A search for the optical and near–infrared counterpart of
the accreting millisecond X–ray pulsar XTE J1751–305

P.G. Jonker1⋆, G. Nelemans1, Z. Wang2, A.K.H. Kong3, D. Chakrabarty2, M. Garcia3,
P.J. Groot4, M. van der Klis5, T. Kerr9, B. Mobasher10, M. Sullivan11, T. Augusteijn6,
B.W. Stappers5,12, P. Challis3, R.P. Kirshner3, J. Hjorth7, A. Delsanti8

1 Institute of Astronomy, Madingley Road, CB3 0HA, Cambridge, UK
2 Department of Physics and Center for Space Research, Massachusetts Institute of Technology, Cambridge, MA 02139, USA
3 Harvard–Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138, USA
4 Department of Astrophysics, University of Nijmegen, P.O.Box 9010, Nijmegen, The Netherlands
5 Astronomical Institute “Anton Pannekoek”, University of Amsterdam, Kruislaan 403, 1098 SJ Amsterdam, The Netherlands
6 Nordic Optical Telescope, Apartado 474, 38700 Santa Cruz de La Palma, Canary Islands, Spain
7 Astronomical Observatory, University of Copenhagen, Juliane Maries Vej 30, DK-2100 Copenhagen Æ, Denmark
8 European Southern Observatory, Alonso de Cordova 3107, Vitacura, Casilla 19001, Santiago 19, CHILE
9 Joint Astronomy Centre, University Park, 660 N. A’ohoku Place, Hilo, Hawaii 96720, USA
10 Space Telescope Science Institute, 3700 San Martin Drive, Baltimore MD 21218, USA
11 Physics Department University of Durham, Science Labs, South Road, Durham, DH1 3LE, UK
12 Dwingeloo, Postbus 2, 7990 AA, Dwingeloo, The Netherlands

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ABSTRACT

We have obtained optical and near–infrared images of the field of the accreting millisecond X–ray pulsar XTE J1751–305. There are no stars in the 0.7" error circle (0.7" is the overall uncertainty arising from tying the optical and X–ray images and from the intrinsic uncertainty in the Chandra X–ray astrometric solution). We derive limiting magnitudes for the counterpart of R > 23.1, I > 21.6, Z > 20.6, J > 19.6, K > 19.2. We compare these upper limits with the magnitudes one would expect for simple models for the possible donor stars and the accretion disk subject to the reddening observed in X–rays for XTE J1751–305 and when put at the distance of the Galactic Centre (8.5 kpc). We conclude that our non–detection does not constrain any of the models for the accretion disk or possible donor stars. Deep, near–infrared images obtained during quiescence will, however, constrain possible models for the donor stars in this ultra–compact system.

Key words: stars: individual (XTE J1751–305) — stars: neutron — X-rays: stars

1 INTRODUCTION

Low–mass X–ray binaries (LMXBs) are typically old (> 10⁸ yr) binary systems in which a low–mass companion star (≤ 1 M⊙) transfers matter to a neutron star or a black hole. The neutron star LMXBs are thought to be among the predecessors of the millisecond radio pulsars. Due to accretion of matter and decay of the magnetic field during the LMXB phase, the neutron star spins–up to millisecond periods (see Bhattacharya 1995 for a review). However, until March 2002 only six LMXBs were known to show pulsations (see Jonker & van der Klis 2001 for a short overview), and only one of them, a transient system, was shown to have a millisecond period (SAX J1808.4–3658; Wijnands & van der Klis 1998 and Chakrabarty & Morgan 1998). In 2002 Markwardt & Swank (2002) (see also Markwardt et al. 2002) discovered millisecond pulsations in two other transient sources, XTE J1751–305 and XTE J0929–314, in outburst.

