intervals is consistent with unity within the uncertainties.
Given the observed suppression pattern as a function of jet rapidity, in which the suppression in \( R_{\mathrm{CP}} \) at fixed \( p_T \) systematically increases at more forward-going rapidities, it is natural to ask if it is possible to find a single relationship between the \( R_{\mathrm{CP}} \) values in the different rapidity intervals which is a function of jet kinematics alone. To test this, the \( R_{\mathrm{CP}} \) values in each rapidity bin were plotted against the quantity \( p_T \times \cosh(y^*) \approx E \), where \( y^* \) is the centre of the rapidity bin and \( E \) is the total energy of the jet. In relativistic kinematics, the total energy of a particle is given by \( E = m_T \cosh(y^*) \), where the transverse mass \( m_T = \sqrt{m^2 + p_T^2} \).

In the kinematic range studied, the mass of the typical jet is sufficiently small relative to its transverse momentum that approximating the transverse mass, \( m_T \), with the \( p_T \) is reasonable. The 0–10\% and 90–100\% \( R_{\mathrm{CP}} \) versus \( p_T \times \cosh(y^*) \) is shown for all ten rapidity ranges in Fig. 6. When plotted against this variable, the \( R_{\mathrm{CP}} \) values in each of the five forward-going rapidities \( y^* > +0.8 \) fall along the same curve, which is approximately linear in the logarithm of \( E \). This trend is also observed in the two most forward of the remaining rapidity intervals \( -0.3 < y^* < +0.8 \), but the \( R_{\mathrm{CP}} \) values at backward rapidities \( y^* < -0.3 \) do not follow this trend. This pattern is also observed in other centrality intervals, albeit with a different slope in \( \ln(E) \) for each centrality interval.

These patterns suggest that the observed modifications may depend on the initial parton kinematics, such as the longitudinal momentum fraction of the parton originating in the proton, \( x_p \). In particular, a dependence on \( x_p \) would explain why the data follow a consistent trend vs. \( p_T \times \cosh(y^*) \) at forward rapidities (where jet production at a given jet energy \( E \) is dominated by \( x_p \sim E/(\sqrt{S}/2) \) partons in the proton) but do not do so at backward rapidities (where the longitudinal momentum fraction of the parton originating in the lead nucleus, \( x_p \), as well as \( x_p \) are both needed to relate the jet and parton kinematics).

By analogy with Fig. 6 where the \( R_{\mathrm{CP}} \) values are plotted versus \( p_T \times \cosh(y^*) \), the \( R_{\mathrm{Pb}} \) values in the four most forward-going bins studied are plotted against this variable in Fig. 7. The \( R_{\mathrm{Pb}} \) values in central and peripheral events are shown separately. Although the systematic uncertainties are larger on \( R_{\mathrm{Pb}} \) than on \( R_{\mathrm{CP}} \), the observed behaviour for jets with \( p_T > 150 \) GeV is consistent with the nuclear modifications depending only on the approximate total jet energy \( p_T \times \cosh(y^*) \).

In central (peripheral) events, the \( R_{\mathrm{Pb}} \) values at forward rapidities are consistent with a rapidity-independent decreasing (increasing) function of \( p_T \times \cosh(y^*) \). Thus, the single trend in \( R_{\mathrm{CP}} \) versus \( p_T \times \cosh(y^*) \) at forward rapidities appears to arise from opposite trends in the central and peripheral \( R_{\mathrm{Pb}} \), both a single function of \( p_T \times \cosh(y^*) \).
The results presented here use the standard Glauber model with fixed $\sigma_{NN}$ to estimate the geometric quantities. The impact of geometric models which incorporate event-by-event changes in the configuration of the proton wavefunction [41] has also been studied. Using the so called Glauber–Grigov Colour Fluctuation model to determine the geometric parameters amplifies the effects seen with the Glauber model. In this model, the suppression in central events and the enhancement in peripheral events would be increased.

10. Conclusions

This paper presents the results of a measurement of the centrality dependence of jet production in $p+\text{Pb}$ collisions at $\sqrt{s_{NN}} = 5.02$ TeV over a wide kinematic range. The data were collected with the ATLAS detector at the LHC and correspond to 27.8 nb$^{-1}$ of integrated luminosity. The centrality of $p+\text{Pb}$ collisions was characterised using the total transverse energy measured in the forward calorimeter on the Pb-going side covering the interval $-4.9 < \eta < -3.2$. The average number of nucleon–nucleon collisions and the mean nuclear thickness factor were evaluated for each centrality interval using a Glauber Monte Carlo analysis.

