Search for direct production of charginos and neutralinos in events with three leptons and missing transverse momentum in $\sqrt{s} = 7$ TeV $pp$ collisions with the ATLAS detector

The ATLAS Collaboration

Abstract

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1. Introduction

 supersymmetry (SUSY)\cite{1,8} postulates the existence of SUSY particles, or “sparticles”, with spin differing by one-half unit with respect to that of their Standard Model (SM) partner. If $R$-parity\cite{10–14} is conserved, the lightest SUSY particle (LSP) is stable and sparticles can only be pair-produced and decay into final states with SM particles and LSPs. Charginos ($\tilde{\chi}^{\pm}_i$, $i = 1, 2$) and neutralinos ($\tilde{\chi}^{0}_j$, $j = 1, 2, 3, 4$) are the mass eigenstates formed from the linear superposition of the SUSY partners of the Higgs and electroweak gauge bosons. These are the Higgsinos, and the winos, zino, and bino, collectively known as gauginos. Naturalness requires $\tilde{\chi}^{\pm}_i$ and $\tilde{\chi}^{0}_j$ (and third-generation sparticles) to have masses in the hundreds of GeV range\cite{13,16}. In scenarios where squark and gluino masses are larger than a few TeV, the direct production of gauginos may be the dominant SUSY process at the Large Hadron Collider (LHC). Charginos can decay into leptonic final states via sneutrinos ($\tilde{\nu}_i$), sleptons ($\tilde{l}_i$) or $W$ bosons ($W\tilde{\chi}^{0}_j$), while unstable neutralinos can decay via sleptons ($\tilde{l}_i$) or $Z$ bosons ($Z\tilde{\chi}^{0}_j$).

This Letter presents a search with the ATLAS detector for the direct production of charginos and neutralinos decaying to a final state with three leptons (electrons or muons) and missing transverse momentum, the latter originating from the two undetected LSPs and the neutrinos. The analysis is based on $4.7\text{ fb}^{-1}$ of proton-proton collision data delivered by the LHC at a centre-of-mass energy $\sqrt{s} = 7\text{ TeV}$ between March and October 2011. The search described here significantly extends the current mass limits on charginos and neutralinos set by ATLAS\cite{14,18} and LEP\cite{21}, where a model-independent lower limit of 103.5 GeV was set at 95% confidence level (CL) on the mass of promptly decaying charginos.

2. Detector Description

 ATLAS\cite{22} is a multipurpose particle detector with forward-backward symmetric cylindrical geometry. It includes an inner tracker (ID) immersed in a 2 T magnetic field providing precision tracking of charged particles for pseudorapidities $|\eta| < 2.5$\cite{1}. Calorimeter systems with either liquid argon or scintillating tiles as the active media provide energy measurements over the range $|\eta| < 4.9$. The muon detectors are positioned outside the calorimeters and are contained in an air-core toroidal magnetic field produced by superconducting magnets with field integrals varying from 1 T·m to 8 T·m. They provide trigger and high-precision tracking capabilities for $|\eta| < 2.4$ and $|\eta| < 2.7$, respectively.

3. New Physics Scenarios

In this analysis, results are interpreted in the phenomenological minimal supersymmetric SM (pMSSM\cite{23}) and in simplified models\cite{24}.

In the pMSSM the mixing for the $\tilde{\chi}^{\pm}_i$ and $\tilde{\chi}^{0}_j$ depends on the gaugino masses $M_1$ and $M_2$, the Higgs mass parameter $\mu$, and $\tan\beta$, the ratio of the expectation values of the two Higgs doublets. The dominant mode for gaugino production leading to three-lepton final states is $\tilde{\chi}^{\pm}_1\tilde{\chi}^{0}_2$ production via the $s$-channel exchange of a virtual gauge boson. Other $\tilde{\chi}^{\pm}_1\tilde{\chi}^{0}_j$ processes contribute a maximum of

\footnote{ ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z-axis along the beam pipe. The x-axis points from the IP to the centre of the LHC ring, and the y-axis points upward. Cylindrical coordinates ($r, \phi$) are used in the transverse plane, $\phi$ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle $\theta$ as $\eta = -\ln(\tan(\theta/2))$.}
20% to three-lepton final states depending on the values of the mass parameters. The right-handed sleptons (including third-generation sleptons) are assumed to be degenerate and have a mass \( m_{\tilde{\ell}_R} = (m_{\tilde{\chi}^0_2} + m_{\tilde{\chi}^+_1})/2 \), set via the right-handed SUSY-breaking slepton mass parameter at the electroweak scale. In these scenarios, decays to sleptons are favoured. The parameter \( \tan\beta \) is set to 6, yielding comparable branching ratios into each slepton generation. The masses of the gluinos, squarks and left-handed sleptons are chosen to be larger than 2 TeV. In order to achieve maximum mixing in the top squark sector the \( \tilde{t}_1 \) is set to be bino-like. Two different scenarios are considered. In the first case, the \( \tilde{\chi}^+_1 \) and \( \tilde{\chi}^0_2 \) are pair-produced and decay via left-handed sleptons, including staus, and sneutrinos of mass \( m_{\tilde{\chi}} = m_{\tilde{\chi}^0_1} = (m_{\tilde{\chi}^0_2} + m_{\chi^+}/2) \) with a branching ratio of 50% each. In the second scenario, the \( \tilde{\chi}^+_1 \) and \( \tilde{\chi}^0_2 \) decay via \( W \), and \( Z \) bosons.

