Search for Pair Production of a Heavy Up-Type Quark Decaying to a W Boson and a b Quark in the Lepton+jets Channel with the ATLAS Detector

ATLAS Collaboration

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A search is presented for production of a heavy up-type quark ($t'$) together with its antiparticle, assuming subsequent decay to a $W$ boson and a $b$ quark, $t'\bar{t'} \to W^+ b W^- \bar{b}$. The search is based on 1.04 fb$^{-1}$ of proton-proton collisions at $\sqrt{s} = 7$ TeV collected by the ATLAS detector at the CERN Large Hadron Collider. Data are analyzed in the lepton+jets final state, characterized by a high transverse momentum isolated electron or muon, high missing transverse momentum and at least three jets. No significant excess of events above the background expectation is observed. A 95% C.L. lower limit of 404 GeV is set for the mass of the $t'$ quark.

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The discovery of the top quark [1] completed the third generation of fundamental fermions in the quark sector of the Standard Model (SM) of particle physics. It is natural to ask whether heavier quarks may exist. These quarks are often present in new physics models aimed at solving the limitations of the SM. For example, models with a fourth generation of heavy chiral fermions could provide new sources of CP violation to explain the matter-antimatter asymmetry in the Universe, and allow for a heavier Higgs boson while remaining consistent with precision electroweak data [2]. The latter is accomplished by keeping a small mass splitting between the heavy up-type quark ($t'$) and the heavy down-type quark ($b'$). Assuming that $m_{t'} - m_{W} < m_{W}$, where $m_{W}$ is the $W$ boson mass, results in the $t'$ quark predominantly decaying to a $W$ boson and a down-type quark $q$ ($q=d, s, b$). Another possibility is the addition of isospin singlets or doublets of vector-like quarks, which appear in many extensions of the SM such as Little Higgs or extra-dimensional models [3]. In both scenarios the $t'$ quark can decay into $Wb$ with a large branching ratio, provided there is a significant mixing with the third generation of quarks, consistent with the existing mass and mixing patterns of the known quarks.

The high center-of-mass energy and integrated luminosity in $pp$ collisions available at the Large Hadron Collider (LHC) offers a unique opportunity to probe these scenarios. At the LHC, these new heavy quarks would be predominantly produced in pairs via the strong interaction for masses below $\sim 1$ TeV, while for larger masses electroweak production of single heavy quarks could become the primary production mechanism, depending on the strength of their interactions with the SM quarks and weak gauge bosons [3].

A search is presented in this Letter for $t'\bar{t'}$ production using $pp$ collision data at $\sqrt{s} = 7$ TeV collected with the ATLAS detector. It is assumed that the $t'$ quark decays exclusively into $Wb$. The lepton+jets final state signature is considered, characterized by a high transverse momentum ($p_T$) isolated electron or muon, high missing transverse momentum ($E_T^{\text{miss}}$) and at least three jets. Similar searches in this channel have been published by the CDF and D0 collaborations [4,5]; the most stringent limits preclude the existence of a $t'$ quark with a mass below 358 GeV at 95% confidence level (C.L.). A search for $t'\bar{t'}$ in the dilepton final state has been performed by the ATLAS collaboration [6], excluding a $t'$ quark with a mass below 350 GeV at 95% C.L. The lepton+jets signature has also been recently exploited by the ATLAS collaboration to search for $b'b' \to W^- W^+ t\bar{t}'$ [7].

The ATLAS detector [8] consists of an inner tracking system surrounded by a superconducting solenoid, electromagnetic and hadronic calorimeters, and a muon spectrometer (MS). The inner detector is immersed in a 2 T axial magnetic field, and consists of pixel and silicon microstrip detectors inside a transition radiation tracker, providing charged particle tracking in the region $|\eta| < 2.5$ [9]. The electromagnetic (EM) calorimeter is based on lead/liquid-argon (LAr). Hadron calorimetry is based on two different detector technologies, with scintillator tiles or LAr as active media, and with either steel, copper, or tungsten as the absorber material. The calorimeters provide coverage up to $|\eta| < 4.9$. The MS consists of superconducting air-core toroids, a system of trigger chambers covering the range $|\eta| < 2.4$, and high-resolution tracking chambers allowing muon momentum measurements within $|\eta| < 2.7$.

