The following full text is a publisher's version.

For additional information about this publication click this link.
http://hdl.handle.net/2066/128947

Please be advised that this information was generated on 2017-09-02 and may be subject to change.
Measurement of the $B \rightarrow J/\psi K^*(892)$ Decay Amplitudes

We present a measurement of the decay amplitudes in $B \to J/\psi K^*(892)$ channels using 20.7 fb$^{-1}$ of data collected at the Y(4S) resonance with the BABAR detector at PEP-II. We measure a $P$-wave fraction $R_+ = (16.0 \pm 3.2 \pm 1.4)\%$ and a longitudinal polarization fraction $(59.7 \pm 2.8 \pm 2.4)\%$. The measurement of a relative phase that is neither 0 nor $\pi$, $\phi_0 = 2.50 \pm 0.20 \pm 0.08$ radians, favors a departure from the factorization hypothesis. Although the decay $B \to J/\psi K \pi$ proceeds mainly via $K^*(892)$, there is also evidence for $K^*(1430)$ and $K \pi$ $S$-wave contributions.

DOI: 10.1103/PhysRevLett.87.241801 PACS numbers: 13.25.Hw, 11.30.Er, 12.15.Hh
The decay $B^0 \to J/\psi K^{*0}$ with $K^{*0} \to K^0 \pi^0$ allows a measurement of the CP violation parameter $\sin 2\beta$ that is theoretically as clean as for $B^0 \to J/\psi K^0_s$ [1]. However, due to the presence of even ($L = 0, 2$) and odd ($L = 1$) orbital angular momenta in the $J/\psi K^*$ system, there can be CP-even and CP-odd contributions to the decay rate. If the information contained in the decay angles is ignored, the measured time-dependent CP asymmetry is reduced by the dilution factor $D_\perp = 1 - 2R_\perp$, where $R_\perp$ is the fraction of the $P$-wave. If the angular information is used, the CP components can be separated [2].

The angular analysis also provides a test of the factorization hypothesis, the validity of which is in question for color-suppressed modes [3,4]. In this scheme, the weak decay is described by a product of $J/\psi$ and $B \to K^*$ hadronic currents, and final state interactions are neglected. If factorization holds, the decay amplitudes should have relative phase 0 or $\pi$.

The decay $B \to J/\psi K^*(892)$ is described by three amplitudes. In the transversity basis [2,5] used by CLEO [6] and CDF [7], the amplitudes $A_0$, $A_1$, and $A_2$ have CP eigenvalues $+1$, $+1$, and $-1$, respectively. $A_0$ corresponds to longitudinal polarization, and $A_1$ and $A_2$, respectively, to parallel and perpendicular transverse polarizations of the vector mesons; $R_\perp$ is defined as $|A_1|^2$. For a $\Delta I = 0$ transition, all $K^* \to K \pi$ channels involve the same amplitudes, and so the data for different decay modes can be combined.

The transversity frame is defined in the $J/\psi$ rest frame. The $K^*$ direction defines the negative $x$ axis. The $K \pi$ decay plane defines the $(x,y)$ plane, with $y$ such that $p_x (K) > 0$. The $z$ axis is the normal to this plane, and the coordinate system is right handed. The transversity angles $\theta_\parallel$ and $\phi_\parallel$ are defined as the polar and azimuthal angles of the positive lepton from the $J/\psi$ decay; $\theta_\perp$ is the $K^*$ helicity angle defined as the angle between the $K$ direction and the direction opposite the $J/\psi$ in the $K^*$ rest frame. The normalized angular distribution $g(\theta_\parallel, \phi_\parallel)$ is

$$
g = \frac{1}{\Gamma} \frac{d^3 \Gamma}{d \cos \theta_\parallel d \cos \theta_\perp d \phi_\parallel} = f_1 |A_0|^2 + f_2 |A_1|^2 + f_3 |A_2|^2 + f_4 \text{Im}(A_0^* A_1^*) + f_5 \text{Re}(A_0^* A_2^*) + f_6 \text{Im}(A_0^* A_2^*) ,
$$

with

$$f_1 = (9/32\pi) \times 2 \cos^2 \theta_K (1 - \sin^2 \theta_\parallel \cos^2 \phi_\parallel) ,
$$

$$f_2 = (9/32\pi) \times \sin^2 \theta_K (1 - \sin^2 \theta_\parallel \sin^2 \phi_\parallel) ,
$$

