Evidence for light-by-light scattering in heavy-ion collisions with the ATLAS detector at the LHC

ATLAS Collaboration†

Light-by-light scattering ($\gamma \gamma \to \gamma \gamma$) is a quantum-mechanical process that is forbidden in the classical theory of electrodynamics. This reaction is accessible at the Large Hadron Collider thanks to the large electromagnetic field strengths generated by ultra-relativistic colliding lead ions. Using 480 μb$^{-1}$ of lead-lead collision data recorded at a centre-of-mass energy per nucleon pair of 5.02 TeV by the ATLAS detector, here we report evidence for light-by-light scattering. A total of 13 candidate events were observed with an expected background of 2.6 ± 0.7 events. After background subtraction and analysis corrections, the fiducial cross-section of the process $\mathrm{Pb} + \mathrm{Pb} (\gamma \gamma) \to \mathrm{Pb}^{\ast n} + \mathrm{Pb}^{\ast m} \gamma \gamma$, for photon transverse energy $E_T > 3 \mathrm{GeV}$, photon absolute pseudorapidity $|\eta| < 2.4$, diphoton invariant mass greater than 6 GeV, diphoton transverse momentum lower than 2 GeV and diphoton acoplanarity below 0.01, is measured to be 70 ± 24 (stat.) ±17 (syst.) nb, which is in agreement with the standard model predictions.

One of the key features of Maxwell’s equations is their linearity in both the sources and the fields, from which follows the superposition principle. This forbids effects such as light-by-light (LbLy) scattering, $\gamma \gamma \to \gamma \gamma$, which is a purely quantum-mechanical process. It was realized in the early history of quantum electrodynamics (QED) that LbLy scattering is related to the polarization of the vacuum$^1$. In the standard model of particle physics, the virtual particles that mediate the LbLy coupling are electrically charged fermions or $W^{\pm}$ bosons. In QED, the $\gamma \gamma \to \gamma \gamma$ reaction proceeds at lowest order in the fine-structure constant ($\alpha_{\text{em}}$) via virtual one-loop box diagrams involving fermions (Fig. 1a), which is an $\mathcal{O}(\alpha_{\text{em}}^3 \approx 3 \times 10^{-9})$ process, making it challenging to test experimentally. Indeed, the elastic LbLy scattering has remained unobserved: even the ultra-intense laser experiments are not yet powerful enough to probe this phenomenon$^2$.

LbLy scattering via an electron loop has been precisely, albeit indirectly, tested in measurements of the anomalous magnetic moment of the electron and muon$^3$ where it is predicted to contribute substantially, as one of the QED corrections$^4$. The $\gamma \gamma \to \gamma \gamma$ reaction has been measured in photon scattering in the Coulomb field of a nucleus (DeBürck scattering) at fixed photon energies below 7 GeV (refs 6–9). The analogous process, where a photon splits into two photons by interaction with external fields (photon splitting), has been observed in the energy region of 0.1–0.5 GeV (ref. 10). A related process involving only real photons, in which several photons fuse to form an electron–positron pair ($e^+e^-$), has been measured in ref. 11. Similarly, the multiphoton Compton scattering, in which up to four laser photons interact with an electron, has been observed$^{12}$.

An alternative way by which LbLy interactions can be studied is by using relativistic heavy-ion collisions. In ‘ultra-peripheral collision’ (UPC) events, with impact parameters larger than twice the radius of the nuclei$^{13,14}$, the strong interaction does not play a role. The electromagnetic (EM) field strengths of relativistic ions scale with the proton number (Z). For example, for a lead (Pb) nucleus with $Z = 82$ the field can be up to $10^{25} \mathrm{V} \mathrm{m}^{-1}$ (ref. 15), much larger than the Schwinger limit$^{16}$ above which QED corrections become important. In the 1930s it was found that highly relativistic charged particles can be described by the equivalent photon approximation (EPA)$^{17–19}$, which is schematically shown in Fig. 1b. The EM fields produced by the colliding Pb nuclei can be treated as a beam of quasi-real photons with a small virtuality of $Q^2 < 1/R^2$, where $R$ is the radius of the charge distribution and so $Q^2 < 10^{-7} \mathrm{GeV}^2$. Then, the cross-section for the reaction $\mathrm{Pb} + \mathrm{Pb} (\gamma \gamma) \to \mathrm{Pb} + \mathrm{Pb} \gamma \gamma$ can be calculated by convolving the respective photon flux with the elementary cross-section for the process $\gamma \gamma \to \gamma \gamma$. Since the photon flux associated with each nucleon scales as $Z^2$, the cross-section is extremely enhanced as compared with proton–proton (pp) collisions.