4. Monte Carlo simulation

Several Monte Carlo (MC) generators are used to simulate SM processes and new physics signals relevant for this analysis. \textsc{Sherpa} [39] is used to simulate diboson processes \( W+Z \) and \( Z+Z \). These include all diagrams leading to three leptons and one neutrino, and to four leptons, respectively, including internal conversions (virtual photons converting into lepton pairs). \textsc{Herwig} [28] is used for \( WW \), while \textsc{MadGraph} [27] is used for the \( t\bar{t}W \), \( t\bar{t}WW \), \( t\bar{t}Z \), \( W\gamma \) and \( Z\gamma \) processes. \textsc{Mcnlo} [31] is chosen for the simulation of single- and pair-production of top quarks, and \textsc{ALPGEN} [29] is used to simulate the phase space for diboson processes. Expected diboson yields are normalised using next-to-leading-order (NLO) QCD predictions obtained with \textsc{MC@NLO} [38, 39]. The top-quark pair-production contribution is normalised to approximate next-to-by-leading-order calculations (NNLO) [32] and the \( t\bar{t}W/W \) contributions are normalised to \( \text{NLO} [33, 34] \). The \( W\gamma \) and \( Z\gamma \) yields are normalised to be consistent with the ATLAS cross-section measurement [35]. The QCD NNLO \textsc{FeWZ} [36, 37] cross-sections are used for normalisation of the inclusive \( W+\text{light-flavour jets} \) and \( Z+\text{light-flavour jets} \). The ratio of the NNLO to \( \text{LO} \) cross-section is used to rescale the \( W+\text{heavy-flavour jets} \) and \( Z+\text{heavy-flavour jets} \) \( \text{LO} \) cross-sections.

The choice of the parton distribution functions (PDFs) depends on the generator. The \textsc{Cteq6l1} [38] PDFs are used with \textsc{MadGraph} and \textsc{ALPGEN}, and the \textsc{Cteq} [38] PDFs with \textsc{Mcnlo} and \textsc{Sherpa}. The \textsc{MrtSmcal} PDF set [40] is used for \textsc{Herwig}.

The pMSSM samples are produced with \textsc{Herwig} and the simplified model samples with \textsc{Herwig++} [41]. The yields of the SUSY samples are normalised to the NLO cross-sections obtained from \textsc{prospino} [12] using the PDF set \textsc{cteq6.6} with the renormalisation/factorisation scales set to the average of the relevant gaugino masses.

Fragmentation and hadronisation for the \textsc{Alpgen} and \textsc{Mcnlo} (\textsc{Madgraph}) samples are performed with \textsc{Herwig} (\textsc{Pythia} [42]), while for \textsc{Sherpa}, these are performed internally. \textsc{Jimmy} [44] is interfaced to \textsc{Herwig} for simulating the underlying event. For all MC samples, the propagation of particles through the ATLAS detector is modelled using \textsc{Geant4} [13, 19]. The effect of multiple proton-proton collisions from the same or different bunch crossings is incorporated into the simulation by overlaying additional minimum bias events onto hard-scatter events using \textsc{Pythia}. Simulated events are weighted to match the distribution of the number of interactions per bunch crossing observed in data.

5. Event Reconstruction and Preselection

The data sample was collected with an inclusive selection of single-lepton and double-lepton triggers. For single-lepton triggers, at least one reconstructed muon (electron) is requested to have transverse momentum \( p_T > 20 \text{ GeV} \) (25 GeV). For di-lepton triggers, at least two leptons are required to be present in the event with transverse energy or momentum above threshold.

The two muons are required to have \( p_T > 12 \text{ GeV} \) for di-muon triggers, and the two electrons to have \( E_T > 17 \text{ GeV} \) for di-electron triggers, while the thresholds for electron-muon triggers are \( E_T > 15 \text{ GeV} \) and \( p_T > 10 \text{ GeV} \). These thresholds are chosen such that the overall trigger efficiency is high, typically in excess of 90%, and independent of the transverse momentum of the triggerable objects within uncertainties.

Events recorded during normal running conditions are analysed if the primary vertex has five or more tracks associated to it. The primary vertex of an event is identified with the high-Sigma \( \Sigma p_T^2 \) of associated tracks.

Electrons must satisfy “tight” identification criteria [47] and fulfill \( |\eta| < 2.47 \) and \( E_T > 10 \text{ GeV} \), where \( E_T \) and \( |\eta| \) are determined from the calibrated clustered energy deposits in the electromagnetic calorimeter and the matched ID track respectively. Muons are reconstructed by combining tracks in the ID and tracks in the muon spectrometer [48]. Reconstructed muons are considered as candidates if they have transverse momentum \( p_T > 10 \text{ GeV} \) and \( |\eta| < 2.4 \).

