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Search for Quantum Black Hole Production in High-Invariant-Mass Lepton+Jet Final States Using $pp$ Collisions at $\sqrt{s} = 8$ TeV and the ATLAS Detector

The ATLAS Collaboration

Abstract

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Quantum black holes (QBHs) $^{1}$ $^{2}$ are predicted in low-scale quantum gravity theories that offer solutions to the mass hierarchy problem of the Standard Model (SM) by lowering the scale of quantum gravity ($M_D$) from the Planck scale ($\sim 10^{16}$ TeV) to a value of about 1 TeV. In models with large extra dimensions such as the Arkani-Hamed–Dimopoulous–Dvali (ADD) model $^{3}$ $^{5}$, only the gravitational field is allowed to penetrate the n extra dimensions, while all SM fields are localized in the usual four-dimensional space-time. QBHs with masses near $M_D$, postulated to conserve total angular momentum, color and electric charge, may decay to two particles $^{2}$ $^{6}$. The behavior of QBHs is distinct from semi-classical black holes that decay via Hawking radiation to a large number of objects $^{7}$. Searches for semi-classical black holes typically require three or more objects $^{8}$ $^{9}$.

The quantum approximations used in the modeling of black hole production are valid when black-hole masses are above a minimal threshold mass, $M_{\text{th}}$, which is taken to be equivalent to the QBH inverse gravitational radius. If the QBHs investigated in this Letter are accessible at the Large Hadron Collider (LHC) $^{10}$, they can produce lepton+jet final states $^{2}$ $^{6}$ motivating this first dedicated search for high-invariant-mass final states with a single electron ($e$) or a single muon ($\mu$), and at least one jet. Two-particle QBH decays to a final state consisting of a lepton and a quark-jet violate lepton and baryon number conservation, producing a distinctive signal for physics beyond the SM. Previous searches for QBHs relied on signatures such as dijet mass distributions $^{11}$ $^{12}$, generic multijet configurations $^{9}$ and photon+jet final states $^{13}$.

The largest QBH cross section for the final states considered is predicted for the collision of two $u$-quarks ($\sigma_{uu}$), which produces charge +4/3 objects with equal branching fractions (BFs) of $\text{BF}_{uu} = 11\%$ to each lepton+jet final state. The next largest cross sections are for charge +1/3 (ud) and $−2/3$ (dd) QBHs with lepton+jet BFs of $\text{BF}_{ud} = 5.7\%$ and $\text{BF}_{dd} = 6.7\%$ $^{10}$. Processes with initial states having antiquarks and heavier sea-quarks are suppressed by at least a factor of 100 and can be neglected. The QBH cross section is a steeply declining function of $M_{\text{th}}$, and has $\Sigma \sigma_{qq} \times \text{BF}_{qq} \approx 8.6 \times 10^5$ fb, $8.9 \times 10^2$ fb and 0.75 fb for $M_{\text{th}}$ of 1 TeV, 3 TeV and 5 TeV, respectively $^{14}$.

The ATLAS detector $^{15}$ includes an inner tracker, covering a pseudorapidity $^{16}$ range $|\eta| < 2.5$, surrounded by a superconducting solenoid providing a 2 T central field. A liquid-argon (LAr) electromagnetic (EM) sampling calorimeter ($|\eta| < 3.2$), a scintillator-tile hadronic calorimeter ($|\eta| < 1.7$), a LAr hadronic calorimeter ($1.4 < |\eta| < 3.2$), and a LAr forward calorimeter ($3.1 < |\eta| < 4.9$) provide the energy measurements. The muon spectrometer consists of tracking chambers covering $|\eta| < 2.7$, and trigger chambers covering $|\eta| < 2.4$, in a magnetic field produced by a system of air-core toroids. Events considered in this analysis are required to have one high-transverse-momentum (high-$p_T$) lepton ($e/\mu$) that passes requirements of the three-level trigger system $^{17}$. The thresholds applied at the third trigger level are 60 GeV and 36 GeV for electrons and muons, respectively. The analysis is based on the complete 2012 data set of $pp$ collisions taken at a center-of-mass energy of $\sqrt{s} = 8$ TeV by the ATLAS detector at the LHC, corresponding to an integrated luminosity of $20.3 \pm 0.6 \text{ fb}^{-1}$ $^{18}$ after data-quality requirements.