In this article, a measurement of LbLy scattering in Pb + Pb collisions at the Large Hadron Collider (LHC) is reported, following the approach recently proposed in ref. 20. The final-state signature of interest is the exclusive production of two photons, $\mathrm{Pb} + \mathrm{Pb} (\gamma \gamma) \to \mathrm{Pb}^{\ast n} + \mathrm{Pb}^{\ast m} \gamma \gamma$, where a possible EM excitation of the outgoing ions$^{21}$ is denoted by (∗). Hence, the expected signature is two photons and no further activity in the central detector, since the Pb$^{\ast n}$ ions escape into the LHC beam pipe. Moreover, it is predicted that the background is relatively low in heavy-ion collisions and is dominated by exclusive dielectron ($\gamma \gamma \to e^+e^-$) production$^{20,22}$. The misidentification of electrons as photons can occur when the electron track is not reconstructed or the electron emits a hard-bremsstrahlung photon. The fiducial cross-section of the process $\gamma \gamma \to \gamma \gamma$ in Pb + Pb collisions is measured, using a data set recorded at a nucleon–nucleon centre-of-mass energy ($\sqrt{s_{\text{NN}}}$) of 5.02 TeV. This data set was recorded with the ATLAS detector at the LHC in 2015 and corresponds to an integrated luminosity of 480 ± 30 μb$^{-1}$. In addition to the measured fiducial cross-section, the significance of the observed number of signal candidate events is given, assuming the background-only hypothesis.

Experimental set-up

ATLAS is a cylindrical particle detector composed of several subdetectors$^{23}$. ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point in the centre of the detector and the z axis along the beam pipe. The x axis points from the interaction point to the centre of the LHC ring, and the y

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axis points upwards. Cylindrical coordinates \((r, \phi)\) are used in the transverse plane, with \(\phi\) being the azimuthal angle around the \(z\) axis. The pseudorapidity \(\eta\) is defined in terms of the polar angle \(\theta\) as \(\eta = -\ln \tan(\theta/2)\).

Angular distance is measured in units of \(\Delta R \equiv \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2}\). The photon or electron transverse energy is \(E_T = E \sin(\theta)\), where \(E\) is its energy. The inner tracking detector (ITD) consists of a silicon pixel system, a silicon microstrip detector and a straw-tube tracker immersed in a 2T magnetic field provided by a superconducting solenoid. The ITD track reconstruction efficiency is estimated in events that, like UPC Pb + Pb events, have a low average track multiplicity. For charged hadrons in the solenoid. The ITD track reconstruction efficiency is estimated in immersed in a 2T magnetic field provided by a superconducting pixel system, a silicon microstrip detector and a straw-tube tracker immersed in a 2T magnetic field provided by a superconducting solenoid. The ITD track reconstruction efficiency is estimated in immersed in a 2T magnetic field provided by a superconducting solenoid. The ITD track reconstruction efficiency is estimated in immersed in a 2T magnetic field provided by a superconducting solenoid. The ITD track reconstruction efficiency is estimated in immersed in a 2T magnetic field provided by a superconducting solenoid.

The photon or electron transverse energy is \(E\) and scintillator-tile active material and steel absorbers for the EM and forward \(|\eta| > 1.7\) hadronic components of the detector, and scintillator-tile active material and steel absorbers for the EM and forward \(|\eta| < 1.7\) hadronic component. The muon spectrometer consists of separate trigger and high-precision tracking chambers measuring the trajectory of muons in a magnetic field generated by superconducting air-core toroids. The ATLAS minimum-bias trigger scintillators (MBTs) consist of scintillator slabs positioned between the ITD and the endcap calorimeters with each side having an outer ring of four slabs segmented in azimuthal angle, covering 2.87 < |\(\eta| < 2.76\) and an inner ring of eight slabs, covering 2.76 < |\(\eta| < 3.86\). The ATLAS zero-degree calorimeters (ZDCs), located along the beam axis at 140 m from the interaction point on both sides, detect neutral particles (including neutrons emitted from the nucleus). The ATLAS trigger system consists of a Level-1 trigger implemented using a combination of dedicated electronics and programmable logic, and a software-based high-level trigger.