“Tagged” leptons are electrons and muons, well separated from each other and from candidate jets. Events containing at least one tagged muon having transverse impact parameter with respect to the primary vertex \( |d_0| > 0.2 \text{ mm} \) or longitudinal impact parameter with respect to the primary vertex \( |z_0| > 1 \text{ mm} \) are rejected to suppress cosmic muon background. “Signal leptons” are tagged leptons for which the scalar sum of the transverse momenta of tracks within a cone of \( \Delta R = \sqrt{(\Delta \phi)^2 + (\Delta \eta)^2} = 0.2 \)
around the lepton candidate, and excluding the lepton candidate track itself, is less than 10% of the lepton $E_T$ for electrons and less than 1.8 GeV for muons. Tracks selected for the electron and muon isolation requirement, defined above, have $p_T > 1$ GeV and are associated to the primary vertex of the event. To suppress leptons originating from secondary vertices, the distance of closest approach of the lepton track to the primary vertex normalised to its uncertainty is required to be small, with $|d_0|/\sigma(d_0) < 6 (3)$ for electrons (muons).

Jets are reconstructed using the anti-$k_t$ algorithm [49] with a radius parameter of $R = 0.4$ using clustered energy deposits calibrated at the electromagnetic scale. The jet energy is corrected to account for the non-compensating nature of the calorimeter using correction factors parameterised as a function of the jet $E_T$ and $\eta$ [50]. The correction factors applied to jets have been obtained from simulation and have been tuned and validated using data. Jets considered in this analysis have $E_T > 20$ GeV, $|\eta| < 2.5$ and a fraction of the jet’s track transverse momenta that can be associated with the primary vertex greater than 0.75. Events containing jets failing the quality criteria described in Ref. [51] are rejected to suppress both SM and beam-induced background. Jets are identified as containing b-hadron decays, and thus called “b-tagged”, using a multivariate technique based on quantities such as the impact parameters of the tracks associated to a reconstructed secondary vertex. The b-tagging algorithm [51] correctly identifies b-quark jets in simulated top-quark decays with an efficiency of 60% and misidentifies jets containing light-flavour quarks and gluons with a rate of < 1%, for jets with $|\eta| < 2.5$ and jet $E_T > 20$ GeV.

The missing transverse momentum, $E_T^{\text{miss}}$, is the magnitude of the vector sum of the transverse momentum or transverse energy of all $p_T > 10$ GeV muons, $E_T > 20$ GeV electrons, $E_T > 20$ GeV jets, and calibrated calorimeter clusters with $|\eta| < 4.9$ not associated to these objects [52].

6. Signal Region Selection

Selected events must contain exactly three signal leptons. As R-parity conserving leptonic decays of $\tilde{\chi}_1^0$ yield same-flavour opposite-sign (SFOS) lepton pairs, the presence of at least one such pair is required. The invariant mass of any SFOS lepton pair must be above 20 GeV to suppress background from low-mass resonances and the missing transverse momentum must satisfy $E_T^{\text{miss}} > 75$ GeV.

Three signal regions are then defined: two “Z-depleted” regions (SR1a and SR1b), with no SFOS pairs having invariant mass within 10 GeV of the nominal Z-boson mass; and a “Z-enriched” one (SR2), where at least one SFOS pair has an invariant mass within 10 GeV of the Z-boson mass. Events in SR1a and SR1b are further required to contain no b-tagged jets to suppress contributions from b-jet-rich background processes, where a lepton could originate from the decay of a heavy-flavor quark. SR1b is designed to increase sensitivity to scenarios characterised by large mass splittings between the heavy gauginos and the LSP by requiring all three leptons to have $p_T > 30$ GeV. In both SR1b and SR2, the transverse mass variable $m_T$ must take values greater than 90 GeV, where $m_T$ is constructed using the $E_T^{\text{miss}}$ and the lepton not included in the lepton pair with invariant mass closest to the nominal Z-boson mass. The $m_T$ requirement is introduced to suppress background from $W/Z$ events. The SR1a/b regions target neutralino decays via intermediate sleptons or via off-shell Z bosons while SR2 targets decays via an on-shell Z boson. Table 1 summarises the selection requirements for the three signal regions.

7. Standard Model Background Estimation

7.1. Reducible Background Processes

Several SM processes contribute to the background in the signal regions. A “reducible” process has at least one “fake” object, that is either a lepton from a semileptonic decay of a heavy-flavour quark or an electron from an isolated photon conversion. The contribution from misidentified light-flavour quark or gluon jets is negligible in the signal regions. The reducible background includes single- and pair-production of top-quarks and WW or W/Z produced in association with jets or photons. The dominant component is the production of top quarks, with a contribution of 1% or less from $Z+$jets. The reducible background is estimated using a “matrix method” similar to that described in Ref. [53].

In this implementation of the matrix method, the signal lepton with the highest $p_T$ or $E_T$ is taken to be real, which is a valid assumption in 99% of the cases, based on simulation. The number of observed events with one or two fakes is then extracted from a system of linear equations relating the number of events with two additional signal or tagged candidates to the number of events with two additional candidates that are either real or fake. The coefficients of the linear equations are functions of the real-lepton identification efficiencies and of the fake-object misidentification probabilities.

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Table 1: The selection requirements for the three signal regions. The \(Z\)-veto (\(Z\)-requirement) rejects (selects) events with \(m_{3\text{LSP}}\) within 10 GeV of the \(Z\) mass (91.2 GeV). The \(m_T\) is calculated from the \(E_T^{\text{miss}}\) and the lepton not forming the best \(Z\) candidate.