$$f_3 = (9/32\pi) \times \sin^2 \theta_K \sin^2 \theta_\parallel ,
$$

$$f_4 = (9/32\pi) \times \sin^2 \theta_K \sin^2 \theta_\parallel \sin \phi_\parallel \cdot \zeta ,
$$

$$f_5 = -(9/32\pi) \times (1/\sqrt{2}) \sin 2\theta_K \sin^2 \phi_\parallel \sin 2\phi_\parallel ,
$$

$$f_6 = (9/32\pi) \times (1/\sqrt{2}) \sin 2\theta_K \sin 2\phi_\parallel \cos \phi_\parallel \cdot \zeta .
$$

When the final state is not a CP eigenstate, $\zeta$ is +1 for $B^+$ and $B^0$, and −1 for $B^-$. For the CP mode $K^0_S \pi^0$, $\zeta (B^0) = -\zeta (\bar{B}^0) = 1/(1 + x_d^2)$, where $x_d = \Delta m_{B^0}/\Gamma_{B^0} \sim 0.73$; however, since flavor is not determined in the present analysis, $\zeta$ averages to zero for this mode. We define the relative phases of the amplitudes as $\phi_\parallel = \text{arg}(A_1/A_0)$ and $\phi_\perp = \text{arg}(A_2/A_0)$.

In this Letter, we present a measurement of the decay amplitudes in the decays $B^0 \to J/\psi K^{*0}$ and $B^+ \to J/\psi K^{*+}$, where the $K^{*0}$ and $K^{*+}$ are reconstructed in the modes $K^0 \pi^0$, $K^+ \pi^−$ and $K^0 \pi^+$, $K^+ \pi^0$, respectively [8]; only $J/\psi$ decays to $e^+ e^−$ and $\mu^+ \mu^−$ are considered. The data sample corresponds to 20.7 fb$^{-1}$ collected at the Y(4S) in 1999–2000 with the BABAR detector at the PEP-II asymmetric $B$ factory, and contains $\sim 22.7 \times 10^6$ $B$ meson pairs.

The BABAR detector is described elsewhere [9]. Charged particle track parameters are obtained from measurements in a 5-layer double-sided silicon vertex tracker and a 40-layer drift chamber located in a 1.5 T magnetic field; both devices provide $dE/dx$ information. Additional charged particle identification (PID) information is obtained from a detector of internally reflected Cherenkov (DIRC) light consisting of quartz bars that carry the light to a volume filled with water, and equipped with 10752 photomultiplier tubes. Electromagnetic showers are measured in a calorimeter (EMC) consisting of 6580 CsI(Tl) crystals. An instrumented flux return (IFR), containing multiple layers of resistive plate chambers, provides $\mu$ identification.

Electrons are identified by requiring that shower shape and energy deposition in the EMC be compatible with those expected for an electron of the measured momentum; $dE/dx$ measurements must also be compatible with the electron hypothesis. Muon candidates must penetrate at least two interaction lengths in the detector, and generate a small number of hits per layer in the IFR. If a muon candidate traverses the EMC, its energy deposition must be consistent with that of a minimum ionizing particle. Kaon candidates must survive a pion veto based on DIRC and $dE/dx$ information.

Charged tracks are required to be in regions of polar angle for which the PID efficiency is well measured. For electrons, muons, and kaons the acceptable ranges are 0.41 to 2.41 rad, 0.3 to 2.7 rad, and 0.45 to 2.5 rad, respectively. $J/\psi$ candidates consist of a pair of identified leptons that form a good vertex. The lepton pair invariant mass must be between 3.06 and 3.14 GeV/c$^2$ for electrons and 2.95 and 3.14 GeV/c$^2$ for muons. This corresponds to a $\pm 3\sigma$ interval for muons, and accounts for the radiative tail due to bremsstrahlung for electrons. $K^0_S$ candidates consist of vertexed pairs of oppositely charged tracks with invariant mass between 489 and 507 MeV/c$^2$. In the plane perpendicular to the beam line, the $K^0_S$ flight length must be greater than 1 mm, and its direction must form an angle with the $K^0_S$ momentum vector in this plane that is less than 2.9 rad.
than 0.2 rad. A photon is defined as a neutral cluster of
energy greater than 30 MeV in the EMC that agrees in
lateral shower shape with an electromagnetic shower. A
π^0 candidate consists of a pair of photons with invariant
mass in the interval 106–153 MeV/c^2. The J/ψ, K^0_S, and
π^0 are constrained to the corresponding nominal masses
[10]. K^+ candidates must have Kπ invariant mass within
100 MeV/c^2 of the nominal K*(892) mass [10].

