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Measurement of long-range multiparticle azimuthal correlations with the subevent cumulant method in pp and p+Pb collisions with the ATLAS detector at the CERN Large Hadron Collider

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A detailed study of multiparticle azimuthal correlations is presented using pp data at √s = 5.02 and 13 TeV, and p+Pb data at √sNN = 5.02 TeV, recorded with the ATLAS detector at the CERN Large Hadron Collider. The azimuthal correlations are probed using four-particle cumulants c_n[4] and flow coefficients v_n[4] = (−c_n[4])^{n/4} for n = 2 and 3, with the goal of extracting long-range multiparticle azimuthal correlation signals and suppressing the short-range correlations. The values of c_n[4] are obtained as a function of the average number of charged particles per event, ⟨N_ch⟩, using the recently proposed two-subevent and three-subevent cumulant methods, and compared with results obtained with the standard cumulant method. The standard method is found to be strongly biased by short-range correlations, which originate mostly from jets with a positive contribution to c_2[4]. The three-subevent method, on the other hand, is found to be least sensitive to short-range correlations. The three-subevent method gives a negative c_2[4], and therefore a well-defined v_2[4], nearly independent of ⟨N_ch⟩, which implies that the long-range multiparticle azimuthal correlations persist to events with low multiplicity. Furthermore, v_2[4] is found to be smaller than the v_2[2] measured using the two-particle correlation method, as expected for long-range collective behavior. Finally, the measured values of v_2[4] and v_2[2] are used to estimate the number of sources relevant for the initial eccentricity in the collision geometry. The results based on the subevent cumulant technique provide direct evidence, in small collision systems, for a long-range collectivity involving many particles distributed across a broad rapidity interval.

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I. INTRODUCTION

The study of azimuthal correlations in high-energy nuclear collisions at the Relativistic Heavy Ion Collider (RHIC) and the Large Hadron Collider (LHC) has been important for understanding the multiparton dynamics of QCD in the strongly coupled nonperturbative regime. One striking observation is the long-range ridge [1–5] in two-particle angular correlations (2PC): an apparent collimated emission of particle pairs with small relative azimuthal angle (∆φ) and large separation in pseudorapidity (∆η). The ridge signature from 2PC is characterized by a Fourier decomposition of the correlation function C(∆φ) ∼ 1 + 2 ∑_n v_n^2 cos(n ∆φ), where v_n denotes the single-particle anisotropy harmonic coefficients. The second-order coefficient v_2 is observed to be the largest, followed by v_3 [3,4]. These coefficients carry information about the collective behavior of the produced system. The ridge was first discovered in nucleus-nucleus (A+A) collisions [1–6], but was later observed in small systems such as proton-nucleus (p+A) collisions [7–11], light-ion–nucleus collisions [12], and more recently in proton-proton (pp) collisions [13–16]. The ridge in large systems, such as central or midcentral A+A collisions, is commonly interpreted as the result of collective hydrodynamic expansion of hot and dense nuclear matter created in the overlap region of the colliding nuclei. Since the formation of an extended region of nuclear matter is not expected in small collision systems such as p+A and pp, the origin of the ridge there could be different from that formed in large collision systems. There remains considerable debate in the theoretical community as to whether the ridge in small systems is of hydrodynamic origin, like it is in A+A collisions [17], or stems from other effects such as initial-state gluon saturation [18].

An important question about the ridge is whether it involves all particles in the event (collective flow) or if it arises merely from correlations among a few particles, due to resonance decays, jets, or multijet production (nonflow). In small systems the contributions from nonflow sources, in particular from jets and dijets, are large. The extraction of a ridge signal using the 2PC method requires a large ∆η gap and careful removal of the significant contribution from dijet production [8–10,14,15,19]. Since collective flow is intrinsically a multiparticle phenomenon, it can be probed more directly using cumulants based on multiparticle correlation techniques [20]. Azimuthal correlations involving four, six, and eight particles have been measured in p+Pb, d+Au, and pp collisions, and a significant v_2 signal has been obtained [11,19,21,22]. One weakness of the standard multiparticle cumulant method is that it does not suppress adequately the nonflow correlations.

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in small systems, which lead to a sign change of \( c_2 \{4 \} \) at smaller values of the charged particle multiplicity, \( N_{\text{ch}} \). Furthermore, the magnitude of \( c_2 \{4 \} \) and the \( N_{\text{ch}} \) value at which the sign change occurs are found to depend sensitively on the exact definition of \( N_{\text{ch}} \) used to categorize the events. These observations suggest that the standard cumulant method, on which several previous measurements in small systems are based, is strongly contaminated by nonflow correlations \([11,19,21,22]\), especially in \( pp \) collisions and low \( N_{\text{ch}} \) region.

Recently an improved cumulant method based on the correlation between particles from different subevents separated in \( \eta \) has been proposed to further reduce the nonflow correlations \([23]\). The effectiveness of this method for suppressing nonflow correlations has been validated using the PYTHIA8 event generator \([24]\), which contains only nonflow correlations.

This paper presents measurements of \( c_2 \{4 \} \) and \( c_3 \{4 \} \) in \( pp \) collisions at \( \sqrt{s} = 5.02 \) and 13 TeV, as well as \( p+\text{Pb} \) collisions at \( \sqrt{s_{\text{NN}}} = 5.02 \) TeV. They are obtained using two- and three-subevent cumulant methods and are compared with the standard cumulant method. The \( c_2 \{4 \} \) cumulant is converted to the corresponding \( v_2 \) coefficient and compared with the results obtained using the two-particle correlation method in Refs. \([10,15]\) to assess the nature of the event-by-event fluctuations of the collective flow in these collisions.