Monte Carlo simulation and theoretical predictions
Several Monte Carlo (MC) samples are produced to estimate background contributions and corrections to the fiducial measurement. The detector response is modelled using a simulation based on GEANT4 (refs 26, 27). The data and MC simulated events are passed through the same reconstruction and analysis procedures.

LbL signal events are generated taking into account box diagrams with charged leptons and quarks in the loops, as detailed in ref 28. The contributions from \(W\)-boson loops are omitted in the calculations since they are mostly important for diphoton masses \(m_{\gamma\gamma} > 2m_W\) (ref 29). The calculations are then convolved with the Pb + Pb EPA spectrum from the STARlight 1.1 MC generator (ref 30). Next, various diphoton kinematic distributions are cross-checked with predictions from ref 20 and good agreement is found.

Diphoton reconstruction is performed using a combination of dedicated electronics and programmable logic, and a software-based high-level trigger. The sources of background considered in this analysis are:\(\gamma\gamma \rightarrow e^+e^-\), central exclusive production (CEP) of photon pairs, exclusive production of quark–antiquark pairs (\(\gamma\gamma \rightarrow q\bar{q}\)) and other backgrounds that could mimic the diphoton event signatures. The \(\gamma\gamma \rightarrow e^+e^-\) background is modelled with STARlight 1.1 (ref 30), in which the cross-section is computed by combining the Pb + Pb EPA with the leading-order formula for \(\gamma\gamma \rightarrow e^+e^-\). This process has been recently measured by the ALICE Collaboration, and a good agreement with STARlight is found (ref 20). The exclusive diphoton final state can be also produced via the strong interaction through a quark loop in the exchange of photons that can couple to form a given final state \(X\). The flux of photons is determined from the Fourier transform of the electromagnetic field of the ion, taking into account the nuclear electromagnetic form factors.

Event selection
Candidate diphoton events were recorded in the Pb + Pb run in 2015 using a dedicated trigger for events with moderate activity in the calorimeter but little additional activity in the entire detector. At Level-1 the total \(E_T\) registered in the calorimeter after noise suppression was required to be between 5 and 200 GeV. Then at the high-level trigger, events were rejected if more than one hit was found in the inner ring of the MBTS (MBTS veto) or if more than 10 hits were found in the pixel detector.

The efficiency of the Level-1 trigger is estimated with \(\gamma\gamma \rightarrow e^+e^-\) events passing an independent supporting trigger. This trigger is designed to select events with mutual dissociation of Pb nuclei and small activity in the ITD. It is based on a coincidence of signals in both ZDCs and a requirement on the total \(E_T\) in the calorimeter below 50 GeV. Event candidates are required to have only two reconstructed tracks and two EM energy clusters. Furthermore, to reduce possible backgrounds, each pair of clusters (cl1, cl2) is required to have a small acoplanarity \((1 - \Delta \phi_{cl1,cl2}/\pi < 0.2)\). The extracted Level-1 trigger efficiency is provided as a function of the...
sum of cluster transverse energies ($E_T^{cl1} + E_T^{cl2}$). The efficiency grows from about 70% at ($E_T^{cl1} + E_T^{cl2}$) = 6 GeV to 100% at ($E_T^{cl1} + E_T^{cl2}$) > 9 GeV. The efficiency is parameterized using an error function fit, which is then used to reweight the simulation. Due to the extremely low noise, very high hit reconstruction efficiency and low conversion probability of signal photons in the pixel detector (around 10%), the uncertainty due to the requirement for minimal activity in the ITD is negligible. The MBTS veto efficiency was studied using $\gamma \gamma \rightarrow e^+e^-$ events (for example, cosmic-ray muons), the transverse momentum of the diphoton system ($p_T^{\gamma \gamma}$) is required to be below 2 GeV. To reduce background from CEP $gg \rightarrow \gamma \gamma$ reactions, an additional requirement on diphoton acoplanarity, $Aco = 1 - \Delta \phi_{\gamma \gamma} / \pi \sim 0.01$, is imposed. This requirement is optimized to retain a high signal efficiency and reduce the CEP background significantly, since the transverse momentum transferred by the photon exchange is usually much smaller than that due to the colour-singlet-state gluons.

