Measurements of the Nuclear Modification Factor for Jets in Pb+Pb Collisions at $\sqrt{s_{NN}} = 2.76$ TeV with the ATLAS Detector

The ATLAS Collaboration

Abstract

Measurements of inclusive jet production are performed in $pp$ and Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV with the ATLAS detector at the LHC, corresponding to integrated luminosities of 4.0 pb$^{-1}$ and 0.14 nb$^{-1}$, respectively. The jets are identified with the anti-$k_t$ algorithm with $R = 0.4$, and the spectra are measured over the kinematic range of jet transverse momentum $32 < p_T < 500$ GeV, and absolute rapidity $|y| < 2.1$ and as a function of collision centrality. The nuclear modification factor, $R_{AA}$, is evaluated and jets are found to be suppressed by approximately a factor of two in central collisions compared to $pp$ collisions. The $R_{AA}$ shows a slight increase with $p_T$ and no significant variation with rapidity.
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Measurements of inclusive jet production are performed in pp and Pb+Pb collisions at \( \sqrt{s_{\mathrm{NN}}} = 2.76 \text{ TeV} \) with the ATLAS detector at the LHC, corresponding to integrated luminosities of 4.0 pb\(^{-1}\) and 0.14 nb\(^{-1}\), respectively. The jets are identified with the anti-\( k_T \) algorithm with \( R = 0.4 \), and the spectra are measured over the kinematic range of jet transverse momentum \( 32 < p_T < 500 \text{ GeV} \), and absolute rapidity \( |y| < 2.1 \) and as a function of collision centrality. The nuclear modification factor, \( R_{AA} \), is evaluated and jets are found to be suppressed by approximately a factor of two in central collisions compared to pp collisions. The \( R_{AA} \) shows a slight increase with \( p_T \) and no significant variation with rapidity.

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Relativistic heavy-ion collisions at the LHC produce a medium of strongly interacting nuclear matter composed of deconfined color charges [1–4]. Hard scattering processes occurring in these collisions produce high transverse momentum (\( p_T \)) partons that propagate through the medium and lose energy, resulting in the phenomenon of “jet quenching.” The partonic energy loss can be probed through measurements of the suppression of jet production rates. The effects of energy loss have been observed through the suppression of single hadrons [5–11] and jets constructed from charged particles [12]. ATLAS has previously reported measurements with fully reconstructed jets [13] by comparing the jet yields in central collisions, where the colliding nuclei have large overlap, to the yields in peripheral collisions. Those results indicate that the rate of jets in Pb+Pb collisions is suppressed by a factor of approximately two in central collisions relative to peripheral collisions. A more sensitive probe of energy loss is provided by measurements of the suppression relative to pp collisions, where there are no quenching effects.

The magnitude of the suppression is expected to depend on both the \( p_T \) dependence of the energy loss as well as the shape of the initial jet production \( p_T \) spectrum [1]. This spectrum becomes increasingly steep at larger values of the jet rapidity [14]. Thus measurements of jet suppression for jets in different intervals of rapidity provide complementary information about the energy loss. Additionally, parton showers initiated by quarks may be quenched differently than gluons [15], and the fraction of quark-initiated jets is expected to increase with rapidity.

Hard scattering rates are enhanced in more central collisions; the larger overlap results in a higher integrated luminosity of partons able to participate in hard scattering processes, and these hard scattering rates are expected to be proportional to the nuclear overlap function, \( T_{AA} \). The suppression is quantified by the nuclear modification factor

\[
R_{AA} = \frac{1}{N_{\text{evt}} \frac{d\sigma}{dp_T dy}} \frac{\frac{d^2\sigma}{dp_T^2 dy}_{\text{central}}}{\frac{d^2\sigma}{dp_T^2 dy}_{\text{AA}}}.
\]

This Letter presents measurements of the inclusive jet \( R_{AA} \) in Pb+Pb collisions at a nucleon–nucleon center-of-mass energy of \( \sqrt{s_{\mathrm{NN}}} = 2.76 \text{ TeV} \). It utilizes Pb+Pb data collected during 2011 corresponding to an integrated luminosity of 0.14 nb\(^{-1}\) as well as data from pp collisions recorded during 2013 at the same center-of-mass energy corresponding to 4.0 pb\(^{-1}\). Results are presented for jets reconstructed in the calorimeter with the anti-\( k_T \) jet-finding algorithm [16] with jet radius parameter \( R = 0.4 \). The contribution of the underlying event (UE) to each jet, assumed to be uncorrelated and additive, was subtracted on a per-jet basis.

