Laser induced spin precession in highly anisotropic granular L10 FePt

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Laser induced spin precession in highly anisotropic granular L1₀ FePt

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The dynamic magnetic properties of a highly anisotropic, granular L1₀ FePt thin film in magnetic fields up to 7 T are investigated using time-resolved magneto-optical Kerr effect measurements. We find that ultrashort laser pulses induce coherent spin precession in the granular FePt sample. Frequencies of spin precession up to 400 GHz are observed, which are strongly field and temperature dependent. The high frequencies can be ascribed to the high value of the magnetocrystalline anisotropy constant $K_u$, leading to large anisotropy fields $H_u$ of up to 10.7 T at 170 K. A Gilbert damping parameter of $\alpha \sim 0.1$ was derived from the lifetimes of the oscillations.

The investigated sample is a granular FePt thin film. It was grown by sputter deposition at a substrate temperature of 550 °C on a set of underlayers with MgO (001) surfaces directly beneath the FePt, resulting in an out of plane easy axis of magnetization (for more information see Ref. 11). The elevated growth temperature leads to the formation of the face centered tetragonal (fct) L1₀ phase rather than the face centered cubic (fcc) A1 phase. The L1₀ phase is chemically ordered (chemical order parameter S = 1), as atomic layers of Fe and Pt alternate along the [001] direction. It has been shown that in contrast to the fcc phase, the fct phase has a high magnetocrystalline anisotropy of $K_u \sim 7 \times 10^7$ erg/cm³, and the higher the chemical ordering in the sample. Structural defects such as mis-oriented grains cause a slight reduction of the overall order in the sample. To ensure magnetically decoupled, separate grains, 35 vol. % of carbon is added as a segregant during the growth process. Carbon does not coalesce with FePt hence forming physically isolated and exchange decoupled grains of roughly 7 nm diameter and a grain size distribution of $\sigma = 16\%$. The magnetic layer is capped by a 3 nm thick protective diamond like carbon (DLC) overcoat, allowing for optical access.

To characterize the static magnetic properties of the sample, the polar magneto-optical Kerr effect was measured at a wavelength of 800 nm. Magnetic hysteresis loops were recorded applying an external magnetic field along the easy axis of magnetization, normal to the sample plane. The resulting hysteresis loops (Fig. 1) reveal a high coercive field ($H_c$) of 5.0 ± 0.1 T at room temperature. Decreasing the temperature to 4.2 K leads to a significant increase in coercive field to 8.2 ± 0.5 T. To saturate the magnetization at 4.2 K, fields greater than $H_z = 11 \pm 0.5$ T are required. This increase

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in \( H_c \) and \( H_r \) can be attributed to an increase in anisotropy at lower temperatures\(^{15}\) as is also confirmed by our measurements of the magnetization dynamics (see below). The rather wide switching field distribution that can be observed in the hysteresis curves can be explained with structural (grain size distribution) and chemical (S < 1) inhomogeneities.\(^{11}\) Another notable effect was observed when applying the external field at a 45° angle with respect to the sample normal. In this configuration, the coercive field is reduced to 3.7 ± 0.1 T. This effect can be attributed to the granular nature of the FePt layer. The individual grains are single domain, exchange decoupled ferromagnets which can be described by a Stoner-Wohlfarth (SW) model.\(^{12}\) For ideal SW particles, the switching field is minimal if applied between the hard and easy axis of magnetization which corresponds to a 45° angle for the here presented sample.

In order to investigate the dynamic magnetic properties, we used an all optical, TR-MOKE pump probe technique described elsewhere.\(^4\) The sample was excited with a 100 fs light pulse at a central wavelength of 800 nm and a fixed fluence of 0.4 mJ/cm\(^2\). It was subsequently probed with a similar pulse 100 times lower in power. During the measurement, an external magnetic field up to 7 T was applied at an angle of 45° with the sample normal. A measurement at room temperature and an applied magnetic field of 7 T is shown in Fig. 2(a). The laser pulse excitation leads to an initial ultrafast demagnetization within 1.5 ps. The point of maximum demagnetization \( t \approx 3.7 \text{T} \). Applying the field at a 45° angle (blue curve) to the easy axis results in a reduced \( H_c = 3.7 \text{T} \). Following the initial demagnetization, damped oscillations are visible during the magnetization relaxation, even more so if the exponential contribution of the relaxation is subtracted (b). (c) TR-MOKE data at 290 K at different applied magnetic fields ranging from 7 T down to 3 T.

FIG. 1. Hysteresis loops acquired using the polar MOKE at 290 K (red curve) and 4.2 K (black curve). Inset (a) shows the polar measurement geometry (red and black curve) while (b) shows the 45° geometry (blue curve). The high coercive field \( H_c = 5 \text{T} \) at 290 K is increased to 8.2 T at 4.2 K. Applying the field at a 45° angle (blue curve) to the easy axis results in a reduced \( H_c = 3.7 \text{T} \).