<table>
<thead>
<tr>
<th>Selection</th>
<th>SR1a</th>
<th>SR1b</th>
<th>SR2</th>
</tr>
</thead>
<tbody>
<tr>
<td>Targeted Intermediate Decay</td>
<td>$l(\ell^\ast)$ or $Z^\ast$</td>
<td>$l(\ell^\ast)$ or $Z^\ast$</td>
<td>on-shell $Z$</td>
</tr>
<tr>
<td>(N) leptons (e, (\mu))</td>
<td>Exactly 3</td>
<td>Exactly 3</td>
<td>Exactly 3</td>
</tr>
<tr>
<td>Lepton charge, flavour</td>
<td>(Z)-veto</td>
<td>(Z)-veto</td>
<td>(Z)-requirement</td>
</tr>
<tr>
<td>(m_{3\text{LSP}})</td>
<td>&gt; 20 GeV</td>
<td>&gt; 75 GeV</td>
<td></td>
</tr>
<tr>
<td>(N) b-jets</td>
<td>0</td>
<td>0</td>
<td>any</td>
</tr>
<tr>
<td>(m_T)</td>
<td>&gt; 90 GeV</td>
<td>&gt; 90 GeV</td>
<td>&gt; 90 GeV</td>
</tr>
<tr>
<td>(p_T) all</td>
<td>&gt; 10 GeV</td>
<td>&gt; 10 GeV</td>
<td>&gt; 10 GeV</td>
</tr>
</tbody>
</table>
The identification efficiency is measured in data using lepton candidates from $Z \rightarrow \ell \ell$ decays. Misidentification probabilities for each relevant fake type (heavy flavour or conversion) and for each reducible background process, parameterised with the lepton $p_T$ and $\eta$, are obtained using simulated events with one signal and two tagged leptons. These misidentification probabilities are then corrected using the ratio (fake scale factor) of the misidentification probability in data to that in simulation obtained from dedicated control samples. For heavy-flavour fakes, the correction factor is measured in a $bb$-dominated control sample. This is defined by selecting events with only one $b$-tagged jet (containing a muon) and a tagged lepton, for which the fake rate is measured. The non-$bb$ background includes top-quark pair production and $W$ bosons produced in association with a $b$-quark. An $E_{\text{miss}}$ requirement of less than 40 GeV suppresses both the $t\bar{t}$ and the $W$ contamination, while requiring $m_t < 40$ GeV reduces the $W$ background. The remaining (small) background is subtracted from data using MC predictions. The fake scale factor for the conversion candidates is determined in a sample of photons radiated from a muon in $Z \rightarrow \mu\mu$ decays. These are selected by requiring $m_{\mu\mu}$ to lie within 10 GeV of the nominal $Z$-boson mass value. A weighted average misidentification probability is then calculated by weighting the corrected type- and process-dependent misidentification probabilities according to the relative contributions in a given signal or validation region, defined below.

### 7.2. Irreducible Background Processes

A background process is considered “irreducible” if it leads to events with three real and isolated leptons, referred to as “real” leptons below. Such processes include diboson ($WZ$ and $ZZ$) and $t\bar{t}W/Z$ production, where the gauge boson may be produced off-mass-shell. The $ZZ$ and $t\bar{t}W/Z$ contribution is determined using the corresponding MC samples, for which lepton and jet selection efficiencies are corrected to account for differences with respect to data.

The largest irreducible background, $WZ$, is determined using a semi-data-driven approach. The $WZ$ background is fit to data in a control region including events with exactly three leptons, one SFOS lepton pair, a $Z$ candidate, $E_{\text{miss}} < 50$ GeV, a $b$-veto, and $m_T > 40$ GeV. The $WZ$ purity in the control region is $\sim 80\%$. Non-$WZ$ backgrounds, both irreducible and reducible, are determined based on simulation or by using the matrix method and subtracted. A $WZ$ normalisation factor $1.25 \pm 0.12$ is obtained in the control region under a background-only hypothesis and used to estimate the $WZ$ background in the validation regions. To obtain the model-independent 95% CL upper limit on the new phenomena cross-section, a fit is performed simultaneously in the $WZ$ control region and in the signal region, with floating $WZ$ normalisation factor and a non-negative signal in the signal region only. This allows the propagation of the uncertainties on the normalisation factor. When setting limits on specific new physics scenarios, the potential signal contamination in the $WZ$ control region is accounted for in the simultaneous fit.

### 8. Background Model Validation

The background predictions have been tested in various validation regions. A region (VR1) dominated by Drell-Yan and $WZ$ events is selected by requiring three signal leptons, at least one SFOS lepton pair, $30 \text{ GeV} < E_{\text{miss}} < 75$ GeV, and a $Z$-boson veto. A reducible-background dominated region (VR2, where top-quark pair-production and decay to two real and one fake lepton is the main contribution) is built by requiring three signal leptons, $E_{\text{miss}} > 50$ GeV and by vetoing SFOS lepton pairs. Finally, a $WZ$-dominated region (VR3) is defined by selecting events with three signal leptons, at least one SFOS lepton pair, a $Z$ candidate, and $50 \text{ GeV} < E_{\text{miss}} < 75$ GeV. The data and predictions are in agreement within the quoted statistical and systematic uncertainties as shown in Table 2.

