Search for supersymmetry in events with three leptons and missing transverse momentum in √s = 7 TeV pp collisions with the ATLAS detector

The ATLAS Collaboration

Abstract

A search for the weak production of charginos and neutralinos decaying to a final state with three leptons (electrons or muons) and missing transverse momentum is presented. The analysis uses 2.06 fb⁻¹ of √s = 7 TeV proton-proton collision data delivered by the Large Hadron Collider and recorded with the ATLAS detector. Observations are consistent with standard model expectations in two signal regions that are either depleted or enriched in Z-boson decays. Upper limits at 95% confidence level are set in R-parity conserving phenomenological minimal supersymmetric and simplified models. For the simplified models, degenerate lightest chargino and next-to-lightest neutralino masses up to 300 GeV are excluded for mass differences from the lightest neutralino up to 300 GeV.
Search for supersymmetry in events with three leptons and missing transverse momentum in $\sqrt{s} = 7$ TeV $pp$ collisions with the ATLAS detector

The ATLAS Collaboration

A search for the weak production of charginos and neutralinos decaying to a final state with three leptons (electrons or muons) and missing transverse momentum is presented. The analysis uses $2.06 \text{fb}^{-1}$ of $\sqrt{s} = 7$ TeV proton-proton collision data delivered by the Large Hadron Collider and recorded with the ATLAS detector. Observations are consistent with standard model expectations in two signal regions that are either depleted or enriched in $Z$-boson decays. Upper limits at 95% confidence level are set in R-parity conserving phenomenological minimal supersymmetric and simplified models. For the simplified models, degenerate lightest chargino and next-to-lightest neutralino masses up to 300 GeV are excluded for mass differences from the lightest neutralino up to 300 GeV.

PACS numbers: 14.80.Ly, 14.80.Nb, 12.60.Jv

Supersymmetric (SUSY) extensions of the standard model (SM) naturally address the gauge hierarchy problem by postulating the existence of SUSY particles, or “sparticles”, with spin differing by one-half unit with respect to that of their SM partner. If R-parity is conserved, sparticles can only be pair-produced and will eventually decay into SM particles and the lightest SUSY particle (LSP) which is stable. Charginos ($\tilde{\chi}_i^\pm$, $i = 1, 2$) and neutralinos ($\tilde{\chi}_j^0$, $j = 1, 2, 3, 4$) are the mass eigenstates formed from the linear superposition of the SUSY partners of the Higgs and electroweak gauge bosons. These are the higgsinos, and the winos, zino, and bino, collectively known as gauginos. Naturalness requires $\tilde{\chi}_1^0$ and $\tilde{\chi}_2^0$ (and third generation sparticles) to have masses in the hundreds of GeV range. In scenarios where first and second generation sfermion masses are larger than few TeVs, the direct production of weak gauginos may be the dominant SUSY process at the Large Hadron Collider (LHC).

Leptonic decays of charginos include sneutrinos ($\tilde{\nu}_\ell \nu$), sleptons ($\tilde{\ell}_\ell \nu$) or $W$ bosons ($W^{\pm(0)} \tilde{\chi}_i^0$), while those of unstable neutralinos include sleptons ($\tilde{\ell}_\ell \nu$) or $Z$ bosons ($Z^{(0)} \tilde{\chi}_i^0$). When both gauginos decay leptonically, a distinctive signature with three leptons and significant missing transverse momentum can be observed, the latter originating from the two undetected LSPs and the neutrinos.

This letter presents the first search with the ATLAS detector for the weak production of charginos and neutralinos decaying to a final state with three leptons (electrons or muons) and missing transverse momentum. The analysis is based on $2.06 \text{fb}^{-1}$ of proton-proton collision data delivered by the LHC at a center-of-mass energy $\sqrt{s} = 7$ TeV between March and August 2011. The search significantly extends the current mass limits on charginos and neutralinos and yields sensitivity in the mass region preferred by naturalness.