The paper is organized as follows. Section II describes the framework for the standard, two-subevent and three-subevent four-particle cumulant methods used in this analysis. Details of the detector, trigger, data sets, as well as event and track selections are provided in Secs. III–V. The correlation analysis and systematic uncertainties are described in Secs. VI and VII, respectively. The measured cumulants from the three data sets are provided in Sec. VIII. A summary is given in Sec. IX.

## II. FOUR-PARTICLE CUMULANTS

The multiparticle cumulant method \([20]\) is used to extract the amplitude of long-range azimuthal correlations of particles produced in high-energy collisions. This method has the advantage of suppressing correlations from jets and dijets, instead of relying on an explicit procedure to correct \( v_n \) harmonics for dijet contributions in the 2PC approach, as done in Refs. \([10,14]\). The framework for the standard cumulant is described in Refs. \([25,26]\), which was recently extended to the case of subevent cumulants in Ref. \([23]\). This paper presents measurements of four-particle cumulants obtained with the standard, two-subevent, and three-subevent methods. The following discussion first describes the standard cumulant method, then describes the two- and three-subevent methods focusing on the differences from the standard method.

The cumulant methods involve the calculation of \( 2k \)-particle azimuthal correlations \( \langle \{2k\}_n \rangle \), and \( 2k \)-particle cumulants, \( c_n \{2k\} \), for the \( n \)-th order flow harmonics. The two- or four-particle azimuthal correlations in one event are evaluated as \([23,25,26]\):

\[
\langle \{2\}_n \rangle = \langle e^{i(\phi_2 - \phi_1)} \rangle = \frac{q_n^2 - \tau_1}{1 - \tau_1},
\]

(1)

\[
\langle \{4\}_n \rangle = \langle e^{i(\phi_2 + \phi_3 - \phi_1 - \phi_2)} \rangle = \frac{q_n^4 - 2\tau_1(\text{Re}[q_{n2}\sqrt{q_n^2} + 2q_n^2] + 2q_n^2) + 8\tau_2\text{Re}[q_{n3}q_n^* + q_n^4 + (2 + q_{n2}) - 6\tau_1]}{1 - 6\tau_1 + 8\tau_2 + 3\tau_1^2 - 6\tau_1},
\]

(2)

where \( \langle \cdot \rangle \) denotes a single-event average over all pairs or quadruplets, respectively. The averages from Eqs. (1) and (2) are expanded into per-particle normalized flow vectors \( q_{n,\tau} \) and factors \( \tau_\ell \) with \( \ell = 1, 2, \ldots \):

\[
q_{n,\tau} = \sum_j w^j \left( e^{i\phi_j} \right), \quad q_{n,\tau} = |q_{n,\tau}|, \quad q_n = q_{n,1},
\]

\[
\tau_\ell = \frac{\sum_j w^{j+1}}{\sum_j w_j^{j+1}},
\]

(3)

where the sum runs over all \( M \) particles in the event and \( w_j \) is a weight assigned to the \( j \)-th particle. This weight is constructed to correct for both detector nonuniformity and tracking inefficiency as explained in Sec. VI. For unit weight \( w_j = 1 \), then \( q_{\text{univ}} = q_{\text{norm}} \), and \( \tau_1 = 1/M^2 \).

The two- and four-particle cumulants are obtained from the azimuthal correlations as:

\[
c_2 \{2\} = \langle \{2\}_n \rangle,
\]

(4)

\[
c_4 \{4\} = \langle \{4\}_n \rangle - 2\langle \{2\}_n \rangle^2,
\]

(5)

where \( \langle \cdot \rangle \) represents a weighted average of \( \langle \{2k\}_n \rangle \) over an event ensemble. In the absence of nonflow correlations, \( c_n \{2k\} \) of the detector, trigger, data sets, as well as event and track selections are provided in Secs. III–V. The correlation analysis and systematic uncertainties are described in Secs. VI and VII, respectively. The measured cumulants from the three data sets are provided in Sec. VIII. A summary is given in Sec. IX.
cording to \(-\eta_{\text{max}} < \eta_a < 0 < \eta_b < \eta_{\text{max}}\), where \(\eta_{\text{max}} = 2.5\) is the maximum \(\eta\) used in the analysis and corresponds to the ATLAS detector acceptance for charged particles. The per-event two- and four-particle azimuthal correlations are then evaluated as:

\[
\langle \{2\}_n \rangle_{a|b} = \langle e^{i(\phi_a^c - \phi_b^c)} \rangle = \Re \{g_{a,b}^c g_{n,h,b}^* \},
\]

and

\[
\langle \{4\}_n \rangle_{2a|2b} = \langle e^{i(\phi_a^c + \phi_b^c - \phi_a^c - \phi_b^c)} \rangle = \frac{\langle g_n^2 - \tau_1 q_{2n}^b \rangle}{\langle g_n^2 - \tau_1 q_{2n}^b \rangle} \frac{\langle g_n^2 - \tau_1 q_{2n}^b \rangle}{(1 - \tau_1)_a(1 - \tau_1)_b},
\]

where the superscript or subscript \(a\) (\(b\)) indicates particles chosen from the subevent \(a\) (\(b\)). Here the four-particle cumulant is defined as:

\[
c_n^{4_{2a|2b}} = \langle \langle \{4\}_n \rangle_{2a|2b} \rangle - 2 \langle \langle \{2\}_n \rangle_{a|b} \rangle.
\]

The two-subevent method should suppress correlations within a single jet (intraget correlations), since each jet usually emits particles into only one subevent.

