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Double occupancy in DMFT and the Dual Boson approach

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We discuss the calculation of the double occupancy using Dynamical Mean-Field Theory (DMFT) in finite dimensions. The double occupancy can be determined from the susceptibility of the auxiliary impurity model or from the lattice susceptibility. The former method typically overestimates, whereas the latter underestimates the double occupancy. We illustrate this for the square-lattice Hubbard model. We propose an approach for which both methods lead to identical results by construction and which resolves this ambiguity. This self-consistent dual boson scheme results in a double occupancy that is numerically close to benchmarks available in the literature.

The double occupancy plays an important role in the study of correlated systems. It is indicative of the Mott metal-insulator transition and of local moment formation. In optical lattice experiments, the double occupancy gives information about the phase. On the theory side, the double occupancy has been used to benchmark approximations. It also enters the calculation of total energies and forces in LDA+DMFT studies of strongly correlated materials. Given this important role, it is worthwhile to look at how the double occupancy is determined. Here, we consider the double occupancy in DMFT and its extension Dual Boson (DB). We compare our results with benchmark results available in the literature.

Over the past two decades, DMFT has become the dominant approximation for strongly correlated electron systems. DMFT solves a self-consistently determined auxiliary single-site problem (the impurity problem) to determine properties of the original lattice problem. Initially, DMFT studies focused on infinite-dimensional systems, where the approximation becomes exact. In this case, the double occupancy is given by the double occupancy of the auxiliary impurity problem. Nowadays, DMFT has become an accepted approximation also for finite-dimensional systems. In these calculations, it is often assumed that the double occupancy is equal to that of the impurity problem. Here we show that this assumption is potentially problematic.

To be concrete, we study the Hubbard model as a prototypical example of a strongly correlated system. It is described by the Hamiltonian

\[ H = -t \sum_{\langle ab \rangle \sigma} c_{a \sigma}^{\dagger} c_{b \sigma} + \sum_{a} U n_{a \uparrow} n_{a \downarrow}, \]  

where \( t \) is the hopping amplitude of an electron with spin \( \sigma \) between nearest-neighbor sites \( a \) and \( b \), and \( U \) is the on-site repulsion between electrons. The local repulsion \( U \) disfavors doubly-occupied sites and hence decreases the double occupancy \( d = \langle n_{\uparrow} n_{\downarrow} \rangle \).

The double occupancy is a local two-particle correlation function. It can be written as the difference between the charge and spin susceptibility. In obvious short-hand notation,

\[ d = \frac{1}{4} \left[ \langle nn \rangle - \langle S_{z} S_{z} \rangle \right] = \frac{1}{4} \left[ X_{\text{loc}}^{\text{ch}} - X_{\text{loc}}^{\text{sz}} + \langle n \rangle \langle n \rangle - \langle S_{z} \rangle \langle S_{z} \rangle \right]. \]  

FIG. 1. Visual illustration of Eq. (3): Shown are example diagrams contributing to the local part of the DMFT susceptibility. The impurity susceptibility \( \chi_{\text{imp}} \) contains the local bubble and local vertex corrections built from local propagators. \( \chi \) contains all the remaining diagrams. These describe processes that start and end at a site \( a \) but involve intermediate scattering on different sites \( b \) and nonlocal propagators.

Here \( X_{\text{loc}}^{\text{ch}} = \langle nn \rangle - \langle n \rangle \langle n \rangle \) and \( X_{\text{loc}}^{\text{sz}} = \langle S_{z} S_{z} \rangle - \langle S_{z} \rangle \langle S_{z} \rangle \) are the equal-time, local correlation functions of the charge density \( n = n_{\uparrow} + n_{\downarrow} \) and the magnetization \( S_{z} = n_{\uparrow} - n_{\downarrow} \), respectively. In DMFT, the charge and spin susceptibilities in turn are either approximated by the respective impurity susceptibility, or by the DMFT lattice susceptibilities. The latter are computed by summing ladder diagrams containing the lattice Green’s function and a local irreducible vertex. The computation of the DMFT double occupancy is ambiguous, because the two-particle impurity correlation functions are not identical to the local lattice correlation functions. Recently, we showed this for the compressibility.