**Performance and validation of photon reconstruction**

Since the analysis requires the presence of low-energy photons, which are not typically used in ATLAS analyses, detailed studies of photon reconstruction and calibration are performed. High-$p_T\gamma\gamma\rightarrow e^+e^-$ production with a final-state radiation (FSR) photon is used for the measurement of the photon PID efficiency. Events with a photon and two tracks corresponding to oppositely charged particles with $p_T > 1$ GeV are required to pass the same trigger as in the diphoton selection or the supporting trigger. The $\Delta R$ between a photon candidate and a track is required to be greater than 0.2 to avoid leakage of the electron clusters from the $\gamma \gamma \rightarrow e^+e^-$ process to the photon cluster. The FSR event candidates are identified using a $p_T^{\gamma \gamma} < 1$ GeV requirement, where $p_T^{\gamma \gamma}$ is the transverse momentum of the three-body system consisting of two charged-particle tracks and a photon. The FSR photons are then used to extract the photon PID efficiency, which is defined as the probability for a reconstructed photon to satisfy the identification criteria. Figure 2a shows the photon PID efficiencies in data and simulation as a function of reconstructed photon $E_T$.

The photon reconstruction efficiency is extracted from data using $\gamma \gamma \rightarrow e^+e^-$ events where one of the electrons emits a hard-bremsstrahlung photon due to interaction with the material of the detector. Events with exactly one identified electron, two reconstructed charged-particle tracks and exactly one photon are studied. The electron $E_T$ is required to be above 5 GeV and the $p_T$
Figure 3 | Kinematic distributions for $\gamma\gamma \rightarrow \gamma\gamma$ event candidates. a, Diphoton acoplanarity before applying the Aco < 0.01 requirement. b, Diphoton invariant mass after applying the Aco < 0.01 requirement. Data (points) are compared to MC predictions (histograms). The statistical uncertainties on the data are shown as vertical bars.

of the track that is unmatched with the electron (trk2) is required to be below 2 GeV. The additional hard-bremstrahlung photon is expected to have $E_\gamma \approx (E_T - p_T^{\text{em}})$. The $p_T^{\text{em}} < 2$ GeV requirement ensures a sufficient $\Delta R$ separation between the expected photon and the second electron, extrapolated to the first layer of the EM calorimeter. The data sample contains 247 $\gamma\gamma \rightarrow e^+e^-$ events that are used to extract the photon reconstruction efficiency, which is presented in Fig. 2b. Good agreement between data and $\gamma\gamma \rightarrow e^+e^-$ MC simulation is observed and the photon reconstruction efficiency is measured with a 5–10% relative uncertainty at low $E_T$ (3–6 GeV).

In addition, a cross-check is performed on $Z \rightarrow \mu^+\mu^-$ events identified in $pp$ collision data from 2015 corresponding to an integrated luminosity of 1.6 fb$^{-1}$. The results support (in a similar way to ref. 42) the choice to use the three shower-shape variables in this photon PID selection in an independent sample of low-$E_T$ photons.

The photon cluster energy resolution is extracted from data using $\gamma\gamma \rightarrow e^+e^-$ events. The electrons from the $\gamma\gamma \rightarrow e^+e^-$ reaction (see Supplementary Information) are well balanced in their transverse momenta, with very small standard deviation, $\sigma_{E_T}^{e_1} - \sigma_{E_T}^{e_2} < 30$ MeV, much smaller than the expected EM calorimeter energy resolution. Therefore, by measuring $(E_T^{e_1} - E_T^{e_2})$ distributions in $\gamma\gamma \rightarrow e^+e^-$ events, one can extract the cluster energy resolution, $\sigma_{E_T}^{e}$. For electrons with $E_T < 10$ GeV, the $\sigma_{E_T}^{e_2}/E_T$ is observed to be approximately 8% both in data and simulation. An uncertainty of $\sigma_{E_T}^{e_2}/\sigma_{E_T}^{e_1}$ is observed to be smaller than the expected EM calorimeter energy resolution and takes into account the differences between $\sigma_{E_T}^{e_2}$ in data and $\sigma_{E_T}^{e_1}$ in simulation.