In the three-subevent cumulant method, the event is divided into three subevents \(a\), \(b\), and \(c\) each covering a unique \(\eta\) range, for example \(-\eta_{\text{max}} < \eta_a < -\eta_{\text{max}}/3\), \(\eta_{\text{min}} < \eta_b < \eta_{\text{max}}/3\), and \(\eta_{\text{max}}/3 < \eta_c < \eta_{\text{max}}\). The four-particle azimuthal correlations and cumulants are then evaluated as:

\[
\langle \langle \{4\}_n \rangle_{2a|b,c} \rangle = \langle e^{i(\phi_a^c + \phi_b^c - \phi_a^c - \phi_b^c)} \rangle = \frac{\langle g_n^2 - \tau_1 q_{2n}^b \rangle}{\langle g_n^2 - \tau_1 q_{2n}^b \rangle} \frac{\langle g_n^2 - \tau_1 q_{2n}^b \rangle}{(1 - \tau_1)_a(1 - \tau_1)_b},
\]

and

\[
c_n^{4_{2a|b,c}} = \langle \langle \{4\}_n \rangle_{2a|b,c} \rangle - 2 \langle \langle \{2\}_n \rangle_{a|b} \rangle \langle \langle \{2\}_n \rangle_{a|c} \rangle,
\]

where \(\langle \langle \{2\}_n \rangle_{a|b} \rangle\) and \(\langle \langle \{2\}_n \rangle_{a|c} \rangle\) are two-particle correlators defined as in Eq. (8). Since the two jets in a dijet event usually produce particles in at most two subevents, the three-subevent method further suppresses nonflow contributions from interjet correlations associated with dijets. To enhance the statistical precision, the \(\eta\) range for subevent \(a\) is also interchanged with that for subevent \(b\) or \(c\), and the resulting three \(c_n^{4_{2a|b,c}}\) values are averaged to obtain the final result.

III. DETECTOR AND TRIGGER

The ATLAS detector [27] provides nearly full solid-angle coverage around the collision point with tracking detectors, calorimeters, and muon chambers, and is well suited for measurement of multiparticle correlations over a large pseudorapidity range.1 The measurements were performed primarily using the inner detector (ID), minimum-bias trigger scintillators (MBTS), and the zero-degree calorimeters (ZDCs). The ID detects charged particles within \(|\eta| < 2.5\) using a combination of silicon pixel detector, a silicon microstrip detector (SCT), and a straw-tube transition radiation tracker, all immersed in a 2 T axial magnetic field [28]. An additional pixel layer, the insertable B-layer (IBL) [29] installed between Run 1 (2010–2013) and Run 2 (2015–2018), is available for the Run-2 data sets. The MBTS, rebuilt before Run 2, detects charged particles within \(2.1 \leq |\eta| \leq 3.9\) using two hodoscopes of counters positioned at \(z = \pm 3.6\) m. The ZDCs are positioned at \(\pm 140\) m from the collision point, and detect neutral particles, primarily neutrons and photons, with \(|\eta| > 8.3\).

The ATLAS trigger system [30] consists of a Level-1 (L1) trigger implemented using a combination of dedicated electronics and programmable logic, and a high-level trigger (HLT) implemented in processors. The HLT reconstructs charged-particle tracks using methods similar to those applied in the offline analysis, allowing high-multiplicity track (HMT) triggers that select events based on the number of tracks with \(p_T > 0.4\) GeV associated with the vertex with the largest number of tracks. The different HMT triggers also apply additional requirements on either the transverse energy \((E_T^\perp)\) in the calorimeters or on the number of hits in the MBTS at L1, and on the number of charged-particle tracks reconstructed by the HLT. The pp and p+Pb data were collected using a combination of the minimum-bias and HMT triggers. More details of the triggers used for the pp and p+Pb data can be found in Refs. [15,31] and Refs. [10,32], respectively.

IV. DATA SETS AND MONTE CARLO SIMULATIONS

This analysis uses integrated luminosities of 28 nb\(^{-1}\) of p+Pb data recorded at \(\sqrt{s_{NN}} = 5.02\) TeV, 0.17 pb\(^{-1}\) of pp data recorded at \(\sqrt{s} = 5.02\) TeV, and 0.9 pb\(^{-1}\) of pp data recorded at \(\sqrt{s} = 13\) TeV, all taken by the ATLAS experiment at the LHC. The p+Pb data were mainly collected in 2013, but also include 0.3 nb\(^{-1}\) data collected in November 2016, which increases the number of events at moderate multiplicity (see Sec. V). During both p+Pb runs, the LHC was configured with a 4 TeV proton beam and a 1.57 TeV per-nucleon Pb beam that together produced collisions at \(\sqrt{s_{NN}} = 5.02\) TeV, with a rapidity shift of 0.465 of the nucleon–nucleon center-of-mass frame towards the proton beam direction relative to the ATLAS rest frame. The direction of the Pb beam is always defined to have negative pseudorapidity. The 5.02 TeV pp data were collected in November 2015. The 13 TeV pp data were collected during several special low-luminosity runs of the LHC in 2015 and 2016.

Monte Carlo (MC) simulated event samples are used to determine the track reconstruction efficiency (Sec. V). The 13 TeV and 5.02 TeV pp data were simulated by the PYTHIA8 MC event generator [24] using the A2 set of tuned parameters with MSTW2008LO parton distribution functions [33]. The HIJING event generator [34] was used to produce p+Pb collisions with the same energy and the same boost of the center-of-mass system as in the data. The detector response was simulated using GEANT4 [35,36] with detector conditions matching those during the data taking. The simulated events
samples reconstructed with the same tracking algorithms and the same track selection requirements. Efficiencies, $\epsilon(\eta, p_T)$, are evaluated as a function of track $\eta$, $p_T$ and the number of reconstructed charged-particle tracks, but averaged over the full range in azimuth. For all collision systems, the efficiency increases by about $4\%$ as $p_T$ increases from 0.3 GeV to 0.6 GeV. Above 0.6 GeV, the efficiency is independent of $p_T$ and reaches $86\%$ (72\%) at $\eta \approx 0$ ($|\eta| > 2$) for $pp$ collisions and $83\%$ (70\%) for $p+Pb$ collisions, respectively. The efficiency is independent of the event multiplicity for $N_{ch}^{rec} > 40$. For lower-multiplicity events the efficiency is smaller by up to a few percent due to broader $d_0^{BL}$ and $z_0$ sin $\theta$ distributions.