The measurements presented here were performed with the ATLAS calorimeter, inner detector, trigger, and data acquisition systems [17, 18]. The calorimeter system consists of a liquid argon (LAr) electromagnetic (EM) calorimeter (\( |\eta| < 3.2 \)), a steel-scintillator sampling hadronic calorimeter (\( |\eta| < 1.7 \)), a LAr hadronic calorimeter (1.5 < \( |\eta| \) < 3.2), and a forward calorimeter (FCal) (\( 3.2 < |\eta| < 4.9 \)). Charged-particle tracks were measured over the range \( |\eta| < 2.5 \) using the inner detector [19], which is composed of silicon pixel detectors in the innermost layers, followed by silicon microstrip detectors and a straw-tube transition-radiation tracker (\( |\eta| < 2.0 \)), all immersed in a 2 T axial magnetic field. The zero-degree calorimeters (ZDCs) are located symmetrically at \( z = \pm 140 \text{ m} \) and cover \( |\eta| > 8.3 \). A ZDC coincidence trigger was defined by requiring a signal consistent with one or more neutrons in each of the calorimeters.

The pp events used in the analysis were selected using the ATLAS jet trigger [20] with multiple values of the trigger \( p_T \) thresholds. During pp data taking, the average number of pp interactions per bunch crossing (pile-up) varied from 0.3 to 0.6. The pp events were required to contain at least one primary vertex, reconstructed from at least two tracks, and jets originating from all such vertices were included in the cross section measurement.

Data from Pb+Pb collisions were recorded using either a minimum-bias trigger or a jet trigger. The minimum-
bias trigger, formed from the logical OR of triggers based on a ZDC coincidence or total transverse energy in the event, is fully efficient in the range of centralities presented here. The jet trigger identified jets by applying the anti-$k_t$ algorithm with $R = 0.2$ with a UE subtraction procedure similar to that applied in the offline analysis. 

The jet trigger selected events having at least one jet with transverse energy $E_T > 20$ GeV at the electromagnetic scale [21]. Event selection and background rejection criteria were applied [22] yielding 53 million and 14 million events in the minimum-bias and jet-triggered samples, respectively.

The centrality of Pb+Pb collisions was characterized by $\Sigma E_T^{FCal}$, the total transverse energy measured in the electromagnetic calorimeter response matrix. Separate response matrices were obtained. The response matrix is generally diagonal, indicating that the response distributions broaden with increasing centrality.

The response matrix is used to unfold the jet spectra. Additional MC samples were used in evaluations of the jet energy resolution (JER) using a method based on the Singular Value Decomposition [34]. The effects of the JER, which receives contributions from both the detector response and UE fluctuations, were evaluated by applying the same procedure to the MC samples as was applied to the data and by matching the resulting reconstructed jets and “generator jets” that are reconstructed from final-state PYTHIA hadrons. For each pair, the $p_T$ of the generator and reconstructed jets were used to populate a detector response matrix. Separate response matrices were obtained for each centrality interval.

The response matrix is generally diagonal, indicating that jets are likely to be reconstructed in the same $p_T$ bin as the generator jets. The average $p_T$ difference between reconstructed and generator jets, is $\lesssim 1\%$, independent of centrality. However, the response distributions broaden at low $p_T$ as the relative JER increases due to the larger UE fluctuations. At $p_T = 200$ GeV, the relative JER is approximately 10% and is independent of centrality. However at $p_T = 40$ GeV, it varies from 20–40% between

<table>
<thead>
<tr>
<th>Centrality [%]</th>
<th>$\langle T_{AA} \rangle$ [mb$^{-1}$]</th>
<th>$\langle N_{part} \rangle$</th>
</tr>
</thead>
<tbody>
<tr>
<td>0–10</td>
<td>$23.45 \pm 0.37$</td>
<td>$356.2 \pm 2.5$</td>
</tr>
<tr>
<td>10–20</td>
<td>$14.43 \pm 0.30$</td>
<td>$260.7 \pm 3.6$</td>
</tr>
<tr>
<td>20–30</td>
<td>$8.73 \pm 0.26$</td>
<td>$85.6 \pm 4.3$</td>
</tr>
<tr>
<td>30–40</td>
<td>$5.04 \pm 0.22$</td>
<td>$129.3 \pm 3.8$</td>
</tr>
<tr>
<td>40–50</td>
<td>$2.74 \pm 0.17$</td>
<td>$85.6 \pm 3.6$</td>
</tr>
<tr>
<td>50–60</td>
<td>$1.33 \pm 0.12$</td>
<td>$53.0 \pm 3.1$</td>
</tr>
<tr>
<td>60–70</td>
<td>$0.50 \pm 0.07$</td>
<td>$30.1 \pm 2.5$</td>
</tr>
<tr>
<td>70–80</td>
<td>$0.24 \pm 0.04$</td>
<td>$15.1 \pm 1.7$</td>
</tr>
<tr>
<td>0–1</td>
<td>$29.04 \pm 0.46$</td>
<td>$400.1 \pm 4.3$</td>
</tr>
<tr>
<td>1–5</td>
<td>$25.62 \pm 0.40$</td>
<td>$377.6 \pm 2.2$</td>
</tr>
<tr>
<td>5–10</td>
<td>$20.50 \pm 0.34$</td>
<td>$330.3 \pm 3.0$</td>
</tr>
<tr>
<td>60–80</td>
<td>$0.41 \pm 0.05$</td>
<td>$22.6 \pm 2.1$</td>
</tr>
</tbody>
</table>