The data set used in this analysis was recorded between March and June 2011 using single electron and muon triggers and includes only events collected under stable beam conditions and for which all detector subsystems were fully operational. The corresponding integrated luminosity is 1.04 fb$^{-1}$. The event selection criteria closely follow those used in recent ATLAS top quark studies, e.g. Ref. [10]. Electron candidates are required to satisfy $p_T > 25$ GeV and $|\eta| < 2.47$, excluding the transition
region $1.37 < |\eta| < 1.52$ between the barrel and endcap EM calorimeters. Muon candidates are required to satisfy $p_T > 20$ GeV and $|\eta| < 2.5$. The $p_T$ threshold requirement ensures that the selected leptons are in the efficiency plateau of the single-lepton triggers. Background from multi-jet production is suppressed by a requirement of $E_T^{\text{miss}} > 35(20) \text{ GeV}$ [11] in the electron (muon) channel, followed by $E_T^{\text{miss}} + m_T > 60$ GeV, where $m_T$ is the transverse mass of the lepton and $E_T^{\text{miss}}$ [12]. The $E_T^{\text{miss}}$ is constructed from the vector sum of all calorimeter cells contained in topological clusters [13], calibrated at the energy scale of the associated high-$(p_T)$ object, and including contributions from selected muons. Further requirements are that there be at least three jets with $p_T > 25$ GeV and $|\eta| < 2.5$, with at least one jet satisfying $p_T > 60$ GeV. Jets are reconstructed with the anti-$k_t$ algorithm [14] with radius parameter $R = 0.4$, from topological clusters of energy deposits in the calorimeters calibrated at the EM scale. These jets are then calibrated to the particle level [15] using a $p_T$- and $\eta$-dependent correction factor derived from simulated events and validated using data. Finally, to further reduce the backgrounds, at least one jet is required to be identified as originating from the hadronization of a $b$ quark ($b$-tagging). This is achieved via an algorithm [16] using multivariate techniques to combine information from the impact parameters of displaced tracks as well as topological properties of secondary and tertiary decay vertices reconstructed within the jet: a working point is used with $\sim 70\%$ efficiency for $b$-quark jets and a rejection factor of $\sim 100$ for jets originating from light quarks ($u, d, s$) or gluons. Events with exactly one electron or one muon, and with exactly three jets or with four or more jets are analyzed separately to take advantage of their different signal-to-background ratio and background composition, as discussed below.

After event selection the main background is $t\bar{t}$ production, followed by the production of a $W$ boson in association with jets ($W$+jets). Smaller contributions arise from multi-jet events, single top quark, $Z$+jets and diboson production. All of the backgrounds which do not involve top quarks are significantly suppressed by the $b$-tagging requirement. Multi-jet events contribute to the selected sample via the misidentification of a jet or a photon as an electron or the presence of a non-prompt lepton, e.g. from a semileptonic $b$- or $c$-hadron decay. The normalization and shape of the multi-jet background kinematic distributions are estimated via data-driven methods [11]. For the $W$+jets background, the shape is estimated from the simulation but the normalization is estimated from the asymmetry between $W^+$+jets and $W^-$+jets production [17] in data. All other backgrounds, as well as the signal, are estimated from the simulation and normalized to their theoretical cross sections. A summary of the background estimates in each of the four channels analyzed, and a comparison with the observed yields in data are presented in Table II, showing a good agreement within the uncertainties.

