Search for Dark Matter Candidates and Large Extra Dimensions in events with a photon and missing transverse momentum in $pp$ collision data at $\sqrt{s} = 7$ TeV with the ATLAS detector

The ATLAS Collaboration

Abstract

Results of a search for new phenomena in events with an energetic photon and large missing transverse momentum in proton-proton collisions at $\sqrt{s} = 7$ TeV are reported. Data collected by the ATLAS experiment at the LHC corresponding to an integrated luminosity of 4.6 fb$^{-1}$ are used. Good agreement is observed between the data and the Standard Model predictions. The results are translated into exclusion limits on models with large extra spatial dimensions and on pair production of weakly interacting dark matter candidates.
Search for Dark Matter Candidates and Large Extra Dimensions in events with a photon and missing transverse momentum in pp collision data at \( \sqrt{s} = 7 \) TeV with the ATLAS detector

The ATLAS Collaboration

Results of a search for new phenomena in events with an energetic photon and large missing transverse momentum in proton-proton collisions at \( \sqrt{s} = 7 \) TeV are reported. Data collected by the ATLAS experiment at the LHC corresponding to an integrated luminosity of 4.6 fb\(^{-1}\) are used. Good agreement is observed between the data and the Standard Model predictions. The results are translated into exclusion limits on models with large extra spatial dimensions and on pair production of weakly interacting dark matter candidates.

Events with an energetic photon and large missing momentum in the final state constitute a clean and distinctive signature in searches for new physics at colliders. In particular, monophoton and monojet final states have been studied \([1-8]\) in the context of searches for supersymmetry and large extra spatial dimensions (LED), aiming to provide a solution to the mass hierarchy problem, and the search for weakly interacting massive particles (WIMPs) as candidates for dark matter (DM).

The Arkani-Hamed, Dimopoulos, and Dvali (ADD) model for LED \([9]\) explains the large difference between the electroweak unification scale \( O(10^2) \) GeV and the Planck scale \( M_{Pl} \sim O(10^{19}) \) GeV by postulating the presence of \( n \) extra spatial dimensions of size \( R \), and defining a fundamental Planck scale in \( 4 + n \) dimensions, \( M_D \), given by \( M_{Pl}^2 \sim M_D^{2+n} R^n \). The extra spatial dimensions are compactified, resulting in a Kaluza-Klein tower of massive graviton modes. At hadron colliders, these graviton modes may escape detection and can be produced in association with an energetic photon or a jet, leading to a monophoton or monojet signature.

The presence of a non-baryonic DM component in the universe is inferred from the observation of its gravitational interactions \([10]\), although its nature is otherwise unknown. A WIMP \( \chi \) with mass \( m_{\chi} \) in the range between 1 GeV and a few TeV is a plausible candidate for DM. It could be detected via its scattering with heavy nuclei \([11]\), the detection of cosmic rays (energetic photons, electrons, positrons, protons, antiprotons, or neutrinos) from \( \chi \chi \) annihilation in astrophysical sources \([12]\), or via \( \chi \bar{\chi} \) pair-production at colliders where the WIMPs do not interact with the detector and the event is identified by the presence of an energetic photon or jet from initial-state radiation. The interaction of WIMPs with Standard Model (SM) particles is assumed to be driven by a mediator with mass at the TeV scale and described using a non-renormalizable effective theory \([12]\) with several operators. The vertex coupling is suppressed by an effective cut-off mass scale \( M_* \sim M/\sqrt{g_1 g_2} \), where \( M \) denotes the mass of the mediator and \( g_1 \) and \( g_2 \) are the couplings of the mediator to the WIMP and SM particles.