The event selection is designed to be efficient for generic lepton+jet final states and is based on leading-order simulated-signal QBH events obtained from the QMBH1.04 generator $^{13}$, followed by parton showering and hadronization using PYTHIA8.165 $^{19}$. The signal generator uses the MSTW2008LO $^{20}$ set of leading-order parton distribution functions (PDFs) with the AU2 underlying-event tune $^{21}$. This Letter assumes the ADD model with $M_{\text{th}} = M_D$, $n = 6$, and the QCD factorization scale for the PDFs set to the inverse gravitational radius $^{14}$. Samples with $M_{\text{th}}$ from 1 TeV to 6 TeV, in steps of 0.5 TeV, are generated for both channels.

Events with a high-$p_T$ lepton and one or more jets can also arise from electroweak (EW) processes includ-
ing vector-boson production with additional jets; diboson (WW, WZ, ZZ), top-quark pair (tt) and single top-quark (t or ¯t) production; and multijet processes including non-prompt leptons from semileptonic hadron decays and jets misidentified as leptons.

The EW background in the signal region (SR) is estimated using Monte Carlo (MC) samples normalized to data in control regions. All MC simulated samples are produced using the ATLAS detector simulation [22] based on GEANT4 [23]. The simulated events are reconstructed in the same manner as the data. The tt and single-top-quark events are simulated with MC@NLO4.06 [24] and AcerMC3.8 [25], respectively; the production of W+jets and Z+jets is simulated using ALPGEN2.14 [26]; and diboson production is simulated with SHERPA1.4.1 [27]. The leading-order CTEQ6L1 PDFs [28] are used for ALPGEN and AcerMC samples while the next-to-leading-order CT10 PDFs [29] are used for the SHERPA and MC@NLO samples. The generators for all samples except dibosons are interfaced to HERWI6.G520 [30, 31] for parton showering and hadronization and to JIMMY4.31 [32] for the underlying-event model. The results of higher-order calculations are used to adjust the relative fractions of the simulated events as in Refs. [33, 34]. Additional inelastic pp interactions, termed pileup, are included in the event simulation so as to match the distribution in the data (on average 21 interactions per bunch crossing).

Electron candidates are identified as localized depositions of energy in the EM calorimeter with $p_T > 130$ GeV and $|\eta| < 2.47$, excluding the barrel–endcap transition region, $1.37 < |\eta| < 1.52$, and matched to a track reconstructed in the tracking detectors. Background from jets is reduced by requiring that the shower profiles are consistent with those of electrons [35]. Isolated electrons are selected by requiring the transverse energy deposited in a cone of radius $\Delta R = \sqrt{\Delta \eta^2 + (\Delta \phi)^2} = 0.3$ centered on the electron cluster, excluding the energy of the electron cluster itself, to be less than $0.0055 \times p_T + 3.5$ GeV after corrections for energy due to pileup and energy leakage from the electron cluster into the cone. This criterion provides nearly constant selection efficiency for signal events over the entire $p_T$ range explored and for all the pileup conditions.

Muon candidates are required to be detected in at least three layers of the muon spectrometer and to have $p_T > 130$ GeV and $|\eta| < 2.4$. Possible background from cosmic rays is reduced by requiring the transverse and longitudinal distances of closest approach to the interaction point to be smaller than 0.2 mm and 1.0 mm, respectively. Signal muons are required to be isolated such that $\sum p_T < 0.05 \times p_T$, where $\sum p_T$ is the sum of the $p_T$ of the other tracks in a cone of radius $\Delta R = 0.3$ around the direction of the muon.

