Fine structure of the lowest Landau level in suspended trilayer graphene

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Magnetotransport experiments on ABC-stacked suspended trilayer graphene reveal a complete splitting of the 12-fold degenerated lowest Landau level, and, in particular, the opening of an exchange-driven gap at the charge neutrality point. A quantitative analysis of distinctness of the quantum Hall plateaus as a function of field yields a hierarchy of the filling factors: $v = 6, 4, 0$ are the most pronounced, followed by $v = 3$, and finally $v = 1, 2, 5$. Apart from the appearance of a $v = 4$ state, which is probably caused by a layer asymmetry, this sequence is in agreement with Hund’s rules for ABC-stacked trilayer graphene.

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The unconventional quantum Hall effects observed in single-layer graphene (SLG), bilayer graphene (BLG), and trilayer graphene (TLG) are a hallmark for the relativistic mass enhancement and the Klein tunneling of quasiparticles. In particular, for trilayer graphene the two possible stacking sequences ABA and ABC lead to a Berry phase of $\pi$ and a zero-field mobility $\mu$ of the order of $10^{14} \text{m}^2/\text{Vs}$.

In a magnetic field, exchange effects and the Zeeman interaction split the Landau levels into Landau subbands. The Landau-level spectrum of trilayer graphene can be described by a 4N-fold degenerate zero-energy level, shared equally between electrons and holes, and fourfold degenerate higher Landau levels for electrons and holes separately. When taking more than only interlayer coupling into account, the situation becomes more complicated. In particular, for trilayer graphene the two possible stacking sequences ABA and ABC lead to a distinct level spectrum of the form $\nu = 1, 2, 4, 6$. In this Rapid Communication we present magnetotransport experiments on ABC-stacked suspended trilayer graphene, where the Berry phase is determined using the Hall resistance of the device before annealing.

In Fig. 1 we show the two-terminal conductance of the sample as a function of the magnetic field $B$. The conductance is centered around zero gate voltage. A quantitative analysis of the data yields a hierarchy of the filling factors: $v = 6, 4, 0$ are the most pronounced, followed by $v = 3$, and finally $v = 1, 2, 5$. Apart from the appearance of a $v = 4$ state, which is probably caused by a layer asymmetry, this sequence is in agreement with Hund’s rules for ABC-stacked trilayer graphene.

We have prepared a suspended TLG sample using an acid free method. Following standard techniques, we first exfoliated flakes from highly oriented pyrolytic graphite and deposited them on a Si/SiO$_2$ substrate covered with a 1.15 $\mu$m thick LOR-A resist (MicroChem Corp.) layer. The TLG flake was then identified by its thickness measured through its optical contrast. Subsequently, two electron beam lithography steps were performed in order to contact the flake with Ti-Au contacts and to remove part of the LOR-A below the graphene flake. The resulting device is a freely suspended bridge, 0.5 $\mu$m wide and 1.3 $\mu$m long, across a trench formed in the LOR-A with two metallic contacts on each side. Carriers in the sample are induced by applying a back-gate voltage $V_B$ on the highly n-doped Si wafer yielding a carrier concentration $n = n V_B$. The lever factor $\alpha \approx 1 \times 10^{14} \text{m}^2/\text{Vs}$ is determined experimentally from the position of the filling factors in Fig. 1 and agrees within a factor of 2 with that deduced from the geometric gate capacitance of the device before annealing.

Measurements were performed at low temperatures and high magnetic fields up to 30 T using a low-frequency lock-in technique with an excitation current $I \leq 1$ nA. The sample was mounted on an in situ tilting stage where the angle $\phi$ between the total magnetic field $B_{tot}$ and the perpendicular component $B_\perp = B \cos(\phi)$ can be controlled independently. $\phi$ was determined using the Hall resistance of a second sample on the same substrate. The device was slowly cooled down to 4.2 K and current annealed by applying a dc bias current up to 3 mA. The local annealing resulted in a high quality sample where the charge neutrality point (CNP) is centered around zero gate voltage.