B mesons are formed from J/ψ and K^+ candidates. For
B \to J/ψ(Kπ^0)^* \cos \theta_K, \cos \theta_B \text{ or } \phi, \text{ event } j, \text{ and } \epsilon(\hat{\omega}) \text{ is the efficiency}
at \hat{\omega}. Rewriting Eq. (1) as \( g = \sum_{i=1}^{6} f_i A_i \), where the
A_i (i = 1, \ldots, 6) \text{ represent } |A_0|^2, |A_1|^2, |A_2|^2, Im(A_1^* A_0), \text{ Re}(A_0^* A_1), \text{ and Im}(A_0^* A_2) \text{,}
the mean efficiency \( (\epsilon) = \int g \epsilon d\hat{\omega} = \sum_{i=1}^{6} \epsilon_i A_i \), where the \( \xi_i = \int f_i \epsilon d\hat{\omega} \) are constants. The signal part of the log
likelihood, \( \ln L_{\text{signal}} = \sum_{i=1}^{N_{\text{obs}}} \ln g_{\text{obs}}(\hat{\omega}_i) \), where \( N_{\text{obs}} \text{ is the number of observed events, becomes } \ln L_{\text{signal}} = \sum_{i=1}^{N_{\text{obs}}} \ln g(\hat{\omega}_i) + \sum_{i=1}^{N_{\text{obs}}} \ln \epsilon(\hat{\omega}_i) - N_{\text{obs}} \ln (\sum_{i=1}^{6} A_i \xi_i) \). Since the \( \epsilon(\hat{\omega}) \) are constants, the second term can be dis-
carded. Only the coefficients \( \xi_i \) are required, and detailed
representation of the acceptance is unnecessary [12].

The coefficients \( \xi_i \) are evaluated with a Monte Carlo
simulation. Separate sets of \( \xi_i \) are used for each channel,
and for \( \ell = e, \mu \). The values of \( \xi_i \) (i = 1, 2, 3) are close
to that of \( \langle \epsilon \rangle \); \( \xi_1 \) is always smallest, especially in channels
involving a π^0, because of the requirement on \( \cos \theta_K \).
The values of \( \xi_i \) (i = 4, 5, 6), which are related to the
interference terms, are compatible with zero.

The angular dependence of combinatorial background
events, \( g_{\text{obs}}(\hat{\omega}) \), is described by a pdf similar to that in Eq. (1)
with amplitudes \( B_i, i = 0, ||, \perp, \) and corresponding
terms \( B_i (i = 1, \ldots, 6) \).

The angular distribution of the (S)CF background is am-
plitude dependent. We correct for the effect of this back-
ground by evaluating the modified values \( \hat{\xi}_i \) of the \( \xi_i \) by

![FIG. 1. Beam-energy substituted mass spectra for the four Kπ modes. The curves are from fits using the \( G(m_{ES}) \) and \( F(m_{ES}) \) functions described in the text.](image)

### TABLE I. Fitted parameter values for the individual Kπ modes. The uncertainties are statistical only.

<table>
<thead>
<tr>
<th>Quantity</th>
<th>( K^0_S \pi^0 )</th>
<th>( K^+ \pi^- )</th>
<th>( K^0_S \pi^+ )</th>
<th>( K^+ \pi^0 )</th>
</tr>
</thead>
<tbody>
<tr>
<td>(</td>
<td>A_1</td>
<td>)</td>
<td>0.65 ± 0.13</td>
<td>0.60 ± 0.04</td>
</tr>
<tr>
<td>(</td>
<td>A_2</td>
<td>)</td>
<td>0.07 ± 0.11</td>
<td>0.17 ± 0.05</td>
</tr>
<tr>
<td>(</td>
<td>A_3</td>
<td>)</td>
<td>0.28 ± 0.14</td>
<td>0.23 ± 0.05</td>
</tr>
<tr>
<td>( \phi_{\perp} \text{ (rad)} )</td>
<td>···</td>
<td>−0.1 ± 0.2</td>
<td>0.0 ± 0.3</td>
<td>−0.4 ± 0.4</td>
</tr>
<tr>
<td>( \phi_{\parallel} \text{ (rad)} )</td>
<td>2.1 ± 0.7</td>
<td>2.5 ± 0.3</td>
<td>2.8 ± 0.4</td>
<td>2.6 ± 0.5</td>
</tr>
</tbody>
</table>
TABLE II. Fitted parameter values for the combined data samples. The first uncertainty is statistical, and the second is systematic. Note that \((\phi_{\perp}, \phi_{\parallel}) \rightarrow (\pi - \phi_{\perp}, -\phi_{\parallel})\) is also a solution.