Jets are constructed from three-dimensional noise-suppressed clusters of calorimeter cells using the anti-$k_t$ algorithm with a radius parameter of 0.4 [36, 37]. Jet energies are corrected for losses in material in front of the active calorimeter layers, detector inhomogeneities, the non-compensating nature of the calorimeter, and pileup. Jet energies are calibrated using MC simulation and the combination of several in-situ techniques applied to data [38, 39]. All jets are required to have $p_T > 50$ GeV and $|\eta| < 2.5$. In addition, the most energetic jet is required to have $p_T > 130$ GeV.

The missing transverse momentum (with magnitude $E_T^{\text{miss}}$), used only in the background estimation, is calculated as the negative of the vectorial sum of calibrated clustered energy deposits in the calorimeters, and is corrected for the momenta of any reconstructed muons [41].

In the electron (muon) channel, events are required to have exactly one electron (muon). Multijet background can be reduced, with minimal loss in signal efficiency, by requiring the average value of $\eta$ for the lepton and leading jet to satisfy $|\langle \eta \rangle| < 1.25$ and the difference between the lepton and leading jet $\eta$ to satisfy $|\Delta \eta| < 1.5$. The signal lepton and jet are mostly back-to-back in $\phi$ and are required to satisfy $|\Delta \phi| > \pi/2$.

The invariant mass ($m_{\text{inv}}$) is calculated from the lepton and highest-$p_T$ jet. The SR is defined by a lower bound on $m_{\text{inv}}$, $m_{\min}$, that accounts for experimental resolution. In the electron channel $m_{\min} = 0.9 M_{\text{th}}$ is used. In the muon channel, the requirement is loosened at high invariant mass, as muon resolution has a term quadratic in $p_T \mu$, resulting in $m_{\min} = [0.95 - 0.05 M_{\text{th}}/1 \text{TeV}] M_{\text{th}}$. A low-invariant-mass control region (LIMCR) is defined with $m_{\min}$ between 400 GeV and 900 GeV, which has a negligible contamination from a potential signal (<2%) for the lowest $M_{\text{th}}$ considered.

The acceptance of the event selection is about 65%, based on generator-level quantities and calculated by imposing selection criteria that apply directly to phase space (lepton/jet $\eta$, lepton/jet $p_T$, $\Delta \eta$, $\Delta \phi$, ($\langle \eta \rangle$), and $m_{\text{inv}}$). All other selection criteria, which in general correspond to event and object quality requirements, are used to calculate the experimental efficiency based on the events included in the acceptance. The experimental efficiency falls from 89(59)% to 81(50)% for masses from 1 TeV to 6 TeV in the electron (muon) channel. The experimental efficiency in the muon channel is lower than that in the electron channel because more stringent requirements are applied to ensure the best possible resolution on $m_{\text{inv}}$. The cumulative signal efficiency is the product of the acceptance and experimental efficiency.

In the electron channel, the multijet background is characterized by small values of $E_T^{\text{miss}}$, while EW background events can have large $E_T^{\text{miss}}$ due to the production of high-momentum neutrinos. The discriminating power of $E_T^{\text{miss}}$ is used to determine the normalization of the two backgrounds to the data in the LIMCR. The multijet template is taken from data in which electron candidates pass relaxed identification criteria but fail the normal
identification selection, and the EW template is taken from MC simulation where electron candidates pass the normal selection. The templates are fit to the $E_T^{\text{miss}}$ distribution in the interval $[0,150]$ GeV in five separate detector-motivated regions of $\eta$, to determine normalization factors for both the multijet and EW backgrounds.

To extrapolate both the multijet and EW background to the SR, functions of the form $p_1 x p_2 + p_3 \ln(x) (1 - x)^p_4$ (with $x = m_{\text{inv}}/\sqrt{s}$ and fit parameters $p_1 - p_4$) are used and the contributions are scaled by the corresponding normalization factor derived in the LIMCR. A simple power-law fit, with $p_3$ and $p_4$ fixed to zero, adequately describes both data and simulation. This is used as the baseline, while $p_3$ and $p_4$ are allowed to vary as part of the evaluation of the systematic uncertainty.