### 9. Systematic uncertainties

Several sources of systematic uncertainty are considered in the signal, control and validation regions. The systematic uncertainties affecting the simulation-based estimates (the yield of the irreducible background, the cross-section weighted misidentification probabilities, the signal yield) include the theoretical cross-section uncertainties due to renormalisation and factorisation scale and PDFs, the acceptance uncertainty due to PDFs, the uncertainty on the luminosity, the uncertainty due to the jet energy scale, jet energy resolution, lepton energy scale, lepton energy resolution, lepton efficiency, $b$-tagging efficiency, mistag probability, and the choice of MC generator. In SR1a, the total uncertainty on the irreducible background is 24%. This is dominated by the uncertainty on the efficiency of the signal region selection for the $WZ$ generator, determined by comparing the nominal yield with that obtained with the HERWIG generator and found to be 20%. The next largest uncertainties are the uncertainty due to

<table>
<thead>
<tr>
<th>Selection</th>
<th>VR1</th>
<th>VR2</th>
<th>VR3</th>
</tr>
</thead>
<tbody>
<tr>
<td>$t\bar{t}$</td>
<td>0.17$\pm$0.14</td>
<td>0.12$\pm$0.10</td>
<td>1.1$\pm$0.9</td>
</tr>
<tr>
<td>$t\bar{t}W$</td>
<td>0.6$\pm$0.5</td>
<td>0.7$\pm$0.5</td>
<td>0.10$\pm$0.08</td>
</tr>
<tr>
<td>$t\bar{t}WW$</td>
<td>0.017$\pm$0.014</td>
<td>0.022$\pm$0.017</td>
<td>0.0023$\pm$0.0019</td>
</tr>
<tr>
<td>$ZZ$</td>
<td>17$\pm$15</td>
<td>0.10$\pm$0.05</td>
<td>3.9$\pm$0.6</td>
</tr>
<tr>
<td>$WZ$</td>
<td>46$\pm$8</td>
<td>0.93$\pm$0.29</td>
<td>98$\pm$12</td>
</tr>
<tr>
<td>Reducible Bkg.</td>
<td>50$\pm$28</td>
<td>13$\pm$7</td>
<td>3.1$\pm$4.7</td>
</tr>
<tr>
<td>Total Bkg.</td>
<td>114$\pm$32</td>
<td>15$\pm$7</td>
<td>106$\pm$13</td>
</tr>
<tr>
<td>Data</td>
<td>126</td>
<td>18</td>
<td>109</td>
</tr>
</tbody>
</table>
the MC generator (16%) and that on the cross-sections (9%) of the non-WZ background. The MC generator uncertainty partially accounts for the cross-section uncertainty, leading to a slight overestimate of the overall uncertainty. All the remaining uncertainties on the irreducible background in this signal region range between 0.5 and 5%. The total uncertainty on the irreducible background in SR1b is slightly larger, at 25%, due to the limited number of simulated events. In SR2, the uncertainty on the irreducible background is 24%, with increased contributions from the jet energy scale and resolution and cross-section uncertainties.

The uncertainty on the reducible background includes the MC uncertainty on the weights for the misidentification probabilities from the sources listed above (up to 10%) and the uncertainty due to the dependence of the misidentification probability on $E_{\text{T}}^\text{miss}$ (0.6–15%). Also included in the uncertainty on the reducible background is the uncertainty on the fake scale factors (10–34%), and that due to the limited number of data events with three tagged leptons, of which at least one is a signal lepton (19–130%). The latter uncertainty is highest in SR2 where the reducible background is very low.

The total uncertainties on the signal yields are 10–20%, where the largest contribution is from the uncertainty on the cross-sections (7%). Signal cross-sections are calculated to NLO in the strong coupling constant using PROSPINO. An envelope of cross-section predictions is defined using the 68% CL ranges of the CTEQ6.6 [20] PDF sets, including the $\alpha_S$ uncertainty and the MSTW [53] PDF sets, together with variations of the factorisation and renormalisation scales by factors of two or one half. The nominal cross-section value is taken to be the midpoint of the envelope and the uncertainty assigned is half the full width of the envelope, following the PDF4LHC recommendations [54].

In all of the above, the value used for the uncertainty on the luminosity is 3.9% [57][58]. Correlations of systematic uncertainties between processes and regions are accounted for.

### 10. Results and Interpretation

The numbers of observed events and the prediction for SM backgrounds in SR1a, SR1b and SR2 are given in Table 2. Distributions of the $E_{\text{T}}^\text{miss}$ in SR1a and SR2 are presented in Fig. 1.

No significant excess of events is found in any of the three signal regions. Upper limits on the visible cross-section, defined as the production cross-section times acceptance times efficiency, of 3.0 fb in SR1a, 0.7 fb in SR1b and 2.0 fb in SR2 are placed at 95% CL with the modified frequentist CL$_{s}$ prescription [60]. All systematic uncertainties and their correlations are taken into account via nuisance parameters in a profile likelihood fit [60]. The corresponding expected limits are 3.0 fb, 0.8 fb and 2.0 fb, respectively.

### Table 3: Expected numbers of events from SM backgrounds and observed numbers of events in data, for 4.7 fb$^{-1}$, in signal regions SR1a, SR1b and SR2.