The rate of falsely reconstructed charged-particle tracks is also estimated and found to be negligibly small in all data sets. This rate decreases with increasing $p_T$, and even at the lowest transverse momenta of 0.2 GeV it is below 1\% of the total number of tracks. Therefore, there is no correction for the presence of these tracks in the analysis.

In the simulated events, the reconstruction efficiency reduces the measured charged-particle multiplicity relative to the generated multiplicity for primary charged particles. The multiplicity correction factor $b$ is used to correct $N_{ch}^{rec}$ to obtain the efficiency-corrected number of charged particles per event, $\langle N_{ch} \rangle = b(N_{ch}^{rec})$. The value of the correction factor is found to be independent of $N_{ch}^{rec}$ in the range used in this analysis. Its value and the associated uncertainties are $b = 1.29 \pm 0.05$ for the 2013 $p+Pb$ collisions and $b = 1.18 \pm 0.05$ for Run-2 $p+Pb$ and $pp$ collisions [37]. Both $c_n(4)$ and $v_n(4)$ are then studied as a function of $\langle N_{ch} \rangle$.

VI. DATA ANALYSIS

The multiparticle cumulants are calculated in three steps using charged particles with $|\eta| < 2.5$. In the first step, the multiparticle correlators $\langle k_{n\mu} \rangle$ from Eqs. (1), (2), (8), (9), and (11) are calculated for each event from particles in one of two $p_T$ ranges, $0.3 < p_T < 3$ GeV and $0.5 < p_T < 5$ GeV.

In the second step, the correlators $\langle k_{2\mu} \rangle$ are averaged over events with the same $N_{ch}^{rec}$, the number of reconstructed charged particles in a given $p_T$ range, to obtain $\langle \langle k_{2\mu} \rangle \rangle$ and $c_n(2k)$ from Eqs. (4), (10), and (12). In a previous study [16], it was observed that the $c_n(2k)$ values varied with the exact definition of $N_{ch}^{rec}$. This is because different definitions of $N_{ch}^{rec}$ lead to different multiplicity fluctuations and therefore different nonflow correlations associated with these multiplicity fluctuations. The observed dependence of $c_n(2k)$ on the definition of $N_{ch}^{rec}$ has been attributed to the change in the nonflow correlations when $N_{ch}^{rec}$ is changed [16].

In order to further test the sensitivity of $c_n(2k)$ to the exact definition of $N_{ch}^{rec}$, four different $p_T$ requirements are used to define $N_{ch}^{rec}$ as follows: when $\langle \langle k_{2\mu} \rangle \rangle$ is calculated in the range $0.3 < p_T < 3$ GeV, $N_{ch}^{rec}$ is evaluated in four different track $p_T$ ranges: $0.3 < p_T < 3$ GeV, $p_T > 0.2$ GeV, $p_T > 0.4$ GeV, and $p_T > 0.6$ GeV. When $\langle \langle k_{2\mu} \rangle \rangle$ is calculated in $0.5 < p_T < 5$ GeV, $N_{ch}^{rec}$ is evaluated in four different track $p_T$ ranges: $0.5 < p_T < 5$ GeV, $p_T > 0.2$ GeV, $p_T > 0.4$ GeV, and $p_T > 0.6$ GeV. In each case, the $c_n(2k)$ value is first calculated for events with the same $N_{ch}^{rec}$, the $c_n(2k)$ values are...
then combined in the broader \(^N_{\text{ch}}\) range of the event ensemble to obtain statistically significant results.

In the third step, the \(c_n(2k)\) and \(v_n(2k)\) values obtained for a given \(N_{\text{ch}}\) are mapped to a given \(\langle N_{\text{ch}}\rangle\), the average number of reconstructed charged particles with \(p_T > 0.4\) GeV. The mapping procedure is necessary so that \(c_n(2k)\) obtained for different \(N_{\text{ch}}\) can be compared using a common \(x\) axis defined by \(\langle N_{\text{ch}}\rangle\). The \(\langle N_{\text{ch}}\rangle\) value is then converted to \(\langle \Delta N_{\text{ch}} \rangle\), the efficiency-corrected average number of charged particles with \(p_T > 0.4\) GeV, as discussed in Sec. V.

In order to account for detector inefficiencies and nonuniformity, particle weights used in Eq. (3) are defined as:

\[
    w_i(\phi, \eta, p_T) = d(\phi, \eta)/e(\eta, p_T).
\]  

(13)

The additional weight factor \(d(\phi, \eta)\) accounts for nonuniformities in the azimuthal acceptance of the detector as a function of \(\eta\). All reconstructed charged particles with \(p_T > 0.2\) GeV are entered into a two-dimensional histogram \(N(\phi, \eta)\), and the weight factor is then obtained as \(d(\phi, \eta) = \langle N(\eta) \rangle / N(\phi, \eta)\), where \(\langle N(\eta) \rangle\) is the track density averaged over \(\phi\) in the given \(\eta\) bin. This procedure removes most \(\phi\)-dependent nonuniformity from track reconstruction for any azimuthal correlation analysis [16].