In this analysis, observations are interpreted in the phenomenological minimal supersymmetric SM (pMSSM and in simplified models. In the pMSSM the masses of the $\tilde{\chi}_i^\pm$ and $\tilde{\chi}_0^0$ depend on the gaugino masses $M_1$ and $M_2$, the Higgs mass parameter $\mu$, and $\tan \beta$, the ratio of the expectation values of the two Higgs doublets. The masses of the gluinos, squarks and left-handed sleptons are chosen to be larger than 1 TeV, while the right-handed sleptons (including third-generation ones) are assumed to be degenerate with $m_{\tilde{\ell}_R} = (m_{\tilde{\chi}_1^0} + m_{\tilde{\chi}_2^0})/2$. In these scenarios, decays to sleptons are favored. In the simplified models, the masses of the relevant particles (mass degenerate wino-like $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$, bino-like $\tilde{\chi}_3^0$, $\tilde{\chi}_4^0$; $\tilde{\ell}_L$) are the only free parameters of the theory. The $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$ are produced via the $s$-channel exchange of a virtual gauge boson and decay via left-handed sleptons, including staus, and sneutrinos of mass $m_{\tilde{\chi}_0} = m_{\tilde{\ell}_R} = (m_{\tilde{\chi}_1^0} + m_{\tilde{\chi}_2^0})/2$ with a branching ratio of 50% each.

ATLAS is a multipurpose particle detector with forward-backward symmetric cylindrical geometry. It includes an inner tracker (ID) immersed in a 2 T magnetic field providing precision tracking of charged particles and the active media provide energy measurements over the pseudorapidities $|\eta| < 4.9$. The muon detectors outside the calorimeters are contained in a toroidal magnetic field produced by air-core superconducting magnets with field integrals varying from 1 to 8 T·m. They provide trigger and high-precision tracking capabilities for $|\eta| < 2.4$ and $|\eta| < 2.7$, respectively. Electrons must satisfy tight identification criteria and fulfill $|\eta| < 2.47$ and $E_T > 10$ GeV, where $|\eta|$ and $E_T$ are determined from the calibrated clustered energy deposits in the electromagnetic calorimeter matched to an ID track. Muons are reconstructed by combining tracks in the ID and tracks in the muon spectrometer. Reconstructed muons are considered as candidates if they have transverse momentum $p_T > 10$ GeV, $|\eta| < 2.4$, and transverse impact parameter with respect to the primary vertex $|d_0| < 0.2$ mm. “Tagged” leptons are electrons and...
muons, well separated from each other and from candidate jets. “Signal” leptons are tagged leptons for which the scalar sum of the tracks’ transverse momenta within \( \Delta R \equiv \sqrt{(\Delta \phi)^2 + (\Delta \eta)^2} < 0.2 \) around the lepton candidate is less than 10% of the \( E_T \) for electrons, and less than 1.8 GeV for muons. Jets are reconstructed from clustered energy deposits calibrated at the electromagnetic scale using the anti-\( k_T \) algorithm \( 28 \) with a radius parameter of 0.4. The jet energy is corrected to account for the non-compensating nature of the calorimeter using correction factors obtained from Monte Carlo (MC) simulation and parameterized as a function of the jet \( E_T \) and \( \eta \). \( 29 \). Jets considered in this analysis have \( E_T > 20 \) GeV and \( |\eta| < 2.8 \). Jets are identified as containing a \( b \)-quark, and thus called “\( b \)-tagged”, using a multivariate technique based on quantities such as the impact parameter of the tracks associated to the secondary vertex, tracks in jet and other jet shape information, consistent with the expected topology of \( b \)-quark decays. The \( b \)-tagging algorithm \( 30 \) correctly identifies \( b \)-quark jets in top decays with an efficiency of 60% and misidentifies jets containing light flavor quarks and gluons with a rate of < 1%, for jets with \( |\eta| < 2.5 \) and jet \( E_T > 20 \) GeV. The missing transverse momentum, \( E_T^{\text{miss}} \), is the magnitude of the vector sum of the transverse momentum or transverse energy of all \( p_T > 10 \) GeV muons, \( E_T > 10 \) GeV electrons, \( E_T > 20 \) GeV jets, and calibrated calorimeter clusters with \( |\eta| < 4.5 \) not associated to these objects. \( 31 \).