Results are presented for the nuclear modification factor $R_{p\text{Pb}}$ with respect to a measurement of the inclusive jet cross-section in $\sqrt{s} = 2.76$ TeV $pp$ collisions corresponding to $4.0 \text{pb}^{-1}$ of integrated luminosity. The $pp$ cross-section was $x_{t}$-interpolated to 5.02 TeV using previous ATLAS measurements of inclusive jet production at 2.76 and 7 TeV. Results are also shown for the central-to-peripheral ratio $R_{CP}$. The centrality-inclusive $R_{p\text{Pb}}$ results for 0–90% collisions indicate only a modest enhancement over the geometric expectation. This enhancement has a weak $p_{T}$ and rapidity dependence and is generally consistent with predictions from the modification of the parton distribution functions in the nucleus, which is small in the kinematic region probed by this measurement.

The results of the $R_{CP}$ measurement indicate a strong centrality-dependent reduction in the yield of jets in central collisions relative to that in peripheral collisions, after accounting for the effects of the collision geometries. In addition, the reduction becomes more pronounced with increasing jet $p_{T}$ and at more forward (downstream proton) rapidities. These two results are reconciled by the centrality-dependent $R_{p\text{Pb}}$ results, which show a suppression in central collisions and enhancement in peripheral collisions, a pattern which is systematic in $p_{T}$ and $y$.

The $R_{CP}$ and $R_{p\text{Pb}}$ measurements at forward rapidities are also reported as a function of $p_{T} \times \cosh(y^*)$, the approximate total jet energy. When plotted this way, the results from different rapidity intervals follow a similar trend. This suggests that the mechanism responsible for the observed effects may depend only on the total jet energy or, more generally, on the underlying parton–parton kinematics such as the fractional longitudinal momentum of the parton originating in the proton.

If the relationship between the centrality intervals and proton–lead collision impact parameter determined by the geometric models is correct, these results imply large, impact parameter-dependent changes in the number of partons available for hard scattering. However, they may also be the result of a correlation between the kinematics of the scattering and the soft interactions resulting in particle production at backward (Pb-going) rapidities [42,43].

Recently, the effects observed here have been hypothesised as arising from a suppression of the soft particle multiplicity in collisions producing high energy jets [44]. Independently, it has also been argued that proton configurations containing a large-$x$ parton interact with nucleons in the nucleus with a reduced cross-section, resulting in the observed modifications [45]. In any case the presence of such correlations would challenge the usual factorisation-based framework for describing hard scattering processes in collisions involving nuclei.

Acknowledgements

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently.

We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWFW and FWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DNRN, DNSRC and Lundbeck Foundation, Denmark; EPLANET, ERC and NSRF, European Union; IN2P3-CNRS, CEA-DSM/IRFU, France; GNSF, Georgia; BMFB, BFG, HGF, and AvH Foundation, Germany; GSRT and NSRF, Greece; ISF, MINERVA, GIF, I-CORE and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRSST, Morocco; FOM and NWO, Netherlands; BRF and RCN, Norway; MNiSW and NCN, Poland; GRICES and FCT, Portugal; MINE/IFA, Romania; MES of Russia and ROSATOM, Russian Federation; JINR; MSTD, Serbia; MSSR, Slovakia; ARRS and MIZŠ, Slovenia; DST/NRF, South Africa; MINECO, Spain; SRC and Wallenberg Foundation, Sweden; SER, SNSF and Cantons of Bern and Geneva, Switzerland; NSC, Taiwan; TAEK, Turkey; STFC, the
Fig. 6. Measured $R_{CP}$ values for $R = 0.4$ jets in 0–10% $p + Pb$ collisions. The panel on the left shows the five rapidity ranges that are the most forward-going, while the panel on the right shows the remaining five. The $R_{CP}$ values at each rapidity are plotted as a function of $p_T \times \cosh(y^*)$, where $(y^*)$ is the midpoint of the rapidity bin. Vertical error bars represent the statistical uncertainty while the boxes represent the systematic uncertainties on the jet yields. The shaded box at the left edge (in the left panel) and right edge (in the right panel) of the $R_{CP} = 1$ horizontal line indicates the systematic uncertainty on $R_{CP}$.

Fig. 7. Measured $R_{pT}$ values for $R = 0.4$ jets in $p + Pb$ collisions displayed for multiple rapidity ranges, showing 0–10% events in the left panel and 60–90% events in the right panel. The $R_{pT}$ at each rapidity is plotted as a function of $p_T \times \cosh(y^*)$, where $(y^*)$ is the midpoint of the rapidity bin. Vertical error bars represent the statistical uncertainty while the boxes represent the systematic uncertainties on the jet yields. The shaded box at the left edge of the $R_{pT} = 1$ horizontal line indicates the systematic uncertainties on $T_{AA}$ and the $pp$ luminosity added in quadrature.