This Letter reports results of the search for new phenomena in monophoton final states, based on \( \sqrt{s} = 7 \) TeV proton-proton collision data corresponding to an integrated luminosity of 4.6 fb\(^{-1}\) collected with the ATLAS detector at the LHC during 2011. The ATLAS detector is described in detail elsewhere \([13]\). The data are collected using a three-level trigger system that selects events with missing transverse momentum greater than 70 GeV. In the analysis, events are required to have \( E_T^{\text{miss}} > 150 \) GeV, where \( E_T^{\text{miss}} \) is computed as the magnitude of the vector sum of the transverse momentum of all noise-suppressed calorimeter topological clusters with \( |\eta| < 4.9 \) \([14, 15]\). A photon is also required with transverse momentum \( p_T > 150 \) GeV and \( |\eta| < 2.37 \), excluding the calorimeter barrel/end-cap transition regions \( 1.37 < |\eta| < 1.52 \) \([13]\). With these criteria, the trigger selection is more than 98% efficient, as determined using events selected with a muon trigger. The cluster energies are corrected for the different response of the calorimeters to hadronic jets, \( \tau \) leptons, electrons or photons, as well as dead material and out-of-cluster energy losses. The photon candidate must pass tight identification criteria \([14]\) and is required to be isolated: the energy not associated with the photon cluster in a cone of radius \( \Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2} = 0.4 \) around the candidate is required to be less than 5 GeV. Jets are defined using the anti-\( k_t \) jet algorithm \([17]\) with the distance parameter set to \( R = 0.4 \). The measured jet \( p_T \) is corrected for detector effects and for contributions from multiple proton-proton interactions per beam bunch crossing (pileup) \([18]\).

Events with more than one jet with \( p_T > 30 \) GeV and \( |\eta| < 4.5 \) are rejected. Events with one jet are retained to increase the signal acceptance and reduce systematic uncertainties related to the modeling of initial-state radiation. The reconstructed photon, \( E_T^{\text{miss}} \) vector and jets (if found) are required to be well separated in the transverse plane with \( \Delta\phi(\gamma, E_T^{\text{miss}}) > 0.4 \), \( \Delta R(\gamma, \text{jet}) > 0.4 \), and \( \Delta \phi(\text{jet}, E_T^{\text{miss}}) > 0.4 \). Additional quality criteria \([19]\) are applied to ensure that jets and photons are not produced by noisy calorimeter cells, and to avoid problematic detector regions. Events with identified electrons

PACS numbers: 14.70.Kv, 13.85.Rm, 13.85.Qk, 14.80.Rt, 14.70.Bh
or muons are vetoed to reject mainly W/Z+jets and W/Z + γ background processes with charged leptons in the final state. Electron (muon) candidates are required to have $p_T > 20$ GeV and $|\eta| < 2.47$ ($p_T > 10$ GeV and $|\eta| < 2.4$), and to pass the medium (combined) criteria \cite{20}. The final data sample contains 116 events, where 88 and 28 events have zero and one jet, respectively.

The SM background to the monophoton signal is dominated by the irreducible $Z(\rightarrow \nu \bar{\nu}) + \gamma$ process, and receives contributions from $W/Z + \gamma$ events with unidentified electrons, muons or hadronic $\tau$ decays, and $W/Z + \gamma$ events with an electron or jet misconstructed as a photon. In addition, the monophoton sample receives small contributions from top-quark, $\gamma \gamma$, diboson ($WW, ZZ, WZ$), $\gamma + $jets, and multi-jet processes.

Background samples of simulated $W/Z + \gamma$ events are generated using ALPGEN \cite{21}, interfaced to HERWIG \cite{22} with JIMMY \cite{23}, and SHERPA \cite{24}, using CTEQ6L1 \cite{25} parton distribution functions (PDFs) and requiring a minimum photon $p_T$ of 40 GeV. Background samples of $W/Z + \gamma$ and $\gamma + $jets are generated using ALPGEN plus HERWIG/JIMMY, with CTEQ6L1 PDFs. Top-quark production samples are generated using MC@NLO \cite{26} and CT10 \cite{27} PDFs, while diboson processes are generated using HERWIG/JIMMY normalized to next-to-leading-order (NLO) predictions with MRST2007 \cite{28} PDFs. Multi-jet and $\gamma \gamma$ processes are generated using PYTHIA 6 \cite{29} with MRST2007 PDFs.

Signal Monte Carlo (MC) samples are generated according to the ADD model using the PYTHIA 8 leading-order (LO) perturbative QCD (pQCD) implementation with default settings, requiring a minimum photon $p_T$ of 80 GeV, and an ATLAS tune for the underlying event (UE) contribution \cite{30} including the CTEQ6L1 PDFs. The number of extra dimensions $n$ is varied from 2 to 6 and values of $M_D$ in the $1 - 2$ TeV range are considered. For consistency with a previous monojet analysis performed in ATLAS \cite{31, 32}, the yields corresponding to CTEQ6.6 \cite{33} PDFs are used, as obtained by reweighting these samples. The samples are normalized to NLO predictions \cite{34}. The LO-to-NLO normalization factors decrease from 1.5 to 1.1 as $n$ increases.