Similarly, the EM cluster energy scale can be studied using the $(E_T^{e_1} + E_T^{e_2})$ distribution. It is observed that the simulation provides a good description of this distribution, within the relative uncertainty of 5% that is assigned to the EM cluster energy-scale modelling.

Background estimation

Due to its relatively high rate, the exclusive production of electron pairs ($\gamma\gamma \rightarrow e^+e^-$) can be a source of fake diphoton events. The contribution from the dielectron background is estimated using $\gamma\gamma \rightarrow e^+e^-$ MC simulation (which gives 1.3 events) and is verified using the following data-driven technique. Two control regions are defined that are expected to be dominated by $\gamma\gamma \rightarrow e^+e^-$ backgrounds. The first control region is defined by requiring events with exactly one reconstructed charged-particle track and two identified photons that satisfy the same preselection criteria as for the signal definition. The second control region is defined similarly to the first one, except exactly two tracks are required ($N_{\text{trk}} = 2$). Good agreement is observed between data and MC simulation in both control regions, but the precision is limited by the number of events in data. A conservative uncertainty of 25% is therefore assigned to the $\gamma\gamma \rightarrow e^+e^-$ background estimation, which reflects the statistical uncertainty of data in the $N_{\text{trk}} = 1$ control region. The contribution from a related QED process, $\gamma\gamma \rightarrow e^+e^-\gamma\gamma$, is evaluated using the MadGraph5_aMC@NLO MC generator and is found to be negligible.

The Aco < 0.01 requirement significantly reduces the CEP $gg \rightarrow \gamma\gamma$ background. However, the MC prediction for this process has a large theoretical uncertainty, hence, an additional data-driven normalization is performed in the region Aco > b, where b is a value greater than 0.01 which can be varied. Three values of b (0.01, 0.02, 0.03) are used, where the central value b = 0.02 is chosen to derive the nominal background prediction and the values b = 0.01 and b = 0.03 to define the systematic uncertainty. The normalization is performed using the condition: $f_{\gamma\gamma}^{\text{diphoton}} = (N_{\text{data}}(\text{Aco} > b) - N_{\text{bg}}(\text{Aco} > b))/N_{\gamma\gamma}^{\text{Diphoton}}(\text{Aco} > b)$, for each value of b, where $N_{\text{data}}$ the number of observed events, $N_{\text{bg}}$ is the expected number of signal events, $N_{\gamma\gamma}^{\text{Diphoton}}$ is the expected background from $\gamma\gamma \rightarrow e^+e^-$ events and $N_{\gamma\gamma}^{\text{Diphoton}}$ is the MC estimate of the expected background from CEP $gg \rightarrow \gamma\gamma$ events. The normalization factor is found to be $f_{\gamma\gamma}^{\text{diphoton}} = 0.5 \pm 0.3$ and the background due to CEP $gg \rightarrow \gamma\gamma$ is estimated to be $f_{\gamma\gamma}^{\text{diphoton}} \times N_{\gamma\gamma}^{\text{Diphoton}}(\text{Aco} < 0.01) = 0.9 \pm 0.5$ events. To verify the CEP $gg \rightarrow \gamma\gamma$ background estimation method, energy deposits in the ZDC are studied for events before the Aco selection. It is expected that the outgoing ions in CEP events predominantly dissociate, which results in the emission of neutrons detectable in the ZDC$^{39}$. Good agreement between the normalized CEP $gg \rightarrow \gamma\gamma$ MC expectation and the observed events with a ZDC signal corresponding to at least 1 neutron is observed in the full Aco range (see Supplementary Information for details).