VII. SYSTEMATIC UNCERTAINTIES

The main sources of systematic uncertainty are related to the detector azimuthal nonuniformity, track selection, track reconstruction efficiency, trigger efficiency, and pileup. Most of the systematic uncertainties enter the analysis through the reconstruction efficiency, trigger efficiency, and pileup. Most events used in the analysis are collected with the HMT triggers with several \(N_{\text{ch}}\) thresholds. In order to estimate the possible bias due to trigger inefficiency as a function of \(\langle N_{\text{ch}}\rangle\), the offline \(N_{\text{ch}}\) requirements are changed such repeating the analysis. The results are mostly consistent with the nominal results within statistical uncertainties. As a cross check, the multiparticle correlations are calculated using a mixed-event procedure, where each particle in a 2k multiplet is selected from a different event with similar \(N_{\text{ch}}\) \((|\Delta N_{\text{ch}}| < 10)\) and similar \(z_{\text{vtx}}\) \((|\Delta z_{\text{vtx}}| < 10 \text{ mm})\). The particle weights defined in Eq. (13) are applied for each particle forming the mixed event. The \(c_2(4)\) signal obtained from the mixed events is less than \(0.2 \times 10^{-6}\) in all data sets.

The systematic uncertainty associated with the track selection is estimated by tightening the \(|d_0|\) and \(|z_0 \sin \theta|\) requirements. For each variation, the tracking efficiency is reevaluated and the analysis is repeated. The maximum differences from the nominal results are observed to be less than \(0.3 \times 10^{-6}\), \(0.2 \times 10^{-6}\), and \(0.1 \times 10^{-6}\) in 5.02 TeV \(pp\), 13 TeV \(pp\), and \(p+Pb\) collisions, respectively.

Previous measurements indicate that the azimuthal correlations (both the flow and nonflow components) have a strong dependence on \(p_T\), but a relatively weak dependence on \(\eta\) [10,15]. Therefore, \(p_T\)-dependent systematic effects in the track reconstruction efficiency could affect \(c_n(2k)\) and \(v_n(2k)\) values. The uncertainty in the track reconstruction efficiency is mainly due to differences in the detector conditions and material description between the simulation and the data. The efficiency uncertainty varies between 1% and 4%, depending on track \(\eta\) and \(p_T\) [15,16]. Its impact on multiparticle cumulants is evaluated by repeating the analysis with the tracking efficiency varied up and down by its corresponding uncertainty as a function of \(p_T\). For the standard cumulant method, which is more sensitive to jets and dijets, the evaluated uncertainty amounts to \((0.1-1.5) \times 10^{-6}\) in \(pp\) collisions and less than \(0.3 \times 10^{-6}\) in \(p+Pb\) collisions for \(\langle N_{\text{ch}}\rangle > 50\). For the two- and three-subevent methods, the evaluated uncertainty is typically less than \(0.3 \times 10^{-6}\) for most of the \(\langle N_{\text{ch}}\rangle\) ranges.

Most events used in the analysis are collected with the HMT triggers with several \(N_{\text{ch}}\) selections as indicated in the figure, which is then mapped to \(\langle N_{\text{ch}}\rangle\), the average number of charged particles with \(p_T > 0.4\) GeV. The error bars and shaded boxes represent the statistical and systematic uncertainties, respectively.

**FIG. 1.** The \(c_2(4)\) values calculated for charged particles with \(0.3 < p_T < 3\) GeV (left) and \(0.5 < p_T < 5\) GeV (right) with the standard cumulant method from the 13 TeV \(pp\) data. The event averaging is performed for \(N_{\text{ch}}\) calculated for various \(p_T\) selections as indicated in the figure, which is then mapped to \(\langle N_{\text{ch}}\rangle\), the average number of charged particles with \(p_T > 0.4\) GeV. The error bars and shaded boxes represent the statistical and systematic uncertainties, respectively.
that the HMT trigger efficiency is at least 50% or 80%. The results are obtained independently for each variation. These results are found to be consistent with each other for the two- and three-subevent methods, and show a small difference for the standard cumulant method in the low \( \langle N_{\text{ch}} \rangle \) region. The nominal analysis is performed using the 50% efficiency selection and the differences between the nominal results and those from the 80% efficiency selection are used as a systematic uncertainty. The change amounts to \((0.1–0.7) \times 10^{-6}\).

In this analysis, a pileup rejection criterion is applied to reject events containing additional vertices. In order to check the impact of residual pileup, the analysis is repeated without the pileup rejection criterion, and no difference is observed. For the 5.02 and 13 TeV \( pp \) data sets, which have relatively high pileup, the data is divided into two samples based on the \( \mu \) value: \( \mu > 0.4 \) and \( \mu < 0.4 \), and the results are compared. The average \( \mu \) values differ by a factor of two between the two samples, and the difference in \( c_2[4] \) is found to be less than \( 0.5 \times 10^{-6} \).

To check the impact of dijet events, where both jets have pseudorapidities close to the boundaries of relevant subevent regions, the three-subevent cumulants are calculated by requiring a \( \Delta \eta = 0.5 \) gap between the adjacent regions. The results are found to be consistent with the nominal result.

The systematic uncertainties from different sources are added in quadrature to determine the total systematic uncertainty. The uncertainty is \((0.1–1) \times 10^{-6}\) for two- and three-subevent methods in the region \( \langle N_{\text{ch}} \rangle > 50 \), where there is a negative \( c_2[4] \) signal. The total systematic uncertainty for the standard method is typically about a factor of two larger.

The systematic uncertainty studies described above are also carried out for \( c_3[4] \), and the absolute uncertainties are found to be smaller than those for \( c_2[4] \), presumably because \( c_3[4] \) is less sensitive to the influence from dijets.