Simulated events corresponding to the $\chi \bar{\chi} + \gamma$ process with a minimum photon $p_T$ of 80 GeV are generated using LO matrix elements from MADGRAPH \cite{35} interfaced to PYTHIA 6 using CTEQ6L1 PDFs. Values for $m_\chi$ between 1 GeV and 1.3 TeV are considered. In this analysis, WIMPs are assumed to be Dirac fermions and the vertex operator is taken to have the structure of a scalar, vector, axial-vector or tensor, corresponding respectively to the operators D1, D5, D8 and D9 in Refs. \cite{12, 36}. These operators correspond to spin-independent (D1 and D5) and spin-dependent (D8 and D9) interactions. The MC samples are passed through a full simulation \cite{35} of the ATLAS detector and trigger system, based on GEANT4 \cite{37}. The simulated events are reconstructed and analyzed with the same analysis chain as the data.

The normalization of the MC predictions for the dominant $W/Z + \gamma$ background processes are set using scale factors determined in a data control sample, resulting in a significant reduction of the background uncertainties. A $\gamma + \mu + E_T^{\text{miss}}$ control sample with an identified muon is defined by inverting the muon veto in the nominal event selection criteria discussed above. According to the simulation, the sample contains a 71% (19%) contribution from $W + \gamma + \mu$ ($Z + \gamma$) processes. This control sample is used to normalize separately the $W + \gamma$ and $Z + \gamma$ MC predictions determined by ALPGEN and SHERPA, respectively. In each case, the scale factor is defined as the ratio of the data to the given MC prediction, after the contributions from the rest of the background processes are subtracted. The scale factors, extracted simultaneously to take into account correlations, are $\kappa(W + \gamma) = 1.0 \pm 0.2$ and $\kappa(Z + \gamma) = 1.1 \pm 0.2$, where statistical and systematic uncertainties are included.

Dedicated studies are performed to determine the probability for electrons or jets to be identified as photons, resulting in data-driven estimates of $W/Z + $jet background contributions. A data sample of Z boson candidates is employed to compute the fraction of electrons from the Z boson decay that are reconstructed as photons. This fraction decreases from 2% to 1% as $p_T$ increases from 150 GeV to 300 GeV, and increases from 1% to 3% as $|\eta|$ increases. These rates are employed to determine the $W(\rightarrow e\nu) + $jets background in the signal region, for which a control data sample selected with the nominal selection criteria and an electron instead of a photon is used. This results in a total $W(\rightarrow e\nu) + $jet background estimation of 14 ± 6 events, where the uncertainty is dominated by the limited size of the control data sample. Control samples enhanced in jets identified as photons are defined using nominal selection criteria with non-isolated photon candidates and/or photon candidates passing a loose selection \cite{16} but not the nominal identification requirements. The ratio of isolated to non-isolated photons passing the nominal identification requirements are used to determine the rate of jets identified as photons in the signal region, after the contribution from $W/Z + \gamma$ processes has been subtracted. This gives an estimate of 4.3 ± 1.9 $W/Z + $jet background events.

The $\gamma + $jet and multi-jet background contributions to the signature of a photon and large $E_T^{\text{miss}}$ originate from the misreconstruction of the energy of a jet in the calorimeter. The direction of $E_T^{\text{miss}}$ vector therefore tends to be aligned with the jet. These background contributions are determined from data using a control sample with the nominal selection criteria and at least one jet with $p_T > 30$ GeV and $\Delta \phi(\text{jet}, E_T^{\text{miss}}) < 0.4$. After the subtraction of electroweak boson and top-quark production processes, a linear extrapolation of the measured
A detailed study of systematic uncertainties on the background predictions has been performed. An uncertainty of 0.3% to 1.5% on the absolute photon energy scale, depending on the photon $p_T$ and $\eta$, translates into a 0.9% uncertainty on the total background prediction. Uncertainties on the simulated photon energy resolution, photon isolation, and photon identification efficiency introduce a combined 1.1% uncertainty on the background yield. Uncertainties on the simulated lepton identification efficiencies introduce a 0.3% uncertainty on the background predictions. The uncertainty on the absolute jet energy scale and jet energy resolution introduce 0.9% and 1.2% uncertainties on the background estimation, respectively. A 10% uncertainty on the absolute energy scale for low $p_T$ jets and unclustered energy in the calorimeter, and a 6.6% uncertainty on the subtraction of pileup contributions, are taken into account. They affect the $E_T^{miss}$ determination and translate into 0.8% and 0.3% uncertainties on the background yield, respectively.