Several MC generators are used to simulate SM processes and SUSY signals relevant for this analysis. \HERWIG (32) is used to simulate diboson processes \((WW^{(*)}, ZZ^{(*)}, WZ^{(*)})\), while \MadGraph (33) is used for the \( t\bar{t}W^{(*)}/Z^{(*)} \) processes. \MCFM (34) is chosen for the simulation of single and pair production of top quarks, while \ALPGEN (35) is used to simulate \( W^{(*)}/Z^{(*)}+\)jets. Expected diboson yields are normalized using next-to-leading order (NLO) QCD predictions obtained with \MCFM (36, 37). The top-quark pair-production contribution is normalized to approximate next-to-next-to-leading (NNLO) order calculations \( 38 \) and the \( t\bar{t}W^{(*)}/Z^{(*)} \) contributions are normalized to NLO results \( 39 \). The QCD NNLO \FEWZ (40, 41) and \MCFM cross-sections are used for NNLO normalization of the \( Z+\)light-flavor jets and \( Z+\)heavy-flavor jets processes, respectively. The choice of the parton distribution functions (PDFs) depends on the generator. \MAST2007 \( 42 \) sets are used for \HERWIG, \CTEQ6L1 (43) with \MadGraph and \ALPGEN, and \CTEQ6.6 (44) with \MCFM. The \pMSSM and simplified model samples are produced with \HERWIG and \HERWIG++ (45), respectively, and the yields are normalized to the NLO cross-sections obtained from \PROSPINO (46) using the PDF set \CTEQ6.6 with the renormalization/factorization scales set to the average of the relevant gaugino masses. Fragmentation and hadronization for the \ALPGEN and \MCFM (\MadGraph) samples are performed with \HERWIG (\PYTHIA (47)). For all MC samples, the propagation of particles through the ATLAS detector is modeled using \GEANT4 (48, 49). The effect of multiple proton-proton collisions from the same or different bunch crossings is incorporated into the simulation by overlaying additional minimum bias events onto hard scatter events using \JIMMY (50). Simulated events are weighted to match the distribution of the mean number of interactions per bunch crossing observed in data.

The data sample was collected with a single-muon trigger \((p_T > 18 \text{ GeV})\) or a single-electron trigger \((E_T > 20 \text{ or } 22 \text{ GeV})\), depending on the instantaneous luminosity. At least one signal lepton is required to have triggered the event and have \( p_T (E_T) \) above 20 GeV \( (25 \text{ GeV})\) for muons \( (\text{electrons})\). Events recorded during normal running conditions are analyzed if at least one of the reconstructed primary vertices has more than four tracks associated to it. Events containing jets with \( |\eta| < 4.9 \) and failing the quality criteria described in Ref. \( 29 \) are rejected to suppress both collisional and non-collisional background. Selected events must contain exactly three signal leptons. As leptonic decays of \( \tilde{\chi}_0^0 \) yield same-flavor opposite-sign (SFOS) lepton pairs, the presence of at least one such pair is required. The invariant mass of any SFOS lepton pair must be above 20 GeV to suppress background from low mass resonances and the missing transverse momentum must satisfy \( E_T^{\text{miss}} > 50 \text{ GeV} \).

Two signal regions are then defined: a “\( Z \)-depleted” region (SR1), with no SFOS pairs having invariant mass within 10 GeV of the nominal \( Z \)-boson mass; and a “\( Z \)-enriched” one (SR2), where at least one SFOS pair has an invariant mass within 10 GeV of the \( Z \)-boson mass. Events in SR1 are further required to contain no \( b \)-tagged jets to suppress contributions from \( b \)-jet-rich backgrounds, where a fake lepton could originate from a heavy-flavor decay. The SR1 and SR2 selections target SUSY events with intermediate slepton or on-mass-shell \( Z \)-boson decays, respectively.