![FIG. 2. The \( c_2[4] \) values calculated for charged particles with \( 0.3 < p_T < 3 \) GeV (left) and \( 0.5 < p_T < 5 \) GeV (right) with the two-subevent cumulant method from the 13 TeV \( pp \) data. The event averaging is performed for \( N_{\text{ch}}^\text{sel} \) calculated for various \( p_T \) selections as indicated in the figure, which is then mapped to \( \langle N_{\text{ch}} \rangle \), the average number of charged particles with \( p_T > 0.4 \) GeV. The error bars and shaded boxes represent the statistical and systematic uncertainties, respectively.](image1)

![FIG. 3. The \( c_2[4] \) values calculated for charged particles with \( 0.3 < p_T < 3 \) GeV (left) and \( 0.5 < p_T < 5 \) GeV (right) with the three-subevent cumulant method from the 13 TeV \( pp \) data. The event averaging is performed for \( N_{\text{ch}}^\text{sel} \) calculated for various \( p_T \) selections as indicated in the figure, which is then mapped to \( \langle N_{\text{ch}} \rangle \), the average number of charged particles with \( p_T > 0.4 \) GeV. The error bars and shaded boxes represent the statistical and systematic uncertainties, respectively.](image2)
FIG. 4. The $c_2[4]$ values calculated for charged particles with $0.3 < p_T < 3$ GeV (left) and $0.5 < p_T < 5$ GeV (right) compared for the three cumulant methods from the 13 TeV $pp$ data. The event averaging is performed for $N_{ch}^{sel}$ calculated for the same $p_T$ range, which is then mapped to $\langle N_{ch} \rangle$, the average number of charged particles with $p_T > 0.4$ GeV. The dashed line indicates the $c_2[4]$ value corresponding to a 4% $v_2$ signal. The error bars and shaded boxes represent the statistical and systematic uncertainties, respectively.

VIII. RESULTS

A. Dependence on the event-class definition

This section presents the sensitivity of $c_2[4]$ to $N_{ch}^{sel}$, which defines the event class used to calculate $\langle \{2\rangle \rangle_s$ and $\langle \{4\rangle \rangle_s$ in Eqs. (10)–(12). The discussion is based on results obtained from the 13 TeV $pp$ data, but the observations for the 5.02 TeV $pp$ and $p+Pb$ data are qualitatively similar.

Figure 1 shows the $c_2[4]$ values obtained using the standard method for four event-class definitions based on $N_{ch}^{sel}$. The $c_2[4]$ values change dramatically as the event-class definition is varied, which, as points out in Ref. [23], reflects different amount of nonflow fluctuations associated with different $N_{ch}^{sel}$. The $c_2[4]$ values for $0.3 < p_T < 3$ GeV become negative when the reference $N_{ch}^{sel}$ is obtained for $p_T > 0.4$ GeV or higher, but the four cases do not converge to the same $c_2[4]$ values. On the other hand, $c_2[4]$ values for $0.5 < p_T < 5$ GeV are always positive, independent of the definition of $N_{ch}^{sel}$. These behaviors suggest that the $c_2[4]$ values from the standard method are strongly influenced by nonflow effects in all $\langle N_{ch} \rangle$ and $p_T$ ranges. Therefore the previously observed negative $c_2[4]$ in $pp$ collisions for $0.3 < p_T < 3$ GeV and $N_{ch}^{sel}$ with $p_T > 0.4$ GeV [19] may be dominated by nonflow correlations instead of long-range collective flow.

Figure 2 shows that the $c_2[4]$ values calculated using the two-subevent method are closer to each other among different event-class definitions. The $c_2[4]$ values decrease gradually with $\langle N_{ch} \rangle$ and become negative for $\langle N_{ch} \rangle > 70$ when $c_2[4]$ is calculated in the range $0.3 < p_T < 3$ GeV range and for $\langle N_{ch} \rangle > 150$ when $c_2[4]$ is calculated in the range $0.5 < p_T < 5$ GeV. Therefore, the $c_2[4]$ values from the two-subevent method are more sensitive to long-range ridge correlations, but nevertheless may still be affected by nonflow effects, especially in the low $\langle N_{ch} \rangle$ region and higher $p_T$.

Figure 3 shows the results from the three-subevent method. For most of the $\langle N_{ch} \rangle$ range, the $c_2[4]$ values are negative, i.e., having the sign expected for long-range ridge correlations.

FIG. 5. The $c_2[4]$ values calculated for charged particles with $0.3 < p_T < 3$ GeV (left) and $0.5 < p_T < 5$ GeV (right) compared for the three cumulant methods from the 5.02 TeV $pp$ data. The event averaging is performed for $N_{ch}^{sel}$ calculated for the same $p_T$ range, which is then mapped to $\langle N_{ch} \rangle$, the average number of charged particles with $p_T > 0.4$ GeV. The dashed line indicates the $c_2[4]$ value corresponding to a 4% $v_2$ signal. The error bars and shaded boxes represent the statistical and systematic uncertainties, respectively.
The $c_2[4]$ values show some sensitivity to the definition of the reference $N_{ch}^{sel}$ but they are close to each other for all definitions in the region $\langle N_{ch} \rangle > 100$. This suggests that the residual nonflow effects may still be important at small $\langle N_{ch} \rangle$, but are negligible at $\langle N_{ch} \rangle > 100$. It is also observed that the $c_2[4]$ values for $0.5 < p_T < 5$ GeV are more negative than those for $0.3 < p_T < 3$ GeV, which is consistent with the observation that the $v_2$ value associated with the long-range collectivity increases with $p_T$ [10,15].