<table>
<thead>
<tr>
<th>Quantity</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>(</td>
<td>A_0</td>
</tr>
<tr>
<td>(</td>
<td>A_1</td>
</tr>
<tr>
<td>(</td>
<td>A_2</td>
</tr>
<tr>
<td>(\phi_{\perp} (\text{rad}))</td>
<td>-0.17 ± 0.16 ± 0.07</td>
</tr>
<tr>
<td>(\phi_{\parallel} (\text{rad}))</td>
<td>2.50 ± 0.20 ± 0.08</td>
</tr>
</tbody>
</table>

including the (S)CF events, in the \(m_{ES}\) signal region, in addition to the signal [12]. In contrast to the \(\xi_i\), the \(\hat{\xi}_i\) depend on the amplitudes used in the simulation, but the maximum effect on the fitted amplitudes is found to be on the order of \(10^{-3}\). The complete log likelihood is

\[
\ln L = \sum_{j=1}^{N_{\text{obs}}} \ln [x G(m_{ES})g(\hat{\omega}_j) + (1 - x) F(m_{ES})g_B(\hat{\omega}_j)] - N_{\text{obs}} \ln \left(\sum_{j=1}^{N_{\text{obs}}} \hat{\xi}_i x |A_i| + (1 - x) |B_i|\right) - N, 
\]

where \(x\) is the fraction of signal integrated over the \(m_{ES}\) range 5.2–5.3 GeV/c². The normalization of \(g\) and \(g_B\) is relaxed in an extended likelihood approach [13], with convergence to the required condition \(a^2 = |A_0|^2 + |A_1|^2 + |A_2|^2 = 1\) imposed through the additional term \(N_{\text{obs}} a^2\) while \(|B_0|^2 + |B_1|^2 + |B_2|^2 = a^2\) holds by construction. The fit parameters are the mean and width of \(G(m_{ES})\); the shape parameter of \(F(m_{ES})\); the fraction \(x\); the signal amplitudes and phases \(|A_0|^2, |A_1|^2, |A_2|^2, \phi_{\perp}\), and \(\phi_{\parallel}\); and the corresponding background amplitudes and phases.

The agreement among the results for the individual decay channels is shown in Table I, while the fit result for the combined sample is summarized in Table II.

A partial representation of the fit is given by the one-dimensional projections of the angular distribution in Fig. 2. As a check of the fit quality, fits were performed to Monte Carlo samples with the angular distribution and number of events observed in the data. The maximum likelihood in the data is 1.35 standard deviations below the mean obtained from the Monte Carlo fits. The probability of obtaining a lower likelihood is 8.8%.

Systematic uncertainties are detailed in Table III. Limited simulation statistics (32 000 events per mode) give rise to a systematic uncertainty in the acceptance and (S)CF corrections (first row). Monte Carlo simulation has been used to estimate uncertainties due to the assumed form for the \(m_{ES}\) angular distributions of the background (second row). In particular, this accounts for any possible absorption of the S(CF) background by the \(F(m_{ES})\) function. The differences between simulated tracking and PID efficiencies and measurements obtained with control samples in the data lead to systematic uncertainties (third row) through their impact on acceptance corrections.