There are two main reasons to search for the companion stars of accretion powered millisecond X–ray pulsars. Accretion powered millisecond X–ray pulsars have most likely accreted a substantial amount of matter as they were spun–up by accretion to a period of milliseconds. A spectroscopical determination of the radial velocity curve of the companion star yields a lower limit to the mass...
of the neutron star; measuring a mass of considerably more than 1.4 \( M_\odot \) for even one neutron star would imply the firm rejection of many proposed equations of state (see the discussion by van Paradijs & McClintock 1995 for an overview of the mass determinations of millisecond radio pulsars in neutron star–white dwarf binaries see Thorsett & Chakrabarty 1999). In order to spectroscopically measure the radial velocity of the companion star, clearly a detection of the object is first required. Secondly, at present there are three known ways to produce compact X–ray binaries (Savonije, de Kool & van den Heuvel 1986; Tutukov & Fedorova 1989; Yungelson, Nelemans & van den Heuvel 2002; Podsiadlowski, Rappaport & Pfahl 2002). Detection of the companion star will allow us to differentiate between these different formation scenarios.

In this Paper we present the results of our search for the optical and near-infrared counterpart of the ultra–compact accreting millisecond X–ray pulsar XTE J1751–305 \((P_{\text{orb}} = 42 \text{ minutes}; \text{Markwardt et al. 2002})\). Early reports on the data presented in this Paper were published by Jonker et al. (2002), Wang et al. (2002), Wang & Chakrabarty (2002) and Kong et al. (2002).

2 OBSERVATIONS AND ANALYSIS

We observed the region around the accurate Chandra X–ray position of the source using the 6.5 m Magellan, the 3.8 m UKIRT, the 3.58 m NTT, and the 1.54 m Danish telescopes. A spectrum of the candidate optical counterpart presented by Kong et al. (2002) was obtained with ESO’s 3.6 m telescope at La Silla using Grism number 12 which covers 601–1032 nm. A log of the observations can be found in Table 1. In case of the near–infrared data the total exposure time is given and the dither pattern, i.e. in case of the UKIRT J band images we took 5 exposures of 60 seconds each, in between each observation the telescope was dithered; this pattern was repeated twice. In the case of the Magellan near–infrared observations we obtained J and Ks images of 60 seconds integration time each. We also obtained 11 J band images of 20 seconds integration each.

The data reduction was performed in IRAF\(^1\). The near–infrared images (J and K band) were sky subtracted (using the sky determined from the dithered images), flatfielded, aligned, and combined to form one image per band per night. The Magellan near–infrared data have been linearised before the reduction. The optical images (R, I, and Z band) were reduced in the standard fashion. Aperture and point spread function fitting photometry were done using the packages APHOT and DAOPHOT in IRAF. The spectrum obtained with ESO’s 3.6 m telescope was extracted using the task APALL in the package SPECRED in IRAF. All nights were photometric except for the nights of April 13 and 14, 2002 when the humidity was very high at Magellan and April 14, 2002 at UKIRT when thin cirrus was present. The J band observations during that night also suffered from a bright sky as they were obtained partially during twilight. Charge overflow from a neighbouring bright star covered the location of XTE J1751–305 on the CCD in our image obtained with the Danish telescope, rendering it impossible to determine whether a source is present or not at the Chandra X–ray location. Hence, we will not consider this observation any further.

We derived an astrometric solution for the optical I band image obtained with the NTT of the field of XTE J1751–305 using the positions of 4 nearby unsaturated stars which appear in the USNO-A1.0 catalogue. The rms of the fit was 0.015”. The typical astrometric error of stars in the USNO–A1.0 catalogue is 0.25” (68 per cent confidence); the uncertainty in the X–ray position is dominated by the Chandra aspect solution (0.6”); 90 per cent confidence; Markwardt et al. 2002). Hence, the overall astrometric uncertainty in the position of the error circle is 0.7” (90 per cent confidence). Next, we tied the astrometric solution of the I band image to the optical R and Z band and to the near–infrared J and K(s) band images by assigning the known position of several stars (from the I band) to the R, Z, J(s), or K band image; this did not increase the error in the astrometry significantly.