Given the relatively small dependence of $c_2[4]$ on the reference $N_{ch}^{sel}$ in the three-subevent method, the remaining discussion focuses on cases where the reference $N_{ch}^{sel}$ is calculated in the same $p_T$ ranges as those used for calculating $c_2[4]$, i.e., $0.3 < p_T < 3$ GeV and $0.5 < p_T < 5$ GeV.

B. Comparison between different cumulant methods

Figures 4–6 show direct comparisons of the results for the standard, two-subevent, and three-subevent methods for $pp$ collisions at $\sqrt{s} = 13$ TeV, $pp$ at $\sqrt{s} = 5.02$ TeV, and $p+Pb$ collisions at $\sqrt{s_{NN}} = 5.02$ TeV, respectively. The results from $5.02$ TeV $pp$ collisions are qualitatively similar to those from the $13$ TeV $pp$ collisions, i.e., the $c_2[4]$ values are smallest for the three-subevent method and largest for the standard method. The same hierarchy between the three methods is also observed in $p+Pb$ collisions, but only for the $\langle N_{ch} \rangle < 100$ region, suggesting that nonflow effects in $p+Pb$ collisions are much smaller than those in $pp$ collisions at comparable $\langle N_{ch} \rangle$. In $p+Pb$ collisions, all three methods give consistent results for $\langle N_{ch} \rangle > 100$. Furthermore, the three-subevent method
gives negative $c_2 \{4 \}$ values in most of the measured $\langle N_{\text{ch}} \rangle$ range.

The comparison of the $c_2 \{4 \}$ values between the three data sets, for the standard and the three-subevent methods, is shown in Figs. 7 and 8. The large positive $c_2 \{4 \}$ values observed in the small $\langle N_{\text{ch}} \rangle$ region in the standard method are likely due to nonflow correlations, since this trend is absent when using the three-subevent cumulant method. In $p+Pb$ collisions, the absolute value of $c_2 \{4 \}$ seems to become smaller for $\langle N_{\text{ch}} \rangle > 200$.

The same analysis is performed for the third-order harmonics. Figures 9 and 10 compare the $c_3 \{4 \}$ values between the three data sets for the standard cumulant method and the three-subevent method. The $c_3 \{4 \}$ values from the three-subevent method are close to zero in all three systems. For the standard method, the positive $c_3 \{4 \}$ values in the small $\langle N_{\text{ch}} \rangle$ region indicate the influence of nonflow correlations, but the influence is not as strong as that for $c_2 \{4 \}$.

Figure 11 shows the $c_3 \{4 \}$ values from $p+Pb$ collisions in the two $p_T$ ranges, obtained with the three-subevent method;
they are zoomed-in version of the \( p+Pb \) data shown in Figs. 8 and 9. Within their large statistical and systematic uncertainties, the values of \( c_1(4) \) are systematically below zero, especially for \( 0.5 < p_T < 5 \text{ GeV} \), where the \( c_1(4) \) values are comparable to \(-0.16 \times 10^{-6}\), corresponding to a \( v_3 \) value of 2]\% as indicated in the figure. The negative \( c_1(4) \) values from the three-subevent method support the existence of long-range multiparticle triangular flow in \( p+Pb \) collisions.

C. Three-subevent flow harmonic \( v_2(4) \)

The harmonic flow coefficients \( v_2(4) \) can be obtained from the measured values of \( c_2(4) \) according to Eq. (7). Figure 12 shows the \( v_2(4) \) values for charged particles with \( 0.3 < p_T < 3 \text{ GeV} \) calculated using the three-subevent method in the three data sets. Results for the higher \( p_T \) range \((0.5 < p_T < 5 \text{ GeV})\) are presented in Fig. 13. The value of \( v_2(4) \) is measured down to \( \langle N_{ch} \rangle \approx 50 \) in \( pp \) collisions and down to \( \langle N_{ch} \rangle \approx 20–40 \) in \( p+Pb \) collisions. The \( v_2(4) \) values are observed to be approximately independent of \( \langle N_{ch} \rangle \) in the measured range in the three data sets: \( 50 < \langle N_{ch} \rangle < 150 \) for 5.02 TeV \( pp \), \( 50 < \langle N_{ch} \rangle < 200 \) for 13 TeV \( pp \), and \( 20 < \langle N_{ch} \rangle < 380 \) for 5.02 TeV \( p+Pb \), respectively. Moreover, the \( p+Pb \) data suggest the value of \( v_2(4) \) is lower for \( \langle N_{ch} \rangle > 200 \), as expected from the similar behavior of \( |c_2(4)| \) in Figs. 7 and 8 at large \( \langle N_{ch} \rangle \).

FIG. 10. The \( c_1(4) \) values calculated for charged particles with \( 0.5 < p_T < 5 \text{ GeV} \) using the standard cumulants (left) and the three-subevent method (right) compared between 5.02 TeV \( pp \), 13 TeV \( pp \), and 5.02 TeV \( p+Pb \). The event averaging is performed for \( N_{ch}^{sel} \) calculated for the same \( p_T \) range, which is then mapped to \( \langle N_{ch} \rangle \), the average number of charged particles with \( p_T > 0.4 \text{ GeV} \). The error bars and shaded boxes represent the statistical and systematic uncertainties, respectively.