In Fig. 1 we show the two-terminal conductance $G$ of our sample as a function of $V_B$ in a perpendicular magnetic field ($\phi = 0$). Before calculating the conductance, we have subtracted a constant background resistance of $550 \Omega$ originating from the finite contact and lead resistance from the measured two-terminal resistance. Using the slope of the dashed line in the figure, $G = n e \mu w / l$, we estimate a zero-field mobility $\mu \approx 8 \text{m}^2/\text{Vs}$ around the CNP. The sample was measured at $B < 3$ T, quantum Hall plateaus at filling factors $v = 4$ and $v = 6$ already start to develop. A further increase of the magnetic field up to 10 T results in the complete lifting of the lowest Landau level and the formation of quantized Hall plateaus at filling factors $v = 5, 3, 2, 1$. The conductance $G$ is the inverse of the resistance $R$, which is determined by a combination of the magnetoresistance...
FIG. 1. (Color online) Conductance traces at $T = 1.3$ K for magnetic fields between 0 and 30 T. The dashed line through the 0 T data is used to estimate the device mobility. The numbers indicate the quantization of $G$ in integer units of $e^2/h$. In this way, plateaus in $\nu$ field $B$, which is plotted in Fig. 2(a) for the measured data in Fig. 1. Motivated by the empiric relation $R_{xx} \propto B \times dR_{xx}/dB$, and in order to accentuate the plateaus more clearly we define therefore a normalized derivative,

$$D = -V_G \frac{dR}{dV_G},$$

(1)

which is plotted in Fig. 2(a) for the measured data in Fig. 1. In this way, plateaus in $R$, originating from $R_{xx}$, result in clear minima at integer filling factors $\nu = ne/hB$ that are related to Shubnikov–de Haas minima in $R_{xx}$. These minima are well pronounced on the electron side only. In a regime where two neighboring Landau levels are still overlapping, the Dingle factor at the oscillation minima scales as $R_D \propto \exp(-B_0/B)$. For higher fields, where the levels are fully separated, $R_D$ saturates and becomes field independent. Filling factors with the largest excitation gap appear first at the lowest $B$, while filling factors corresponding to smaller gaps appear at higher $B$. Quantitatively, using the above equation, we define an onset field $B_0$ where $37\% (1/e)$ of the maximum oscillation amplitude is reached.

In Fig. 2(b) we plot the position of the minima in $D$ as a function of gate voltage (proportional to $\nu$) for magnetic fields between 0 and 7 T. The lines indicate the expected linear behavior for the given magnetic fields between 0 and 7 T. Curves are shifted upwards to appear in the 5 T trace in Fig. 2(a); due to the limited visibility. Therefore, we will focus our analysis on the electron side only.

In Fig. 2(c) we plot the amplitude $A_\nu(B,T)$, a period $P$, and a phase $\psi$. Depending on the corresponding gap at a given filling factor $\nu$, the amplitude is different for different $\nu$. $A_\nu(B,T)$ contains a temperature dependent term $R_T$ and a field dependent Dingle term $R_D$. In order to concentrate on the field dependence alone, we have performed all measurements at a constant temperature $T = 1.3$ K, i.e., leaving $R_T$ constant for all measurements.

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In Fig. 2(c) we plot the amplitude $A_\nu$ [defined as the distance from the oscillation minimum to the average of the two neighboring maxima—see the arrow in Fig. 2(a)] as a function of $1/B$ for the different filling factors. The solid lines indicate the slope of the data points and determine the values of $B_0$ summarized in Table I. The $\nu = 5$ minimum just starts to appear in the 5 T trace in Fig. 2(a); due to the limited
gate voltage range we were not able to perform a quantitative analysis of its amplitude $A_0(B)$. As stated above, the values of $B_0$ in Table I describe a hierarchy of the filling factors. The most pronounced filling factor is found to be $v = 6$, separating the lowest 12-fold degenerate Landau level of ABC-stacked TLG from the first Landau level. It confirms that our sample is indeed TLG.\textsuperscript{5–8}

The subsequent filling factors are related to the lifting of the degeneracy of the lowest Landau level: first $v = 3$ and at higher fields $v = 1, 2, 5$, a sequence which is in agreement with Hund’s rules of ABC-TLG.\textsuperscript{28,29} However, filling factor $v = 4$ is also observed experimentally in this sequence whereas it is not predicted by Hund’s rules for ABC-TLG. We attributed its appearance to a layer asymmetry caused by an external electric field of the back gate or local inhomogeneities.\textsuperscript{13}