<table>
<thead>
<tr>
<th>Selection</th>
<th>SR1a</th>
<th>SR1b</th>
<th>SR2</th>
</tr>
</thead>
<tbody>
<tr>
<td>$t\bar{t}Z$</td>
<td>0.06±0.05</td>
<td>0.025±0.023</td>
<td>0.6±0.5</td>
</tr>
<tr>
<td>$t\bar{t}W$</td>
<td>0.36±0.29</td>
<td>0.10±0.08</td>
<td>0.09±0.08</td>
</tr>
<tr>
<td>$t\bar{t}WW$</td>
<td>0.010±0.008</td>
<td>0.0023±0.0019</td>
<td>0.004±0.004</td>
</tr>
<tr>
<td>$Z\bar{Z}$</td>
<td>0.67±0.21</td>
<td>0.09±0.08</td>
<td>0.34±0.17</td>
</tr>
<tr>
<td>$WZ$</td>
<td>15.5±2.9</td>
<td>1.05±0.28</td>
<td>9.3±2.1</td>
</tr>
<tr>
<td>Reducible Bkg.</td>
<td>10.5</td>
<td>0.35±0.34</td>
<td>0.5±0.5</td>
</tr>
<tr>
<td>Total Bkg.</td>
<td>25±6</td>
<td>1.6±0.5</td>
<td>10.9±2.4</td>
</tr>
<tr>
<td>Data</td>
<td>24</td>
<td>0</td>
<td>11</td>
</tr>
</tbody>
</table>

SR1a and SR1b provide the best sensitivity for the pMSSM scenarios; in particular SR1a (SR1b) targets scenarios with small (large) mass splitting between the heavy gauginos and the LSP. The limits are calculated using the signal region providing the best expected limit for each of the model points. The uncertainties on the signal cross-section are not included in the limit calculation but their impact on the observed limit is shown.

The main features in the exclusion limits shown in Fig. 2 as a function of the three parameters $M_1$, $M_2$ and $\mu$ can be explained in broad terms as follows. For a given value of $M_1$, for example $M_1 = 100$ GeV in Fig. 2(a), the production cross-section decreases as $M_2$ and $\mu$ increase, which explains why limits become less stringent when both $M_2$ and $\mu$ take high values. In general, the sensitivity is reduced in the region at low $M_2$ and high $\mu$, due to the small mass splitting between the $\tilde{\chi}^0_1$ and the $\tilde{\chi}^0_2$. When $\mu$ is greater than $M_1$ and $M_2$, which is true for example in the rightmost part of the exclusion plots for $M_1 = 100$ GeV (Fig. 2(a)) and $M_1 = 140$ GeV (Fig. 2(b)), the mass of the gauginos does not depend on $\mu$ and the sensitivity remains constant as a function of $\mu$. On the contrary, in a large section of the plane shown for $M_1 = 250$ GeV (Fig. 2(c)), the condition that $\mu$ should be greater than $M_1$ is not fulfilled and the resulting limits on the same plane become less stringent. Additionally, the reduced reach at high $M_2$ and low $\mu$ for $M_1 = 140$ GeV can be explained in terms of smaller cross-section values and smaller mass splittings in that section of the parameter space. The difference between expected and observed limits seen in the upper right corner of the $M_1 = 100$ GeV exclusion plot, where SR1b has the best sensitivity, is explained by the observed under-fluctuation in data with respect to SM predictions.
The value of \( \tan \beta \) does not have a significant impact on 
\[ \sigma(pp \rightarrow \tilde{\chi}^\pm_1 \tilde{\chi}^0_1) \times BR(\tilde{\chi}^\pm_1 \tilde{\chi}^0_1 \rightarrow l\nu \tilde{\chi}^0_2 \ell \tilde{\chi}^0_1) \],
which decreases by 10% if \( \tan \beta \) is raised from 6 to 10.

The results obtained in signal regions SR1a and SR1b are combined with results from the relevant signal region in the ATLAS two-lepton search (SR-m\(_{T2}\)) [6]. The fits are performed on the combined likelihood function from SR-m\(_{T2}\) with SR1a, and from SR-m\(_{T2}\) with SR1b. The combination yielding the highest expected sensitivity is selected for optimal exclusions in the pMSSM planes (Fig. 3). The uncertainties are profiled in the likelihood and correlations between channels and processes are taken into account. An improvement in the sensitivity for \( M_1 = 250 \text{ GeV} \) and small values of \( M_2 \) is seen when results from the three-lepton and the two-lepton analyses are combined.

Region SR1b provides the best sensitivity to the sim-
The colour coding is the same as that in Figure 2. The various limits are as described in Figure 2. Combination of the results from this search with those of the two-lepton pMSSM, 2+3 leptons. The results from this search with those of the two-lepton pMSSM, 2+3 leptons. The results from this search with those of the two-lepton pMSSM, 2+3 leptons. 

Figure 3: Observed and expected 95% CL limit contours for chargino and neutralino production in the pMSSM for $M_1 = 100$ GeV (a), $M_1 = 140$ GeV (b) and $M_1 = 250$ GeV (c). Contours from the combination of the results from this search with those of the two-lepton ATLAS search in [61]. The various limits are as described in Figure 2. The colour coding is the same as that in Figure 2.