Several SM processes contribute to the background in SR1 and SR2. A background process is considered “irreducible” if it leads to events with three real and isolated leptons, referred to as “real” leptons below. These include diboson \((WZ^{(*)} \text{ and } ZZ^{(*)})\) and \( t\bar{t}W/Z^{(*)} \) production, where the gauge boson may be produced off-mass-shell. Their contribution is determined using the corresponding MC samples, for which lepton and jet selection efficiencies are corrected to account for differences with respect to data. A “reducible” process has at least one “fake” object, that is either a lepton from a semileptonic decay of a heavy-flavor quark or an electron from an isolated photon conversion. The contribution from misidentified light-flavor quarks is negligible. The reducible background includes single- and pair-production of top-quark and \( WW^{(*)} \text{ or } W^{(*)}/Z^{(*)} \) produced in association with jets or photons. The dominant component is the production of top quarks, with a contribution of 1%
or less from \(Z^{(*)}+\text{jets}\) production. The reducible background is estimated using a “matrix method” similar to that described in Ref. \[51\].

In this implementation of the matrix method, the signal lepton with the highest \(p_T\) or \(E_T\) is taken to be real, which is a valid assumption in 99% of the cases, based on MC studies. The number of observed events with one or two fakes is then extracted from a system of linear equations linking the number of events with two additional signal or tagged candidates to the number of events with two additional candidates that are either real or fake. The coefficients of the linear equations are functions of the real lepton identification efficiencies and of the fake object misidentification probabilities. The identification efficiency is measured in data using lepton candidates from \(Z \to \ell\ell\) decays.

Misidentification probabilities for each fake type (heavy flavor, conversion) and for each reducible background process are obtained using simulated events with one signal and two tagged leptons. These misidentification probabilities are then corrected using the ratio (fake scale factor) of the misidentification probability in data and that in MC simulation obtained in dedicated control samples. For heavy flavor fakes, the correction factor is measured in \(b\bar{b}\) events, while for conversion fakes it is determined in a sample of photons radiated from a muon in \(Z \to \mu\mu\) decays. A weighted average misidentification probability is then calculated by weighting the corrected type- and process-dependent misidentification probabilities according to the process cross-section.

An additional source of background is due to events with two signal leptons and one virtual photon converting into two muons, one with \(p_T\) above 10 GeV. The contribution from events in which both muons from the virtual photons have \(p_T\) above 10 GeV is negligible due to the requirement on the dilepton pair invariant mass. For events with only one muon above threshold, an upper limit of 0.5 ± 0.5 in SR1 and of 0.7 ± 0.7 in SR2 is obtained from data as follows. The number of observed events with exactly two signal leptons and \(E_T^{\text{miss}} > 50\) GeV is rescaled by the probability that any of the signal leptons could have radiated the converted photon. This probability is measured in events with \(E_T^{\text{miss}} < 50\) GeV as the ratio of number of events with three signal muons with trilepton invariant mass within 10 GeV of the nominal \(Z\) boson mass to the number of events with two signal muons having the dilepton mass in the same mass window.

The matrix method has been tested using MC events and shown to be accurate within 2%. The background predictions have been validated in a region dominated by \(Z^{(*)}+\text{jets}\) production (VR1: three signal leptons, 30 < \(E_T^{\text{miss}}\) < 50 GeV) and in one dominated by top pair-production (VR2: three signal leptons, SFOs lepton pairs vetoed, \(E_T^{\text{miss}} > 50\) GeV). The data and predictions are in agreement within the quoted statistical and systematic uncertainties as shown in Table \[1\].