The dependence of the predicted $W/Z + \gamma$ backgrounds on the lepton selection criteria is 20%.

Uncertainties due to the choice of PDFs and the variation of the renormalization and factorization scales in the $W/Z$ backgrounds 137$^{+18}_{-17}$ depending on the photon $p_T$ increases. Variations from top-quark, $\gamma\gamma$, and diboson production processes, determined using MC samples, are small. Finally, non-collision backgrounds are negligible.

Typical event selection efficiencies of $\epsilon \sim 75\%$ are found in simulated ADD and WIMP signal samples.

<table>
<thead>
<tr>
<th>Background source</th>
<th>Prediction ± (stat.) ± (syst.)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$Z(\rightarrow \nu\nu) + \gamma$</td>
<td>93 ± 16 ± 8</td>
</tr>
<tr>
<td>$Z/\gamma^*(\rightarrow \ell^+\ell^-) + \gamma$</td>
<td>0.4 ± 0.2 ± 0.1</td>
</tr>
<tr>
<td>$W(\rightarrow \ell\nu) + \gamma$</td>
<td>24 ± 5 ± 2</td>
</tr>
<tr>
<td>$W/Z + jets$</td>
<td>18 − ± 6</td>
</tr>
<tr>
<td>Top</td>
<td>0.07 ± 0.07 ± 0.01</td>
</tr>
<tr>
<td>$WW, WZ, ZZ, \gamma\gamma$</td>
<td>0.3 ± 0.1 ± 0.1</td>
</tr>
<tr>
<td>$\gamma$+jets and multi-jet</td>
<td>1.0 − ± 0.5</td>
</tr>
<tr>
<td>Total background</td>
<td>137 ± 18 ± 9</td>
</tr>
</tbody>
</table>

| Events in data (4.6 fb$^{-1}$) | 116 |

TABLE I: The number of events in data compared to the SM predictions, including statistical and systematic uncertainties. In the case of $W/Z$ + jets, $\gamma$+jets and multi-jet processes a global uncertainty is quoted.

FIG. 1: The measured $E_T^{miss}$ distribution (black dots) compared to the SM (solid lines), SM+ADD (dashed lines), and SM+WIMP (dotted lines) predictions, for two particular ADD and WIMP scenarios.

The results are translated into 95% CL limits on the parameters of the ADD model. The typical $A \times \epsilon$ of the selection criteria is $20.0 \pm 0.4$(stat.) $\pm 1.6$(syst.)%, approximately independent of $n$ and $M_D$. Experimental uncertainties related to the photon, jet and $E_T^{miss}$ scales and resolutions, the photon reconstruction, the trigger efficiency, the pileup description, and the luminosity introduce a 6.8% uncertainty on the signal yield. Uncertainties related to the modeling of the initial- and final-state gluon radiation translate into a 3.5% uncertainty on the ADD signal yield. Systematic uncertainties due to PDFs result in a 0.8% to 1.4% uncertainty on the signal $A \times \epsilon$ and a 4% to 11% uncertainty on the signal cross section, increasing as $n$ increases. Variations of the renormalization and factorization scales by factors of two and one-half introduce a 0.6% uncertainty on the signal $A \times \epsilon$ and an uncertainty on the signal cross section that decreases from 9% to 5% as $n$ increases.
ure 2 shows the expected and observed 95% CL lower limits on $M_D$ as a function of $n$, as determined using the $CL_s$ method and considering uncertainties on both signal and SM background predictions. Values of $M_D$ below 1.93 TeV ($n = 2$), 1.83 TeV ($n = 3$ or 4), 1.86 TeV ($n = 5$), and 1.89 TeV ($n = 6$) are excluded at 95% CL. The observed limits decrease by 3% to 2% after considering the $−1\sigma$ uncertainty from PDFs, scale variations, and parton shower modeling in the ADD theoretical predictions (dashed lines in Figure 2). These results improve upon previous limits on $M_D$ from LEP and Tevatron experiments [1, 3]. In this analysis, no weights are applied for signal events in the phase space region with $\hat{s} > M_D^2$, which is sensitive to the unknown ultraviolet behavior of the theory. For $M_D$ values close to the observed limits, the visible signal cross sections decrease by 15% to 75% as $n$ increases when truncated samples with $\hat{s} < M_D^2$ are considered. This analysis probes a kinematic range for which the model predictions are defined but ambiguous.