FIG. 11. The \( c_1(4) \) values calculated for charged particles with \( 0.3 < p_T < 3 \text{ GeV} \) (left) or \( 0.5 < p_T < 5 \text{ GeV} \) (right) with the three-subevent cumulant method for the \( p+Pb \) data. The event averaging is performed for \( N_{ch}^{sel} \) calculated for various \( p_T \) selections as indicated in the figure, which is then mapped to \( \langle N_{ch} \rangle \), the average number of charged particles with \( p_T > 0.4 \text{ GeV} \). The dashed line indicates the \( c_1(4) \) value corresponding to a 2]\% \( v_3 \) signal. The error bars and shaded boxes represent the statistical and systematic uncertainties, respectively.
The values of $v_2[4]$ presented in Figs. 12 and 13 are also compared to the values of $v_2[2]$ obtained from the 2PC measurements [10,15] where the nonflow effects are estimated using low-multiplicity events ($\langle N_{ch} \rangle < 20$) and then subtracted. The subtraction was performed either by a template fit, which includes the pedestal level from the $\langle N_{ch} \rangle < 20$ events, or by a peripheral subtraction, which sets the pedestal level by a zero-yield at minimum (ZYAM) procedure [6]. The peripheral subtraction explicitly assumes that the most peripheral events do not contain any long-range correlations [15], and so $v_2$ is forced to be zero at the corresponding $\langle N_{ch} \rangle$ value, which biases $v_2$ to a lower value in other multiplicity ranges.

D. Dependence on the number of sources in the initial state

Figures 12 and 13 show that the $v_2[4]$ values are smaller than the $v_2[2]$ values extracted using the template-fit method in both the $pp$ and $p+Pb$ collisions. In various hydrodynamic models for small collision systems [38,39], this difference can be interpreted as the influence of event-by-event flow fluctuations associated with the initial state, which is closely related to the effective number of sources $N_s$ for particle production in the transverse density distribution of the initial state [39]:

$$\frac{v_2[4]}{v_2[2]} = \left( \frac{4}{3 + N_s} \right)^{1/4} \quad \text{or} \quad N_s = \frac{4v_2[2]^4}{v_2[4]^4} - 3. \quad (14)$$

Figure 14 shows the extracted values of $N_s$ as a function of $\langle N_{ch} \rangle$ in 13 TeV $pp$ and 5.02 $p+Pb$ collisions, estimated using charged particles with $0.3 < p_T < 3$ GeV and $0.5 < p_T < 5$ GeV. It is observed that the $N_s$ value increases with $\langle N_{ch} \rangle$ in $p+Pb$ collisions, reaching $N_s \sim 20$ in the highest multiplicity class, and it is consistent between the two $p_T$ ranges.

In the model framework in Refs. [38,39], the values of $|c_{2[4]}|$ and $v_2[4]$ are expected to decrease for large $N_s$, which is compatible with the presented results. The slight decreases of $|c_{2[4]}|$ shown in Figs. 7 and 8 for $p+Pb$ collisions are compatible with the model predictions. The results for 13 TeV $pp$ collisions cover a limited $\langle N_{ch} \rangle$ range compared to $p+Pb$, but agree with $p+Pb$ collisions in this range.

FIG. 12. The $v_2[4]$ values calculated for charged particles with $0.3 < p_T < 3$ GeV using the three-subevent method in 5.02 TeV $pp$ (left), 13 TeV $pp$ (middle), and 5.02 TeV $p+Pb$ collisions (right). They are compared to $v_2$ obtained from the 2PC analyses [10,15] where the nonflow effects are removed by a template fit procedure (solid circles) or with a fit after subtraction with a ZYAM assumption (peripheral subtraction, open circles). The error bars and shaded boxes represent the statistical and systematic uncertainties, respectively.

FIG. 13. The $v_2[4]$ values calculated for charged particles with $0.5 < p_T < 5$ GeV using the three-subevent method in 5.02 TeV $pp$ (left), 13 TeV $pp$ (middle), and 5.02 TeV $p+Pb$ collisions (right). They are compared to $v_2$ obtained from the 2PC analyses [10,15] where the nonflow effects are removed by a template fit procedure (solid circles) or with a fit after subtraction with a ZYAM assumption (peripheral subtraction, open circles). The error bars and shaded boxes represent the statistical and systematic uncertainties, respectively.
nonflow fluctuations associated with the event class chosen for correlation. In general, it is easy to obtain incorrect results from nonflow correlations instead of a long-range collective flow.\[16,19\], based on the standard method, may be dominated by particles used to form the event classes used for averaging. This provides direct evidence that the ridge is indeed a long-range final-state hydrodynamic collective effect.

The subevent cumulant technique and the new results for the presence of long-range multiparticle azimuthal correlations in broad \( \langle N_{ch} \rangle \) ranges in pp and \( p+Pb \) collisions, and these long-range multiparticle correlations persist even in events with rather low multiplicity of \( \langle N_{ch} \rangle \sim 40 \). The \( c_2[4] \) values are consistent with zero in \( pp \) collisions, but are systematically below zero in \( p+Pb \) collisions, compatible with the presence of significant long-range multiparticle triangular flow in \( p+Pb \) collisions.

The single-particle harmonic coefficient \( v_2[4] = (c_2[4])^{1/2} \) is calculated and compared with \( v_2[2] \) obtained previously using the two-particle correlation method, where the nonflow contributions were estimated and subtracted. The magnitude of \( v_2[4] \) is smaller than that for \( v_2[2] \), as expected for a long-range final-state hydrodynamic collective effect. The ratio of \( v_2[4] \) to \( v_2[2] \) is used, in a model-dependent framework, to infer the number of particle-emitting sources in the initial-state geometric configuration. The number of sources extracted within this framework is found to increase with \( \langle N_{ch} \rangle \) in \( p+Pb \) collisions.

The subevent cumulant technique and the new results provide direct evidence that the ridge is indeed a long-range collective phenomenon involving many particles distributed across a broad rapidity interval. The results of \( v_2[4] \) and its dependence on \( p_T \) and \( \langle N_{ch} \rangle \), largely free from nonflow effects, can be used to understand the space-time dynamics and the properties of the medium created in small collision systems.

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