TABLE I. The $\langle T_{AA} \rangle$ and $\langle N_{part} \rangle$ values and their uncertainties in each centrality bin.
peripheral and central collisions. The unfolding is most sensitive in this region and the range of jet $p_T$ used in the unfolding was chosen separately in each centrality interval to be as low as possible while maintaining stability in the unfolding procedure. The statistical covariance of each unfolded spectrum was evaluated using the pseudo-experiment procedure described in Ref. [13]. Systematic uncertainties in the unfolding procedure were evaluated by varying the choice of regularization parameter used in the unfolding.

The effects of any inefficiency in the jet reconstruction, including inefficiency introduced by the UE jet rejection requirement, were corrected for by a multiplicative correction applied after unfolding. This factor, obtained from the MC sample, is unity for $p_T > 100$ GeV and reaches a maximum of 1.3 in the most central collisions at the lowest $p_T$. For values larger than unity, an uncertainty of 0.5% was assigned to this correction based on the comparison of the jet reconstruction efficiency with respect to track jets between the data and MC sample.

Uncertainties on the JER and JES have been evaluated using data-driven techniques in $pp$ collisions [21, 35]. A systematic uncertainty of 1.5% on the JES was assigned to account for potential differences, not described by the MC simulations, between the two data-taking periods. This value was obtained by comparing the calorimetric response with respect to the $p_T$ of matched track jets in $pp$ and peripheral Pb+Pb collisions.

A centrality-dependent uncertainty on the JES due to differences between $pp$ and Pb+Pb in the partonic composition of jets and in their fragmentation was estimated with the PYQUEN sample. The jet response in that sample was found to differ by up to 1% from that in the PYTHIA sample. The magnitude of this variation was checked with a similar study using track jets to compare central and peripheral Pb+Pb data. The uncertainty was taken to be 1% in the most central collisions with the uncertainty decreasing in more peripheral collisions.

The impacts of the JER and JES uncertainties on the spectra were assessed by constructing new response matrices with a systematically varied relationship between the reconstructed and generator jet kinematics and repeating the unfolding. Correlations in the JES and JER uncertainties across the $pp$ and Pb+Pb samples were accounted for in the propagation of the uncertainties to the $R_{AA}$. 

FIG. 1. The double differential jet cross section in $pp$ collisions as a function of $p_T$ in different rapidity bins (scaled by successive powers of $10^2$). The statistical and systematic uncertainties are indicated by the error bars (too small to be seen on this scale) and shaded bands, respectively. The points and horizontal error bars indicate the $p_T$ bin center and width, respectively.

FIG. 2. The per-event jet yield in Pb+Pb collisions, multiplied by $1/\langle T_A^A \rangle$, as a function of $p_T$ (scaled by successive powers of $10^2$). The upper panel shows the 0–2.1 rapidity range in different centrality intervals. The lower panel shows the 0–10% centrality interval in different rapidity ranges. The statistical and systematic uncertainties are indicated by the error bars (too small to be seen on this scale) and shaded bands, respectively. The points and horizontal error bars indicate the $p_T$ bin center and width, respectively. The solid and dashed lines represent the $pp$ jet cross section for the same rapidity interval scaled by the same factor.
Uncertainties on the $T_{AA}$ and integrated luminosity affect the overall normalization of the yields and thus are independent of jet $p_T$ and rapidity. The uncertainties on $\langle T_{AA} \rangle$ vary between 1% and 10% in the most central and peripheral collisions, respectively, with the full set of values given in Table I. The uncertainty on the integrated luminosity is estimated to be 3.1%. It is determined, following the same methodology as that detailed in Ref. [36], from a calibration of the luminosity scale derived from beam-separation scans performed during the 2.76 TeV operation of the LHC in 2013.