Monte Carlo (MC) samples of $t\bar{t}$ and single top quark background are generated using MC@NLO v3.41 [13], assuming a top quark mass of 172.5 GeV, using the CTEQ6.6 set of parton distribution functions (PDF) [19], and are normalized to the approximate next-to-next-to-leading-order (NNLO) theoretical cross sections [20, 21]. Samples of $W/Z$+jets and $tW$+jets background are generated using ALPGEN v2.13 [22] and the CTEQ6L1 PDF set [19]. The $Z$+jets background is normalized to the NNLO theoretical cross section [23], while the $W$+jets background normalization is extracted from data. Both MC@NLO and ALPGEN are interfaced to HERWIG v6.5 [24] to model the parton shower and fragmentation, while JIMMY [25] is used to simulate the underlying event. The diboson backgrounds are modeled using HERWIG v6.5 and normalized to their NLO theoretical cross sections [26]. The signal is modeled using PYTHIA 6.421 [27]. Signal samples are generated for a range of masses, $m_{t\bar{t}}$, from 250 to 500 GeV in steps of 50 GeV and are normalized to the approximate NNLO theoretical cross sections [20] using the CTEQ6.6 PDF. The MC samples generated using HERWIG or PYTHIA use the MRST2007 LO* PDF set [28]. All MC samples include multiple $pp$ interactions and are processed through a full simulation [29] of the detector geometry and response using GEANT4 [30], and the same reconstruction software as the data. Simulated events are corrected to match the object identification efficiencies and resolutions determined in data control samples. The total signal detection efficiency, considering both lepton flavors and jet multiplicities analyzed, ranges from $5.2\%$ for $m_{t\bar{t}} = 250$ GeV to $17.3\%$ for $m_{t\bar{t}} = 500$ GeV.

This analysis uses the reconstructed heavy quark mass ($m_{\text{reco}}$) as the primary discriminating variable. In the case of events with $\geq 4$ jets, $m_{\text{reco}}$ is estimated by performing a kinematic likelihood fit [17] to the $t\bar{t} \rightarrow W^+W^-\ell\nu$ hypothesis, imposing the constraints that $t'$ and $\ell'$ have the same mass, and that the mass of the lepton-neutrino system, as well as that of a jet pair, equals the nominal $W$ boson mass. The final state objects considered are the lepton, $E_T^{\text{miss}}$ and the four jets with highest $p_T$. Among all possible jet-parton permutations, the one yielding the highest likelihood value after maximization over the fit parameters is kept. In the case of events with exactly three jets, $m_{\text{reco}}$ is taken to be the invariant mass of the three-jet system. In order to ensure a robust background prediction in the tail of the $m_{\text{reco}}$ distribution, a dynamic binning scheme is adopted; starting from the high side and low side of the distributions, bins are merged until the statistical uncertainty in the sum of the background predictions in that bin drops below 5%.

Systematic uncertainties affecting the normalization and shape of the $m_{\text{reco}}$ distribution are estimated for...
both signal and background, taking into account correlations among processes as well as channels. The dominant sources of uncertainty arise from the modeling of the $t\bar{t}$ background. The uncertainties on the $t\bar{t}$ background come from the theoretical uncertainty on the cross section ($\pm 0.6\%$) as well as the effects on both normalization and shape of the $m_{\text{reco}}$ distribution from a number of sources; these are uncertainties on the fragmentation model (based on the comparison of HERWIG and PYTHIA fragmentations), on the NLO event generator (based on the comparison of MC@NLO and POWHEG \cite{31}) and on the top quark mass (taken to be $\pm 1$ GeV).

The uncertainty on the jet energy scale affects the normalization of signal (2–12%) and backgrounds (5–30%) modeled through the simulation, as well as the shape of their $m_{\text{reco}}$ distributions.

Uncertainties on the modeling of initial- and final-state QCD radiation (ISR/FSR), evaluated by varying corresponding generator parameters, are considered as correlated between the $t\bar{t}$ background and the $t\bar{t}l^0$ signal.