Royal Society and Leverhulme Trust, United Kingdom; DOE and NSF, United States of America.

The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA) and in the Tier-2 facilities worldwide.

References

ATLAS Collaboration


1 Department of Physics, University of Adelaide, Adelaide, Australia
2 Physics Department, SUNY Albany, Albany, NY, United States
3 Department of Physics, University of Alberta, Edmonton, AB, Canada
4 (a) Department of Physics, Ankara University, Ankara; (b) Department of Physics, Gazi University, Ankara; (c) Division of Physics, TOBB University of Economics and Technology, Ankara;
5 Turkish Atomic Energy Authority, Ankara, Turkey

6 IAPP, CNRS/IN2P3 and Université de Savoie, Annecy-le-Vieux, France
7 High Energy Physics Division, Argonne National Laboratory, Argonne, IL, United States
8 Department of Physics, University of Arizona, Tucson, AZ, United States
9 Department of Physics, The University of Texas at Arlington, Arlington, TX, United States
10 Physics Department, University of Athens, Athens, Greece
11 Physics Department, National Technical University of Athens, Zografou, Greece
12 Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
13 Institut de Física d'Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain
14 (a) Institute of Physics, University of Belgrade, Belgrade; (b) Vinca Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia
15 Department for Physics and Technology, Bergen, Bergen, Norway
16 Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, United States

17 Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland
18 School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
19 (a) Department of Physics, Bogazici University, Istanbul; (b) Department of Physics, Dogus University, Istanbul; (c) Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey
20 (a) INFN Sezione di Bologna; (b) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
21 Physikalisches Institut, University of Bonn, Bonn, Germany
22 Department of Physics, Boston University, Boston, MA, United States
23 Department of Physics, Brandeis University, Waltham, MA, United States
24 (a) Universidade Federal do Rio de Janeiro COPPE/EEF, Rio de Janeiro; (b) Electrical Circuits Department, Federal University of Juiz de Fora (UFJF), Juiz de Fora; (c) Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei; (d) Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil
25 Physics Department, Brookhaven National Laboratory, Upton, NY, United States
26 (a) National Institute of Physics and Nuclear Engineering, Bucharest; (b) National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca; (c) University Politehnica Bucharest, Bucharest; (d) West University in Timisoara, Timisoara, Romania
27 Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
28 Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
29 Department of Physics, Carleton University, Ottawa, ON, Canada
30 CERN, Geneva, Switzerland
31 Enrico Fermi Institute, University of Chicago, Chicago, IL, United States
32 (a) Departamento de Fisica, Pontificia Universidad Catolica de Chile, Santiago; (b) Departamento de Fisica, Universidad Técnica Federico Santa Maria, Valparaiso, Chile
33 (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (b) Department of Modern Physics, University of Science and Technology, Anhui; (c) Department of Physics, Nanjing University, Jiangsu;
34 (a) School of Physics, Shandong University, Shandong; (b) Physics Department, Shanghai Jiao Tong University, Shanghai, China
35 Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France
36 Nevis Laboratory, Columbia University, Irvington, NY, United States
37 Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
38 (a) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati; (b) Dipartimento di Fisica, Università della Calabria, Rende, Italy
39 (a) ACH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; (b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
40 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland
41 Physics Department, Southern Methodist University, Dallas, TX, United States
42 Physics Department, University of Texas at Dallas, Richardson, TX, United States
43 DESY, Hamburg and Zeuthen, Germany
44 Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
45 Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
46 Department of Physics, Duke University, Durham, NC, United States
47 SUPA – School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
48 INFN Laboratori Nazionali di Frascati, Frascati, Italy
49 Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
50 Section de Physique, Université de Genève, Geneva, Switzerland
51 (a) INFN Sezione di Genova; (b) Dipartimento di Fisica, Università di Genova, Genova, Italy
52 (a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi; (b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
53 (a) Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
54 (a) Division of Physics, TOBB University of Economics and Technology, Ankara;
55 Turkish Atomic Energy Authority, Ankara, Turkey

56 Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France
57 Department of Physics, Hampton University, Hampton, VA, United States
58 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, United States
59 (a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (c) ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
60 Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
61 Department of Physics, Indiana University, Bloomington, IN, United States
62 Institut für Astron- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
63 University of Iowa, Iowa City, IA, United States
64 Department of Physics and Astronomy, Iowa State University, Ames, IA, United States
65 Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
66 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
67 Graduate School of Science, Kobe University, Kobe, Japan
68 Faculty of Science, Kyoto University, Kyoto, Japan
69 Kyoto University of Education, Kyoto, Japan
144 SLAC National Accelerator Laboratory, Stanford, CA, United States