Several sources of systematic uncertainty are considered in the signal regions. The systematic uncertainties affecting the MC based estimates (irreducible background yield, misidentification probabilities, signal yield) include the theoretical cross section uncertainty due to scale and PDFs, the acceptance uncertainty due to PDFs, jet energy scale, jet energy resolution, lepton energy scale, lepton energy resolution, lepton efficiency, \(b\)-tagging efficiency, event quality selection, and the uncertainty on the luminosity. In SR1, the total uncertainty on the irreducible background is 17%. This includes the uncertainty on the acceptance due to PDFs (14%), that on the theoretical cross section due to scale and PDFs (7%), and that from the limited number of simulated events (10%), while all the remaining uncertainties on the irreducible background in this signal region range between 0.5-5%. The total uncertainty on the reducible background is 29%. This includes an uncertainty on the object misidentification probability of 10-30% from the sources listed above. The uncertainty from the dependence of the misidentification probability on \(E_T^{\text{miss}}\) (0.4–17%) and the uncertainty on the fake scale factors (10–50%) are also included in the total, together with the uncertainty from the limited number of data events with three tagged leptons, of which at least one is a signal lepton. The total uncertainties on the signal cross-section range between 10-15%. These include uncertainties due to the renormalization and factorization scale, \(\alpha_s\), and PDFs. The maximum uncertainty obtained from either the CTEQ6.6 or the MSTW \[52\] PDF set is used. In SR2, the values of systematic uncertainties are similar to those obtained in SR1. The only exceptions are the uncertainty from the limited number of simulated events (4%) and the uncertainty on the reducible background (52%). In all of the above, the value used for the uncertainty on the luminosity is 3.7%.

The numbers of observed events and the prediction for SM backgrounds in SR1 and SR2 are reported in Table \[1\]. The probability that the background fluctuates to the observed number of events or higher is calculated in the frequentist approach and found to be 19% in SR1 and 6% in SR2. The distributions of the \(E_T^{\text{miss}}\) in the two signal regions are presented in Fig. \[1\]. The yield in SR1 for one of the simplified model scenarios \((m_{\chi_1^0}, m_{\chi_2^0}, m_{\chi_1^\pm} = 250, 250, 175, 100\text{ GeV})\) is also shown for illustration purposes.

No significant excess of events is found in either signal region. Upper limits on the visible production cross-section of 9.9 fb in SR1 and 23.8 fb in SR2 are placed at 95% confidence level (CL) with the modified frequentist CLs prescription \[53\]. No corrections for the effects of experimental resolution, acceptance and efficiency are applied. All systematic uncertainties and their correlations are taken into account via nuisance parameters. The corresponding expected limits are 7.1 fb and 14.1 fb, respectively. In SR2, the observed limit on the visible
TABLE I. Expected numbers of events from SM backgrounds (Bkg.) and observed numbers of events in data, for 2.06 fb$^{-1}$, in control regions VR1 and VR2, and in signal regions SR1 and SR2. Both statistical and systematic uncertainties are included.

<table>
<thead>
<tr>
<th>Selection</th>
<th>VR1</th>
<th>VR2</th>
<th>SR1</th>
<th>SR2</th>
</tr>
</thead>
<tbody>
<tr>
<td>$t\bar{t}W^{(<em>)}/Z^{(</em>)}$</td>
<td>1.4±1.1</td>
<td>0.7±0.6</td>
<td>0.4±0.3</td>
<td>2.7±2.1</td>
</tr>
<tr>
<td>$ZZ^{(*)}$</td>
<td>6.7±1.5</td>
<td>0.03±0.04</td>
<td>0.7±0.2</td>
<td>3.4±0.8</td>
</tr>
<tr>
<td>$WZ^{(*)}$</td>
<td>61±11</td>
<td>0.4±0.2</td>
<td>11±2</td>
<td>58±11</td>
</tr>
<tr>
<td>Reducible Bkg.</td>
<td>56±35</td>
<td>14±9</td>
<td>14±4</td>
<td>7.5±3.9</td>
</tr>
<tr>
<td>Total Bkg.</td>
<td>125±37</td>
<td>15±9</td>
<td>26±5</td>
<td>72±12</td>
</tr>
<tr>
<td>Data</td>
<td>122</td>
<td>12</td>
<td>32</td>
<td>95</td>
</tr>
</tbody>
</table>

FIG. 1. $E_{T}^{miss}$ distributions for events in signal regions SR1 (left) and SR2 (right). The error band includes both statistical and systematic uncertainty, while the errors on the data points are statistical only. The SUSY reference point used in SR1 is described in the text.