While the normalization is obtained from the asymmetry measurement, the uncertainties on the normalization of the $W+$jets background are derived from measurements of $W+2$ jets dominated data samples and take into account the uncertainty on the heavy-flavor content of the samples as well as the extrapolation to higher jet multiplicities. The total uncertainty on the $W+$jets normalization is 50% and 70% for events with exactly 3 jets and $\geq 4$ jets, respectively. Uncertainties on the shape of the $m_{\text{reco}}$ distribution for the $W+$jets background are estimated by varying the choices of the matching scale (from 15 to 10 GeV) and the factorization scale (from $\mu_R^2 = m_W^2 + \sum p_{T,jet}^2$ to $\mu_R^2 = m_W^2 + p_{T,W}^2$) in ALPGEN.

Table I. Number of events observed compared to the background expectation after final event selection in each of the four channels considered. Also shown are the expected signal yields assuming $m_{t'} = 400$ GeV. The quoted uncertainties are prior to the fit to data and include both statistical and systematic contributions, taking into account correlations among processes.

<table>
<thead>
<tr>
<th>Channel</th>
<th>$e^+3$ jets</th>
<th>$\mu^+3$ jets</th>
<th>$e^{+4}$ jets</th>
<th>$\mu^{+4}$ jets</th>
</tr>
</thead>
<tbody>
<tr>
<td>$t\bar{t}$</td>
<td>$2320 \pm 460$</td>
<td>$3000 \pm 630$</td>
<td>$4470 \pm 920$</td>
<td>$5900 \pm 1200$</td>
</tr>
<tr>
<td>$W+$jets</td>
<td>$1440 \pm 790$</td>
<td>$2200 \pm 1200$</td>
<td>$830 \pm 580$</td>
<td>$1160 \pm 790$</td>
</tr>
<tr>
<td>$Z+$jets</td>
<td>$92 \pm 53$</td>
<td>$118 \pm 62$</td>
<td>$86 \pm 56$</td>
<td>$83 \pm 46$</td>
</tr>
<tr>
<td>Single top</td>
<td>$382 \pm 68$</td>
<td>$554 \pm 94$</td>
<td>$262 \pm 70$</td>
<td>$325 \pm 79$</td>
</tr>
<tr>
<td>Dibosons</td>
<td>$28 \pm 7$</td>
<td>$37 \pm 11$</td>
<td>$12 \pm 5$</td>
<td>$17 \pm 5$</td>
</tr>
<tr>
<td>Multi-jet</td>
<td>$520 \pm 520$</td>
<td>$550 \pm 550$</td>
<td>$320 \pm 320$</td>
<td>$340 \pm 340$</td>
</tr>
<tr>
<td>Total prediction</td>
<td>$4800 \pm 1000$</td>
<td>$6500 \pm 1500$</td>
<td>$6000 \pm 1100$</td>
<td>$7800 \pm 1400$</td>
</tr>
<tr>
<td>Data</td>
<td>$4533$</td>
<td>$6421$</td>
<td>$6145$</td>
<td>$8149$</td>
</tr>
<tr>
<td>$p_T$ (400 GeV)</td>
<td>$20.0 \pm 3.3$</td>
<td>$21.0 \pm 3.6$</td>
<td>$102.0 \pm 10.5$</td>
<td>$98.1 \pm 11.1$</td>
</tr>
</tbody>
</table>
event migration to different jet multiplicities when varying uncertainties such as jet energy scale or ISR/FSR. In addition to the jet multiplicity spectrum, the jet energy scale affects the peak position of the $m_{\text{reco}}$ spectrum for $t\bar{t}$ background, and can be constrained owing to the small uncertainty on the measured top quark mass \cite{X}. Nuisance parameters associated with smaller systematic uncertainties (e.g. lepton identification/trigger) are only weakly constrained.