145 (a) Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Košice, Slovak Republic

146 (a) Department of Physics, University of Cape Town, Cape Town; (b) Department of Physics, University of Johannesburg, Johannesburg; (c) School of Physics, University of the Witwatersrand, Johannesburg, South Africa

147 (a) Department of Physics, Stockholm University; (b) The Oskar Klein Centre, Stockholm, Sweden

148 Physics Department, Royal Institute of Technology, Stockholm, Sweden

149 Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook, NY, United States

150 Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom

151 School of Physics, University of Sydney, Sydney, Australia

152 Institute of Physics, Academia Sinica, Taipei, Taiwan

153 Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel

154 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

155 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

156 International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan

157 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan

158 Department of Physics, Tokyo Institute of Technology, Tokyo, Japan

159 Department of Physics, University of Toronto, Toronto, ON, Canada

160 (a) TRIUMF, Vancouver, BC; (b) Department of Physics and Astronomy, York University, Toronto, ON, Canada

161 Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan

162 Department of Physics and Astronomy, Tufts University, Medford, MA, United States

163 Centro de Investigaciones, Universidad Antonio Nariño, Bogota, Colombia

164 Department of Physics and Astronomy, University of California Irvine, Irvine, CA, United States

165 INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; (c) Istituto di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy

166 Department of Physics, University of Illinois, Urbana, IL, United States

167 Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

168 Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain

169 Department of Physics, University of British Columbia, Vancouver, BC, Canada

170 Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada

171 Department of Physics, University of Warwick, Coventry, United Kingdom

172 Waseda University, Tokyo, Japan

173 Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

174 Department of Physics, University of Wisconsin, Madison, WI, United States

175 Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

176 Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany

177 Department of Physics, Yale University, New Haven, CT, United States

178 Yerevan Physics Institute, Yerevan, Armenia

179 Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France

\[ \text{Also at Department of Physics, King’s College London, London, United Kingdom.} \]

\[ \text{Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.} \]

\[ \text{Also at Novosibirsk State University, Novosibirsk, Russia.} \]

\[ \text{Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.} \]

\[ \text{Also at TRIUMF, Vancouver, BC, Canada.} \]

\[ \text{Also at Department of Physics, California State University, Fresno, CA, United States.} \]

\[ \text{Also at Department of Physics, University of Fribourg, Fribourg, Switzerland.} \]

\[ \text{Also at Tomsk State University, Tomsk, Russia.} \]

\[ \text{Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.} \]

\[ \text{Also at Università di Napoli Parthenope, Napoli, Italy.} \]

\[ \text{Also at Institute of Particle Physics (IPP), Canada.} \]

\[ \text{Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia.} \]

\[ \text{Also at Chinese University of Hong Kong, China.} \]

\[ \text{Also at Louisiana Tech University, Ruston, LA, United States.} \]

\[ \text{Also at Institucio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.} \]

\[ \text{Also at Department of Physics, The University of Texas at Austin, Austin, TX, United States.} \]

\[ \text{Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.} \]

\[ \text{Also at CERN, Geneva, Switzerland.} \]

\[ \text{Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan.} \]

\[ \text{Also at Manhattan College, New York, NY, United States.} \]

\[ \text{Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.} \]

\[ \text{Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France.} \]

\[ \text{Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.} \]

\[ \text{Also at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France.} \]

\[ \text{Also at School of Physical Sciences, National Institute of Science Education and Research, Bhubaneswar, India.} \]

\[ \text{Also at Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy.} \]

\[ \text{Also at Moscow Institute of Physics and Technology (MIPT), Dolgoprudny, Russia.} \]

\[ \text{Also at Section de Physique, Université de Genève, Geneva, Switzerland.} \]

\[ \text{Also at International School for Advanced Studies (SISSA), Trieste, Italy.} \]

\[ \text{Also at Department of Physics and Astronomy, University of South Carolina, Columbia, SC, United States.} \]

\[ \text{Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.} \]

\[ \text{Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow, Russia.} \]

\[ \text{Also at National Research Nuclear University MEPhI, Moscow, Russia.} \]

\[ \text{Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.} \]

\[ \text{Also at Department of Physics, Oxford University, Oxford, United Kingdom.} \]

\[ \text{Also at Department of Physics, Nanjing University, Jiangsu, China.} \]
Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.

Also at Department of Physics, The University of Michigan, Ann Arbor, MI, United States.

Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa.

Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia.

* Deceased.