To emphasize their difference, we can decompose the susceptibility into the impurity susceptibility \( \chi_{\text{imp}} \) and a
this means that the DMFT susceptibility (blue triangles) overestimates \(X^{sz}\) and thus, according to (2), underestimates \(d\). This happens particularly at large \(U\), where the tendency towards antiferromagnetism is more pronounced and hence the strongest deviations occur. In fact, we find that the DMFT susceptibility results in a negative double occupancy at sufficiently large values of \(U\), which is clearly unphysical. The double occupancy of the auxiliary impurity, on the other hand, gets close to the DCA results at large \(U\). As \(U\) increases, the electrons localize and nonlocal corrections are less important.

We now turn our attention to the self-consistent DB (sc-DB) approach\(^{14,18}\). The hallmark of this approach is the self-consistency condition\(^{19}\),

\[
X_{\text{loc},\omega} = \chi_{\omega}^{\text{imp}},
\]

which resolves the above ambiguity by construction.\(^{20}\) This self-consistency condition is achieved by introducing a frequency-dependent interaction \(\Lambda_\omega\) to the auxiliary impurity model. The effect of this interaction is twofold: First of all, the auxiliary impurity problem and the associated \(\chi^{\text{imp}}\) are different from their DMFT values. Second, the lattice susceptibility has to be calculated as \(X^{-1} = X_{\text{DMFT}}^{-1} + \Lambda\), where \(X_{\text{DMFT}}\) is the usual DMFT susceptibility obtained from the two-particle ladder, and the equation should be understood in momentum and frequency space. It resembles the RPA expression for the susceptibility, where \(X_{\text{DMFT}}\) plays the role of the bare susceptibility and \(-\Lambda\) the role of the interaction.\(^{21}\) We determine a self-consistent field both for the density and the spin \((S_z)\) channel, as they differ in general.

The self-consistent dual boson approach described here shares certain characteristics with the two-particle self-consistent approach\(^{22–24}\) (TPSC). Both approaches determine an effective interaction for the spin and charge channel according to a self-consistency condition. Note however, that the original TPSC approach takes the RPA as its starting point, sc-DB starts from DMFT. Strong correlation effects are included in DB from the start. Furthermore, the effective interaction in TPSC is static. The Moriya-esque \(\lambda\) correction\(^{25}\) in DΓA\(^{26}\) is also somewhat similar in spirit. It, too, is a correction to the susceptibility that is used to impose self-consistency. Like in TPSC, the Moriya-esque correction is static. Therefore, in both cases, these can only fix certain sum-rules. The effective interaction in sc-DB, on the contrary, may be regarded as a frequency-dependent \(\Lambda\)-correction, which removes the ambiguity of calculating local susceptibilities.

The purple crosses in Fig. 2 show the result of an sc-DB calculation with self-consistently determined frequency-dependent interactions \(\Lambda_\omega\) in the density and spin \((S_z)\) channel. The results are much closer to the DCA benchmark than both DMFT approaches. This shows that taking into account feedback of collective excitations onto the effective impurity problem can lead to a significant improvement.

We note that it was difficult to converge our scheme for \(U > 8\). This coincides with the parameter region where
The retardation interactions appear as a difference between the screened and bare interaction for different values of $U$. The retarded interaction is largest in the charge channel. This behavior is also observed in TPSC. The Hubbard repulsion strongly suppresses charge fluctuations. At $U = 8$, the charge channel develops a small maximum at intermediate frequency. This occurs at the typical energy scale of charge fluctuations, i.e., $\omega_n \approx U$, however the precise physical origin of this maximum is currently unknown.

The dimensionality of the model plays an important role in DMFT-based studies. To illustrate the effect of the dimensionality, we have performed similar calculations in a 3d simple cubic lattice. The results are shown in Fig. 4. A direct comparison is difficult by the change in energy scales that occurs when changing the dimension. To overcome this, we use the $Z$-factor, $Z = (1 - d \Re \Sigma_n / d \omega)^{-1}$, obtained in DMFT, to indicate the importance of interaction effects on the one-particle level. The renormalization of the double occupancy by nonlocal processes, $d_{\text{sc-DB}} / d_{\text{DMFT}}$, is shown as a function of $Z$. The figure clearly shows that the sc-DB susceptibility contains important nonlocal corrections. Although it is clear that other factors than dimensionality also play a role, Fig. 4 also suggests that the nonlocal renormalization of the double occupancy is less important in the 3d system. This result matches our physical expectation that nonlocal correlation effects are stronger in lower dimensions. In fact, vertex corrections to the susceptibility vanish in infinite dimension, $X^{\text{local}} = 0$, and $d = d_{\text{DMFT}}$ in DMFT, so no ambiguity occurs and there is no nonlocal renormalization of the double occupancy.