The cross-section is less stringent than the expected limit because of the upwards fluctuation in the number of events in data with respect to the expected background. SR1 provides better sensitivity in the parameter space considered and the limits are interpreted in simplified models and pMSSM scenarios with $M_{1}=100$ GeV and $\tan\beta=6$ (Fig. 2). The chosen $M_{1}$ value leads to a sizable mass splitting between $\tilde{\chi}_{1}^{\pm}$ and $\tilde{\chi}_{0}^{0}$ and therefore to a large acceptance. The value of $\tan\beta$ does not have a significant impact on $\sigma(pp\to\tilde{\chi}_{i}^{\pm}\tilde{\chi}_{j}^{0}) \times BR(\tilde{\chi}_{i}^{\pm}\tilde{\chi}_{j}^{0}\to\ell\ell\tilde{l}\bar{\nu})$, which varies by ~10% if $\tan\beta$ is raised to 10.

In the simplified models, degenerate $\tilde{\chi}_{1}^{\pm}$ and $\tilde{\chi}_{0}^{0}$ masses up to 300 GeV are excluded for large mass differences from the $\tilde{\chi}_{1}^{0}$. Care has to be taken when interpreting the simplified model limit in the context of a pMSSM scenario, where the mass of the sneutrino is lighter than the mass of the left-handed slepton, leading to higher lepton momenta from chargino decays and to a change in the branching ratios of the $\tilde{\chi}_{1}^{0}$.

FIG. 2. Observed and expected 95% CL limit contours for chargino and neutralino production in the pMSSM (upper) and simplified model (lower) scenarios. For the simplified models, the 95% CL upper limit on the production cross-section is also shown. Interpolation is used to account for the discreteness of the signal grids.

In summary, results from the first ATLAS search for the weak production of charginos and neutralinos decaying to a final state with three leptons (electrons or muons) and missing transverse momentum are reported. The analysis is based on 2.06 fb$^{-1}$ of proton-proton collision data delivered by the LHC at $\sqrt{s}=7$ TeV. No significant excess of events is found in data, where upwards fluctuations of less than 2-sigma are observed. The null result is interpreted in pMSSM and in simplified models. For the simplified models with intermediate sleptons considered in this paper, degenerate lightest chargino and next-to-lightest neutralino masses are excluded up to 300 GeV for mass differences to the lightest neutralino up to 300 GeV.

ACKNOWLEDGEMENTS

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently.

We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWF, Austria;
The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CERN/IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA) and in the Tier-2 facilities worldwide.

The polar angle θ is defined in terms of the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ as η = − ln tan(θ/2).

[26] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z-axis along the beam pipe. The x-axis is from the IP to the centre of the LHC ring, and the y-axis points upward. Cylindrical coordinates (R, φ) are used in the transverse plane, φ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ as η = − ln tan(θ/2).