Similarly, 90% CL upper limits on the pair production cross section of dark matter WIMP candidates are obtained. The $A \times \epsilon$ of the selection criteria are typically $11.0 \pm 0.2$(stat.) $\pm 1.6$(syst.$)$% for the D1 operator, $18.0 \pm 0.3$(stat.) $\pm 1.4$(syst.$)$% for the D5 and D8 operators, and $23.0 \pm 0.3$(stat.) $\pm 2.1$(syst.$)$% for the D9 operator, with a moderate dependence on $m_\chi$. Experimental uncertainties, as discussed above, translate into a 6.6% uncertainty on the signal yields. Theoretical uncertainties on initial- and final-state gluon radiation introduce a 3.5% to 10% uncertainty on the signal yields. The uncertainties related to PDFs result in 1.0% to 8.0% and 5.0% to 30% uncertainties on the signal $A \times \epsilon$ and cross section, respectively. Variations of the renormalization and factorization scales lead to a change of 1.0% to 2.0% and 8.0% in the signal $A \times \epsilon$ and cross section, respectively. In the case of the D1 (D5) spin-independent operator, values of $M_\chi$ below 31 GeV and 5 GeV (585 GeV and 156 GeV) are excluded at 90% CL for $m_\chi$ equal to 1 GeV and 1.3 TeV, respectively. Values of $M_\chi$ below 585 GeV and 100 GeV (794 GeV and 188 GeV) are excluded for the D8 (D9) spin-dependent operator for $m_\chi$ equal to 1 GeV and 1.3 TeV, respectively. These results can be translated into upper limits on the nucleon-WIMP interaction cross section using the prescription in Refs. [12, 39]. Figure 3 shows 90% CL upper limits on the nucleon-WIMP cross section as a function of $m_\chi$. In the case of the D1 (D5) spin-independent interaction, nucleon-WIMP cross sections above $2.7 \times 10^{-39}$ cm$^2$ and $5.8 \times 10^{-34}$ cm$^2$ ($2.2 \times 10^{-39}$ cm$^2$ and $1.7 \times 10^{-36}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ ($2.2 \times 10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ ($2.2 \times 10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ ($2.2 \times 10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ ($2.2 \times 10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ ($2.2 \times 10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ ($2.2 \times 10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ ($2.2 \times 10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ ($2.2 \times 10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ ($2.2 \times 10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively. Spin-dependent interactions cross sections in the range $7.6 \times 10^{-41}$ cm$^2$ to $3.4 \times 10^{-37}$ cm$^2$ ($2.2 \times 10^{-41}$ cm$^2$ to $2.7 \times 10^{-38}$ cm$^2$) are excluded at 90% CL for $m_\chi = 1$ GeV and $m_\chi = 1.3$ TeV, respectively.

In summary, we report results on the search for new phenomena in events with an energetic photon and large missing transverse momentum in proton-proton collisions at $\sqrt{s} = 7$ TeV at the LHC, based on ATLAS data corresponding to an integrated luminosity of 4.6 fb$^{-1}$. The measurements are in agreement with the SM predic-
The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA) and in the Tier-2 facilities worldwide.

INFIN Sezione di Bologna; (b) Dipartimento di Fisica, Università di Bologna, Bologna, Italy

Department of Physics, Boston University, Boston MA, United States of America

Department of Physics, Brandeis University, Waltham MA, United States of America

Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; (b) Federal University of Juiz de Fora (UFJF), Juiz de Fora; (c) Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei; (d) Instituto de Física, Universidade de Sao Paulo, Sao Paulo, Brazil

Physics Department, Brookhaven National Laboratory, Upton NY, United States of America

National Institute of Physics and Nuclear Engineering, Bucharest; (b) University Politehnica Bucharest, Bucharest; (c) West University in Timisoara, Timisoara, Romania

Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina

Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom

Department of Physics, Carleton University, Ottawa ON, Canada

CERN, Geneva, Switzerland

Enrico Fermi Institute, University of Chicago, Chicago IL, United States of America