In Figure 1 (top panel) we show the R band image from our NTT observations. Separately, in Figure 1 (bottom panel) we show the near–infrared UKIRT (K band) image. From the astrometry it is clear that each of the two candidate counterparts (both indicated with two arrows in the top panel) is just outside the 90 per cent confidence error circle. We derived an astrometric solution for the optical J band image to the optical R and Z band and to the near–infrared J and K(s) band images by assigning the known position of several stars (from the I band) to the R, Z, J(s), or K band image; this did not increase the error in the astrometry significantly.

We determined 3 \( \sigma \) upper limits on the presence of a star in the error circle of each of the images obtained under photometric conditions. To do this we added a simulated artificial star (created using the point spread function of the stars in the image) at the position of the error circle. The magnitude of this extra star was varied and measured using the standard photometric tasks. We define the 3 \( \sigma \) limiting magnitude as the magnitude at which the error in the magnitude of such an artificial star is \( \sim 0.3 \) magnitudes (an 0.3 magnitude error is equivalent to a \( \sim 30 \) per cent error on the flux measurement, i.e. a three sigma detection; see also Hulleman et al. 2000). For all the images we need \( \sim 3–5 \) trials to cover the error circle (the number of trials is different for the different images; it varies as a function of the seeing). We did not take the number of trials into account when calculating the upper limits (the error introduced by this is \( \sim 0.1 \) magnitude). The error on the determined limiting magnitude is also \( \sim 0.1 \) magnitude. Together with the uncertainty in the photometric zero–point (at most 0.1 magnitude) this yields an uncertainty of \( \sim 0.2 \) magnitude in the limiting magnitudes.

For the NTT Z band image we converted the (Sloan) z’ magnitude of the observed standard star (PG0918+029D) to the UKIRT UFTI Z band using the transformation given on the UKIRT web page. However, since the NTT and the UKIRT Z band differ, and there is no cross-calibration of these two Z bands, in doing so we introduced an uncertainty in the Z band magnitudes which can be more than 0.2 magnitudes. Therefore, the Z band upper limit is given

\(^{1}\) IRAF is distributed by the National Optical Astronomy Observatories
Table 1. Log of the observations. MJD and start time refer to the time at the start of the first observation.