Figure 1 shows a comparison of the post-fit $m_{\text{reco}}$ distribution between data and the background prediction for the combined $e/\mu + 3$ jets and $e/\mu + \geq 4$ jets channels. The fitted parameters are typically within one standard deviation of their nominal values and their uncertainties are consistent with expectations based on pseudo-experiments. Several additional studies were performed to check the integrity of the fitting procedure. The likelihood was verified to be parabolic near the minimum for each of the fitted parameters and to yield reasonable fit uncertainties; the lack of sensitivity to the assumed $p_T$ and $\eta$ correlation of the jet energy scale uncertainty was verified.

In the absence of any significant data excess, either in the $e$+jets or $\mu$+jets channels individually or in their combination, 95% C.L. upper limits on the $t\bar{t}$ production cross section are derived using the $CL_s$ method \cite{Y}, which employs the LLR test-statistic described above. Pseudo-experiments are generated under both the signal-plus-background ($s+b$) and background-only (b) hypotheses, taking into account per-bin statistical fluctuations of the total predictions according to Poisson statistics, as well as Gaussian fluctuations in the signal and background expectations describing the effect of systematic uncertainties. The fraction of $s+b$ and b pseudo-experiments with LLR larger than the median or observed LLR defines $CL_s$ and $CL_b$ for the expected or observed limits, respectively. Signal cross sections for which $CL_s = CL_s+b/CL_b < 0.05$ are deemed excluded at the 95% C.L.

The resulting observed and expected upper limits on the $t\bar{t}$ production cross section are shown in Fig. 2 as a function of the $t'$ mass, compared to the theoretical prediction, assuming a $BR(t' \rightarrow Wb) = 1$. As a result, an observed (expected) 95% C.L. lower limit of 404 (394) GeV on the mass of the $t'$ quark is derived.

In summary, a search for $t\bar{t}$ production has been performed in the lepton+jets final state under the assumption $BR(t' \rightarrow Wb) = 1$. No significant excess of events in the tail of the $m_{\text{reco}}$ distribution was found, resulting in an observed lower limit of $m_{t'} > 404$ GeV at 95% C.L. This represents the most stringent limit to date. This limit is also directly applicable to a down-type vector-like quark with electric charge of $-4/3$ decaying into a $W$ boson and a b quark \cite{Z}.

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FIG. 2. Observed (solid line) and expected (dashed line) 95% C.L. upper limits on the $t'\bar{t}'$ cross section as a function of the $t'$ mass. The surrounding shaded bands correspond to the 1 and 2 standard deviations (s.d.) around the expected limit. The thin line shows the theoretical prediction including its 1 s.d. uncertainty band. The shaded area is the mass region previously excluded by the CDF experiment. 