Low-$p_T$ dijet events can produce multiple $n\pi^0$ mesons, which could potentially mimic diphoton events. The event selection requirements are efficient in rejecting such events, and based on studies performed with a supporting trigger, the background from hadronic processes is estimated to be 0.3 ± 0.3 events. MC studies show that the background from $\gamma\gamma \rightarrow \rho\rho$ processes is negligible.

Exclusive neutral two-meson production can be a potential source of background for LbL events, mainly due to their back-to-back topology being similar to that of the CEP $gg \rightarrow \gamma\gamma$ process. The cross-section for this process is calculated to be below 10% of the CEP $gg \rightarrow \gamma\gamma$ cross-section$^{44,45}$ and it is therefore considered to
give a negligible contribution to the signal region. The contribution from bottomonia production (for example, $\gamma\gamma \rightarrow \eta_b \gamma \gamma$ or $\gamma \mathrm{Pb} \rightarrow \gamma \gamma \eta_b \rightarrow 3\gamma$) is calculated using parameters from refs 46, 47 and is found to be negligible.

The contribution from other fake diphoton events (for example those induced by cosmic-ray muons) is estimated using photons that fail to satisfy the longitudinal shower-shape requirement. The total background due to other fake photons is found to be $0.1 \pm 0.1$ events. As a further cross-check, additional activity in the region $p_T^{\gamma} < 2$ GeV, no events with muon activity are found, which is compatible with the abovementioned estimate of $0.1 \pm 0.1$.

The contribution from UPC events where both nuclei emit a bremsstrahlung photon is estimated using calculations from ref. 13, and is found to be negligible for photons with $|\eta| < 2.4$ and $E_\gamma > 3$ GeV.

### Results

Photon kinematic distributions for events satisfying the selection criteria are shown in Fig. 3. The shape of the diphoton acoplanarity distribution for $\gamma\gamma \rightarrow e^+e^-$ events in Fig. 3a reflects the trajectories of the electron and positron in the detector magnetic field, before they emit hard photons in their collisions with the ITD material. In total, 13 events are observed in data whereas 7.3 signal events and 6.3 background events are expected. In general, good agreement between data and MC simulation is observed. The effect of sequential selection requirements on the number of events selected is shown in Table 1, for each of the data, signal and background samples.

To quantify an excess of events over the background expectation, a test statistic based on the profile likelihood ratio $C$ is used. The $p$ value for the background-only hypothesis, defined as the probability for the background to fluctuate and give an excess of events as large or larger than that observed in the data, is found to be $5 \times 10^{-6}$. The $p$ value can be expressed in terms of Gaussian tail probabilities, which, given in units of standard deviation ($\sigma$), corresponds to a significance of $4.4\sigma$. The expected $p$ value and significance (obtained before the fit of the signal-plus-background hypothesis to the data and using standard model predictions from ref. 28) are $8 \times 10^{-3}$ and $3.8\sigma$, respectively.

The cross-section for the $\mathrm{Pb} + \mathrm{Pb}$ ($\gamma\gamma \rightarrow \gamma\gamma$) process is measured in a fiducial phase space defined by the photon transverse energy $E_\gamma > 3$ GeV, photon absolute pseudorapidity $|\eta| < 2.4$, diphoton invariant mass greater than 6 GeV, diphoton transverse momentum lower than 2 GeV and diphoton acoplanarity below 0.01. Experimentally, the fiducial cross-section is given by

$$\sigma_{\text{fid}} = \frac{N_{\text{data}} - N_{\text{bkgr}}}{C \times \int L \, dt}$$

(1)

where $N_{\text{data}}$ is the number of selected events in data, $N_{\text{bkgr}}$ is the expected number of background events and $\int L \, dt$ is the integrated luminosity.

### Conclusion

In summary, this article presents evidence for the scattering of LbyL in quasi-real photon interactions from 480 $\mu$b$^{-1}$ of ultra-peripheral $\mathrm{Pb} + \mathrm{Pb}$ collisions at $\sqrt{s_{NN}} = 5.02$ TeV by the ATLAS experiment at the LHC. The statistical significance against the background-only hypothesis is found to be 4.4 standard deviations. After background subtraction and analysis corrections, the fiducial cross-section for the $\mathrm{Pb} + \mathrm{Pb}$ ($\gamma\gamma \rightarrow \gamma\gamma$) process was measured and is compatible with standard model predictions.