<table>
<thead>
<tr>
<th>Telescope</th>
<th>Instrument</th>
<th>Observation date and start time (UT)</th>
<th>MJD (UTC)</th>
<th>Filters (except J, Ks)</th>
<th>Exposure Time (s)</th>
<th>Airmass</th>
<th>Seeing (arcseconds)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Danish</td>
<td>DFOSC</td>
<td>08-04-2002 08:51</td>
<td>52372.36870</td>
<td>R</td>
<td>900</td>
<td>1.01</td>
<td>~1.2</td>
</tr>
<tr>
<td>Magellan</td>
<td>Classic-CAM</td>
<td>09-04-2002 08:02</td>
<td>52373.33472</td>
<td>Ks</td>
<td>60</td>
<td>1.05</td>
<td>~0.65</td>
</tr>
<tr>
<td>Magellan</td>
<td>Classic-CAM</td>
<td>09-04-2002 07:32</td>
<td>52373.31388</td>
<td>J</td>
<td>280 (11x20+60)</td>
<td>1.05</td>
<td>~0.65</td>
</tr>
<tr>
<td>Magellan</td>
<td>MagIC</td>
<td>13-04-2002 07:43</td>
<td>52377.32153</td>
<td>I</td>
<td>240 (3x)</td>
<td>1.05</td>
<td>~0.6</td>
</tr>
<tr>
<td>Magellan</td>
<td>MagIC</td>
<td>14-04-2002 08:30</td>
<td>52378.35416</td>
<td>I</td>
<td>30 (1x), 240 (18x)</td>
<td>1.00-1.03</td>
<td>~0.7</td>
</tr>
<tr>
<td>UKIRT</td>
<td>UFTI</td>
<td>14-04-2002 14:41</td>
<td>52378.61201</td>
<td>J</td>
<td>600 (2x5x60)</td>
<td>1.6</td>
<td>~0.8</td>
</tr>
<tr>
<td>UKIRT</td>
<td>UFTI</td>
<td>14-04-2002 14:55</td>
<td>52378.61611</td>
<td>K</td>
<td>600 (2x5x60)</td>
<td>1.6</td>
<td>~0.65</td>
</tr>
<tr>
<td>UKIRT</td>
<td>UFTI</td>
<td>18-04-2002 14:47</td>
<td>52382.61623</td>
<td>J</td>
<td>600 (2x5x60)</td>
<td>1.6</td>
<td>~0.65</td>
</tr>
<tr>
<td>UKIRT</td>
<td>UFTI</td>
<td>18-04-2002 15:06</td>
<td>52382.62925</td>
<td>K</td>
<td>600 (2x5x60)</td>
<td>1.6</td>
<td>~0.5</td>
</tr>
<tr>
<td>NTT</td>
<td>SUSI2</td>
<td>18-04-2002 08:33</td>
<td>52382.35682</td>
<td>R</td>
<td>600</td>
<td>1.00</td>
<td>~0.8</td>
</tr>
<tr>
<td>NTT</td>
<td>SUSI2</td>
<td>18-04-2002 08:23</td>
<td>52382.34950</td>
<td>I</td>
<td>600</td>
<td>1.00</td>
<td>~0.9</td>
</tr>
<tr>
<td>NTT</td>
<td>SUSI2</td>
<td>18-04-2002 08:12</td>
<td>52382.34185</td>
<td>Z</td>
<td>600</td>
<td>1.01</td>
<td>~0.8</td>
</tr>
<tr>
<td>3.6M (spec)</td>
<td>EFOSC</td>
<td>01-05-2002 07:19</td>
<td>52395.30470</td>
<td>Gr#12 0.7&quot; slit</td>
<td>1500</td>
<td>1.01</td>
<td>~0.9</td>
</tr>
</tbody>
</table>

Table 2. Upper limits (3 σ) on the presence of a star at the position of the Chandra error circle.

<table>
<thead>
<tr>
<th>Telescope</th>
<th>Observation date</th>
<th>Limiting magnitude</th>
</tr>
</thead>
<tbody>
<tr>
<td>Magellan</td>
<td>April, 9, 2002</td>
<td>J &gt;19.6</td>
</tr>
<tr>
<td>Magellan</td>
<td>April, 9, 2002</td>
<td>K &gt;18.2</td>
</tr>
<tr>
<td>UKIRT</td>
<td>April, 18, 2002</td>
<td>J &gt;19.5</td>
</tr>
<tr>
<td>UKIRT</td>
<td>April, 18, 2002</td>
<td>K &gt;19.2</td>
</tr>
<tr>
<td>NTT</td>
<td>April, 18, 2002</td>
<td>R&gt;23.1</td>
</tr>
<tr>
<td>NTT</td>
<td>April, 18, 2002</td>
<td>&gt;21.6</td>
</tr>
<tr>
<td>NTT</td>
<td>April, 18, 2002</td>
<td>Z&gt;20.6</td>
</tr>
</tbody>
</table>

as reference only; it should be considered approximate. The upper limits are given in Table 2.

We searched for variability in the I band magnitude for the candidate counterpart proposed by Wang & Chakrabarty (2002) (star 1) by comparing the magnitudes in the Magellan and the NTT observations. The I band magnitude of star 1 is consistent with being the same during the observations. Unfortunately, the I band magnitude of the candidate counterpart proposed by Kong et al. (2002) (star 2 in Figure 1) could not be determined due to the presence of the nearby star. However, differential photometry showed that star 2 was not variable in the Magellan Sloan i’ band images taken a day apart.