Figure 4: Observed and expected 95% CL limit contours for chargino and neutralino production in the simplified model scenario with intermediate slepton decay (a) and intermediate gauge boson decay (b). The colour coding is the same as that in Figure 2. For scenarios with intermediate slepton decay (with no intermediate slepton decay) the interpretation is shown in Fig. 4 (a). In the simplified model with gauginos decaying via gauge bosons (Fig. 4 (b)). The signal region SR1a has the best sensitivity for small mass differences between the heavy and light neutralinos, while SR2 is sensitive to decays of $\tilde{\chi}_2^0$ into on-shell $Z$ bosons.

11. Summary

Results from a search for direct production of charginos and neutralinos in the final state with three leptons (elect-
trons or muons) and missing transverse momentum are reported. The analysis is based on 4.7 fb$^{-1}$ of proton-proton collision data delivered by the LHC at $\sqrt{s} =$ 7 TeV and collected by ATLAS. No significant excess of events is found in data. The null result is interpreted in the pMSSM and simplified models. For the pMSSM, an improvement in the sensitivity for $M_1 = 250$ GeV and small values of $M_2$ is seen when results from this analysis are combined with those from the corresponding two-lepton ATLAS search. For the simplified models with intermediate slepton decays, degenerate $\tilde{\chi}_1^0$ masses up to 500 GeV are excluded for large mass differences from the $\tilde{\chi}_1^0$. The analysis presented here also has sensitivity to direct gaugino production with decays via gauge bosons.

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Nevis Laboratory, Columbia University, Irvington NY, United States of America
Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark
AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland
The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland
Physics Department, Southern Methodist University, Dallas TX, United States of America
Physics Department, University of Texas at Dallas, Richardson TX, United States of America
DESY, Hamburg and Zeuthen, Germany
Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
Institut für Kern-und Teilchenphysik, Technical University Dresden, Dresden, Germany
Department of Physics, Duke University, Durham NC, United States of America
SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
INFN Laboratori Nazionali di Frascati, Frascati, Italy
Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
Section de Physique, Université de Genève, Geneva, Switzerland
INFN Sezione di Genova; (b)Dipartimento di Fisica, Università di Genova, Genova, Italy
E. Andronikashvili Institute of Physics, Tbilisi State University, Tbilisi; (b)High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
II Physikalisches Institut, Justus-Liebig-Universität Gießen, Gießen, Germany
SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
Laboratoire de Physique Subatomique et de Cosmologie, Université Joseph Fourier and CNRS/IN2P3 and Institut National Polytechnique de Grenoble, Grenoble, France
Department of Particle Physics and Cosmology, Harvard University, Cambridge MA, United States of America
Kochhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (b)Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (c)ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
Department of Physics, Indiana University, Bloomington IN, United States of America
Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
University of Iowa, Iowa City IA, United States of America
Department of Physics and Astronomy, Iowa State University, Ames IA, United States of America
Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
Graduate School of Science, Kobe University, Kobe, Japan
Faculty of Science, Kyoto University, Kyoto, Japan
Kyoto University of Education, Kyoto, Japan
Department of Physics, Kyushu University, Fukuoka, Japan
Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
Physics Department, Lancaster University, Lancaster, United Kingdom
INFN Sezione di Lecce; (b)Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
Department of Physics, Jozef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia
School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
Department of Physics, Royal Holloway University of London, Surrey, United Kingdom
Department of Physics and Astronomy, University College London, London, United Kingdom
Laboratoire de Physique Nucléaire et de Hautes Énergies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
Fysiska institutionen, Lunds universitet, Lund, Sweden
Departamento de Física Teórica C-15, Universidad Autonoma de Madrid, Madrid, Spain
Institut für Physik, Universität Mainz, Mainz, Germany
School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
Department of Physics, University of Massachusetts, Amherst MA, United States of America
Department of Physics, McGill University, Montreal QC, Canada
School of Physics, University of Melbourne, Victoria, Australia
Department of Physics, The University of Michigan, Ann Arbor MI, United States of America
Department of Physics and Astronomy, Michigan State University, East Lansing MI, United States of America
Department of Physics and Astronomy, University of Delaware, Newark DE, United States of America
B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus
National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus
Department of Physics, Massachusetts Institute of Technology, Cambridge MA, United States of America
Group of Particle Physics, University of Montreal, Montreal QC, Canada
P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia
Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia
Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
Nagasaki Institute of Applied Science, Nagasaki, Japan
Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
INFN Sezione di Napoli; Dipartimento di Scienze Fisiche, Università di Napoli, Napoli, Italy
Department of Physics and Astronomy, University of New Mexico, Albuquerque NM, United States of America
Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
Department of Physics, Northern Illinois University, DeKalb IL, United States of America
Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
Department of Physics, New York University, New York NY, United States of America
Ohio State University, Columbus OH, United States of America
Faculty of Science, Okayama University, Okayama, Japan
Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK, United States of America
Palacký University, RCPTM, Olomouc, Czech Republic
Center for High Energy Physics, University of Oregon, Eugene OR, United States of America
LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
Graduate School of Science, Osaka University, Osaka, Japan
Department of Physics, University of Oslo, Oslo, Norway
Department of Physics, Oxford University, Oxford, United Kingdom
INFN Sezione di Pavia; Dipartimento di Fisica, Università di Pavia, Pavia, Italy
Department of Physics, University of Pennsylvania, Philadelphia PA, United States of America
Petersburg Nuclear Physics Institute, Gatchina, Russia
INFN Sezione di Pisa; Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA, United States of America
Laboratorio de Instrumentacao e Fisica Experimental de Particulas - LIP, Lisboa, Portugal; Departamento de Fisica Teorica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain
Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic
Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic
Czech Technical University in Prague, Praha, Czech Republic
State Research Center Institute for High Energy Physics, Protvino, Russia
Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
Physics Department, University of Regina, Regina SK, Canada
Ritsumeikan University, Kusatsu, Shiga, Japan
INFN Sezione di Roma I; Dipartimento di Fisica, Università La Sapienza, Roma, Italy
INFN Sezione di Roma Tor Vergata; Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
INFN Sezione di Roma Tre; Dipartimento di Fisica, Università Roma Tre, Roma, Italy
Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca; Centre National de l’Energie des Sciences Techniques Nucléaires, Rabat; Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech; Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; Faculté des sciences, Université Mohammed V-Agdal, Rabat, Morocco
DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique), Gif-sur-Yvette, France
Santa Cruz Institute for Particle Physics, University of California, Santa Cruz, Santa Cruz CA, United States of America

