The following full text is a publisher's version.

For additional information about this publication click this link.
http://hdl.handle.net/2066/112779

Please be advised that this information was generated on 2018-02-03 and may be subject to change.
High field study of the magnetic phase transition of $\text{URu}_2\text{Si}_2$

S.A.M. Mentink$^a$, U. Wyder$^b$, J.A.J. Perenboom$^b$, A. de Visser$^c$, A.A. Menovsky$^{c,d}$, G.J. Nieuwenhuys$^d$, J.A. Mydosh$^d$, T.E. Mason$^a$

$^a$Department of Physics, University of Toronto, Toronto, Ont., Canada M5S 1A7
$^b$Nijmegen High Magnetic Field Laboratory, Nijmegen, The Netherlands
$^c$Van der Waals–Zeeman Laboratory, University of Amsterdam, The Netherlands
$^d$Kamerlingh Onnes Laboratory, Leiden University, The Netherlands

Abstract

We investigate the phase transition of $\text{URu}_2\text{Si}_2$ at $T_0 = 17.5 \text{ K}$ by thermal expansion and heat capacity in magnetic fields, $H$, up to 25 T. At $T_0$ an energy gap of $\Delta \approx 115 \text{ K}$ opens, which decreases to 65 K in 25 T. The $H$ dependence of $T_0$ scales with $\Delta$, whereas a different $H$ dependence and energy scale was observed for the ordered dipole moment. The anomalous $H$ dependence of the thermal expansion coefficients may discriminate between the various scenarios for the primary order parameter of $\text{URu}_2\text{Si}_2$.

Keywords: $\text{URu}_2\text{Si}_2$; Antiferromagnetism; High magnetic field

The origin of the putative magnetic phase transition [1] in the heavy-fermion superconductor $\text{URu}_2\text{Si}_2$ remains elusive [2, 3]. Since the entropy released in the phase transition is very large compared to the antiferromagnetically ordered moments of $\mu_{\text{ord}} = 0.04 \mu_B$/U-atom [4], purely dipolar ordering can be ruled out. However, no definitive identification of the primary order parameter has yet been made. It has been suggested that the transition may involve multipolar ordering of uranium ions with a $5f^2$ configuration, with the implication that the small moment arises from an admixture of higher lying $5f^3$ or $5f^4$ states [5, 6]. Previously, we have demonstrated [7] by resistivity measurements up to 25 T that the field dependence of $T_0$ scales directly with the energy gap $\Delta$, which opens at $T_0$, following $\Delta/\Delta^{H=0} \propto T_0/T_0^{H=0} \propto 1 - (H/H_0)^2$. Here the critical fields for $T_0$ and $\Delta$, are equal: $\mu_0H_0 = 40 \pm 2 \text{ T}$. Neutron diffraction [8] indicated a substantial reduction (25%) of $\mu_{\text{ord}}$ with modest applied fields up to 8 T $||c$. The extrapolated critical field for the disappearance of $\mu_{\text{ord}}$ of $14.5 \pm 0.3 \text{ T}$ is much smaller than that for $T_0$ and $\Delta$, thereby suggesting that there are two different energy scales relevant to the magnetic phase transition of $\text{URu}_2\text{Si}_2$ [7]. Hence the application of high fields might be the appropriate experimental method towards discriminating between the various scenarios put forth to explain this transition.

Here we extend the field range over which the energy gap can be reliably determined up to 25 T by analyzing the exponential $T$ dependence of the thermal expansion, $\alpha$, along the main $a$ and $c$ axes. We also calculate the uniaxial pressure dependence of $T_0$ at the various applied fields by means of the Ehrenfest relation. The anomalous behavior observed above about 15 T is discussed to stimulate further work in this high-field region of the phase diagram of $\text{URu}_2\text{Si}_2$ [9].