(a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; (b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile

(a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (b) Department of Modern Physics, University of Science and Technology of China, Anhui; (c) Department of Physics, Nanjing University, Jiangsu; (d) School of Physics, Shandong University, Shandong, China

Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France

Nevis Laboratory, Columbia University, Irvington NY, United States of America

Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark

(a) INFN Gruppo Collegato di Cosenza; (b) Dipartimento di Fisica, Università della Calabria, Arcavata di Rende, Italy

AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland

The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland

Physics Department, Southern Methodist University, Dallas TX, United States of America

Physics Department, University of Texas at Dallas, Richardson TX, United States of America

DESY, Hamburg and Zeuthen, Germany

Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany

Institut für Kern- und Teilchenphysik, Technical University Dresden, Dresden, Germany

Department of Physics, Duke University, Durham NC, United States of America

SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom

INFN Laboratori Nazionali di Frascati, Frascati, Italy

Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany

Section de Physique, Université de Genève, Geneva, Switzerland

(a) INFN Sezione di Genova; (b) Dipartimento di Fisica, Università di Genova, Genova, Italy

(a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi; (b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia

II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany

SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom

II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany

Laboratoire de Physique Subatomique et de Cosmologie, Université Joseph Fourier and CNRS/IN2P3 and Institut National Polytechnique de Grenoble, Grenoble, France

Department of Physics, Hampton University, Hampton VA, United States of America

Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge MA, United States of America

(a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (c) ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany

Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan

Department of Physics, Indiana University, Bloomington IN, United States of America

Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria

University of Iowa, Iowa City IA, United States of America

Department of Physics and Astronomy, Iowa State University, Ames IA, United States of America
Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
Graduate School of Science, Kobe University, Kobe, Japan
Faculty of Science, Kyoto University, Kyoto, Japan
Kyoto University of Education, Kyoto, Japan
Department of Physics, Kyushu University, Fukuoka, Japan
Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
Physics Department, Lancaster University, Lancaster, United Kingdom
Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia
School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
Department of Physics, Royal Holloway University of London, Surrey, United Kingdom
Department of Physics and Astronomy, University College London, London, United Kingdom
Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
Fysiska institutionen, Lunds universitet, Lund, Sweden
Departmentamento de Física Teórica C-15, Universidad Autonoma de Madrid, Madrid, Spain
Institut für Physik, Universität Mainz, Mainz, Germany
School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
Department of Physics, University of Massachusetts, Amherst MA, United States of America
Department of Physics, McGill University, Montreal QC, Canada
School of Physics, University of Melbourne, Victoria, Australia
Department of Physics, The University of Michigan, Ann Arbor MI, United States of America
Department of Physics and Astronomy, Michigan State University, East Lansing MI, United States of America
(a)INFN Sezione di Milano; (b)Dipartimento di Fisica, Università di Milano, Milano, Italy
B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus
National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus
Department of Physics, Massachusetts Institute of Technology, Cambridge MA, United States of America
Group of Particle Physics, University of Montreal, Montreal QC, Canada
P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia
Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
Moscow Engineering and Physics Institute (MEPhI), Moscow, Russia
Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia
Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
Nagasaki Institute of Applied Science, Nagasaki, Japan
Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
(a)INFN Sezione di Napoli; (b)Dipartimento di Scienze Fisiche, Università di Napoli, Napoli, Italy
Department of Physics and Astronomy, University of New Mexico, Albuquerque NM, United States of America
Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
Department of Physics, Northern Illinois University, DeKalb IL, United States of America
Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
Department of Physics, New York University, New York NY, United States of America
Ohio State University, Columbus OH, United States of America
Faculty of Science, Okayama University, Okayama, Japan
Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK, United States of America
Department of Physics, Oklahoma State University, Stillwater OK, United States of America
Palacký University, RCP TM, Olomouc, Czech Republic
Center for High Energy Physics, University of Oregon, Eugene OR, United States of America
LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
Graduate School of Science, Osaka University, Osaka, Japan
$^{ah}$ Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom

$^{ai}$ Also at Department of Physics, Oxford University, Oxford, United Kingdom

$^{aj}$ Also at Institute of Physics, Academia Sinica, Taipei, Taiwan

$^{ak}$ Also at Department of Physics, The University of Michigan, Ann Arbor MI, United States of America

$^{al}$ Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa

* Deceased