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[2] See e.g. B. Holdom et al., PMC Physics A 3, 4 (2009) and references therein.
[9] Pseudorapidity is defined as $\eta = -\ln[\tan(\theta/2)]$, where $\theta$ is the polar angle relative to the beam direction and $\phi$ is the azimuthal angle in the plane transverse to the beam direction.
[12] The transverse mass is defined by the formula $m_T = \sqrt{2p_T E_T^{miss}(1 - \cos \Delta\phi)}$, where $p_T$ is the lepton and $\Delta\phi$ is the azimuthal angle separation between the lepton and $E_T^{miss}$.
101 (a)INFN Sezione di Napoli; (b)Dipartimento di Scienze Fisiche, Università di Napoli, Napoli, Italy
102 Department of Physics and Astronomy, University of New Mexico, Albuquerque NM, United States of America
103 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
104 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
105 Department of Physics, Northern Illinois University, DeKalb IL, United States of America
106 Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
107 Department of Physics, New York University, New York NY, United States of America
108 Ohio State University, Columbus OH, United States of America
109 Faculty of Science, Okayama University, Okayama, Japan
110 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK, United States of America
111 Department of Physics, Oklahoma State University, Stillwater OK, United States of America
112 Palacký University, RCPTM, Olomouc, Czech Republic
113 Center for High Energy Physics, University of Oregon, Eugene OR, United States of America
114 LAL, Univ. Paris-Sud and CNRS/IN2P3, Orsay, France
115 Graduate School of Science, Osaka University, Osaka, Japan
116 Department of Physics, University of Oslo, Oslo, Norway
117 Department of Physics, Oxford University, Oxford, United Kingdom
118 (a)INFN Sezione di Pavia; (b)Dipartimento di Fisica, Università di Pavia, Pavia, Italy
119 Department of Physics, University of Pennsylvania, Philadelphia PA, United States of America
120 Petersburg Nuclear Physics Institute, Gatchina, Russia
121 (a)INFN Sezione di Pisa; (b)Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
122 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA, United States of America
123 (a)Laboratorio de Instrumentacao e Fisica Experimental de Partículas - LIP, Lisboa, Portugal; (b)Departamento de Fisica Teorica y del Cosmos and CAPE, Universidad de Granada, Granada, Spain
124 Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic
125 Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic
126 Czech Technical University in Prague, Praha, Czech Republic
127 State Research Center Institute for High Energy Physics, Protvino, Russia
128 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
129 Physics Department, University of Regina, Regina SK, Canada
130 Ritsumeikan University, Kusatsu, Shiga, Japan
131 (a)INFN Sezione di Roma I; (b)Dipartimento di Fisica, Università La Sapienza, Roma, Italy
132 (a)INFN Sezione di Roma Tor Vergata; (b)Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
133 (a)INFN Sezione di Roma Tre; (b)Dipartimento di Fisica, Università Roma Tre, Roma, Italy
134 (a)Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca; (b)Centre National de l’Energie des Sciences Techniques Nucleaires, Rabat; (c)Faculté des Sciences Semlalia, Université Cadi Ayyad, LPEA-Marrakech; (d)Faculté des Sciences, Université Mohammed Premier and LPTPM, Oujda; (e)Faculté des Sciences, Université Mohammed V- Agdal, Rabat, Morocco
135 DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat a l’Energie Atomique), Gif-sur-Yvette, France
136 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA, United States of America
137 Department of Physics, University of Washington, Seattle WA, United States of America
138 Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
139 Department of Physics, Shinshu University, Nagano, Japan
140 Fachbereich Physik, Universität Siegen, Siegen, Germany
141 Department of Physics, Simon Fraser University, Burnaby BC, Canada
142 SLAC National Accelerator Laboratory, Stanford CA, United States of America
143 (a)Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava; (b)Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
144 (a)Department of Physics, University of Johannesburg, Johannesburg; (b)School of Physics, University of the Witwatersrand, Johannesburg, South Africa
145 (a)Department of Physics, Stockholm University; (b)The Oskar Klein Centre, Stockholm, Sweden
146 Physics Department, Royal Institute of Technology, Stockholm, Sweden
Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook NY, United States of America

Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom

School of Physics, University of Sydney, Sydney, Australia

Institute of Physics, Academia Sinica, Taipei, Taiwan

Department of Physics, Technion: Israel Inst. of Technology, Haifa, Israel

Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan

Department of Physics and Technology, Tokyo Metropolitan University, Tokyo, Japan

Department of Physics, University of Sussex, Brighton, United Kingdom

School of Physics, University of Sydney, Sydney, Australia

Institute of Physics, Academia Sinica, Taipei, Taiwan

Department of Physics, Technion: Israel Inst. of Technology, Haifa, Israel

Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan

Department of Physics and Technology, Tokyo Metropolitan University, Tokyo, Japan