In the muon channel, the multijet and EW backgrounds can be discriminated on the basis of the transverse impact parameter ($d_0$) distribution of the muon since multijet background is dominated by jets containing charm and bottom hadrons decaying to muons while EW backgrounds are dominated by prompt muons. The template for the EW background is selected using $Z$-boson decays to two muons while the template for multijet background is taken from muons that fail the isolation requirement. Both templates are taken from data. The templates are fit to the $d_0$ distribution in the interval $[-0.1, +0.1]$ mm to determine the normalization factors. The fraction of multijet background, $0.046 \pm 0.005$, is neglected when extrapolating the background in SR. The procedure for extrapolating the EW background to the SR is the same as for the electron channel.

The background estimate in the SR, shown in Fig. 1, was not compared to data until the final fit method and parameters were fixed. The hatched area in Fig. 1 shows the total uncertainty in the background estimate, which is dominated by the systematic uncertainties. In extracting the limits, the fits described above are used to extrapolate the background into the high invariant-mass region.

The systematic uncertainties on the background are evaluated as a function of $m_{\text{inv}}$, and are dominated by uncertainties on the fits used to extrapolate the background to the highest $m_{\text{inv}}$, uncertainties on PDFs, and the choice of MC generator. Systematic uncertainties due to the choice of fitting functions are evaluated by fitting the $m_{\text{inv}}$ spectrum with parameters $p_3$ and $p_4$ free and taking the difference between these fits and the fits with $p_3$ and $p_4$ fixed to zero. Additionally, SHERPA samples are used instead of ALPGEN and the fits are repeated. The uncertainty in the PDFs is estimated using a set of 44 PDF eigenvectors for CTEQ6.6 [43]. For each of the 44 sets, the background fits are repeated and the extrapolated backgrounds are estimated. To estimate the uncertainty in the multijet background in the electron channel, an alternative selection of background-enriched data events, based on photons, is used. The systematic uncertainties from the simulation of the detector response are associated with the jet and electron energy scales and resolutions, the muon momentum scale and resolution, and the trigger requirement. The combined uncertainty in the background prediction ranges from 16% (1 TeV) to 100% (6 TeV) for the electron channel and from 50% (1 TeV) to 170% (6 TeV) for the muon channel. Background systematic uncertainties for $M_{M_{\text{TH}}} = 5$ TeV are given in Table I.

Uncertainties on the signal efficiency in each of the mass bins are associated with the requirements on $\Delta y$, $\Delta \phi$, $\langle n \rangle_{\text{jet}}$, $m_{\text{inv}}$, and isolation. In addition, uncertainties on the detector simulation, mentioned above for the background, as well as the uncertainty in luminosity are taken into account. The combined uncertainty in the signal efficiency from these sources ranges from 3.5% at 1 TeV to
TABLE I. Breakdown of relative systematic uncertainties on the SM background for the threshold mass $M_{th} = 5$ TeV. The uncertainties are added in quadrature to obtain the total uncertainty.

<table>
<thead>
<tr>
<th>Source</th>
<th>Electron+jet %</th>
<th>Muon+jet %</th>
</tr>
</thead>
<tbody>
<tr>
<td>Lepton reconstruction, scale and resolution</td>
<td>29 (4.3)</td>
<td>18 (2.7)</td>
</tr>
<tr>
<td>Jet reconstruction, scale and resolution</td>
<td>31 (4.7)</td>
<td>19 (2.8)</td>
</tr>
<tr>
<td>Multijet modeling</td>
<td>23 (3.4)</td>
<td>14 (2.0)</td>
</tr>
<tr>
<td>PDF</td>
<td>35 (5.1)</td>
<td>29 (4.1)</td>
</tr>
<tr>
<td>Fit</td>
<td>49 (7.1)</td>
<td>35 (4.9)</td>
</tr>
<tr>
<td>Total</td>
<td>100 (9.9)</td>
<td>87 (7.4)</td>
</tr>
</tbody>
</table>