The \(K \pi \) \(S\)-wave systematic uncertainty (fourth row) is obtained as follows. Although the \(K \pi \) mass distribution for \(B \rightarrow J/\psi K \pi\) is dominated by the \(K^*\) [Fig. 3(a)], a significant number of candidates are at higher mass with
a clear peak at \( \sim 1.4 \text{ GeV}/c^2 \). The states in this region that couple strongly to \( \pi K \) are the \( K_0^*(1430) \) and the \( K_L^*(1430) \) [10]. Since it has width \( \sim 300 \text{ MeV}/c^2 \), the \( K_0^*(1430) \) alone would yield significantly more events above and below the peak than are observed. The \( K_L^*(1430) \) alone describes the high mass region but, when combined with the \( K^*(892) \) tail, yields too few events in the 1.1--1.3 \text{ GeV}/c^2 range [Fig. 3(b)]. This suggests a significant S-wave contribution, in which case the ratio \( J/\psi \) has a helicity angle distribution \( \sin^2(\theta_{J/\psi}) \). The observed behavior [Fig. 3(c)] agrees with this conjecture. This, together with the absence of the S-wave above 1.5 \text{ GeV}/c^2, is consistent with the mass dependence of the S-wave \( \pi K \) scattering amplitude [14]. If the \( \pi K \) S-wave in \( B \) decay behaves like this, a coherent S-wave amplitude should also be present in the \( K^*(892) \) region; \( S-P \) interference should occur, which, if ignored, can affect the P-wave amplitudes extracted from the data.

The effect of the S-wave in the \( K^*(892) \) region has been estimated by including a scalar term in the total amplitude. This yields a more complicated angular distribution \( g_S \), with ten \( f_i \) functions. A fit of \( g_S \) to the data in the 1.1--1.3 \text{ GeV}/c^2 region yields an S-wave fraction of \( (62 \pm 9)\% \), in agreement with the failure of a \( P- \) and D-wave fit to describe the mass spectrum. By repeating the analysis using \( g_S \), we find the S-wave contribution in the \( K^*(892) \) region to be \( (1.2 \pm 0.7)\% \). The differences in the P-wave results with and without S-wave are taken as estimates of systematic uncertainty (Table III, fourth row) since, with the present statistics, the presence of the S-wave in the \( K^*(892) \) region cannot be confirmed.

Table IV compares our results to those of CLEO [6] and CDF [7]. They are consistent, but the present measurement is significantly more precise. Longitudinal polarization is seen to dominate and the P-wave intensity is small. If \( \sin2\beta \) were measured in the \( B \rightarrow J/\psi K^0_\ell q^0 \) channel from the decay-time information only, the value of the dilution from the present measurement, \( D_\perp = 0.68 \pm 0.07 \), would contribute a 10% uncertainty.

Finally, we find that \( |\phi_\| | \) differs significantly from \( \pi \). This agrees with the CDF measurement, and indicates a departure from the factorization of the hadronic currents. In addition, there is evidence that S- and D-wave amplitude contributions are necessary for a description of the \( \pi K \) mass spectrum from \( B \rightarrow J/\psi K^* \) decay.

We thank Luis Oliver for many enlightening discussions. We are grateful for the excellent luminosity and machine conditions provided by our PEP-II colleagues. The collaborating institutions wish to thank SLAC for its support and kind hospitality. This work is supported by DOE and NSF (USA), NSERC (Canada), IHEP (China), CEA and CNRS-IN2P3 (France), BMBF (Germany), INFN (Italy), NFR (Norway), MIST (Russia), and PPARC (United Kingdom).

Individuals have received support from the Swiss NSF, A. P. Sloan Foundation, Research Corporation, and Alexander von Humboldt Foundation.

**Table IV.** Comparison with other experiments. Statistical and systematic uncertainties are added in quadrature.

| | \( |A_\parallel|^2 \) | \( |A_\perp|^2 \) | \( \phi_\perp \) | \( \phi_\parallel \) |
|---|---|---|---|---|
|CLEO [6] | 0.52 \( \pm \) 0.08 | 0.16 \( \pm \) 0.09 | \(-0.11 \pm 0.46\) | 3.00 \( \pm \) 0.37 |
|CDF [7] | 0.59 \( \pm \) 0.06 | 0.13 \( \pm \) 0.011 | \(-0.56 \pm 0.54\) | 2.16 \( \pm \) 0.47 |
|BABAR | 0.60 \( \pm \) 0.04 | 0.16 \( \pm \) 0.03 | \(-0.17 \pm 0.17\) | 2.50 \( \pm \) 0.22 |

*Also with: Università di Perugia, Perugia, Italy.
† Also with Università della Basilicata, Potenza, Italy.
[8] Throughout the paper, the use of charge conjugate modes is implied wherever relevant.