Department of Physics, University of Sussex, Brighton, United Kingdom

School of Physics, University of Sydney, Sydney, Australia

Institute of Physics, Academia Sinica, Taipei, Taiwan

Department of Physics, Technion: Israel Inst. of Technology, Haifa, Israel

Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan

Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan

Department of Physics, Tokyo Institute of Technology, Tokyo, Japan

Department of Physics, University of Toronto, Toronto ON, Canada

(a) TRIUMF, Vancouver BC; (b) Department of Physics and Astronomy, York University, Toronto ON, Canada

Institute of Pure and Applied Sciences, University of Tsukuba, 1-1-1 Tennodai, Tsukuba, Ibaraki 305-8571, Japan

Science and Technology Center, Tufts University, Medford MA, United States of America

Centre de Investigaciones, Universidad Antonio Narino, Bogota, Colombia

Department of Physics and Astronomy, University of California Irvine, Irvine CA, United States of America

(a) INFN Gruppo Collegato di Udine; (b) ICTP, Trieste; (c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy

Department of Physics, University of Illinois, Urbana IL, United States of America

Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNMT), University of Valencia and CSIC, Valencia, Spain

Department of Physics, University of British Columbia, Vancouver BC, Canada

Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada

Waseda University, Tokyo, Japan

Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

Department of Physics, University of Wisconsin, Madison WI, United States of America

Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany

Department of Physics, Yale University, New Haven CT, United States of America

Yerevan Physics Institute, Yerevan, Armenia

Domaine scientifique de la Doua, Centre de Calcul CNRS/IN2P3, Villeurbanne Cedex, France

Faculty of Science, Hiroshima University, Hiroshima, Japan

Also at Laboratorio de Instrumentacao e Física Experimental de Partículas - LIP, Lisboa, Portugal

Also at Faculdade de Ciencias and CFNUL, Universidade de Lisboa, Lisboa, Portugal

Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom

Also at TRIUMF, Vancouver BC, Canada

Also at Department of Physics, California State University, Fresno CA, United States of America

Also at Novosibirsk State University, Novosibirsk, Russia

Also at Fermilab, Batavia IL, United States of America

Also at Department of Physics, University of Coimbra, Coimbra, Portugal

Also at Università di Napoli Parthenope, Napoli, Italy

Also at Institute of Particle Physics (IPP), Canada

Also at Department of Physics, Middle East Technical University, Ankara, Turkey

Also at Louisiana Tech University, Ruston LA, United States of America

Also at Department of Physics and Astronomy, University College London, London, United Kingdom

Also at Group of Particle Physics, University of Montreal, Montreal QC, Canada

Also at Department of Physics, University of Cape Town, Cape Town, South Africa

Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany

Also at Manhattan College, New York NY, United States of America

Also at School of Physics, Shandong University, Shandong, China

Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
\textsuperscript{u} Also at School of Physics and Engineering, Sun Yat-sen University, Guanzhou, China
\textsuperscript{v} Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan
\textsuperscript{w} Also at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat a l’Energie Atomique), Gif-sur-Yvette, France
\textsuperscript{x} Also at Section de Physique, Université de Genève, Geneva, Switzerland
\textsuperscript{y} Also at Departamento de Física, Universidade de Minho, Braga, Portugal
\textsuperscript{z} Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, United States of America
\textsuperscript{aa} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary
\textsuperscript{ab} Also at California Institute of Technology, Pasadena CA, United States of America
\textsuperscript{ac} Also at Institute of Physics, Jagiellonian University, Krakow, Poland
\textsuperscript{ad} Also at LAL, Univ. Paris-Sud and CNRS/IN2P3, Orsay, France
\textsuperscript{ae} Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
\textsuperscript{af} Also at Department of Physics, Oxford University, Oxford, United Kingdom
\textsuperscript{ag} Also at Institute of Physics, Academia Sinica, Taipei, Taiwan
\textsuperscript{ah} Also at Department of Physics, The University of Michigan, Ann Arbor MI, United States of America
\textsuperscript{ai} Also at Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
\textsuperscript{*} Deceased