FIG. 2. (a) TR-MOKE time trace on FePt at 290 K, 7 T. Following the initial demagnetization, damped oscillations are visible during the magnetization relaxation, even more so if the exponential contribution of the relaxation is subtracted (b). (c) TR-MOKE data at 290 K at different applied magnetic fields ranging from 7 T down to 3 T.
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occupy 65 vol. % of the film and due to the high $K_u$ the anisotropy energy will be significantly higher than the demagnetizing energy. Using spherical coordinates, the equilibrium angles for the magnetization $\mathbf{M}$, $\theta_{eq}$, and $\varphi_{eq}$, that minimize $U$ are then given by

$$
\sin(\theta_{eq} - \theta_h) = \frac{4\pi M - 2K_u/M}{2H}, \quad \cos \varphi_{eq} = 0,
$$

with $\theta_h$ is the polar angle of the applied magnetic field $\mathbf{H}$ relative to the sample normal. For the magnetization $\mathbf{M}$, time dependent azimuthal and polar angles $\theta$, $\varphi \sim e^{i\omega t}$, and small oscillation amplitudes around the equilibrium are assumed. This yields the following formula for ferromagnetic resonance from Eq. (1):

$$
\frac{\omega^2}{\gamma^2} = H_1 \cdot H_2,
$$

$$
H_1 = H \cos(\theta_{eq} - \theta_h) - \frac{(4\pi M - 2K_u/M) \cos^2 \theta_{eq}}{2},
$$

$$
H_2 = H \cos(\theta_{eq} - \theta_h) - \frac{(4\pi M - 2K_u/M) \cos 2\theta_{eq}}{2}.
$$

Using the measured saturation magnetization for this sample of 950 emu/cm$^3$ at room temperature, the fit yields values for the anisotropy fields and the g-factor as shown in Table I. The anisotropy field found at room temperature is not equal to the coercive field, as theory would predict for ideal Stoner-Wohlfarth particles. Here, however, a great ensemble of grains is measured. The grains in the FePt sample have slightly differing shapes and are spread in size as well

as in crystalline alignment, thus showing merely near-SW behaviour. We find that the anisotropy field $H_{an} = 2K_u/M$ increases with decreasing temperature. It is known for L1$_0$ FePt that with decreasing temperature, $K_u$ increases stronger than $M$. If one extrapolates the fit towards zero applied field (Fig. 3(a)), the FMR frequency does not drop to zero but reaches a value of $\sim 240$ GHz at room temperature ($\sim 280$ GHz at 170 K). As can be seen from Eqs. (4) to (5), the magnitude of this zero field frequency is determined by the demagnetizing field $4\pi M$ and the anisotropy field $2K_u/M$ acting against each other. At room temperature, the demagnetizing field equals $\sim 1.2$ T, which is small compared to the anisotropy field of 8.9 T. This leads to the conclusion that the strong magnetocrystalline anisotropy $K_u$ is responsible for the high frequencies of spin precession. As $K_u$ has a stronger temperature dependence than $M$, the shift in the frequency observed between the measurements at 290 K and 170 K can likewise be ascribed to the increase of $K_u$ with decreasing temperature.

The light induced excitation of coherent spin precession can be explained phenomenologically. Initially, the magnetization points along the out of plane easy axis (Fig. 3(a-I)). After applying an external field $\mathbf{H}$, the magnetization aligns along the effective field at an angle $\theta_{eq}$ (see Eq. (3)) between the easy axis and the direction of the applied field (Fig. 3(a-II)). Upon pump arrival, two possible processes can initiate the above seen coherent spin precession: An ultrafast demagnetization, decreasing the Zeeman and demagnetization energy, or an ultrafast change of the magnetocrystalline anisotropy $K_u$. In both cases, the effective field $\mathbf{H}_{eff}$ is tilted out of its equilibrium position (Fig. 3(a-III)), so it does not lie parallel to $\mathbf{M}$ anymore. This creates a torque that acts on the magnetization, thereby starting the precessional motion around $\mathbf{H}_{eff}$ which also relaxes back to its initial position as the energy is dissipated (Fig. 3(a-IV)).

Furthermore, by evaluating the lifetimes $\tau_L$ of the oscillations the damping parameter $\alpha$ can be extracted, using the following relationship derived from the LLG equation assuming that $\alpha \ll 1$:

$$
\alpha = \frac{2}{\tau_L \gamma (H_1 + H_2)}.
$$

A damping factor of $\alpha \approx 0.1$ is obtained for measurements at 290 K as is shown in Fig. 3(b). No significant change in $\alpha$ above the experimental error is observed upon reducing the temperature to 170 K. Also, the damping does not change notably as a function of applied magnetic field. This suggests that there is only little (if any) extrinsic contribution to the measured Gilbert damping parameter. Indeed, strong external magnetic fields as applied here are known to suppress extrinsic damping. Similarly relatively high values for the

<table>
<thead>
<tr>
<th>Temperature (K)</th>
<th>g-factor</th>
<th>$H_{an}$ (T)</th>
</tr>
</thead>
<tbody>
<tr>
<td>290</td>
<td>2.2 ± 0.1</td>
<td>8.9 ± 0.3</td>
</tr>
<tr>
<td>170</td>
<td>2.2 ± 0.1</td>
<td>10.7 ± 0.3</td>
</tr>
</tbody>
</table>
intrinsic Gilbert damping parameter were reported on continuous FePt films with lower magnetocrystalline anisotropy.\textsuperscript{21–23} In summary, our here presented results show that coherent spin precession can be excited in a highly anisotropic, granular L1\textsubscript{0} FePt films using ultrashort light pulses of 100 fs. The high anisotropy fields up to more than 10 T found in the here investigated sample lead to FMR frequencies in the THz range. Frequencies of that magnitude have never before been observed in ferromagnets, and are of considerable relevance to the writing speed in magnetic data storage: Considering conventional damped gyroscopic switching, where the switching field is applied parallel to the easy axis, the switching time is inversely proportional to the frequency of magnetization precession. Lowering the temperature increases the magnetocrystalline anisotropy K\textsubscript{u} which leads to a substantial increase in the observed precession frequency. Furthermore, a Gilbert damping parameter of $\alpha \sim 0.1$ is found, similar to results for continuous films of L1\textsubscript{0}.

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