The slit we used to obtain a spectrum using ESO’s 3.6 m telescope at La Silla had a width of 0.7”. The slit orientation was such that both star 1 and 2 were in the slit as well as most of the position marked by the error circle. However, due to the seeing of ~0.9” during these observations the spectrum of star 2 may have been contaminated by light of the bright nearby star. In order to investigate this we converted the magnitudes of this bright nearby star (R=18.3, I=16.2) and those of star 2 (R=22.6, I=19.2; we note that the uncertainty on these magnitudes is large due to the non–photometric conditions at the time of the Magellan I band observations; for comparison the best estimate of the magnitude of star 1 was 20.2) to flux densities and we compared the amplitude of the spectral energy distribution according to these R and I band flux densities with the amplitude of the flux calibrated spectrum. The fact that the flux density of the spectrum is higher than that of star 2 shows that the spectrum will have been contaminated with light from the bright nearby star. We note that in converting the magnitudes to fluxes we assumed that star 2 was not variable. We show the spectrum in Figure 2. We also label the position of Hα, although it is unclear whether the donor star of XTE J1751–305 contains hydrogen or not. The spectrum is featureless except for the atmospheric absorption feature near 7613 Angstrom and the band at ~9300 Angstrom. The spectrum of star 1 was not detected, rendering further support to the conclusion that the detected light from the position of star 2 was dominated by the bright nearby star.

3 DISCUSSION

We have obtained optical and near–infrared images of the field of the accreting millisecond X–ray pulsar XTE J1751–305. Two possible counterparts have been investigated but since they both fall outside the 90 per cent confidence Chandra error circle we conclude that neither the optical nor the near–infrared counterpart was detected. The star closest to the error circle was suggested to be the counterpart by Kong et al. (2002) (star 2 in Figure 1). We placed upper limits on the presence of a star in the error circle in the R, I, Z, J, and K band.

To evaluate the constraints these upper limits can place on the properties of the system, we consider the possible evolutionary states for XTE J1751–305. To arrive at an approximately forty minute orbital period X–ray bi-
Figure 1. Top panel: The R band image (∼27′′ × 14′′; 10 minutes integration; North is up, East to the left) of the region of XTE J1751–305 obtained with the 3.58 m NTT. The X-ray error circle for the location of the source is overplotted (0.7′′; 90 per cent confidence). The double set of arrows indicate the location of the possible counterparts proposed by Wang & Chakrabarty (2002; star 1) and Kong et al. (2002; star 2). Bottom panel: The UKIRT K band image obtained on April 18, 2002 (∼27′′ × 14′′; 10 minutes integration; North is up, East is to the left) of the region of XTE J1751–305. The Chandra error circle is overplotted. The vertical black strips in the centre of the image are artifacts introduced when combining the dithered images.

there are essentially three evolutionary paths. The first starts from a detached white dwarf – neutron star binary, which is brought into contact by angular momentum loss due to gravitational wave radiation. For a discussion of this scenario for ultra–compact X-ray binaries see Yungelson, Nelemans & van den Heuvel (2002). The mass of the donor in this case would be ∼0.02 M⊙ and the mass transfer rate ∼10⁻¹¹ M⊙yr⁻¹. Recently, Bildsten (2002) discussed XTE J1751–305 and the similar accreting millisecond X-ray pulsar XTE J0929–314 and calculated models for hot white dwarf donors, rather than the previously used zero–temperature models. He finds a luminosity of the order of 10⁻³ L⊙ and an effective temperature of ∼6000 K for the donors.

It is also possible to form a forty minute binary from a helium star that transfers matter to a neu-
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Figure 2. Spectrum of the Southernmost source identified with the marks in Figure 1 (star 2; to the East of the bright neighbour), obtained with ESO’s 3.6 m telescope at La Silla on May 1, 2002. The line connects the magnitudes of this star in the R and I band converted to flux densities at the central wavelength of the R and I band filter.

