Novel Phases in the Field Induced Spin Density Wave State in (TMTSF)$_2$PF$_6$

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Magnetoresistance measurements on the quasi one-dimensional organic conductor (TMTSF)$_2$PF$_6$ performed in magnetic fields $B$ up to 16 T, temperatures $T$ down to 0.12 K and under pressures $P$ up to 14 kbar have revealed new phases on its $P-B-T$ phase diagram. We found a new boundary which subdivides the field induced spin density wave (FISDW) phase diagram into two regions. We showed that a low-temperature region of the FISDW diagram is characterized by a hysteretic behavior typical for the first order transitions, as observed in a number of studies. In contrast to the common believe, in high temperature region of the FISDW phase diagram, the hysteresis and, hence, the first order transitions were found to disappear. Nevertheless, sharp changes in the resistivity slope are observed both in the low and high temperature domains indicating that the cascade of transitions between different subphases exists over all range of the FISDW state. We also found that the temperature dependence of the resistance (at a constant $B$) changes sign at about the same boundary. We compare these results with recent theoretical models.

Layered organic compounds tetramethiltetraselenafulvalene (TMTSF)$_2$X, where the anion $X$ is ClO$_4$, AsF$_6$, PF$_6$ etc are unique material systems with a very rich phase diagram (for a review see Refs. [1–4]). Conduction vallene (TMTSF)$_2$X is F ermi wave vector, $\mathbf{q}$, and $\mathbf{F}$ is F is the size of the electron-hole symmetry, in momentum space. Under the assumption of the electron-state whose period determines a nesting vector in the $\mathbf{q}$, first quenches the superconducting state and, further induces a cascade of phase transitions between FISDW states accompanied by the quantum Hall effect [8,9].

A so called ‘standard’ model was suggested [1,2,7] to explain the metal-SDW transition in magnetic field. Later it was developed into a ‘Quantized Nesting Model’ in Refs. [8–11] to describe a cascade of the first order transitions between different FISDW sub-phases. According to this model, electrons condense in the SDW state whose period determines a nesting vector in the momentum space. Under the assumption of the electron-hole symmetry, $x$–component of the nesting vector $Q_x$ is quantized as

$$Q_x = 2k_F - N\frac{eBb}{h},$$

where $k_F$ is Fermi wave vector, $b$ is the size of the elementary cell in $y$– direction, and $N$ is an integer. As magnetic field varies, $N$ changes by an integer, causing step-like changes in the nesting vector, which result in the sequence of the first-order phase transitions. According to the recent analysis [12], however the electron-hole symmetry in the SDW state is not fulfilled unless $N = 0$ in Eq. (1). As a result, (i) the nesting vector is not strictly quantized and (ii) the step-like changes in the nesting vector may disappear above a certain temperature transforming into oscillations. Correspondingly, as temperature increases, the first order transitions with $\Delta N \approx 1$ may disappear, whereas FISDW state still persists. Thus, in contrast to the ‘Quantized Nesting Model’ which predicts the first order phase transition to exist over the whole range of temperatures where FISDW develops, the ‘novel model’ predicts the first order phase transitions may disappear above a certain temperature, $T_0$. The latter possibility depends on the parameter $h\omega_c/(2\pi k_F T_0)$ [12], where $\omega_c$ is the cyclotron frequency.

In order to verify the theoretical predictions of the two models above, we studied temperature dependence of the magnetoresistance in (TMTSF)$_2$PF$_6$ at various pressures. Specifically, we measured, at different pressures, a temperature evolution of the hysteresis intrinsic to magnetoresistance traces of $R(B)$. We observed that the hysteresis indeed disappears above a temperature $T_0$ whereas FISDW state still persists. We found such behavior to manifest itself over the whole explored range of the existence of the FISDW. According to our results, the total $P-B-T$ phase diagram of the FISDW state can be subdivided into two domains, the ‘low-$T$’ domain where the first order phase transitions between FISDW sub-phases take place, and the ‘high-$T$’ domain where the transitions between the FISDW states do not exhibit first order behavior. This observation is in agreement with the ‘novel model’; in the latter case, the ‘low $T$-phase’ is treated as a ‘quantum FISDW’ state with step-like changes in the
nesting vector, whereas the ‘high $T$-phase’ is treated as
the ‘semi-classical FISDW’ state where the nesting vector
oscillates. We also found that as temperature decreases
and crosses the boundary between the two domains, the
behavior of the resistivity changes qualitatively, from or-
dinary insulating-like ($dR/dT < 0$) through the ‘high-
$T$’-domain, to the metallic-like ($dR/dT > 0$) through the
‘low-$T$’-domain. Around the $T = T_0$-boundary, $R(T)$ ex-
hibits a maximum.