Additionally, a field-induced gap opens at the CNP, as theoretically expected for ABC-stacked TLG but not for ABA-stacked TLG.\textsuperscript{12,28,30} We analyze the nature of this gap in more detail by focusing on the diverging resistance at the CNP. Already at zero magnetic field a strong activated temperature dependence is observed [see Fig. 3(a)]. It can be described by $R_{\text{CNP}} \propto \exp(\Delta_0/k_BT)$, with a gap size $\Delta_0 \approx 0.38$ meV. In a magnetic field, $R_{\text{CNP}}$ grows rapidly with increasing $B/T$, suggesting an increase of the relevant gap. For $B/T$ up to 0.5 T/K the roughly exponential behavior of $R_{\text{CNP}}$ suggests an Arrhenius-activated transport with a field-enhanced gap $\Delta(B) = \Delta_0 + \gamma B$ with $\gamma = 1.1$ meV/T. This value is one order of magnitude larger than the bare Zeeman gap $\Delta = \mu_B B$ (0.116 meV at 1 T). We therefore suggest that an exchange-enhanced mechanism is responsible for the field-induced gap opening.

\begin{table}[h]
\centering
\caption{Field values $B_0$ characterizing the strength of the plateaus at $v = 6, 4, 3, 2,$ and 1.}
\begin{tabular}{ccc}
\hline
\textbf{$v$} & 6 & 4 & 3 & 2 & 1 \\
\hline
\textbf{$B_0$ (T)} & 1.4 ± 0.5 & 2.2 ± 0.5 & 4 ± 1 & 14 ± 2 & 14 ± 2 \\
\hline
\end{tabular}
\end{table}

The scenario of an exchange-driven gap at the CNP is further supported by experiments in tilted magnetic fields shown in Fig. 3(b). The resistance at the CNP decreases as a function of $B_\parallel/T$ in the roughly exponential behavior $R_{\text{CNP}} \propto \exp\left(-\Delta_0/k_BT\right)$. The solid lines show the expected resistance decrease, $R \propto \exp\left(-2\mu_B B/k_BT\right)$, scaling with the bare Zeeman energy. (b) Proposed scenario for the behavior in a parallel magnetic field: A gap $\Delta_0$ present at 0 T closes due to spin splitting in both energy levels. An exchange mechanism prevents a further decrease of the gap for parallel fields above 2 T, and a finite gap $\Delta_{\text{ex}} = 0.25$ meV remains.

In conclusion, magnetotransport experiments on a two-probe suspended ABC-stacked trilayer graphene sample show the full lifting of the 12-fold degeneracy in the lowest Landau level. Performing a quantitative analysis on the distinctness of the related quantum Hall plateaus, we have determined an order for the appearance of the corresponding filling factors: $v = 6, 4, 0$ appear first, followed by $v = 3$, and finally $v = 1, 2.5$. Furthermore, we have studied the opening of a gap at the CNP. Already at zero magnetic field we observe a gap $\Delta_0 = 0.38$ meV, which can be partly closed by the Zeeman effect in a parallel magnetic field. In contrast, in a perpendicular magnetic field, the gap at $v = 0$ was shown to increase linearly with field and to be an order of magnitude larger than the bare Zeeman gap. These facts point to a spin unpolarized ground state with an exchange-driven gap.

FIG. 3. (Color online) (a) $R_{\text{CNP}}$ in a magnetic field perpendicular to the two-dimensional electron system plotted as a function of $B/T$ for different temperatures. (b) $R_{\text{CNP}}$ at 1.3 K in tilted magnetic fields plotted as a function of the perpendicular field component $B_\perp$.

FIG. 4. (Color online) (a) Resistance at the CNP in a parallel magnetic field plotted as a function $B/T$. The solid lines show the expected resistance decrease, $R \propto \exp\left(-2\mu_B B/k_BT\right)$, scaling with the bare Zeeman energy. (b) Proposed scenario for the behavior in a parallel magnetic field: A gap $\Delta_0$ present at 0 T closes due to spin splitting in both energy levels. An exchange mechanism prevents a further decrease of the gap for parallel fields above 2 T, and a finite gap $\Delta_{\text{ex}} = 0.25$ meV remains.
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