Fig. 4 also shows that the nonlocal renormalization is small in weakly interacting systems, with $Z \approx 1$. According to Eq. 3, nonlocal vertex corrections are responsible for this nonlocal renormalization, and the vertex is small at small $U$. On the other hand, at very large $U$, the electrons would be very localized and as a result, nonlocal vertex corrections are again small. The uptick between $U = 7$ to $U = 8$ in Fig. 4 is a first sign of this. This behav-
ior is similar to the eigenvalue analysis in dual fermion. The corrections to the local auxiliary impurity are largest at intermediate $U$.

In Fig. 5, we show the double occupancy similar to Fig. 2, away from half-filling, at constant $U = 2$ and lower temperature $\beta = 8$ as a function of $\langle n \rangle$. For clarity, we show the double occupancy normalized by the density squared. For a noninteracting system, $d = \langle n \rangle^2$ / 4, so the deviation from 1/4 shows the interaction effects. At this lowered temperature, the half-filled system has strong antiferromagnetic fluctuations that inhibit sc-DB calculations. As in Fig. 2, both the DCA and sc-DB results lie between the results obtained from the DMFT lattice susceptibility and from the auxiliary impurity model. Again the impurity susceptibility over- and the lattice susceptibility underestimate the double occupancy. The largest inconsistency between the two occurs actually towards lower filling, similar to what happens for the compressibility. While the DCA benchmarks only go down to $\langle n \rangle = 0.3$, these results suggest that the region at lower fillings still contains interesting physics.

In conclusion, we have proposed a scheme which resolves the ambiguity in the calculation of the double occupancy inherent to DMFT in finite dimensions. DMFT is not self-consistent on the two-particle level, and as a result, the single- and two-particle quantities are not compatible, as exemplified here by the difference between the double occupancy from the Galitskii-Migdal formula and from the susceptibility. This ambiguity extends to the determination of total energies and forces from DMFT.

The self-consistent dual boson approach gives results which are in good agreement with benchmarks in the literature. We have further shown that the DMFT impurity double occupancy typically overestimates the double occupancy, while the occupancy determined from the DMFT lattice susceptibility underestimates it significantly. For large interaction, this can even lead to unphysical negative results. We have seen that in two-dimensional systems and at a moderate $Z$-factor of 0.5, the results may differ by 10%-20%. The double occupancies determined from the DMFT lattice and impurity susceptibilities differ strongly also away from half filling. The discrepancy may have a significant effect on total energy and force calculations and hence the determination of equilibrium positions of the atoms in realistic LDA+DMFT calculations. This should be kept in mind when interpreting the results. In order to avoid unphysical results we recommend to approximate the double occupancy by the impurity double occupancy in LDA+DMFT. The accuracy of this approach however remains to be determined.

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Most papers are not very explicit in how they calculate the double occupancy. An alternative expression is
\[ d = \int_{\nu} 2T/U \Sigma_{\nu} G_{\nu}, \]
based on the Galitskii-Migdal formula for the total energy. In infinite dimensions and in the DMFT approximation, the self-energy is local, and the formula simplifies to
\[ d = \int_{\nu} 2T/U \Sigma_{\nu} G_{\nu}^{\text{loc}}. \]
This yields exactly the same double occupancy as the impurity problem, since the impurity problem is solved exactly.

To emphasize the analogy between magnetic and charge susceptibilities, we define magnetization without a factor \( \frac{1}{2} \).


Without the self-consistency condition and in absence of nonlocal interaction, the dual boson approach reduces to DMFT, in particular for the Hubbard model.

By introducing an additional interaction \( \Lambda \) to the impurity part of the problem, the lattice model is left with an interaction \( -\Lambda \).