[27] ATLAS Collaboration, JHEP 12, 60 (2010)
121 Petersburg Nuclear Physics Institute, Gatchina, Russia
122 (a)INFN Sezione di Pisa; (b)Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
123 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA, United States of America
124 (a)Laboratorio de Instrumentacao e Fisica Experimental de Particulas - LIP, Lisboa, Portugal; (b)Departamento de Fisica Teorica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain
125 Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic
126 Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic
127 Czech Technical University in Prague, Praha, Czech Republic
128 State Research Center Institute for High Energy Physics, Protvino, Russia
129 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
130 Physics Department, University of Regina, Regina SK, Canada
131 Ritsumeikan University, Kusatsu, Shiga, Japan
132 (a)INFN Sezione di Roma I; (b)Dipartimento di Fisica, Università La Sapienza, Roma, Italy
133 (a)INFN Sezione di Roma Tor Vergata; (b)Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
134 (a)INFN Sezione di Roma Tre; (b)Dipartimento di Fisica, Università Roma Tre, Roma, Italy
135 (a)Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca; (b)Centre National de l’Energie des Sciences Techniques Nucleaires, Rabat; (c)Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech; (d)Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; (e)Faculty of sciences, Mohammed V-Agdal University, Rabat, Morocco
136 DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat a l’Energie Atomique), Gif-sur-Yvette, France
137 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA, United States of America
138 Department of Physics, University of Washington, Seattle WA, United States of America
139 Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
140 Department of Physics, Shinshu University, Nagano, Japan
141 Fachbereich Physik, Universität Siegen, Siegen, Germany
142 Department of Physics, Simon Fraser University, Burnaby BC, Canada
143 SLAC National Accelerator Laboratory, Stanford CA, United States of America
144 (a)Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava; (b)Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
145 (a)Department of Physics, University of Johannesburg, Johannesburg; (b)School of Physics, University of the Witwatersrand, Johannesburg, South Africa
146 (a)Department of Physics, Stockholm University; (b)The Oskar Klein Centre, Stockholm, Sweden
147 Physics Department, Royal Institute of Technology, Stockholm, Sweden
148 Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook NY, United States of America
149 Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
150 School of Physics, University of Sydney, Sydney, Australia
151 Institute of Physics, Academia Sinica, Taipei, Taiwan
152 Department of Physics, Technion: Israel Inst. of Technology, Haifa, Israel
153 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
154 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
155 International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan
156 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
157 Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
158 Department of Physics, University of Toronto, Toronto ON, Canada
159 (a)TRIUMF, Vancouver BC; (b)Department of Physics and Astronomy, York University, Toronto ON, Canada
160 Institute of Pure and Applied Sciences, University of Tsukuba,1-1-1 Tennodai,Tsukuba, Ibaraki 305-8571, Japan
161 Science and Technology Center, Tufts University, Medford MA, United States of America
162 Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia
163 Department of Physics and Astronomy, University of California Irvine, Irvine CA, United States of America
164 (a)INFN Gruppo Collegato di Udine; (b)ICTP, Trieste; (c)Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy
165 Department of Physics, University of Illinois, Urbana IL, United States of America
Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain

Department of Physics, University of British Columbia, Vancouver BC, Canada

Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada

Department of Physics, University of Warwick, Coventry, United Kingdom

Waseda University, Tokyo, Japan

Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

Department of Physics, University of Wisconsin, Madison WI, United States of America

Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany

Department of Physics, Yale University, New Haven CT, United States of America

Yerevan Physics Institute, Yerevan, Armenia

Domaine scientifique de la Doua, Centre de Calcul CNRS/IN2P3, Villeurbanne Cedex, France

Also at Laboratorio de Instrumentação e Física Experimental de Partículas - LIP, Lisboa, Portugal

Also at Facultade de Ciencias and CFNUL, Universidade de Lisboa, Lisboa, Portugal

Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom

Also at TRIUMF, Vancouver BC, Canada

Also at Department of Physics, California State University, Fresno CA, United States of America

Also at Novosibirsk State University, Novosibirsk, Russia

Also at Fermilab, Batavia IL, United States of America

Also at Department of Physics, University of Coimbra, Coimbra, Portugal

Also at Università di Napoli Parthenope, Napoli, Italy

Also at Institute of Particle Physics (IPP), Canada

Also at Department of Physics, Middle East Technical University, Ankara, Turkey

Also at Louisiana Tech University, Ruston LA, United States of America

Also at Department of Physics and Astronomy, University College London, London, United Kingdom

Also at Group of Particle Physics, University of Montreal, Montreal QC, Canada

Also at Department of Physics, University of Cape Town, Cape Town, South Africa

Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany

Also at Manhattan College, New York NY, United States of America

Also at School of Physics, Shandong University, Shandong, China

Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France

Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China

Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan

Also at Dipartimento di Fisica, Università La Sapienza, Roma, Italy

Also at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Énergie Atomique), Gif-sur-Yvette, France

Also at Section de Physique, Université de Genève, Geneva, Switzerland

Also at Departamento de Física, Universidade de Minho, Braga, Portugal

Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, United States of America

Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary

Also at California Institute of Technology, Pasadena CA, United States of America

Also at Institute of Physics, Jagiellonian University, Krakow, Poland

Also at LAL, Univ. Paris-Sud and CNRS/IN2P3, Orsay, France

Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom

Also at Department of Physics, Oxford University, Oxford, United Kingdom

Also at Institute of Physics, Academia Sinica, Taipei, Taiwan

Also at Department of Physics, The University of Michigan, Ann Arbor MI, United States of America

* Deceased