* Corresponding author.
The experiments have been performed on annealed single crystalline URu$_2$Si$_2$. Thermal expansion up to 25 T was measured with a capacitive method, using RuO$_2$ and capacitance thermometry. The heat capacity was recorded up to 16 T using a pulse method. The magnetic field was only applied along the easy c axis. Fig. 1 displays $\alpha$ parallel to the $a$ and $c$ axes, recorded at fields of 0, 10, 14, 18 and 25 T. The $H = 0$ data are in good agreement with literature [10]. Cooling through $T_0$, the $a$ axis contracts, while the $c$ axis length increases. Interestingly, and unexpectedly for a regular magnetic ordering, the anomaly grows and remains very sharp upon increasing $H$. Below $T_0$, the data can be well described by an (dominating) exponential $T$ dependence, in addition to the linear and cubic terms for the electronic and phonon contributions:

$$\alpha = \alpha_0 + AT + BT^3 + C \exp \left(-\frac{\Delta}{T}\right).$$  \hspace{1cm} (1)

In Fig. 2 the gap as derived from the volume expansion, $\Delta V = \frac{1}{2}(2\alpha_a + \alpha_c)$, is shown versus $H$. The magnitude of $\Delta V = 0$ is the same as obtained from infrared reflectance, vacuum tunneling, nonlinear susceptibility, heat capacity and neutron scattering. $\Delta$ decreases quadratically with $H$, with a critical field $\mu_0H_0 = 37 \pm 2$ T, consistent with Ref. [7].

![Fig. 1. Thermal expansion of URu$_2$Si$_2$ along the $a$ and $c$ axes, with magnetic fields along the $c$ axis; $\mu_0H = 0$ ( ), 10 ( ), 14 ( ), 18 ( ), and 25 T ( ).](image)

![Fig. 2. The energy gap of URu$_2$Si$_2$ versus magnetic field parallel to the $c$ axis. The extrapolated critical field equals 37 ± 2 T.](image)

The $T_0(H)$ phase boundary found in Ref. [7] has been verified by a collection of magnetostriction curves [9]. From the ratio of the jumps at $T_0$ of thermal expansion and specific heat, $c$ (not shown), the initial uniaxial pressure dependence can be calculated through the Ehrenfest relation:

$$\left(\frac{\partial T_0}{\partial P_{a,c}}\right)_{B,T'} = \frac{V_m \Delta \alpha_{a,c}}{\Delta(c/T)}. \hspace{1cm} (2)$$

The $c/T$ jumps at $T_0$ vary slowly from 0.30 J/mole-K$^2$ at $H = 0$ to 0.33 J/mole-K$^2$ at 16 T [9]. By using $V_m = 4.95 \times 10^{-5}$ m$^3$/mole, and assuming that $\Delta(c/T)$ at 18 T and 25 T does not change significantly from the 16 T value, we can derive the pressure derivatives as shown in Fig. 3. The pressure dependences are almost $H$-independent up to about 15 T, but increase strongly at high fields. Note that the anisotropy, reflected in the directly measured $c/a$ ratio (not shown), changes in a similar way [9], which is a further indication that the above assumption on $c/T$ is reasonable. The question arises whether this anomalous high-field behavior is due to the diminishing gap, allowing for cross-gap excitations, or perhaps related to the vanishing ordered moment. The observed decrease of $\gamma = c/T$ at 2 K upon increasing $H$ [9] makes the first possibility unlikely.

We have presented the first thermal expansion experiments on URu$_2$Si$_2$ in magnetic fields up to
25 T. While confirming the equivalent field dependence of $T_0$ and $\Delta$ found in resistivity, the gap derived from $\alpha_T$ equals that found in tunneling and neutron scattering, and decreases by 45% at 25 T. The $c/a$ ratio and, possibly also the uniaxial pressure dependence of $T_0$, have qualitatively different behaviors in low fields ($\mu_0H \leq 15$ T) and high fields ($\mu_0H > 15$ T). Such behavior would be consistent with the presence of an extra energy scale of about 15 T, proposed to relate to the disappearance of the staggered moment. More experiments in this field range are necessary to substantiate such speculations and could eventually enable us to identify the primary order parameter of URu$_2$Si$_2$.

We acknowledge W.J.L. Buyers, M.B. Walker and H. Amitsuka for discussions. This work was partially supported by CIAR and NSERC of Canada, the Dutch Foundation FOM, and the European Large Installation Program.

References

[9] A full account of this work will be published elsewhere.