**Experimental.** Three samples (of a typical size $2 \times 0.8 \times
0.3\, \text{mm}^3$) were grown by a conventional electrochemical
technique. Measurements were made using either four
Ohmic contacts formed at the $a - b$ plane or eight con-
tacts at two $a - c$ planes; in all cases $25\mu \text{m}$ Pt-wires were
attached by a graphite paint to the sample along the most
dominating direction $a$. The sample and a manganin pressure
gauge were inserted into a Teflon cylinder placed inside a nonmagnetic 18 mm o.d. pressure
chamber, in a bore of a 16 T superconducting
magnet. For all measurements, the magnetic field was
applied along the least conducting direction, $z$, of the
crystal. Sample resistance was measured by four probe
ac technique at 132 Hz frequency, with current 1-4 $\mu\text{A}$
to avoid nonlinear effects. The out-of phase component
of the measured voltage was found to be negligible in all
measurements, indicating Ohmic contacts to the sample.

The sample temperature was varied slowly, at a rate
less than 0.25 K/min in order to avoid breaking the sam-
ple. The measured changes in the sample resistance were
fully reproducible during the full run of measurements
including temperature sweeps; this indicated that the
sample quality did not change. The magnetoresistance
was measured either at a constant $T$ and varying mag-
netic field $B$, or at a constant $B$ and varying $T$. Sample
temperature was determined by RuO$_2$ resistance ther-
ometer with a pre-calibrated magnetoresistance. Mea-
surements were done in magnetic fields up to 16 T and
for temperatures in the range from 1.4 to 30 K (mainly)
and down to 0.12 K (partly). The most detailed results
were obtained for pressures 7, 8, 10 and 14 kbar.

Figure 1 shows magnetoresistance traces measured (a)
at $P = 10$ kbar in the temperature range 0.6-4.2 K and
(b) at 8 kbar, 0.12 K. In agreement with earlier observa-
tions [1], when magnetic field exceeds the critical value
(which is 0.16 T in our case), the superconductivity is
quenched and the sample resistance starts gradually in-
creasing. Further, this smooth dependence transforms
into step-like changes in $R$. As temperature decreases,
the step-like changes become steeper and appear at pro-
gressively lower fields. This behaviour is also consistent
with earlier observations [1,5,6,7] and is interpreted as
transitions between different sub-phases in FISDW
[7,11,12]. This interpretation is further supported by the
hysteresis between $R(B)$ traces for the field ramping up

and down, which is clearly seen in Fig. 1. The hysteresis
is also consistent with earlier observations [1] and signals
the onset of the first order phase transitions.

![FIG. 1. Magnetoresistance $R_{xx}$ vs magnetic field $B_z$: a) For $P = 10$ kbar and for eight temperatures (indicated for each curve). The curves are shifted vertically, for clar-
ity (the offset values are indicated on the left side). The inset magnifies the temperature evolution of the hysteresis regions of $R(B)$-curves near $B = 15$ T. b) For $P = 8$ kbar
and $T = 0.12$ K.]

As temperature increases, the hysteresis weakens and
tends to disappear as illustrated in the inset to Fig. 1 a.
Nevertheless, the steps in $R(B)$ persist to higher temper-
atures, being therefore non- or at least partly correlated
with the hysteresis. In order to quantify the hysteresis
strength, we calculated the maximal width of the hyste-
resis loop $\Delta B$ for each curve and plotted it in Fig. 2
as a function of temperature; in this determination, the
$L/R$ time constant of the magnet was carefully measured
and taken into account.

For three groups of the data in Fig. 2 the hysteresis
width decreases linearly with temperature and vanishes
at a certain temperature $T_0$; above $T = T_0$ it remains
equal to zero. The falling part of these dependences
were fitted with linear curves (solid lines), which appear
to have the same slope. We plotted linear curves with
the same slope through other single data points (dashed
tines) in order to estimate $T_0$ for all transitions at differ-
ent pressure values.

Measurements at two other pressures, 7 and 14 kbar
have shown qualitatively similar results. At $P = 7$ kbar
the steps (transitions) shift to lower fields and persist
up to higher temperatures. The hysteresis, $\delta B$, is bigger than that at 8 and 10 kbar and disappears at slightly higher temperature. At $P = 14$ kbar, the trend is opposite: $T_0$ becomes lower than that for 10 kbar.

Figure 3 represents the results of temperature sweeps taken at six fixed magnetic fields for $P = 7$ kbar. Starting from high temperatures, the resistance increases as $T$ decreases, then passes through a maximum at a certain temperature $T_{\text{max}}$ and further decreases towards low temperatures. Triangles mark the onset of the phase transition at each curve. The similar $T$–dependences measured at $P = 10$ and 14 kbar were qualitatively similar to those shown in Fig. 3 but shifted to lower temperatures.