Department of Physics, University of Washington, Seattle WA, United States of America

Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom

Department of Physics, Shinshu University, Nagano, Japan

Fachbereich Physik, Universität Siegen, Siegen, Germany

Department of Physics, Simon Fraser University, Burnaby BC, Canada

SLAC National Accelerator Laboratory, Stanford CA, United States of America

(a) Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic

(a) Department of Physics, University of Johannesburg, Johannesburg; (b) School of Physics, University of the Witwatersrand, Johannesburg, South Africa

(a) Department of Physics, Stockholm University; (b) The Oskar Klein Centre, Stockholm, Sweden

Physics Department, Royal Institute of Technology, Stockholm, Sweden

Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook NY, United States of America

Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom

School of Physics, University of Sydney, Sydney, Australia

Institute of Physics, Academia Sinica, Taipei, Taiwan

Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel

Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel

Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece

International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan

Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan

Department of Physics, Tokyo Institute of Technology, Tokyo, Japan

Department of Physics, University of Toronto, Toronto ON, Canada

(a) TRIUMF, Vancouver BC; (b) Department of Physics and Astronomy, York University, Toronto ON, Canada

Institute of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan

Science and Technology Center, Tufts University, Medford MA, United States of America

Centro de Investigaciones, Universidad Antonio Nariño, Bogota, Colombia

Department of Physics and Astronomy, University of California Irvine, Irvine CA, United States of America

(a) INFN Gruppo Collegato di Udine; (b) ICTP, Trieste; (c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy

Department of Physics, University of Illinois, Urbana IL, United States of America

Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain

Department of Physics, University of British Columbia, Vancouver BC, Canada

Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada

Department of Physics, University of Warwick, Coventry, United Kingdom

Waseda University, Tokyo, Japan

Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

Department of Physics, University of Wisconsin, Madison WI, United States of America

Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany

Department of Physics, Yale University, New Haven CT, United States of America

Yerevan Physics Institute, Yerevan, Armenia

Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France

(a) Also at Laboratorio de Instrumentacao e Fisica Experimental de Particulas - LIP, Lisboa, Portugal

(b) Also at Faculdade de Ciencias and CFNUL, Universidade de Lisboa, Lisboa, Portugal

(c) Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom

(d) Also at TRIUMF, Vancouver BC, Canada

(e) Also at Department of Physics, California State University, Fresno CA, United States of America

(f) Also at Novosibirsk State University, Novosibirsk, Russia
Also at Fermilab, Batavia IL, United States of America
Also at Department of Physics, University of Coimbra, Coimbra, Portugal
Also at Department of Physics, UASLP, San Luis Potosi, Mexico
Also at Università di Napoli Parthenope, Napoli, Italy
Also at Institute of Particle Physics (IPP), Canada
Also at Department of Physics, Middle East Technical University, Ankara, Turkey
Also at Louisiana Tech University, Ruston LA, United States of America
Also at Dep Fisica and CEFITEC of Faculdade de Ciencias e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal
Also at Department of Physics and Astronomy, University College London, London, United Kingdom
Also at Group of Particle Physics, University of Montreal, Montreal QC, Canada
Also at Department of Physics, University of Cape Town, Cape Town, South Africa
Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany
Also at Manhattan College, New York NY, United States of America
Also at School of Physics, Shandong University, Shandong, China
Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
Also at School of Physics and Engineering, Sun Yat-sen University, Guanzhou, China
Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan
Also at Dipartimento di Fisica, Università La Sapienza, Roma, Italy
Also at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat a l’Energie Atomique), Gif-sur-Yvette, France
Also at Section de Physique, Université de Genève, Geneva, Switzerland
Also at Departamento de Física, Universidade de Minho, Braga, Portugal
Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, United States of America
Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary
Also at California Institute of Technology, Pasadena CA, United States of America
Also at Institute of Physics, Jagiellonian University, Krakow, Poland
Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
Also at Nevis Laboratory, Columbia University, Irvington NY, United States of America
Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
Also at Department of Physics, Oxford University, Oxford, United Kingdom
Also at Institute of Physics, Academia Sinica, Taipei, Taiwan
Also at Department of Physics, The University of Michigan, Ann Arbor MI, United States of America
Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa
* Deceased