3.9% at 6 TeV for the electron channel and from 3.6% at 1 TeV to 5.6% at 6 TeV for the muon channel. The cumulative efficiency, shown in Table II, is taken from the signal MC simulation for charge +4/3 QBHs. The differences in the efficiency between the charge +4/3 state and the other charged states are much smaller than the uncertainties mentioned above and are neglected. The effect of the 0.65% uncertainty in the LHC beam energy is to change the QBH production cross section. Since the QBH cross section is nearly constant in $M_{th}/\sqrt{s}$ this is effectively an uncertainty in $M_{th}$ and has a negligible effect on the limits.

The observed numbers of events and the expected backgrounds, shown in Table II, are in agreement within the total uncertainty. There is no evidence for any excess. Upper limits on $\Sigma\sigma_{qq} \times BF_{qq}$ for the production of QBHs above $M_{th}$ are determined in the interval 1–6 TeV assuming lepton universality and using the CLs method, which is designed to give conservative limits in cases where the observed background fluctuates below the expected values. The statistical combination of the channels employs a likelihood function constructed as the product of Poisson probability terms describing the total number of events observed in each channel. Systematic uncertainties are incorporated as nuisance parameters into the likelihood through their effect on the mean of the Poisson functions and through convolution with their assumed Gaussian distributions. Correlations between channels are taken into account.

Figure 2 shows the 95% confidence level (C.L.) combined lepton+jet upper limit on the cross section times branching fraction for the production of QBHs as a function of $M_{th}$. Above 3.5 TeV, the limit is 0.18 fb. For the $n = 6$ QBH model assumed in this Letter, the 95% C.L. lower limit on $M_{th}$ is 5.3 TeV. For $n = 2$, and all other model assumptions the same, the 95% C.L. lower limit on $M_{th}$ is 4.7 TeV. Treating the channels separately, the 95% C.L. upper limit on the electron (muon)+jet $\Sigma\sigma_{qq} \times BF_{qq}$ above 3.5 TeV is 0.27 (0.49) fb, and the $n = 6$ lower limit on $M_{th}$ is 5.2 (5.1) TeV.

In conclusion, a first search for two body lepton+jet final states with large invariant mass has been performed using 20.3 fb$^{-1}$ of $pp$ collisions recorded at $\sqrt{s} = 8$ TeV with the ATLAS detector at the LHC. In the invariant-mass region above 1 TeV the observed events are consistent with data-driven extrapolated backgrounds from the low-invariant-mass control region. Above 3.5 TeV the ex-
pected background drops below one event and the 95% C.L. upper limit on the electron (m\text{e}n\text{e})+\text{-jet} \Sigma \sigma_{\text{QBH}} \times BF_{\text{QBH}} is 0.27 (0.49) fb. Assuming lepton universality, the 95% C.L. upper limit on the sum of the product of QBH lepton+jet production cross sections and branching fractions is 0.18 fb.

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[16] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the z-axis along the beam pipe. The x-axis points from the IP to the center of the LHC ring, and the y-axis points upward. The pseudorapidity is defined in terms of the polar angle \theta as \eta = -\ln \tan(\theta/2). Cylindrical coordinates (r, \phi) are used in the transverse plane, \phi being the azimuthal angle around the beam pipe, referred to the x-axis.


Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
Department of Physics, Indiana University, Bloomington IN, United States of America
Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
University of Iowa, Iowa City IA, United States of America
Department of Physics and Astronomy, Iowa State University, Ames IA, United States of America
Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
Graduate School of Science, Kobe University, Kobe, Japan
Faculty of Science, Kyoto University, Kyoto, Japan
Kyoto University of Education, Kyoto, Japan
Department of Physics, Kyushu University, Fukuoka, Japan
Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
Physics Department, Lancaster University, Lancaster, United Kingdom
Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia
School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
Department of Physics, Royal Holloway University of London, Surrey, United Kingdom
Department of Physics and Astronomy, University College London, London, United Kingdom
Louisiana Tech University, Ruston LA, United States of America
Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
Fysiska institutionen, Lunds universitet, Lund, Sweden
Departamento de Física Teorica C-15, Universidad Autonoma de Madrid, Madrid, Spain
Institut für Physik, Universität Mainz, Mainz, Germany
School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
Department of Physics, University of Massachusetts, Amherst MA, United States of America
Department of Physics, McGill University, Montreal QC, Canada
School of Physics, University of Melbourne, Victoria, Australia
Department of Physics, The University of Michigan, Ann Arbor MI, United States of America
Department of Physics and Astronomy, Michigan State University, East Lansing MI, United States of America
(a) INFN Sezione di Milano; (b) Dipartimento di Fisica, Università di Milano, Milano, Italy
B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus
National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus
Department of Physics, Massachusetts Institute of Technology, Cambridge MA, United States of America
Group of Particle Physics, University of Montreal, Montreal QC, Canada
P.N. Lebedev Institute of Physics, Moscow, Russia
Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
Moscow Engineering and Physics Institute (MEPhI), Moscow, Russia
D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
Nagasaki Institute of Applied Science, Nagasaki, Japan
Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
(a) INFN Sezione di Napoli; (b) Dipartimento di Scienze Fisiche, Università di Napoli, Napoli, Italy
Department of Physics and Astronomy, University of New Mexico, Albuquerque NM, United States of America
Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
Department of Physics, Northern Illinois University, DeKalb IL, United States of America
Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
Department of Physics, New York University, New York NY, United States of America
Ohio State University, Columbus OH, United States of America
Faculty of Science, Okayama University, Okayama, Japan
Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK, United States of America
Department of Physics, Oklahoma State University, Stillwater OK, United States of America
Palacký University, RCPTM, Olomouc, Czech Republic
Center for High Energy Physics, University of Oregon, Eugene OR, United States of America
LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
Graduate School of Science, Osaka University, Osaka, Japan
Department of Physics, University of Oslo, Oslo, Norway
Department of Physics, Oxford University, Oxford, United Kingdom
(a) INFN Sezione di Pavia; (b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
Department of Physics, University of Pennsylvania, Philadelphia PA, United States of America
Petersburg Nuclear Physics Institute, Gatchina, Russia
(a) INFN Sezione di Pisa; (b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA, United States of America
(a) Laboratorio de Instrumentación e Física Experimental de Partículas - LIP, Lisboa, Portugal; (b) Departamento de Física Teórica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain
Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic
Czech Technical University in Prague, Praha, Czech Republic
Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic
State Research Center Institute for High Energy Physics, Protvino, Russia
Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
Physics Department, University of Regina, Regina SK, Canada
Ritsumeikan University, Kusatsu, Shiga, Japan
(a) INFN Sezione di Roma I; (b) Dipartimento di Fisica, Università La Sapienza, Roma, Italy
(a) INFN Sezione di Roma Tor Vergata; (b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
(a) INFN Sezione di Roma Tre; (b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
(a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca; (b) Centre National de l’Énergie des Sciences Techniques Nucléaires, Rabat; (c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech; (d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; (e) Faculté des sciences, Université Mohammed V-Agdal, Rabat, Morocco
DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Énergie Atomique et aux Énergies Alternatives), Gif-sur-Yvette, France
Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA, United States of America
Department of Physics, University of Washington, Seattle WA, United States of America
Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