$B - T$ phase diagram in Fig. 4 summarizes the results of all measurements, at $P = 7$ kbar (the main panel) and at $P = 10$ kbar (the inset). The closed squares depict the onsets of the steps in $R(B)$ obtained from magnetic field sweeps at fixed temperatures and triangles are for the temperature sweeps $R(T)$ at fixed field.

In addition to the data taken directly, the open squares show the lower temperature data, $T = 0.12$ K, taken at $P = 8$ kbar which has been recalculated to correspond to the data at $P = 7$ kbar (main panel) and to 10 kbar (inset). In this procedure, the data for 8 kbar were shifted in magnetic field according to the pressure coefficient $d(B^{-1})/dP = -0.015 T^{-1}\text{kbar}^{-1}$ which we determined from the higher temperature data at $T = 1.4$ K for $P = 7$, 10, and 14 kbar. The hysteresis width is obtained from Fig. 2 and is depicted by the split lines.

The above three features, (i) the existence of the hysteresis in $R(B)$ at low temperature, (ii) its disappearance above a certain temperature $T_0$ and (iii) the persistence of the steps in $R(B)$ to temperatures higher than $T_0$, are observed in our experiments for several transitions (see Figs. 1, 2). It seems likely that these features are generic also to other transitions in the FISDW part of the phase diagram and that the hysteresis for higher $N$-values was not observed in our measurements just because $T_0$ for these transitions is lower than our lowest accessible temperature, 1.4 K (for the majority of measurements).
domains vs temperature. The hysteresis domains for different transitions collapse above $T = T_0$; this was determined for seven transitions and the corresponding points are denoted with open circles. The separatrices $T_0$ thus subdivide the phase boundaries into the two regions: the low temperature domain ($T < T_0$) of the hysteretic behaviour and the high temperature domain ($T > T_0$) where the FISDW transitions develop without a hysteresis. The disappearance of the hysteresis with rising temperature at one fixed pressure was mentioned earlier but to the best of our knowledge no studies of this effect followed. The subdivision of the phase diagram is qualitatively consistent with the ‘novel model’ for the FISDW [12], where the low-temperature domain corresponds to the ‘quantum FISDW’ sub-phase and the high temperature domain corresponds to the ‘semiclassical FISDW’. According to the model, in the former sub-phase the transitions between different FISDW states take place (i) with a hysteresis behaviour and the high temperature domain ($T > T_0$) where the FISDW transitions develop without a hysteresis. The disappearance of the hysteresis with rising temperature at one fixed pressure was mentioned earlier [1] but to the best of our knowledge no studies of this effect followed. The subdivision of the phase diagram is qualitatively consistent with the ‘novel model’ for the FISDW [12], where the low-temperature domain corresponds to the ‘quantum FISDW’ sub-phase and the high temperature domain corresponds to the ‘semiclassical FISDW’. According to the model, in the former sub-phase the transitions between different phases take place with jumps in the nesting vector, are of the first order, and are accompanied by hysteresis of various physical quantities. In the latter sub-phase, the transitions between different phases take place without any jumps in the nesting vector, and are therefore not first order phase transitions.

The $B - T$ phase diagram on the main panel in Fig. 4 is plotted for pressure of 7 kbar. For $P = 10$ kbar (inset), the phase diagram looks qualitatively similar but the new phase boundary is shifted to lower temperature by about $(0.5 \sim 1)$ K.

The dashed line in Fig. 4 connects the $T_0$—values (circles). In order to check whether the dashed line splits not only the phase boundaries but the overall $B - T$ parameter space, we plotted onto the same phase diagram the $B, T$-coordinates of the maxima on $R(T)$ curves from $T$—sweeps similar to that shown in Fig. 3.

The maxima in $R(T)$ dependences in the FISDW regime were observed earlier [20,21] and were associated with a non-linear temperature dependence of the $\sigma_{xy}$ caused by the transition from the QHE to metallic regime [20,22,21]. In this interpretation, the maxima should correspond to the onset of the QHE regime. On the other hand, the hysteresis behaviour if associated with jumps in the nesting vector [12], should manifest in the QHE regime only. Therefore, $T_{\text{max}}$ is expected to be $\leq T_0$. It is surprising that in our case $T_{\text{max}}$ almost coincide with $T_0$ for different pressures, even though $T_0$ varies with pressure.

To summarize, from studies of the temperature and magnetic field dependences of the resistivity of the quasi-1D organic conductor, we found that its $P - B - T$ phase diagram splits in two domains, where the transitions between different FISDW states take place (i) with a hysteresis as the first order phase transitions (for low temperatures), and (ii) without hysteresis (for high temperatures). This result is not expected within the ‘Quantized nesting model’ and is consistent with the recent suggestion by Lebed [12] that the period of the spin structure in FISDW state can be either partially quantized or not quantized at all. We experimentally plotted the new phase boundary where such behavior occurs. We found that the $P - B - T$ boundary separating the two domains almost coincides with a line at which the $dR/dT$ changes sign.

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