The third evolutionary scenario producing ultra–compact binaries involves a main sequence star close to core hydrogen exhaustion that starts mass–transfer to a neutron star [Tutukov, de Kool & van den Heuvel 1986; Tutukov & Fedorova 1989]. Such a scenario goes through the forty minute period range twice. The first time, the system goes through the orbital period of 40 minutes while the orbital period decreases, then the donor is a luminous, hot helium star of about 0.6 M⊙. The mass transfer rate at that stage is of the order of 10−8 M⊙ yr−1. After having reached a period minimum of about 10 minutes, the system returns to longer periods with a dim, semi–degenerate helium star donor, not unlike the low–mass white dwarfs in the first scenario.

The X–ray observations already place some constraints on the system. The mass–function for the companion (1.3×10−6 M⊙; Markwardt et al. 2002) suggests that the companion is a low–mass star. If the companion were to be a 0.6 M⊙ helium star, the inclination should be less than 1.6°. The two pre period–minimum models of [Podsiadlowski, Rappaport & Pfahl 2002] with donor masses of 0.14 and 0.094 M⊙ would imply inclinations less than 6° and 9°, respectively. Markwardt et al. (2002) infer a mass transfer rate of 2.1×10−11 M⊙ yr−1, assuming a neutron star radius of 10 km, a distance of 10 kpc and a recurrence time of 3.8 yr. Even though there is quite some uncertainty in the distance and recurrence time, the high mass transfer rate as expected from a helium star donor seems to be unlikely, even more so, because at such high rates the system would be expected to be a persistent X–ray source [Tsugawa & Osaki 1997].

Next, we estimate the expected absolute magnitudes of the different donor star and accretion disk models in the optical and near–infrared. Since the source was still in outburst when we obtained our observations it is likely that the accretion disk dominates in the optical and near–infrared.

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We assumed an absolute visual magnitude of the disk of 4.5 based on the observed absolute magnitude of 4U 1916–053 of 5.3, which has a similar orbital period, but allowing for the higher X-ray luminosity of XTE J1751–305 (assuming that the disk luminosity scales with the square root of the X-ray luminosity see van Paradijs & McClintock 1995), and the fact that 4U 1916–053 has a high inclination reducing its absolute visual magnitude. Since these ultra-compact systems have to have small accretion disks, they are expected to be hot (see van Paradijs & McClintock 1995). We estimated the absolute magnitudes for disks modelling them as simple blackbodies of 10000 and 30000 K. The results are shown in Figure 3 (dashed lines in the top panel).

The absolute magnitudes of the possible donor stars in a forty minute binary, again assuming simple blackbody spectral energy distributions, are calculated for the hot white dwarf scenario proposed by Bildsten (2002) and for the two pre period–minimum models of Podsiadlowski, Rappaport & Pfahl (2002). We did not consider their post–period minimum model because according to Bildsten (2002) the companions in X–ray binaries will never cool down so much. Also, such a donor star would be very similar to the hot white dwarf donor model.

For the hot helium star donor, we used a model having a luminosity of 100 L⊙ and T = 60000 K (see Tutukov & Fedorova 1989). Interestingly, the absolute magnitudes of this model fall on top of the 30000 K disc model. The absolute magnitudes of all four models are shown in Figure 3 (top panel).

To compare the theoretical models with our upper limits, we assumed XTE J1751–305 is near the Galactic Centre, with a distance modulus of 14.65 magnitudes. Furthermore, we inferred an absorption (A_V = 5.6), based on the measured N_H = 10^{22} cm^{-2} (Miller et al. 2002) and calculated the absorption in the other bands according to the relations found by Rieke & Lebofsky (1985). The results are shown in Figure 3 (bottom panel) where the single symbols denote the upper limits. Unfortunately, our upper limits do not constrain any of the models. Further deep near–infrared imaging could, however, start to rule out several possible models for the counterpart of the accreting millisecond X–ray pulsar XTE J1751–305.

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