The analysis is mostly limited by the amount of data available and the lower limit on transverse energy for reconstructed photons ($E_\gamma > 3$ GeV), below which more signal is expected. Advancements on these two points would also allow for reconstruction of low-mass mesons decaying into two photons, which in turn could be used to improve detector calibration. The heavy-ion data yield is expected to double at the end of 2018 (and again increase tenfold after

### Table 1 | The number of events accepted by the sequential selection requirements for data, compared with the number of background and signal events expected from the simulation.

<table>
<thead>
<tr>
<th>Selection</th>
<th>$\gamma\gamma \rightarrow e^+e^-$</th>
<th>CEP $gg \rightarrow \gamma\gamma$</th>
<th>Hadronic fakes</th>
<th>Other fakes</th>
<th>Total background</th>
<th>Signal</th>
<th>Data</th>
</tr>
</thead>
<tbody>
<tr>
<td>Preselection</td>
<td>74</td>
<td>4.7</td>
<td>6</td>
<td>19</td>
<td>104</td>
<td>9.1</td>
<td>105</td>
</tr>
<tr>
<td>$N_{\text{tag}} = 0$</td>
<td>4.0</td>
<td>4.5</td>
<td>6</td>
<td>19</td>
<td>33</td>
<td>8.7</td>
<td>39</td>
</tr>
<tr>
<td>$p_T^{\gamma} &lt; 2$ GeV</td>
<td>3.5</td>
<td>4.4</td>
<td>3</td>
<td>1.3</td>
<td>12.2</td>
<td>8.5</td>
<td>21</td>
</tr>
<tr>
<td>Acceptance ($p_\text{T}$)</td>
<td>0.01</td>
<td>0.9</td>
<td>0.3</td>
<td>0.1</td>
<td>0.7</td>
<td>0.7</td>
<td>1.5</td>
</tr>
<tr>
<td>Uncertainty</td>
<td>0.3</td>
<td>0.5</td>
<td>0.3</td>
<td>0.7</td>
<td>1.5</td>
<td>1.5</td>
<td>2.6</td>
</tr>
</tbody>
</table>

The signal simulation is based on calculations from ref. 28. In addition, the uncertainties on the expected number of events passing all selection requirements are given.

### Table 2 | Summary of systematic uncertainties.

<table>
<thead>
<tr>
<th>Source of uncertainty</th>
<th>Relative uncertainty</th>
</tr>
</thead>
<tbody>
<tr>
<td>Trigger</td>
<td>5%</td>
</tr>
<tr>
<td>Photon reco. efficiency</td>
<td>12%</td>
</tr>
<tr>
<td>Photon PID efficiency</td>
<td>16%</td>
</tr>
<tr>
<td>Photon energy scale</td>
<td>7%</td>
</tr>
<tr>
<td>Photon energy resolution</td>
<td>11%</td>
</tr>
<tr>
<td>Total</td>
<td>24%</td>
</tr>
</tbody>
</table>

The table shows the relative systematic uncertainty on detector correction factor $C$ broken into its individual contributions. The total is obtained by adding them in quadrature.
LHC Run 4, scheduled to start in 2026), which would significantly reduce the statistical uncertainty. Future upgrades of ATLAS, such as extended tracking acceptance from $|\eta| < 2.5$ to $|\eta| < 4.0$, will further improve this.

**Data availability.** The experimental data that support the findings of this study are available in HEPData with the identifier http://dx.doi.org/10.17182/hepdata.77761.

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**References**


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Author contributions
All authors have contributed to the publication, being variously involved in the design and the construction of the detectors, in writing software, calibrating subsystems, operating the detectors and acquiring data, and finally analysing the processed data. The ATLAS Collaboration members discussed and approved the scientific results. The manuscript was prepared by a subgroup of authors appointed by the collaboration and subject to an internal collaboration-wide review process. All authors reviewed and approved the final version of the manuscript.

Additional information
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Competing financial interests
The authors declare no competing financial interests.

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