The total systematic uncertainty on the $pp$ cross sections is dominated by the JES uncertainty, which is as large as 15%. For the $Pb+Pb$ jet yields this uncertainty is also dominant and in the most central collisions is 22%. In the $R_{AA}$, much of this uncertainty cancels. However the dominant contribution is due to the JES in most centrality and rapidity intervals and is typically 10%. The uncertainties due to the unfolding are generally a few per cent, but for some $p_T$ values near the upper and lower limits included in the measurement the contributions from this source are as large as 15%. The contributions of the JER to the total uncertainty on $R_{AA}$ are less than 3% except in the most central collisions at low $p_T$ where they are as large as 10%. In the most peripheral bins the $\langle T_{AA} \rangle$ uncertainties that affect the overall normalization are the dominant contribution.

The $pp$ differential jet cross sections are shown in Fig. 1 for the following absolute rapidity ranges: 0–0.3, 0.3–0.8, 0.8–1.2, 1.2–2.1 and 2.1–2.7. These results are consistent with a previous measurement with fewer events [37]. The differential per-event jet yield in $Pb+Pb$ collisions, multiplied by $1/\langle T_{AA} \rangle$, is shown in Fig. 2, in selected rapidity and centrality bins in the lower and upper panels, respectively. The dashed lines represent the $pp$ jet cross sections for that same rapidity bin; the jet suppression is evidenced by the fact that the jet yields fall below these lines.

The jet $R_{AA}$ as a function of $p_T$ is shown in Fig. 3 for different ranges in collision centrality and jet rapidity. The $R_{AA}$ is observed to increase weakly with $p_T$, except in the most peripheral collisions. In the 0–10% and $|y| < 2.1$ centrality and rapidity intervals, which have the smallest statistical uncertainty, the $R_{AA}$ is 0.47 at $p_T \sim 55$ GeV and rises to 0.56 at $p_T \sim 350$ GeV. These distributions were fit, accounting for the pointwise correlations in the uncertainties, to the functional form $a \ln(p_T) + b$. The slope parameter was found to be significantly above zero in all but the most peripheral collisions. The magnitude and weak increase of the $R_{AA}$ in central collisions are described quantitatively by recent theoretical calculations [38, 39]. The results of this measurement are consistent with measurements of the jet central-to-peripheral ratio [13], although in those measurements the uncertainties are too large to infer any significant $p_T$ dependence.

The rapidity dependence of the $R_{AA}$ is shown in the top panel of Fig. 4 for jets with $80 < p_T < 100$ GeV for three centrality bins. The $R_{AA}$ shows no significant rapidity dependence over the $p_T$ and rapidity ranges presented in this measurement. The $\langle N_{\text{part}} \rangle$ dependence...
is shown in the bottom panel of Fig. 4 for jets in the same \( p_T \) interval and with \( |y| < 2.1 \). The \( R_{AA} \) decreases smoothly from the most peripheral collisions (smallest \( \langle N_{\text{part}} \rangle \) values) to central collisions, where it reaches a minimal value of approximately 0.4 in the most central 1% of collisions. A similar \( \langle N_{\text{part}} \rangle \) dependence is observed for jets in different ranges of \( p_T \) and rapidity.

In summary, this Letter presents measurements of inclusive jet production in pp and Pb+Pb collisions over a wide range in \( p_T \), rapidity and centrality. The jet nuclear modification factor, \( R_{AA} \), obtained from these measurements shows a weak rise with \( p_T \), with a slope that varies with collision centrality. No significant slope is observed in the most peripheral collisions. The \( R_{AA} \) decreases gradually with increasing \( \langle N_{\text{part}} \rangle \). At forward rapidity, the increasing steepness of the jet production spectrum is expected to result in more suppression of the jet yields. In this kinematic region, the production is increasingly dominated by quark jets, which may lose less energy than gluon jets [15]. The observed lack of rapidity dependence in the \( R_{AA} \) places constraints on relative energy loss for quark and gluon jets in theoretical descriptions of jet quenching.

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[17] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the z-axis along the beam pipe. The \( x \)-axis points from the IP to the center of the LHC ring, and the \( y \)-axis points upward. Cylindrical coordinates \( (r, \phi) \) are used in the transverse plane, \( \phi \) being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle \( \theta \) as \( \eta = -\ln \tan(\theta/2) \).


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