Department of Physics, Shinshu University, Nagano, Japan
Fachbereich Physik, Universität Siegen, Siegen, Germany
Department of Physics, Simon Fraser University, Burnaby BC, Canada
SLAC National Accelerator Laboratory, Stanford CA, United States of America
(a) Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
(a) Department of Physics, University of Cape Town, Cape Town; (b) Department of Physics, University of Johannesburg, Johannesburg; (c) School of Physics, University of the Witwatersrand, Johannesburg, South Africa
(a) Department of Physics, Stockholm University; (b) The Oskar Klein Centre, Stockholm, Sweden
Physics Department, Royal Institute of Technology, Stockholm, Sweden
Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook NY, United States of America
Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
School of Physics, University of Sydney, Sydney, Australia
Institute of Physics, Academia Sinica, Taipei, Taiwan
Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel
Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan
Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
(a) TRIUMF, Vancouver BC; (b) Department of Physics and Astronomy, York University, Toronto ON, Canada
Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan
Department of Physics and Astronomy, Tufts University, Medford MA, United States of America
Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia
Department of Physics and Astronomy, University of California Irvine, Irvine CA, United States of America
(a) INFN Gruppo Collegato di Udine; (b) ICTP, Trieste; (c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy
Department of Physics, University of Illinois, Urbana IL, United States of America
Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNMTM), University of Valencia and CSIC, Valencia, Spain
Department of Physics, University of British Columbia, Vancouver BC, Canada
Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada
Department of Physics, University of Warwick, Coventry, United Kingdom
Waseda University, Tokyo, Japan
Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel
Department of Physics, University of Wisconsin, Madison WI, United States of America
Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany
Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany
Department of Physics, Yale University, New Haven CT, United States of America
Yerevan Physics Institute, Yerevan, Armenia
Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France
Also at Department of Physics, King’s College London, London, United Kingdom
Also at Laboratorio de Instrumentacão e Física Experimental de Partículas - LIP, Lisboa, Portugal
Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
Also at Faculdade de Ciencias and CFNUL, Universidade de Lisboa, Lisboa, Portugal
Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
Also at TRIUMF, Vancouver BC, Canada
Also at Department of Physics, California State University, Fresno CA, United States of America
Also at Novosibirsk State University, Novosibirsk, Russia
Also at Department of Physics, University of Coimbra, Coimbra, Portugal
Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
Also at Università di Napoli Parthenope, Napoli, Italy
Also at Institute of Particle Physics (IPP), Canada
Also at Department of Physics, Middle East Technical University, Ankara, Turkey
Also at Dep Fisica and CEFITEC of Faculdade de Ciencias e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal
Also at Department of Physics and Astronomy, Michigan State University, East Lansing MI, United States of America
Also at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece
Also at Louisiana Tech University, Ruston LA, United States of America
Also at Institucio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain
Also at Department of Physics, University of Cape Town, Cape Town, South Africa
Also at CERN, Geneva, Switzerland
Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan
Also at Manhattan College, New York NY, United States of America
Also at Institute of Physics, Academia Sinica, Taipei, Taiwan
Also at School of Physics and Engineering, Sun Yat-sen University, Guanzhou, China
Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan
Also at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
Also at School of Physical Sciences, National Institute of Science Education and Research, Bhubaneswar, India
Also at Dipartimento di Fisica, Università La Sapienza, Roma, Italy
Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia
Also at Section de Physique, Université de Genève, Geneva, Switzerland
Also at Departamento de Fisica, Universidade de Minho, Braga, Portugal
Also at Department of Physics, The University of Texas at Austin, Austin TX, United States of America
Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary
Also at DESY, Hamburg and Zeuthen, Germany
Also at International School for Advanced Studies (SISSA), Trieste, Italy
Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, United States of America
Also at Faculty of Physics, M.V.Lomonosov Moscow State University, Moscow, Russia
Also at Physics Department, Brookhaven National Laboratory, Upton NY, United States of America
Also at Moscow Engineering and Physics Institute (MEPhI), Moscow, Russia
Also at Department of Physics, Oxford University, Oxford, United Kingdom
Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany
Also at Department of Physics, The University of Michigan, Ann Arbor MI, United States of America
Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa
* Deceased